## Exploring the QCD phase diagram with Fluctuations and correlations

A. Bzdak, VK, N. Strodthoff: arXiv:1607.07375
A. Bzdak, VK, V. Skokov: arXiv:1612.05128
A. Bzdak, VK: arXiv:1707.02640
A. Bzdak, D. Oliinychenko, J. Steinheimer, VK: arXiv:1804.04463
A. Bzdak, VK: arXiv:1810.01913
A. Bzdak and V.K: arXiv:1811.04456

## The phase diagram



Increase chemical potential by lowering the beam energy
In reality, we add baryons (nucleons) from target and projectile to mid-rapidity

## What we know about the Phase Diagram



## What we "hope" for



## Is there a critical point?



## Cumulants and phase structure



What we always see....


What it really means....
" $\mathrm{T}_{\mathrm{c}}$ " $\sim 160 \mathrm{MeV}$

## Derivatives



## How to measure derivatives

At $\mu=0$ :

$$
\begin{gather*}
Z=\operatorname{tr} e^{-\hat{E} / T+\mu / T \hat{N}_{B}} \\
\langle E\rangle=\frac{1}{Z} \operatorname{tr} \hat{E} e^{-\hat{E} / T+\mu / T \hat{N}_{B}}=-\frac{\partial}{\partial 1 / T} \ln (Z) \\
\left\langle(\delta E)^{2}\right\rangle=\left\langle E^{2}\right\rangle-\langle E\rangle^{2}=\left(-\frac{\partial}{\partial 1 / T}\right)^{2} \ln (Z)=\left(-\frac{\partial}{\partial 1 / T}\right)\langle E\rangle \\
\left\langle(\delta E)^{n}\right\rangle=\left(-\frac{\partial}{\partial 1 / T}\right)^{n-1}\langle E\rangle
\end{gather*}
$$

Cumulants of Energy measure the temperature derivatives of the EOS
Cumulants of Baryon number measure the chem. pot. derivatives of the EOS

## Cumulants of (Baryon) Number

$K_{n}=\frac{\partial^{n}}{\partial(\mu / T)^{n}} \ln Z=\frac{\partial^{n-1}}{\partial(\mu / T)^{n-1}}\langle N\rangle$
$K_{1}=\langle N\rangle, K_{2}=\langle N-\langle N\rangle\rangle^{2}, K_{3}=\langle N-\langle N\rangle\rangle^{3}$

Cumulants scale with volume (extensive): $\quad K_{n} \sim V$

Volume not well controlled in heavy ion collisions

Cumulant Ratios: $\quad \frac{K_{2}}{\langle N\rangle}, \frac{K_{3}}{K_{2}}, \frac{K_{4}}{K_{2}}$

## Simple model

Change degrees of freedom at phase transition

$$
\langle N\rangle=\operatorname{dof}(\mu) e^{\mu / T} \int d^{3} p e^{-E / T}
$$





## Close to $\mu=0$

$$
F=F(r), \quad r=\sqrt{T^{2}+a \mu^{2}}
$$

a ~ curvature of critical line


$$
\left.\frac{\partial^{2}}{\partial \mu^{2}} F(T, \mu)\right|_{\mu=0}=\frac{a}{T} \frac{\partial}{\partial T} F(T, \mu=0) \sim\langle E\rangle
$$

Needs higher order cumulants (derivatives) at $\mu \sim 0$

# Latest STAR result on net-proton cumulants 

X. Luo, NPA 956 (2016) 75

$\mathrm{K}_{4} / \mathrm{K}_{2}$ follows expectation for $\mathrm{CP}, \mathrm{K}_{3} / \mathrm{K}_{2}$ no so much..... URQMD totally fails to get trend for $\mathrm{K}_{4} / \mathrm{K}_{2}$ !

## Let's take the preliminary STAR data at face value

## Further insights: Correlations

Cumulants $\quad K_{n}=\frac{\partial^{n}}{\partial \hat{\mu}^{n}} P / T^{4}$
$K_{2}=\langle N-\langle N\rangle\rangle^{2}=\left\langle(\delta N)^{2}\right\rangle$
$\rho_{2}\left(p_{1}, p_{2}\right)=\rho_{1}\left(p_{1}\right) \rho_{1}\left(p_{2}\right)+C_{2}\left(p_{1}, p_{2}\right), \quad \mathrm{C}_{2}$ : Correlation Function

$$
\begin{aligned}
& K_{3}=\left\langle(\delta N)^{3}\right\rangle \\
& \begin{aligned}
\rho_{3}\left(p_{1}, p_{2}, p_{3}\right)= & \left.\rho_{1}\left(p_{1}\right) \rho_{1}\left(p_{2}\right) \rho_{1}\left(p_{3}\right)+\rho_{1}\left(p_{1}\right) \underline{C_{2}\left(p_{2}, p_{3}\right.}\right)+\rho_{1}\left(p_{2}\right) \underline{C_{2}\left(p_{1}, p_{3}\right)} \\
& +\rho_{1}\left(p_{3}\right) \underline{C_{2}\left(p_{1}, p_{2}\right)}+\underline{C_{3}\left(p_{1}, p_{2}, p_{3}\right)}
\end{aligned}
\end{aligned}
$$

