

# Inhomogeneous phases in the 3+1-dimensional mean-field Nambu-Jona-Lasinio model on the lattice

Laurin Pannullo, Marc Wagner, Marc Winstel

Goethe University Frankfurt

---

HFHF Theory Retreat, 16.09.2022



# Don't investigate

## Inhomogeneous phases in the 3+1-dimensional mean-field Nambu-Jona-Lasinio model on the lattice

Laurin Pannullo, Marc Wagner, Marc Winstel

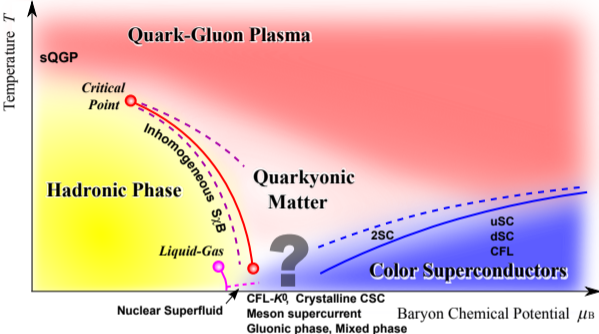
Goethe University Frankfurt

---

HFHF Theory Retreat, 16.09.2022



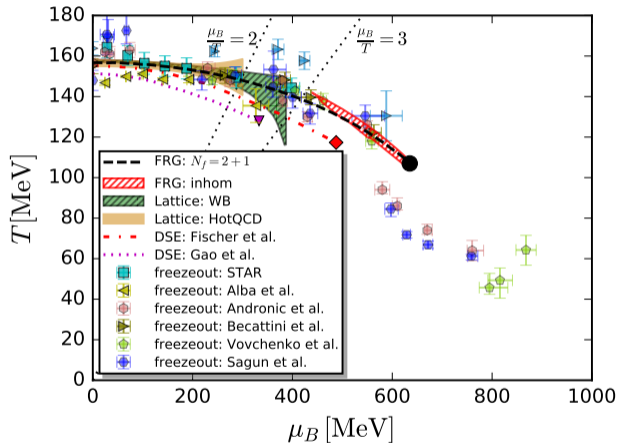
# Motivation I



- QCD phase diagram – a plot full of conjectures
- What goes on at finite or large  $\mu_B$ ?

[ K. Fukushima, T. Hatsuda, *Reports on Prog. Phys.* **74** (2011) ]

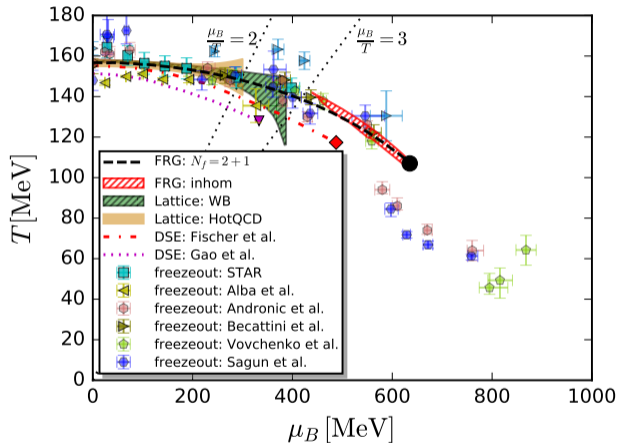
# Motivation I



[ W.-j. Fu *et al.*, *Phys. Rev. D.* **101** (2020) ]

- QCD phase diagram – a plot full of conjectures
- What goes on at finite or large  $\mu_B$ ?
  - Do first principal calculations  
⇒ very hard / impossible

# Motivation I



[ W.-j. Fu *et al.*, *Phys. Rev. D.* **101** (2020) ]

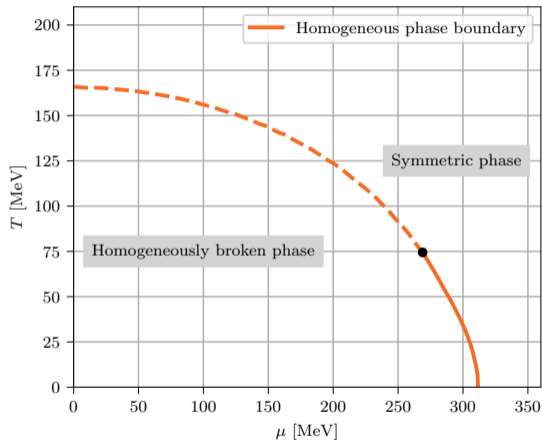
- QCD phase diagram – a plot full of conjectures
- What goes on at finite or large  $\mu_B$ ?
  - Do first principal calculations  
⇒ very hard / impossible
  - Use models of QCD  
⇒ a lot easier; questionable physical relevance of predictions

# Motivation II

- Nambu-Jona-Lasinio model

$$\mathcal{L} = \bar{\psi} (\not{\partial} + \gamma_0 \mu + \sigma + i\gamma_5 \boldsymbol{\tau} \cdot \boldsymbol{\pi}) \psi + \frac{\sigma^2 + \boldsymbol{\pi}^2}{4G}$$

- phase diagram that resembles our QCD expectations



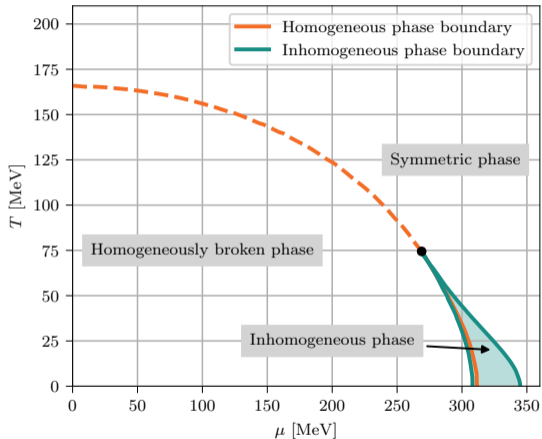
[ D. Nickel, *Phys. Rev. D.* **80** (2009) ]

# Motivation II

- Nambu-Jona-Lasinio model

$$\mathcal{L} = \bar{\psi} (\not{\partial} + \gamma_0 \mu + \sigma + i\gamma_5 \boldsymbol{\tau} \cdot \boldsymbol{\pi}) \psi + \frac{\sigma^2 + \boldsymbol{\pi}^2}{4G}$$

- phase diagram that resembles our QCD expectations
- features an inhomogeneous phase (IP) – a phase with a **space-dependent chiral condensate**



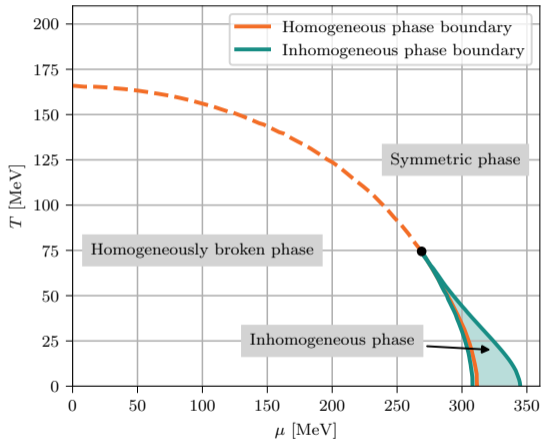
[ D. Nickel, *Phys. Rev. D.* **80** (2009) ]

# Motivation II

- Nambu-Jona-Lasinio model

$$\mathcal{L} = \bar{\psi} (\not{\partial} + \gamma_0 \mu + \sigma + i\gamma_5 \boldsymbol{\tau} \cdot \boldsymbol{\pi}) \psi + \frac{\sigma^2 + \boldsymbol{\pi}^2}{4G}$$

- phase diagram that resembles our QCD expectations
- features an inhomogeneous phase (IP) – a phase with a **space-dependent chiral condensate**
- several problems with this result
  - Mean-field  
⇒ no bosonic quantum fluctuations
  - non-renormalizable model  
⇒ results may depend on the regularization



[ D. Nickel, *Phys. Rev. D.* **80** (2009) ]



