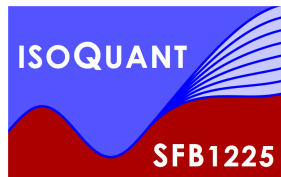


# Event-by-event fluctuations of identified particles in heavy-ion collisions

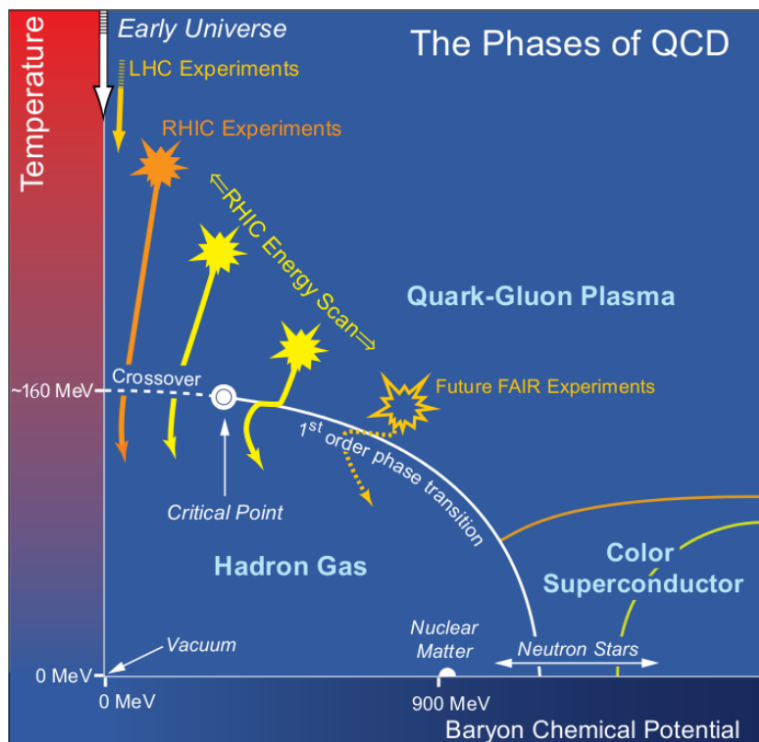
Alice Ohlson  
Universität Heidelberg

Hirschegg, 17 January 2019



# Fluctuations in heavy ion collisions

- Event-by-event fluctuations of particle multiplicities are used to study properties and phase structure of strongly-interacting matter
  - Fluctuations grow in the region near a phase transition and/or critical point
    - can we observe signs of criticality?



## Critical opalescence in CO<sub>2</sub>

J.V. Sengers, A.L Sengers, Chem. Eng. News, June 10, 1968, 104–118, 1968



$T > T_c$

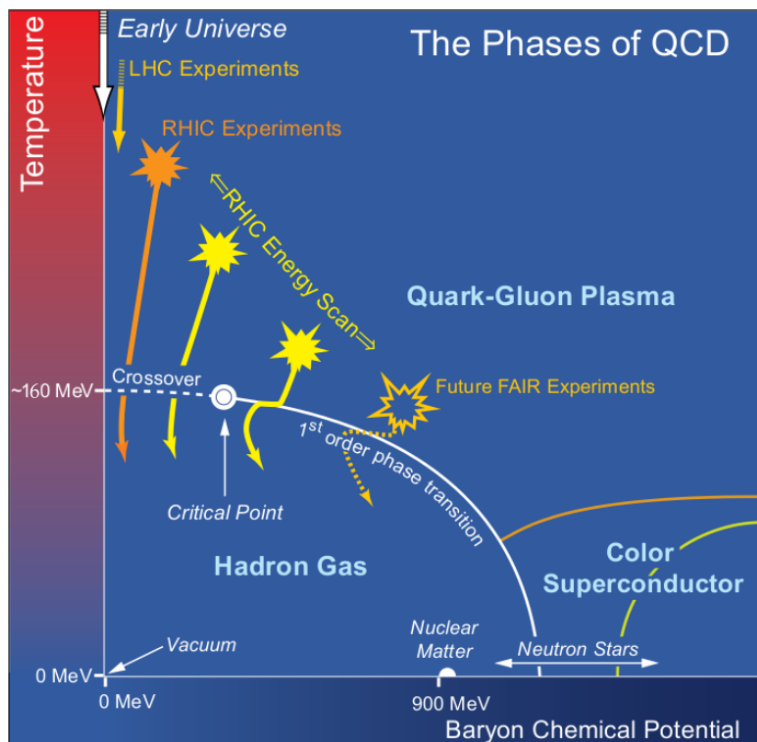
$T \gtrsim T_c$

$T \lesssim T_c$

$T < T_c$

# Fluctuations in heavy ion collisions

- Event-by-event fluctuations of particle multiplicities are used to study properties and phase structure of strongly-interacting matter