## From Cumulants to Correlations (no anti-protons)

Defining integrated correlations function a.k.a factorial cumulants
$C_{n}=\int d p_{1} \ldots d p_{n} C_{n}\left(p_{1}, \ldots, p_{n}\right)$
Simple Algebra leads to relation between correlations $\mathrm{C}_{\mathrm{n}}$ and $\mathrm{K}_{\mathrm{n}}$
$C_{2}=-K_{1}+K_{2}$,
$C_{3}=2 K_{1}-3 K_{2}+K_{3}$,
$C_{4}=-6 K_{1}+11 K_{2}-6 K_{3}+K_{4},$.
or vice versa
$K_{2}=\langle N\rangle+C_{2}$
$K_{3}=\langle N\rangle+3 C_{2}+C_{3}$
$K_{4}=\langle N\rangle+7 C_{2}+6 C_{3}+C_{4}$

## Preliminary Star Data (X. Luo, PoS Cpod 2014 (019))



Significant four particle correlations!

Four particle correlation dominate $\mathrm{K}_{4}$ for central collisions at 7.7 GeV

$$
\begin{aligned}
K_{2} & =\langle N\rangle+C_{2} \\
K_{3} & =\langle N\rangle+3 C_{2}+C_{3} \\
K_{4} & =\langle N\rangle+7 C_{2}+6 C_{3}+C_{4}
\end{aligned}
$$

## Correlations



## Rapidity dependence

$C_{k}(\Delta Y)=\int_{\Delta Y} d y_{1} \ldots d y_{k} \rho_{1}\left(y_{1}\right) \ldots \rho_{1}\left(y_{k}\right) c_{k}\left(y_{1}, \ldots, y_{k}\right)$
Assume: $\quad \rho_{1}(y) \simeq$ const.
short range correlations:

$$
\begin{aligned}
& c_{k}\left(y_{1}, \ldots, y_{k}\right) \sim \delta\left(y_{1}-y_{2}\right) \ldots \delta\left(y_{k-1}-y_{k}\right) \\
& C_{k}(\Delta Y) \sim \Delta Y \rightarrow K_{k} \sim \Delta Y
\end{aligned}
$$

Long range correlations:
$c_{k}\left(y_{1}, \ldots, y_{k}\right)=$ const .

$$
\begin{aligned}
& C_{k}(\Delta Y) \sim(\Delta Y)^{k} \sim\langle N\rangle^{k} \\
& \quad \Rightarrow K_{n}=K_{n}(\langle N\rangle)
\end{aligned}
$$

## Long range correlations

$$
\begin{aligned}
& C_{k}=\langle N\rangle^{k} c_{k} \\
& c_{k}=\text { const. } \Rightarrow K_{n}=K_{n}(\langle N\rangle)
\end{aligned}
$$



## Energy dependence



Note: anti-protons are non- negligible above 19.6 GeV Data are protons only

## Can we understand these correlations?

- Two particle correlations can be understood by simple

Glauber model + Baryon number conservation



Four particle correlations are orders of magnitudes larger in the data Also seen in URQMD calculations by He et al. PLB774 (2017) 623

Need to assume the $\sim 40 \%$ of protons come from 8 -nucleon cluster in order to get magnitude right!

## URQMD



He, Luo PLB774 (2017) 623

## Latest STAR result on net-proton cumulants <br> X. Luo, NPA 956 (2016) 75


$\mathrm{K}_{4} / \mathrm{K}_{2}$ above baseline $\mathrm{K}_{3} / \mathrm{K}_{2}$ below baseline

## Shape of probability distribution

$$
\begin{aligned}
& K_{3}<\langle N\rangle \quad K_{3}=\langle N-\langle N\rangle\rangle^{3} \\
& K_{4}>\langle N\rangle \\
& K_{4}=\langle N-\langle N\rangle\rangle^{4}-3\langle N-\langle N\rangle\rangle^{2}
\end{aligned}
$$