# Motivation III

- **Main goal:** Investigate IPs in models for QCD beyond the mean-field approximation
  - via FRG  $\Rightarrow$  see Lennart Kurth's talk yesterday
  - via lattice Monte-Carlo (MC) simulations  $\Rightarrow$  our approach

# Motivation III

- **Main goal:** Investigate IPs in models for QCD beyond the mean-field approximation
  - via FRG  $\Rightarrow$  see Lennart Kurth's talk yesterday
  - via lattice Monte-Carlo (MC) simulations  $\Rightarrow$  our approach
- **But first:**
  - feasibility of lattice MC simulations of these models at finite  $\mu$   
 $\Rightarrow$  there have been MC simulations in the past [ S. Hands, D. N. Walters, *Phys. Rev. D.* **69** (2004) ]

# Motivation III

- **Main goal:** Investigate IPs in models for QCD beyond the mean-field approximation
  - via FRG  $\Rightarrow$  see Lennart Kurth's talk yesterday
  - via lattice Monte-Carlo (MC) simulations  $\Rightarrow$  our approach
- **But first:**
  - feasibility of lattice MC simulations of these models at finite  $\mu$   
 $\Rightarrow$  there have been MC simulations in the past [ S. Hands, D. N. Walters, *Phys. Rev. D.* **69** (2004) ]
  - explore regulator dependence of IPs
  - suitability of lattice regularizations for IPs

# Motivation III

- **Main goal:** Investigate IPs in models for QCD beyond the mean-field approximation
    - via FRG  $\Rightarrow$  see Lennart Kurth's talk yesterday
    - via lattice Monte-Carlo (MC) simulations  $\Rightarrow$  our approach
  - **But first:**
    - feasibility of lattice MC simulations of these models at finite  $\mu$   
 $\Rightarrow$  there have been MC simulations in the past [ S. Hands, D. N. Walters, *Phys. Rev. D.* **69** (2004) ]
    - explore regulator dependence of IPs
    - suitability of lattice regularizations for IPs
- $\Rightarrow$  Simple setup: **Stability analysis** of the 3 + 1-dimensional NJL model in **mean-field**

# Parameter fixing of the NJL model

coupling  $G$

regulator  $\Lambda$ , e.g., Pauli-Villars mass, lattice spacing

$$\bar{\psi} (\not{\partial} + \gamma_0 \mu + \sigma + i\gamma_5 \boldsymbol{\tau} \cdot \boldsymbol{\pi}) \psi + \frac{\sigma^2 + \boldsymbol{\pi}^2}{4G}$$

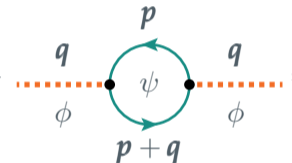
Pion-decay constant  $f_\pi = 88 \text{ MeV}$

Constituent quark mass  $M_0 \sim \langle \bar{\psi} \psi \rangle_{\mu=T=0}$   
in range of 150-500 MeV

[ S. P. Klevansky, *Rev. Mod. Phys.* **64** (1992) ] [ S. Hands, D. N. Walters, *Phys. Rev. D.* **69** (2004) ]

# Stability analysis

- Apply inhomogeneous perturbations  $\delta\tilde{\phi}(\mathbf{q})$  to homogeneous fields
- Curvature of the action in direction  $\delta\tilde{\phi}(\mathbf{q})$  is given by the **bosonic two-point function**  $\Gamma_{\phi}^{(2)}(\mathbf{q})$
- Simple quantity in the mean-field approximation

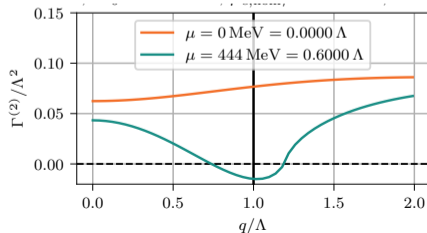
$$\Gamma_{\phi}^{(2)}(\mathbf{q}) = \frac{1}{2G} + \text{diagram} = \frac{1}{2G} + \int \frac{d^4p}{(2\pi)^4} \text{Tr} [X_{\phi} S(p) X_{\phi} S(p + q)]$$


- negative values indicate instability for mode  $q$

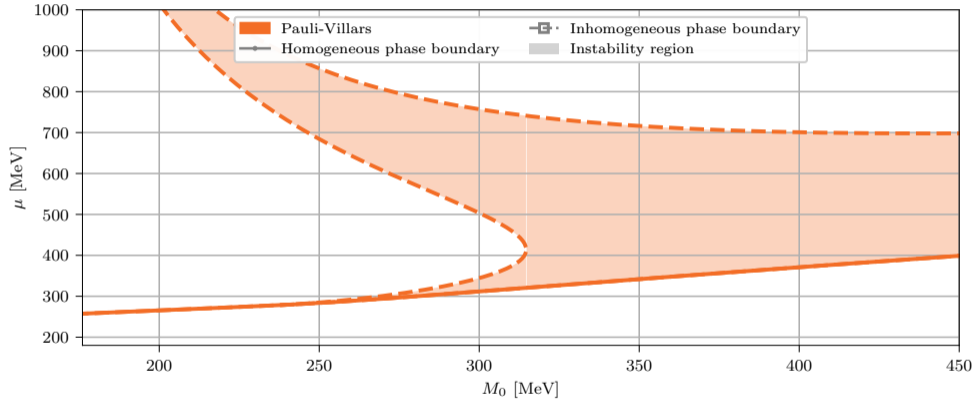
[ M. Buballa *et al.*, *The Eur. Phys. J. Special Top.* **229** (2020) ]

[ A. Koenigstein *et al.* (2021) ]

[ M. Buballa *et al.*, *Phys. Rev. D.* **103** (2021) ]

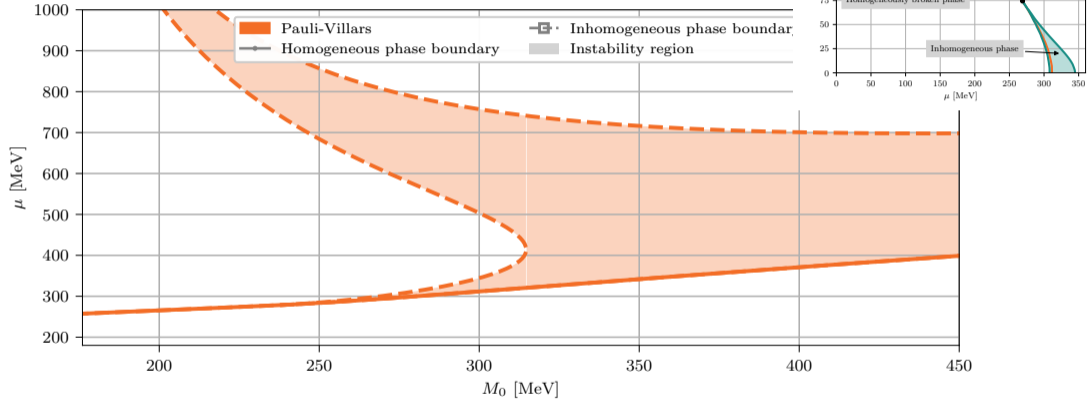


# Quark mass scan



- Most common regularization scheme: **Pauli-Villars** (polynomial suppression of large momenta) [ D. Nickel, *Phys. Rev. D.* **80** (2009) ]