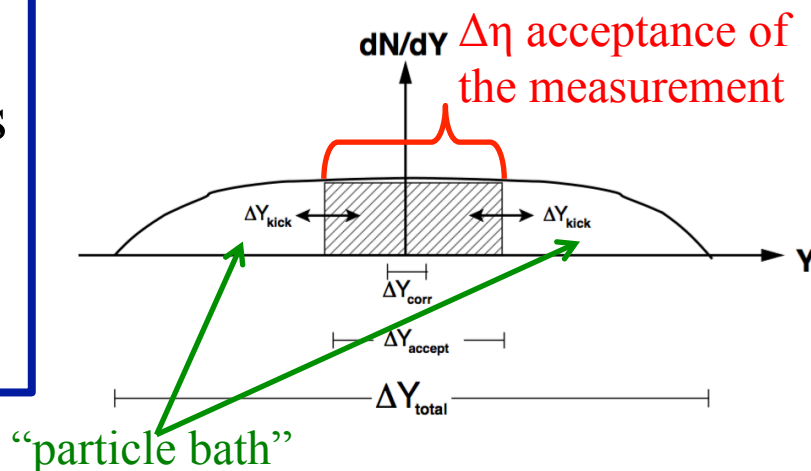
- Fluctuations grow in the region near a phase transition and/or critical point
  - can we observe signs of criticality?
- Fluctuations of conserved charges can be related to susceptibilities calculable in lattice QCD
  - precision test of LQCD at  $\mu_B \approx 0$

# Connecting theory to experiment

- Thermodynamic susceptibilities  $\chi$ 
  - describe the response of a thermalized system to changes in external conditions, fundamental properties of the medium
  - can be calculated within lattice QCD
  - within the Grand Canonical Ensemble, are related to event-by-event fluctuations of the number of conserved charges

Theory:  
susceptibilities

$$\chi_n^B = \frac{\partial^n (P/T^4)}{\partial (\mu_B/T)^n}$$



Experiment:  
moments of  
particle  
multiplicity  
distributions

$$\Delta N_B = N_B - N_{\bar{B}}$$



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- Thermodynamic susceptibilities  $\chi$ 
  - describe the response of a thermalized system to changes in external conditions, fundamental properties of the medium
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Theory:  
susceptibilities

$$\chi_n^B = \frac{\partial^n (P/T^4)}{\partial (\mu_B/T)^n}$$

$$\langle \Delta N_B \rangle = VT^3 \chi_1^B$$

$$\langle (\Delta N_B - \langle \Delta N_B \rangle)^2 \rangle = VT^3 \chi_2^B = \sigma^2$$

$$\langle (\Delta N_B - \langle \Delta N_B \rangle)^3 \rangle / \sigma^3 = \frac{VT^3 \chi_3^B}{(VT^3 \chi_2^B)^{3/2}} = S$$

$$\langle (\Delta N_B - \langle \Delta N_B \rangle)^4 \rangle / \sigma^4 - 3 = \frac{VT^3 \chi_4^B}{(VT^3 \chi_2^B)^2} = \kappa$$

Experiment:  
moments of  
particle  
multiplicity  
distributions

$$\Delta N_B = N_B - N_{\bar{B}}$$

# Connecting theory to experiment

- Thermodynamic susceptibilities  $\chi$ 
  - describe the response of a thermalized system to changes in external conditions, fundamental properties of the medium
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Theory:  
susceptibilities

$$\chi_n^B = \frac{\partial^n (P/T^4)}{\partial (\mu_B/T)^n}$$

$$\langle \Delta N_B \rangle = VT^3 \chi_1^B$$

$$\langle (\Delta N_B - \langle \Delta N_B \rangle)^2 \rangle = S \sigma^2$$

$$K \sigma^2 = \chi_4^B / \chi_2^B$$

$$\langle (\Delta N_B - \langle \Delta N_B \rangle)^3 \rangle / \sigma^3 = \frac{VT^3 \chi_3^B}{(VT^3 \chi_2^B)^{3/2}} = S$$

$$\langle (\Delta N_B - \langle \Delta N_B \rangle)^4 \rangle / \sigma^4 - 3 = \frac{VT^3 \chi_4^B}{(VT^3 \chi_2^B)^2} = K$$