## Simple two component model



Weight of small component: $\sim 0.3 \%$

## Simple two component model

Difficult to see in the real data with efficiency $\varepsilon=0.65$


## Two component model

$$
\begin{aligned}
& P(N)=(1-\alpha) P_{(a)}(N)+\alpha P_{(b)}(N) \\
& \bar{N}=\left\langle N_{(a)}\right\rangle-\left\langle N_{(b)}\right\rangle \\
& C_{2}=C_{2}^{(a)}-\alpha\left\{\bar{C}_{2}-(1-\alpha) \bar{N}^{2}\right\} \\
& C_{3}=C_{3}^{(a)}-\alpha\left\{\bar{C}_{3}+(1-\alpha)\left[(1-2 \alpha) \bar{N}^{3}-3 \bar{N} \bar{C}_{2}\right]\right\} \\
& C_{4}=C_{4}^{(a)}-\alpha\left\{\bar{C}_{4}-(1-\alpha)\left[\left(1-6 \alpha+6 \alpha^{2}\right) \bar{N}^{4}-6(1-2 \alpha) \bar{N}^{2} \bar{C}_{2}+4 \bar{N} \bar{C}_{3}+3\left(\bar{C}_{2}\right)^{2}\right]\right\} \\
& \bar{C}_{n}=C_{n}^{(a)}-C_{n}^{(b)}
\end{aligned}
$$

For Poisson, $\mathrm{C}_{(\mathrm{a})}, \mathrm{C}_{(\mathrm{b})}=0$

Fit to STAR data: $\quad\left\langle N_{(a)}\right\rangle \simeq 40,\left\langle N_{(b)}\right\rangle \simeq 25, \alpha \simeq 0.003$

## Two component model

$P(N)=(1-\alpha) P_{(a)}(N)+\alpha P_{(b)}(N)$
$\bar{N}=\left\langle N_{(a)}\right\rangle-\left\langle N_{(b)}\right\rangle>0$
For $\mathrm{P}_{(\mathrm{a})}, \mathrm{P}_{(\mathrm{b})}$ Poisson, or (to good approximation) Binomial
$C_{n}=(-1)^{n} K_{n}^{B} \bar{N}^{n} \quad n \geq 2$
$K_{n}^{B}$ : Cumulant of Bernoulli distribution
$\alpha \ll 1, K_{n}^{B}=\alpha \Rightarrow C_{n} \simeq \alpha(-1)^{n} \bar{N}^{n}$
$\Rightarrow\left|C_{n}\right| \sim\langle N\rangle^{n}$ as seen by STAR ( i.e. "infinite" correlation length)
predict: $\quad \frac{C_{4}}{C_{3}}=\frac{C_{5}}{C_{4}}=\frac{C_{n+1}}{C_{n}}=-\bar{N} \quad \bar{N} \simeq 15$
Clear and falsifiable prediction: $\quad C_{5} \approx-2650 \quad C_{6} \approx 41000$

## Two component model is Statistics "friendly"



Based on 144393 events (same as STAR 0-5\% at 7.7 GEV)

## This model can be tested RIGHT NOW!

Model prediction:

$$
\begin{array}{ll}
C_{5}=-2645(1 \pm 0.14), & C_{6}=40900(1 \pm 0.18),
\end{array} \quad \text { Efficiency }
$$

$$
\begin{aligned}
& C_{5}=-307(1 \pm 0.31), \quad C_{6}=3085(1 \pm 0.41), \\
& C_{7}=-30155(1 \pm 0.61), \quad C_{8}=271492(1 \pm 1.06),
\end{aligned}
$$

Efficiency
UN-corrected

Based on 144393 events (same as STAR 0-5\% at 7.7 GEV)

## Double the acceptance




Should be visible in raw (unfolded) data

## Speculation




## Simple two component model



Analyse data for $\mathrm{N}_{\mathrm{p}}<20$

- Is flow etc different?
- "Inspect by eye ( $<1 \%$ of all events)


## The return of the camel?



## Net-baryon distribution consisted with Lattice QCD



For details, see arXiv:1810.01913

## Summary

- Fluctuations sensitive to phase structure:
- measure "derivatives" of EOS
- Cumulants contain information about correlations
- Preliminary STAR data:
- Significant four particle correlations at 7.7 and 11.5 GeV
- Fluctuations of system size ( $\mathrm{N}_{\mathrm{part}}$ ) and stopping
- May explain 2-particle correlations
-Fail to reproduce the magnitude of 3- and 4- particle correlations
- 3 and 4 particle correlations are HUGE!
- "Bi-Modal" distribution works
- Can be tested RIGHT NOW by STAR.
- Net-baryon number distribution consistent with lattice
- Deviation from Skellam is very small!