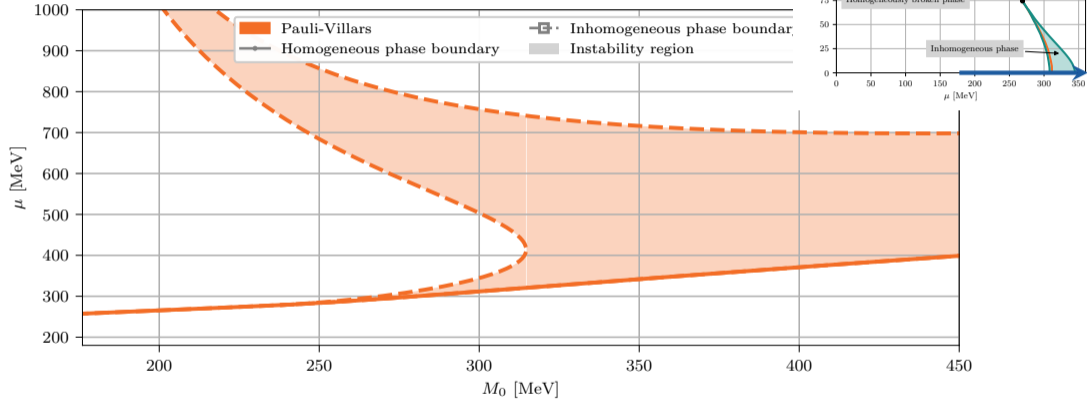
# Quark mass scan



- Most common regularization scheme: **Pauli-Villars** (polynomial suppression of large momenta) [ D. Nickel, *Phys. Rev. D.* **80** (2009) ]

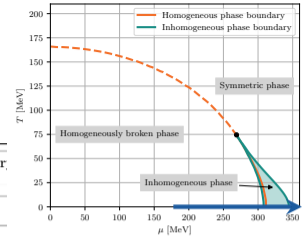
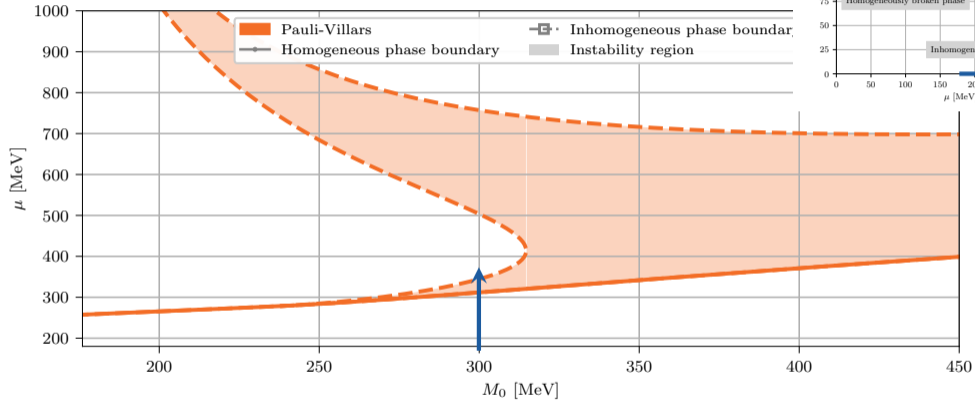


# Quark mass scan



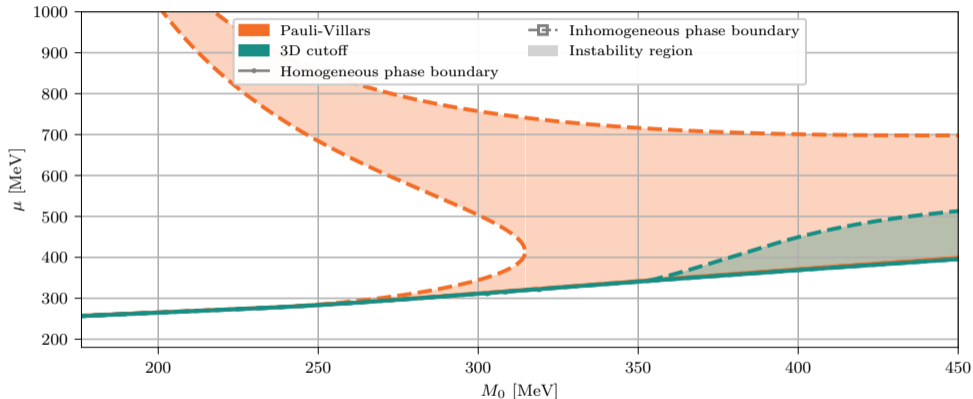
- Most common regularization scheme: **Pauli-Villars** (polynomial suppression of large momenta) [ D. Nickel, *Phys. Rev. D.* **80** (2009) ]

# Quark mass scan



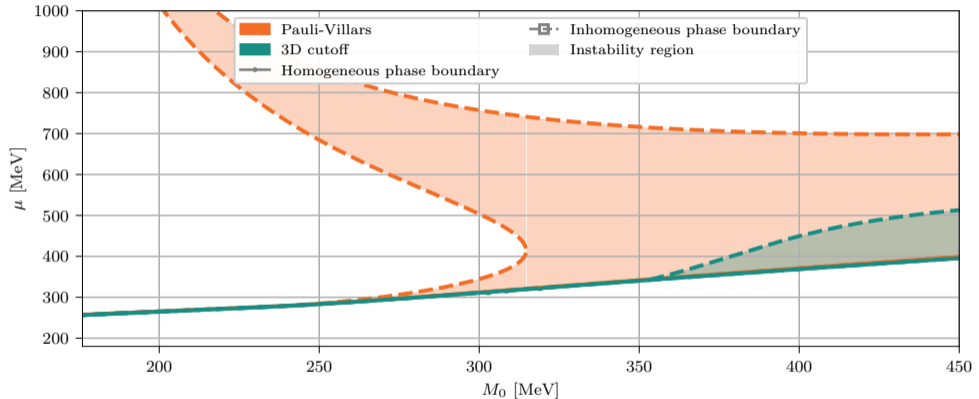
- Most common regularization scheme: **Pauli-Villars** (polynomial suppression of large momenta) [ D. Nickel, *Phys. Rev. D.* **80** (2009) ]

# Quark mass scan



- 3D Cutoff (restriction of spatial loop momenta  $|\mathbf{p}| < \Lambda$ )
- Similar homogeneous phase boundary, but vastly different instability region

# Quark mass scan

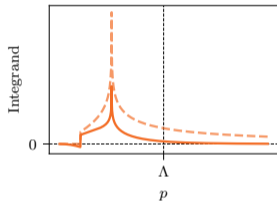


- General problem with cutoff schemes? Explicit breaking of translational invariance?
- Or general problem with results in this non-renormalizable model?

# Lattice regularizations I

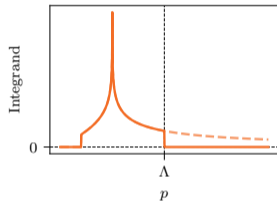
## Pauli-Villars

polynomial suppression



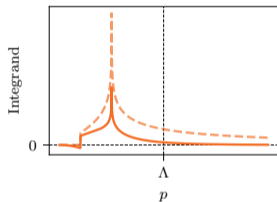
## Momentum Cutoff

hard suppression



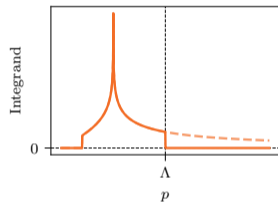
# Lattice regularizations I

Pauli-Villars  
polynomial suppression



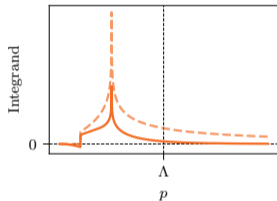
Lattice?

Momentum Cutoff  
hard suppression



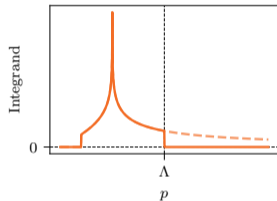
# Lattice regularizations I

Pauli-Villars  
polynomial suppression



Lattice?

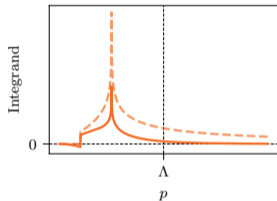
Momentum Cutoff  
hard suppression



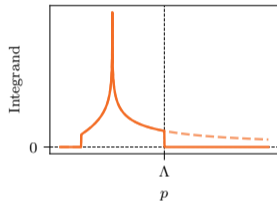
- something like a cutoff  
⇒ maximum momentum  $\pi/a$
- but also periodic dispersion relation  
⇒ can give rise to doublers

# Lattice regularizations I

Pauli-Villars  
polynomial suppression

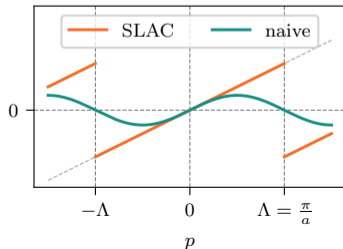


Momentum Cutoff  
hard suppression



Lattice?