Experiment:  
moments of  
particle  
multiplicity  
distributions

$$\Delta N_B = N_B - N_{\bar{B}}$$

# Connecting theory to experiment

- Thermodynamic susceptibilities  $\chi$ 
  - describe the response of a thermalized system to changes in external conditions, fundamental properties of the medium
  - can be calculated within lattice QCD
  - within the Grand Canonical Ensemble, are related to event-by-event fluctuations of the number of conserved charges

Theory:  
fixed volume,  
particle bath in  
GCE

$$\langle \Delta N_B \rangle \neq VT^3 \chi_1^B$$

$$\langle (\Delta N_B - \langle \Delta N_B \rangle)^2 \rangle \neq VT^3 \chi_2^B = \sigma^2$$

$$\langle (\Delta N_B - \langle \Delta N_B \rangle)^3 \rangle / \sigma^3 \neq \frac{VT^3 \chi_3^B}{(VT^3 \chi_2^B)^{3/2}} = S$$

$$\langle (\Delta N_B - \langle \Delta N_B \rangle)^4 \rangle / \sigma^4 - 3 \neq \frac{VT^3 \chi_4^B}{(VT^3 \chi_2^B)^2} = K$$

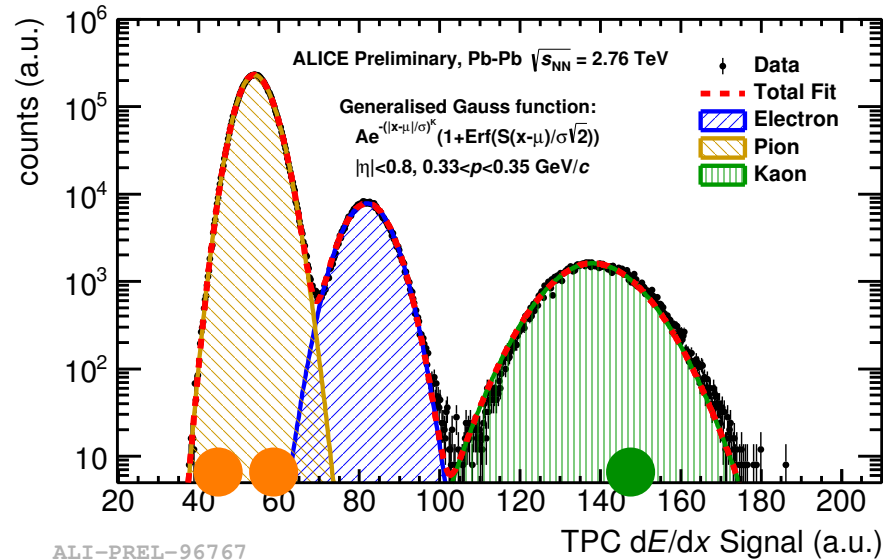
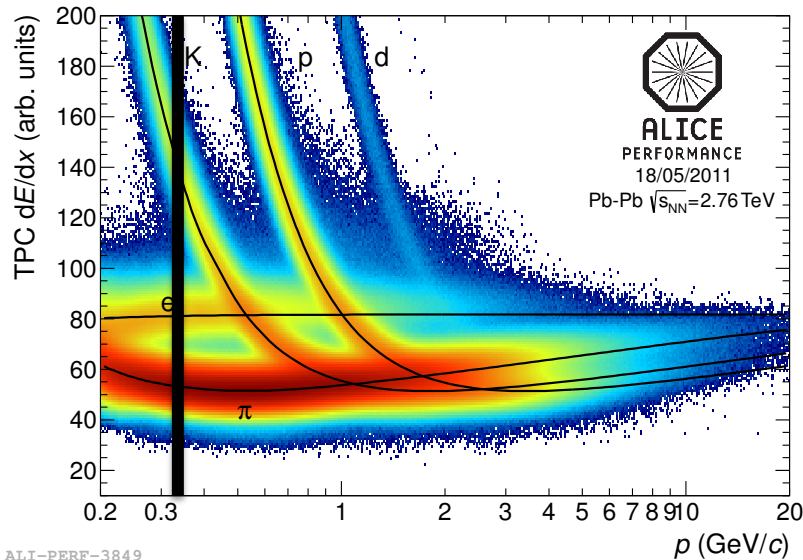
Experiment:  
event-by-event  
volume fluctuations,  
global conservation  
laws