## Thank You

## Net-baryon multiplicity distribution

Utilize of cluster expansion model of Vovchenko et al arXiv:1711.01261
Virial expansion: $\quad \frac{P}{T^{4}}=\frac{1}{V T^{3}} \ln (Z)=\sum_{k=0}^{\infty} p_{k}(T) \cosh \left(k \hat{\mu}_{B}\right)$
Cluster model: $\quad p_{k}=f\left(p_{1}, p_{2}\right) ; k>2$
Lattice QCD: $\quad p_{1}, p_{2} \quad$ Vovchenko et al, arXiv:1708.02852



Figure from:
arXiv:1807.06472

## Net-baryon multiplicity distribution

Virial expansion:

$$
\frac{P}{T^{4}}=\frac{1}{V T^{3}} \ln (Z)=\sum_{k=0}^{\infty} p_{k}(T) \cosh \left(k \hat{\mu}_{B}\right)
$$

$$
Z=\exp \left[V T^{3} \sum_{k=0}^{\infty} p_{k}(T) \cosh \left(k \hat{\mu}_{B}\right)\right]=z_{0}+2 \sum_{\mathcal{N}=1}^{\infty} z_{\mathcal{N}} \cosh \left(\mathcal{N} \hat{\mu}_{B}\right)
$$

Multiplicity distribution:

$$
P(\mathcal{N})=\frac{z_{\mathcal{N}} e^{\hat{\mu}_{B} \mathcal{N}}}{Z}
$$

$$
\hat{\mu}_{B} \rightarrow i \bar{\mu}_{B}
$$

$$
P(\mathcal{N})=\frac{1}{\pi} \int_{0}^{\pi} d \bar{\mu}_{B} \cos \left(\mathcal{N} \bar{\mu}_{B}\right) \frac{\exp \left[V T^{3} \sum_{k=1}^{\infty} p_{k}(T) \cos \left(k \bar{\mu}_{B}\right)\right]}{\exp \left[V T^{3} \sum_{k=1}^{\infty} p_{k}(T)\right]}
$$

LHC



Top RHIC








## Almasi at al Model



## Latest STAR result on net-proton cumulants <br> X. Luo, NPA 956 (2016) 75 <br> 


$\mathrm{K}_{4} / \mathrm{K}_{2}$ follows expectation, $\mathrm{K}_{3} / \mathrm{K}_{2}$ no so much..... URQMD totally fails to get trend for $\mathrm{K}_{4} / \mathrm{K}_{2}$ !

## Correlations near the critical point

M. Stephanov, 0809.3450, PRL 102

Scaling of Cumulants $\mathrm{K}_{\mathrm{n}}$ with correlation length $\xi$
$K_{2} \sim \xi^{2}, K_{3} \sim \xi^{4.5}, K_{4} \sim \xi^{7}$
Cumulants from Correlations

$$
\begin{aligned}
& K_{2}=\langle N\rangle+C_{2} \\
& K_{3}=\langle N\rangle+3 C_{2}+C_{3} \\
& K_{4}=\langle N\rangle+7 C_{2}+6 C_{3}+C_{4}
\end{aligned}
$$

Consequently:

$$
C_{2} \sim \xi^{2}, C_{3} \sim \xi^{4.5}, \quad C_{4} \sim \xi^{7}
$$

Correlations $\mathrm{C}_{\mathrm{n}}$ pick up the most divergent pieces of cumulants $\mathrm{K}_{\mathrm{n}}$ !

## Reduced correlation function

Reduced correlation function
$c_{k}=\frac{\int \rho_{1}\left(y_{1}\right) \cdots \rho_{1}\left(y_{k}\right) c_{k}\left(y_{1}, \ldots, y_{k}\right) d y_{1} \cdots d y_{k}}{\int \rho_{1}\left(y_{1}\right) \cdots \rho_{1}\left(y_{k}\right) d y_{1} \cdots d y_{k}}$

$$
C_{k}=\langle N\rangle^{k} c_{k}
$$

Independent sources such as resonances, cluster, $\mathrm{p}+\mathrm{p}$ :
$c_{k} \sim \frac{\left\langle N_{s}\right\rangle}{\langle N\rangle^{k}} \sim \frac{1}{\langle N\rangle^{k-1}}$
For example two particle correlations:
$c_{2} \sim \frac{\text { Number of sources }}{\text { Number of all pairs }}=\frac{\text { Number of correlated pairs }}{\text { Number of all pairs }}=\frac{1}{\langle N\rangle}$

## Centrality dependence



## Centrality dependence






## Preliminary Star data are consistent with long range correlations


7.7 GeV
central

19.6 GeV central