- something like a cutoff  
⇒ maximum momentum  $\pi/a$
- but also periodic dispersion relation  
⇒ can give rise to doublers





# Lattice regularizations II

- Wilson fermions – explicit breaking of chiral symmetry at finite  $a$  **X**
- Domain wall fermions – too expensive for now **X**

# Lattice regularizations II

- Wilson fermions – explicit breaking of chiral symmetry at finite  $a$  **X**
- Domain wall fermions – too expensive for now **X**
- SLAC – exact chiral symmetry, but problematic jump at the edge of the Brioullin zone **?**

# Lattice regularizations II

- Wilson fermions – explicit breaking of chiral symmetry at finite  $a$  **X**
- Domain wall fermions – too expensive for now **X**
- SLAC – exact chiral symmetry, but problematic jump at the edge of the Brioullin zone **?**
- naive fermions – doubler mixing, but can be cured in the continuum limit **?**
  - ⇒ problems with temporal doublers **X**
  - ⇒ use Hybrid discretization with SLAC in the temporal direction and naive fermions in the spatial directions **?**

# Lattice regularizations II

- Wilson fermions – explicit breaking of chiral symmetry at finite  $a$  **X**
- Domain wall fermions – too expensive for now **X**
- SLAC – exact chiral symmetry, but problematic jump at the edge of the Brioullin zone **?**
- naive fermions – doubler mixing, but can be cured in the continuum limit **?**
  - ⇒ problems with temporal doublers **X**
  - ⇒ use Hybrid discretization with SLAC in the temporal direction and naive fermions in the spatial directions **?**
- staggered fermions – *less* doubler mixing, but can be cured in the continuum limit; not full chiral symmetry **?**

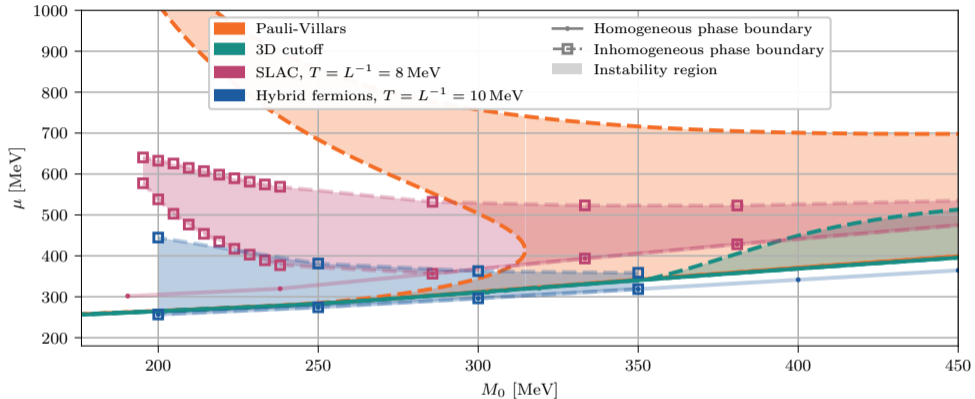
# Lattice regularizations II

- Wilson fermions – explicit breaking of chiral symmetry at finite  $a$  **X**
- Domain wall fermions – too expensive for now **X**
- **SLAC** – exact chiral symmetry, but problematic jump at the edge of the Brioullin zone **?**
- naive fermions – doubler mixing, but can be cured in the continuum limit **?**
  - ⇒ problems with temporal doublers **X**
  - ⇒ use **Hybrid** discretization with SLAC in the temporal direction and naive fermions in the spatial directions **?**
- **staggered** fermions – *less* doubler mixing, but can be cured in the continuum limit; not full chiral symmetry **?**

Finished results with these

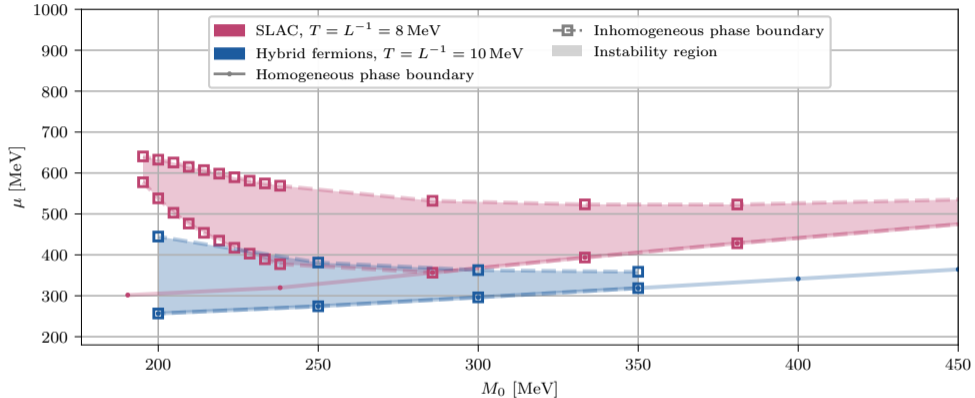
In progress

# Quark mass scan



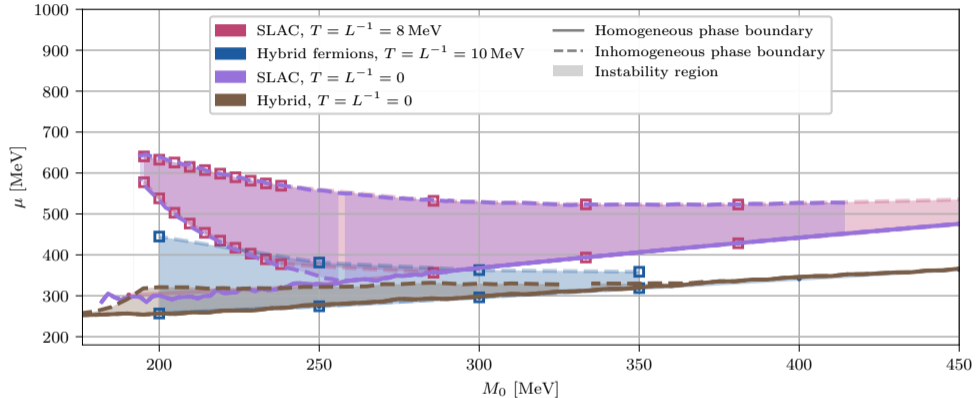
- Even more conflicting results between regularizations

# Quark mass scan



- Lattice results at finite space time volume  
⇒ mean-field stability analysis so simple that we can go on an **infinite lattice**
- SLAC unaffected, but Hybrid much smaller instability

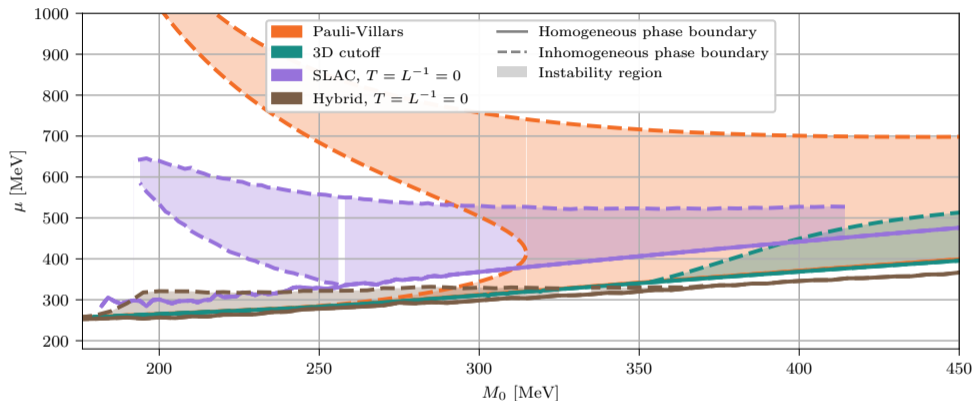
# Quark mass scan



- Lattice results at finite space time volume  
⇒ mean-field stability analysis so simple that we can go on an **infinite lattice**
- SLAC unaffected, but Hybrid much smaller instability