# Experimental Challenges

1. Event-by-event particle identification
2. Event-by-event efficiency correction

We know how to correct the first moments,  
but what about the higher moments?

# The challenge: event-by-event PID



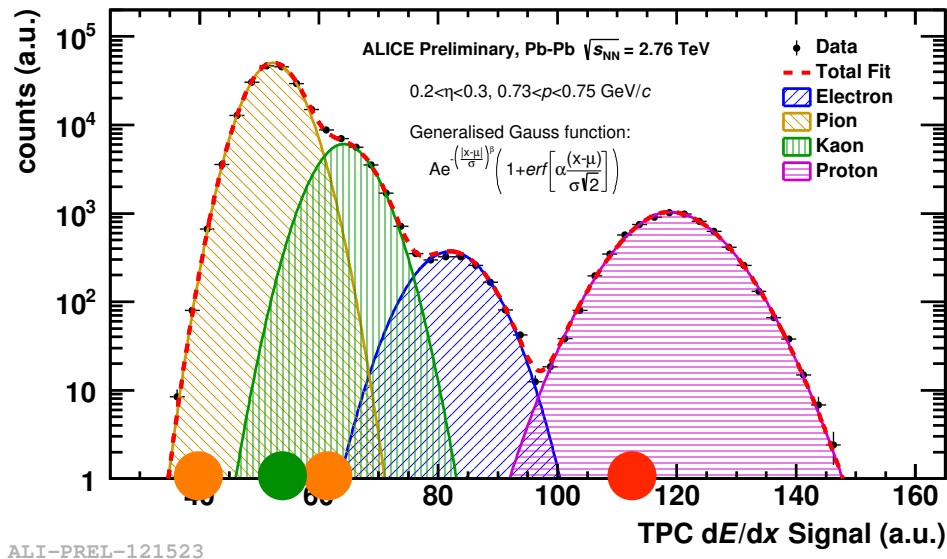
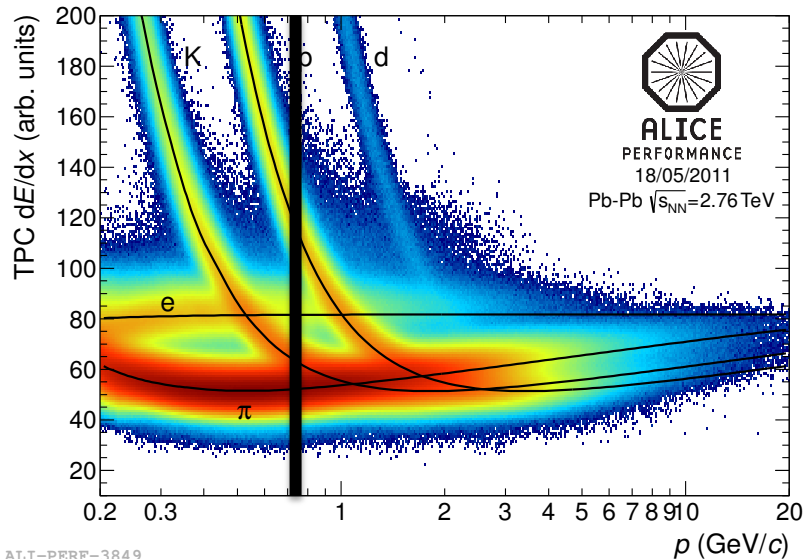
- Traditional method:

- count number of pions ( $N_\pi$ ), kaons ( $N_K$ ), protons ( $N_p$ ) in each event

$$N_p = \sum_i^{\#tracks} \begin{cases} 1 & \text{particle } i \text{ is a proton} \\ 0 & \text{particle } i \text{ is not a proton} \end{cases}$$

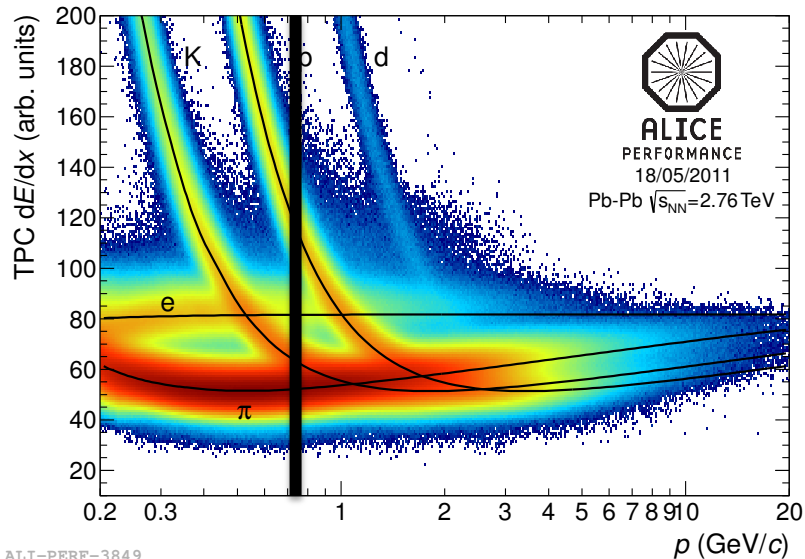
- find moments of distributions of  $N_\pi, N_K, N_p, \dots$

# Traditional method

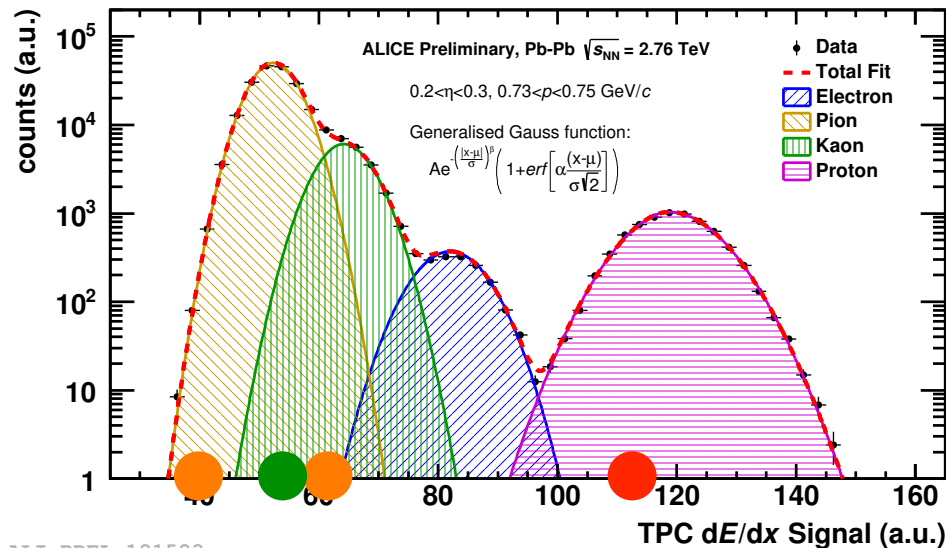


- What if PID is unclear?
  - use other detector information or reject phase space bin
  - results in lower efficiency

# Identity method



ALI-PERF-3849



ALI-PREL-121523

- As a function of the PID variable  $m$ , determine probability  $w$  that particle is of a given species

- Calculate event-by-event sum of weights  $W_p$ ,

$$W_K, W_p, \dots \quad W_p = \sum_i^{\# \text{tracks}} w_p(m_i)$$

- Using knowledge of inclusive  $m$  distributions, unfold moments of  $W$  distributions to get moments of  $N$
- Contamination is accounted for, full phase space can be used

M. Gazdzicki et al., PRC 83 (2011) 054907, arXiv:1103.2887 [nucl-th]

M. I. Gorenstein, PRC 84, (2011) 024902, arXiv:1106.4473 [nucl-th]

A. Rustamov et al., PRC 86 (2012) 044906, arXiv:1204.6632 [nucl-th]

M. Arslanodk and A. Rustamov, arXiv: 1807.06370 [hep-ex]

# Efficiency corrections: several ideas

- Simple scaling of moments using HIJING and/or AMPT
- Correction of factorial moments assuming binomial track loss