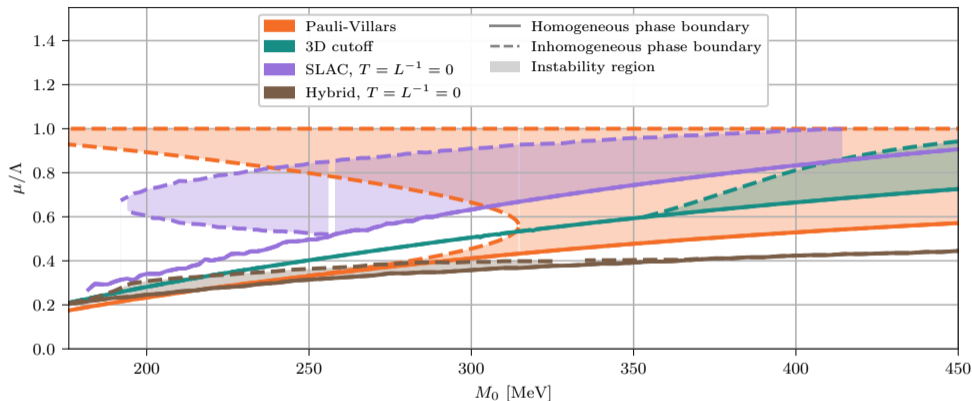


# Quark mass scan



- Lattice results at finite space time volume  
⇒ mean-field stability analysis so simple that we can go on an **infinite lattice**
- SLAC unaffected, but Hybrid much smaller instability

# Quark mass scan

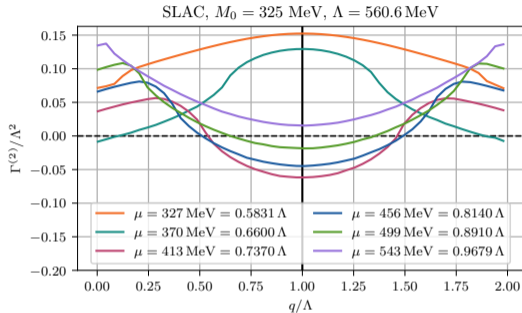
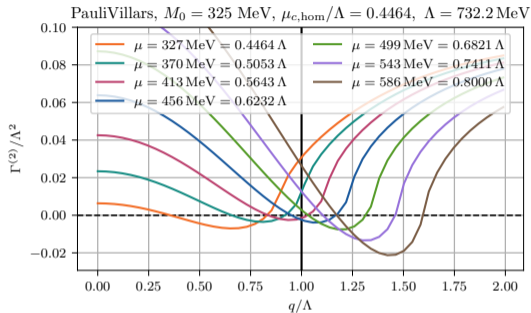


- Plot in cutoff units reveals that chemical potentials are in the order of the cutoff !
- IPs vanish when moving to larger cutoffs

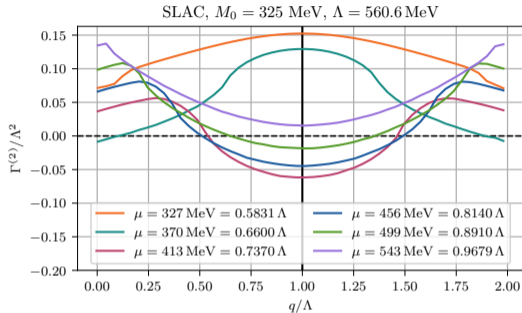
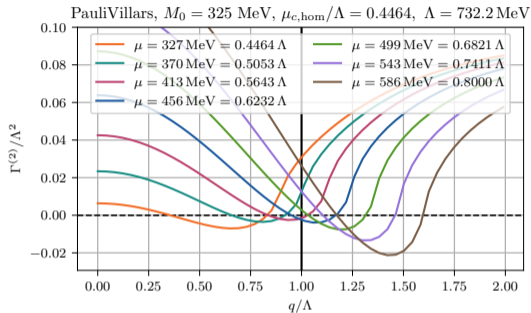
# A closer look at SLAC I

- Shouldn't SLAC be quite similar to the 3D cutoff results? After all SLAC fermions have the continuum dispersion relation...
- Yes, but they also have something like a doubler due to the discontinuity at the edge of the Brillouin zone
- This discontinuity could be probed by the bosonic field !  
⇒ But this is actually not the main problem here.

# A closer look at SLAC II - Two-point functions

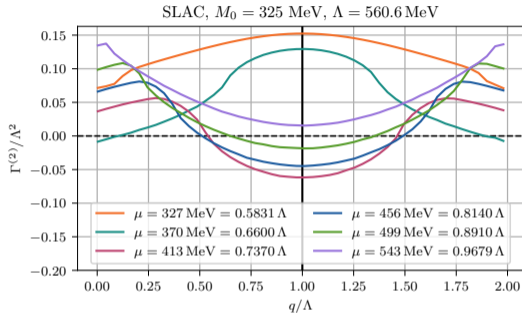
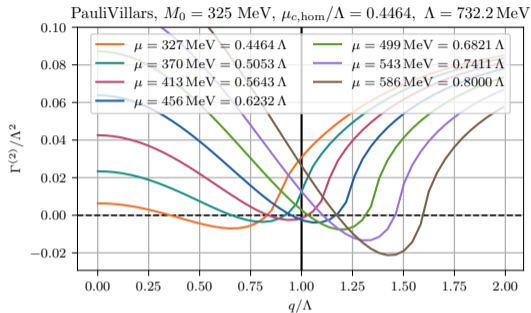


# A closer look at SLAC II - Two-point functions



- SLAC has non-zero slope at  $q = 0$
- caused by the jump at the edge of the Brillouin zone
- minimum runs into maximum  $q$

# A closer look at SLAC II - Two-point functions

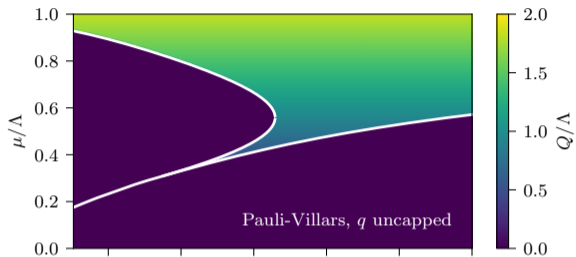


- SLAC has non-zero slope at  $q = 0$
- caused by the jump at the edge of the Brillouin zone
- minimum runs into maximum  $q$

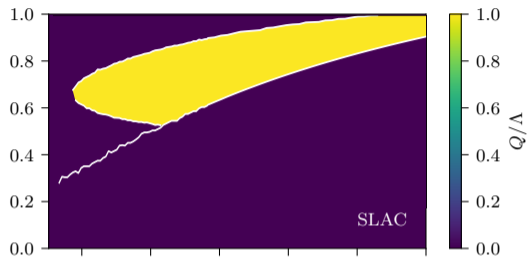
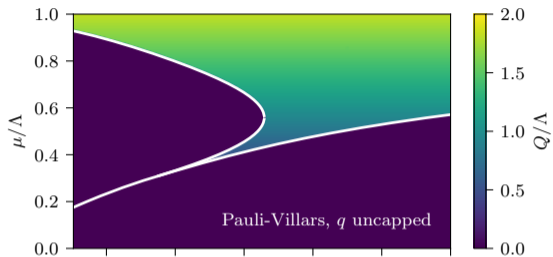
⇒ Take a look at  $Q$  in the  $M_0 - \mu$ -plane

$$Q = \begin{cases} \operatorname{argmin}_q \Gamma^{(2)} & \text{if } \min_q \Gamma^{(2)} < 0 \\ 0 & \text{else} \end{cases}$$