A. Bzdak and V. Koch,  
Phys. Rev. C86, 044904 (2012),  
arXiv:1206.4286 [nucl-th].

A. Bzdak and V. Koch,  
Phys. Rev. C91, 027901 (2015),  
arXiv:1312.4574 [nucl-th].

## – extension to Identity Method

C. Pruneau, Phys. Rev. C96 (2017) 054902,  
arXiv:1706.01333 [physics.data-an]

- Correction using moments of detector response matrix

T. Nonaka et al., Nucl. Inst. Meth. A 906 (2018) 10,  
arXiv:1805.00279 [physics.data-an]

- Full unfolding of moments

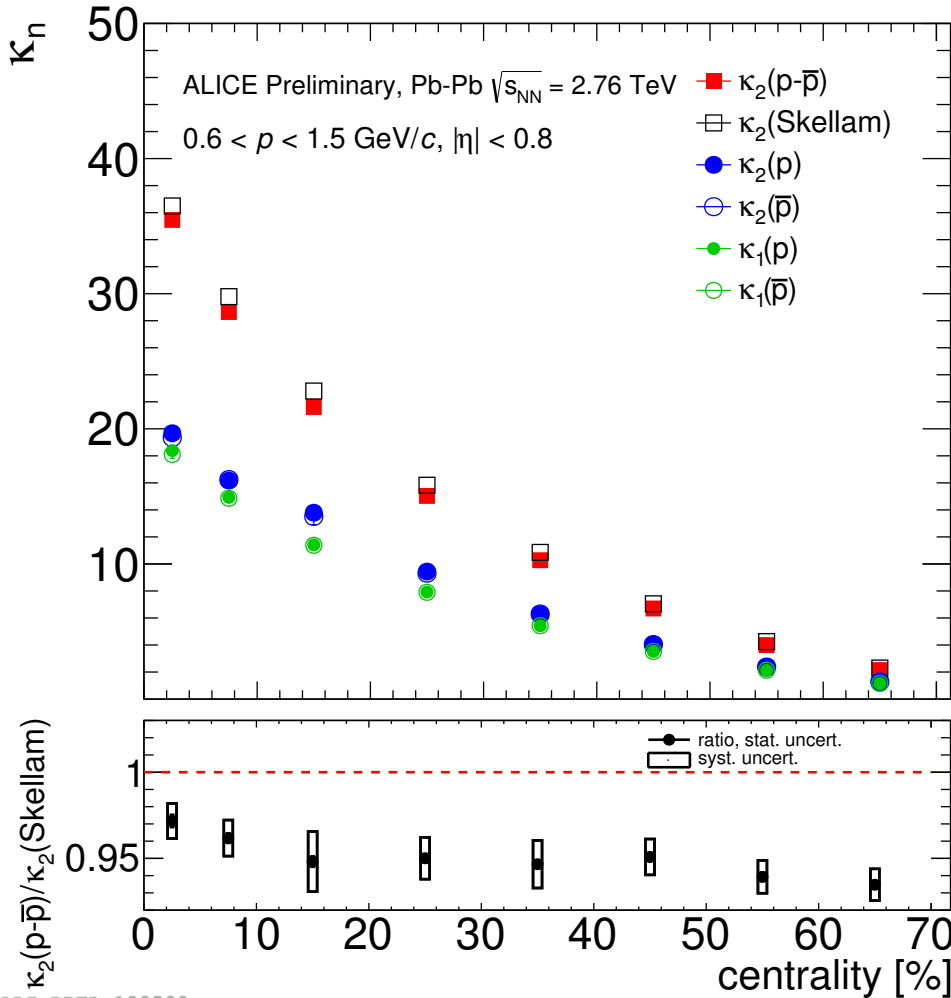
All correction methods rely on different assumptions,  
which must be assessed and tested carefully!