# Q scan

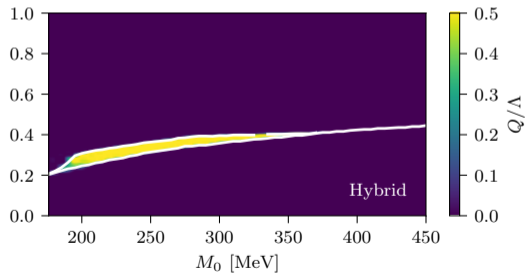
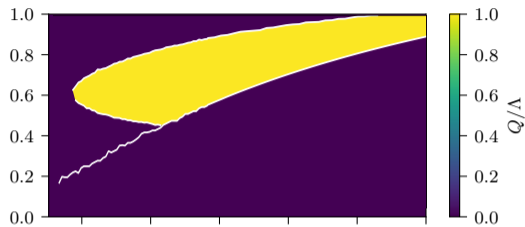
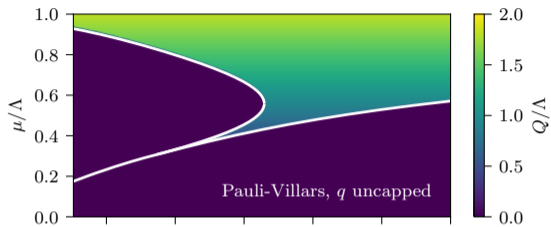


# Q scan

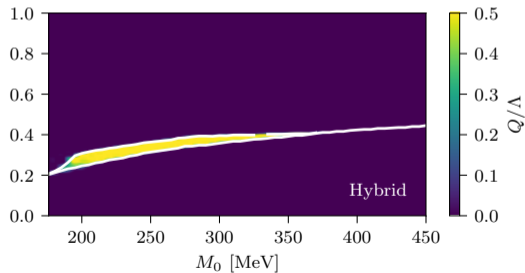
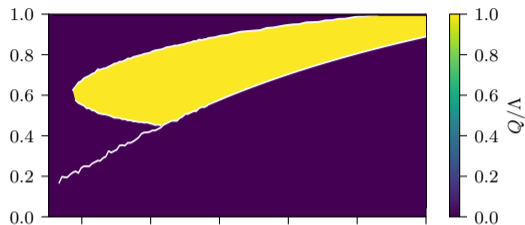
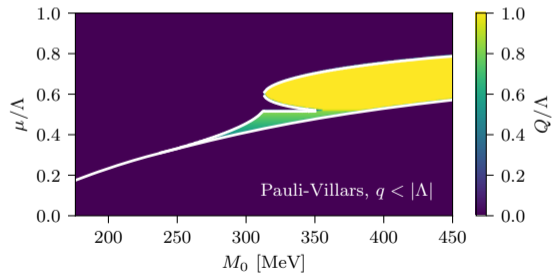
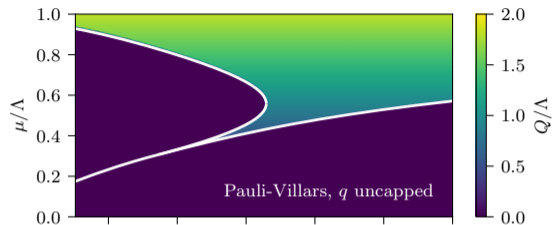




# Q scan



# Q scan



# Summary and Outlook

## Summary:

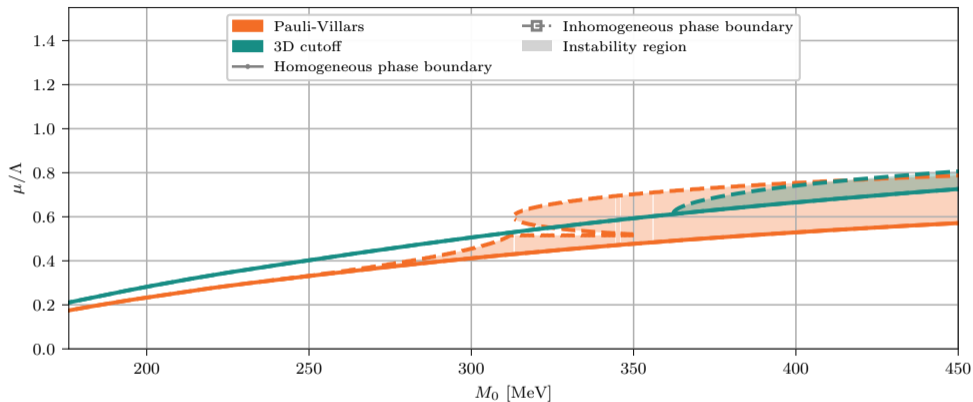
- The parameter regions are very unfavorable - especially for the lattice.
- A *straightforward* lattice investigation of inhomogeneous phases in the  $3 + 1$ -dimensional NJL model is most likely pointless.

## Outlook:

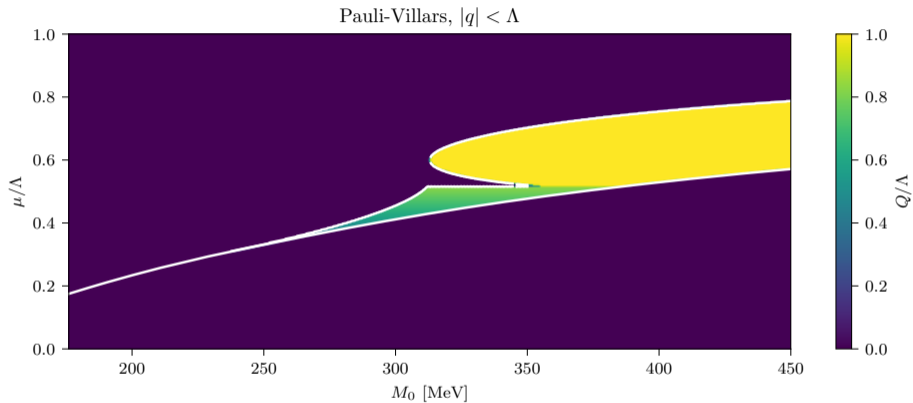
- Making the 3D Cutoff RG consistent yielded promising results
  - ⇒ exploring a similar treatment of the lattice discretizations
  - ⇒ Most likely not applicable in real world
- Finish results with staggered fermions
- Investigate the Quark-Meson model as it might have more favorable parameter regions and is 'renormalizable'

# Appendix

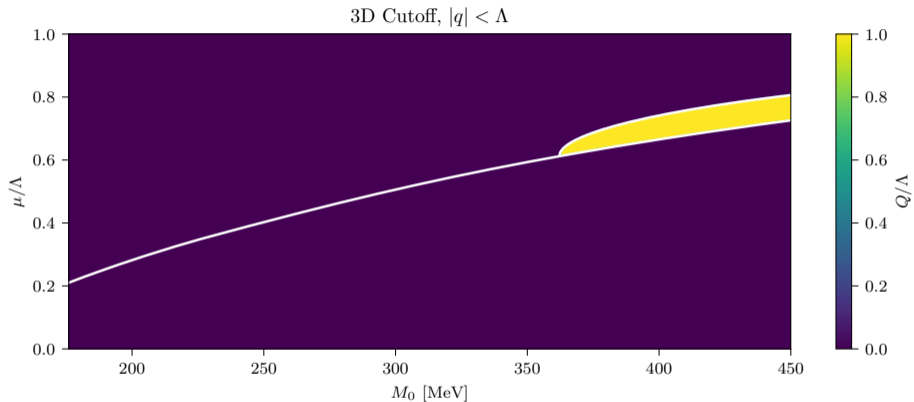
# Continuum capped $q$



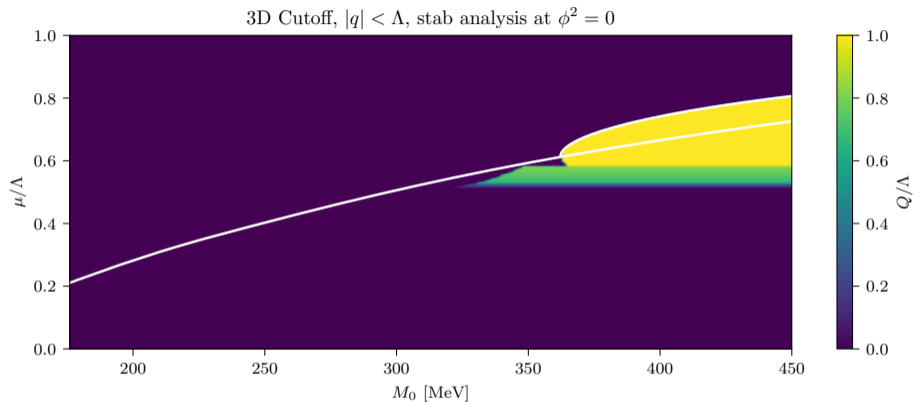
# Continuum capped $q$



# Continuum capped $q$

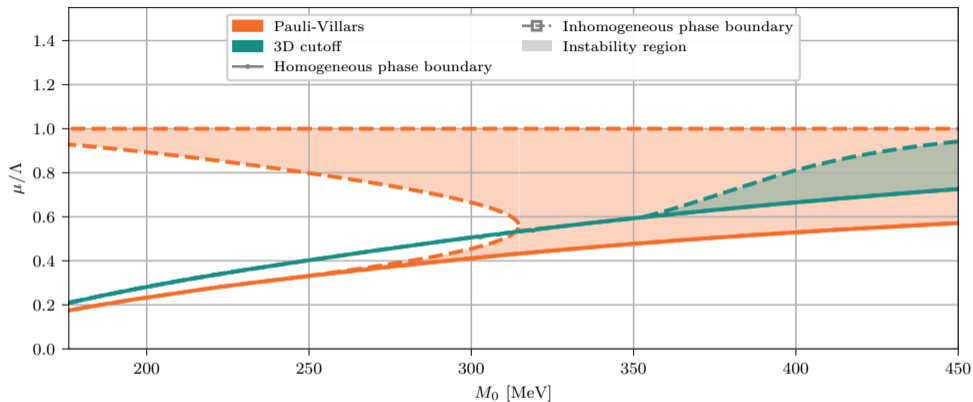


# Continuum capped $q$

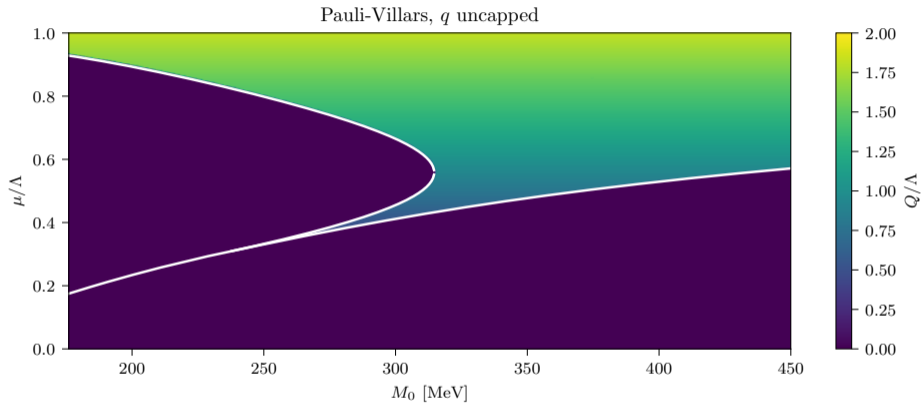




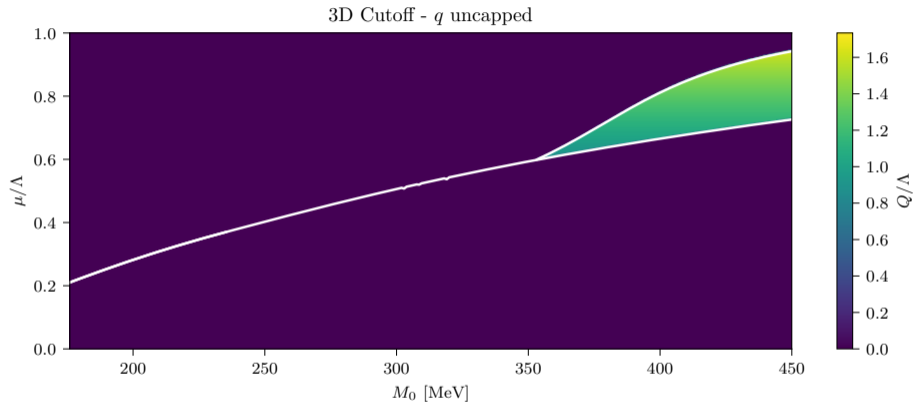
# Continuum Uncapped $q$



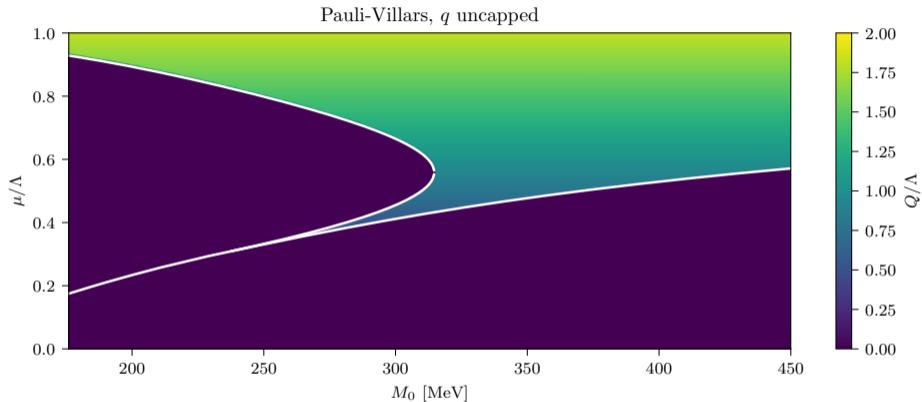
# Continuum Uncapped $q$



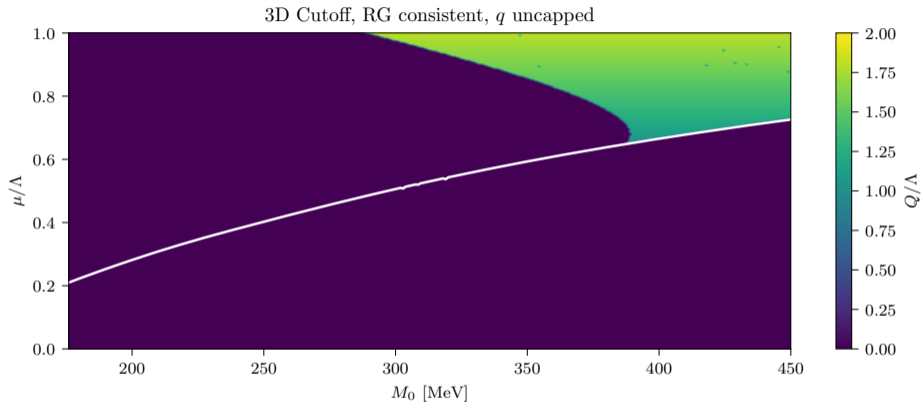
# Continuum Uncapped $q$



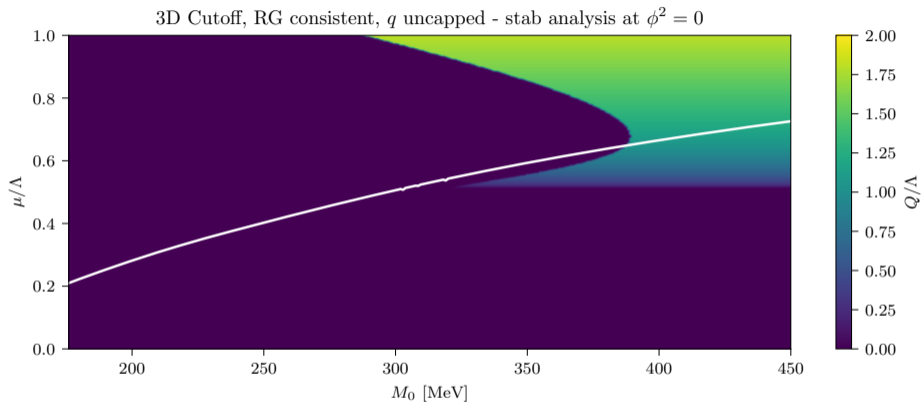
# RG consistent cutoff



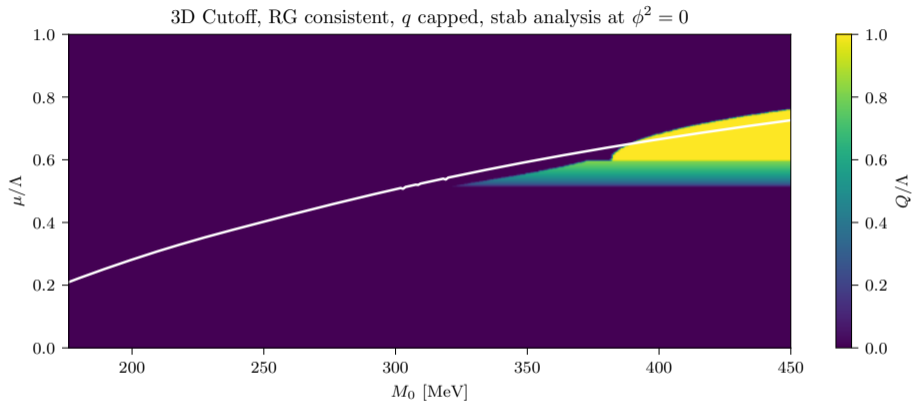
# RG consistent cutoff



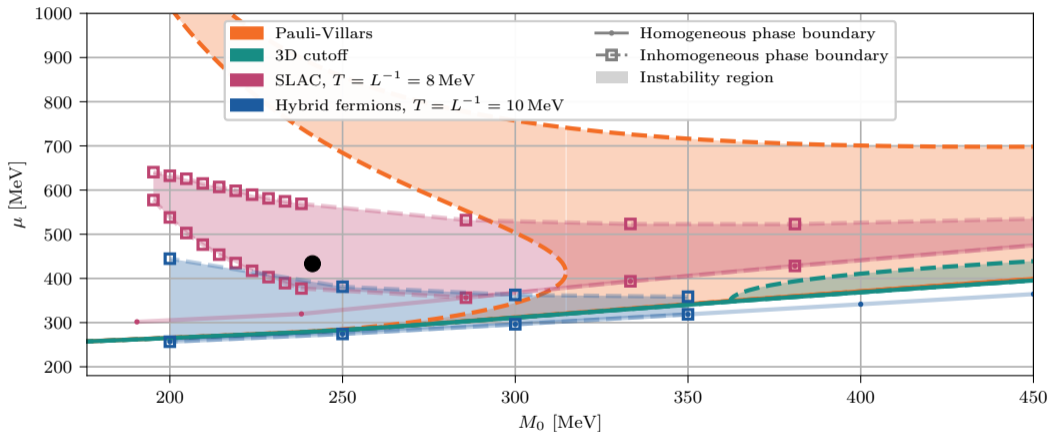
# RG consistent cutoff



# RG consistent cutoff



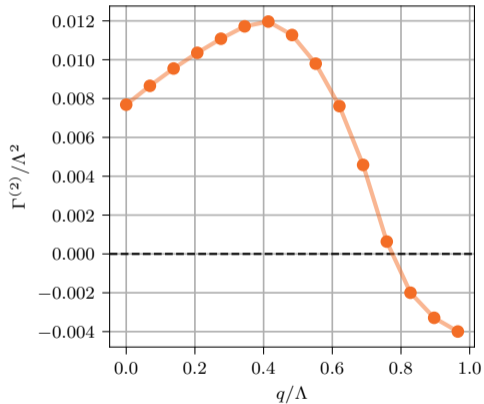
# A closer look at SLAC IV - Minimum configuration



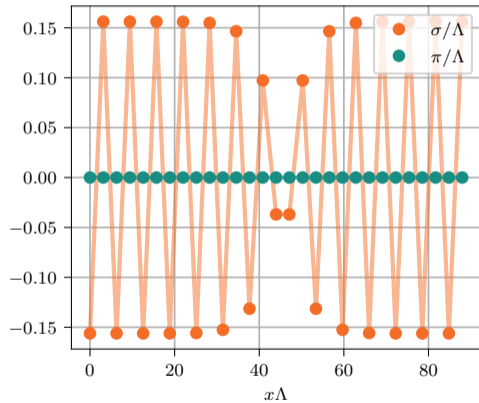


# A closer look at SLAC IV - Minimum configuration

SLAC,  $T \approx 8$  MeV,  $\mu = 442.79$  MeV,  $M_0 = 238$  MeV



SLAC,  $T \approx 8$  MeV,  $\mu = 442.79$  MeV,  $M_0 = 238$  MeV

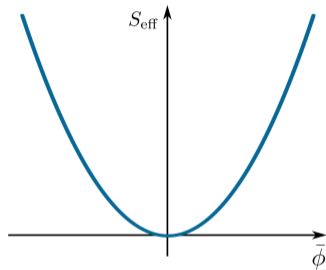


# Stability analysis

- Homogeneous fields

$$\phi(x) = \bar{\phi}$$

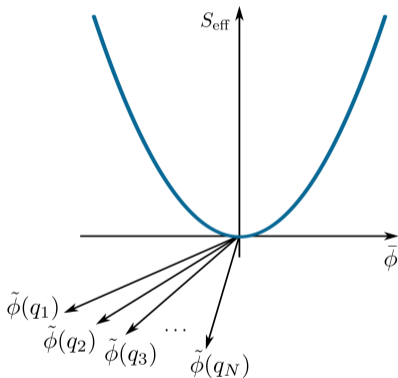
- Minimum is easy to obtain.



# Stability analysis

- In general fields have full space dependence

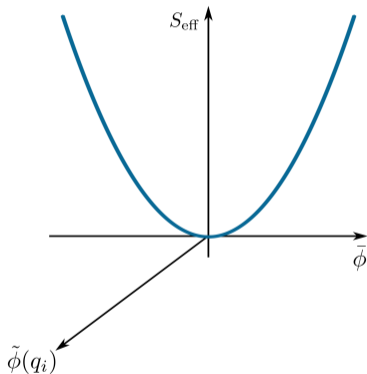
$$\begin{aligned}\phi(x) &= \bar{\phi} + \phi_s(x) \\ &= \bar{\phi} + \sum_j \tilde{\phi}_s(q_j) e^{ixq_j}\end{aligned}$$



# Stability analysis

- In general fields have full space dependence

$$\begin{aligned}\phi(x) &= \bar{\phi} + \phi_s(x) \\ &= \bar{\phi} + \sum_j \tilde{\phi}_s(q_j) e^{ixq_j}\end{aligned}$$

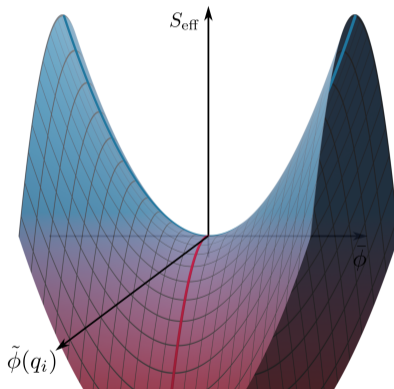


# Stability analysis

- In general fields have full space dependence

$$\begin{aligned}\phi(x) &= \bar{\phi} + \phi_s(x) \\ &= \bar{\phi} + \sum_j \tilde{\phi}_s(q_j) e^{ixq_j}\end{aligned}$$

- Former homogeneous minimum might only be **saddle point**
- Full dependence of  $S_{\text{eff}}$  on  $\phi(x)$  extremely difficult or impossible



# Stability analysis

- Consider only inhomogeneous perturbations

$$\begin{aligned}\phi(x) &= \bar{\phi} + \delta\phi_s(x) \\ &= \bar{\phi} + \sum_j \delta\tilde{\phi}_s(q_j) e^{ixq_j}\end{aligned}$$

- investigate curvature at homogeneous minimum