# Second moments from the LHC

# Net-proton fluctuations

A. Rustamov for ALICE, QM2017



ALI-PREL-122590

$$\kappa_1(p) = \langle N_p \rangle \quad \kappa_2(p) = \left\langle \left( N_p - \langle N_p \rangle \right)^2 \right\rangle$$

$$\kappa_2(p - \bar{p}) = \left\langle \left( N_p - N_{\bar{p}} - \langle N_p - N_{\bar{p}} \rangle \right)^2 \right\rangle$$

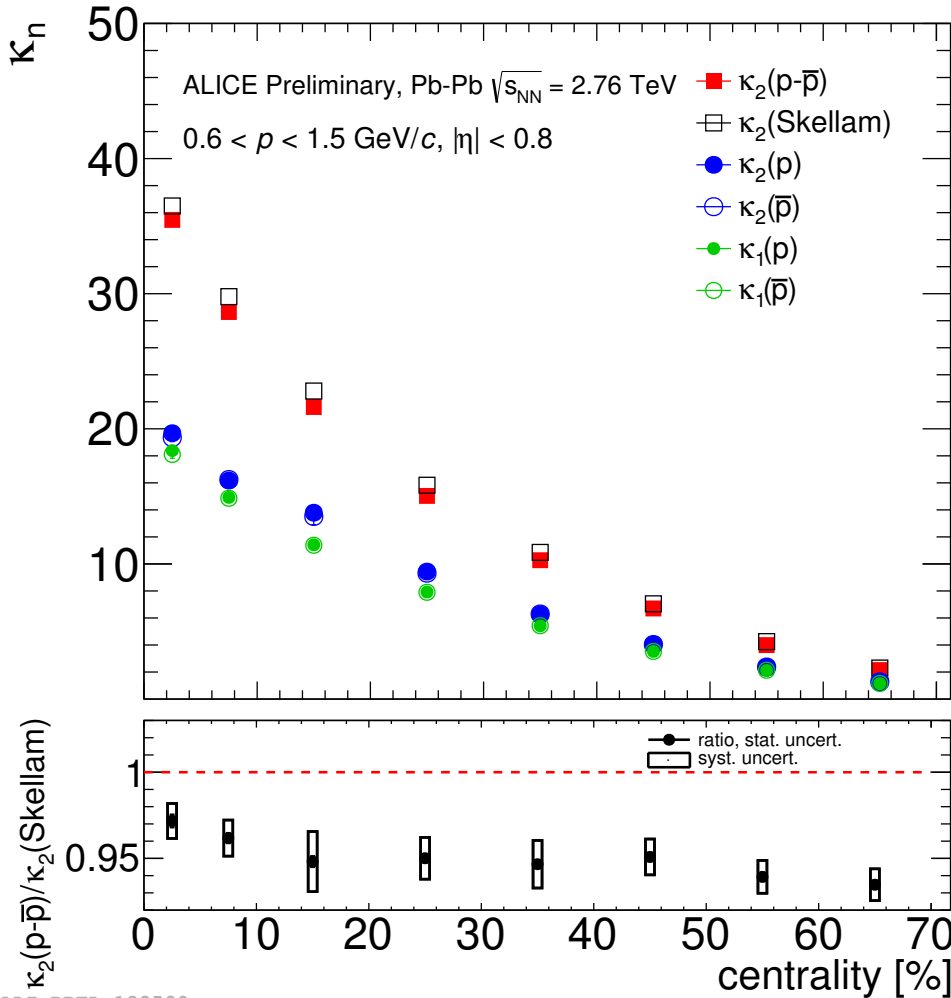
$$= \kappa_2(p) + \kappa_2(\bar{p}) - 2 \underbrace{\left( \langle N_p N_{\bar{p}} \rangle - \langle N_p \rangle \langle N_{\bar{p}} \rangle \right)}_{\text{correlation term}}$$

- If multiplicity distributions of protons and anti-protons are Poissonian and uncorrelated  
 $\rightarrow$  Skellam distribution for net-protons

$$\kappa_2(\text{Skellam}) = \kappa_1(p) + \kappa_1(\bar{p})$$

# Net-proton fluctuations

A. Rustamov for ALICE, QM2017



ALI-PREL-122590

$$\kappa_1(p) = \langle N_p \rangle \quad \kappa_2(p) = \left\langle \left( N_p - \langle N_p \rangle \right)^2 \right\rangle$$

$$\kappa_2(p - \bar{p}) = \left\langle \left( N_p - N_{\bar{p}} - \langle N_p - N_{\bar{p}} \rangle \right)^2 \right\rangle$$

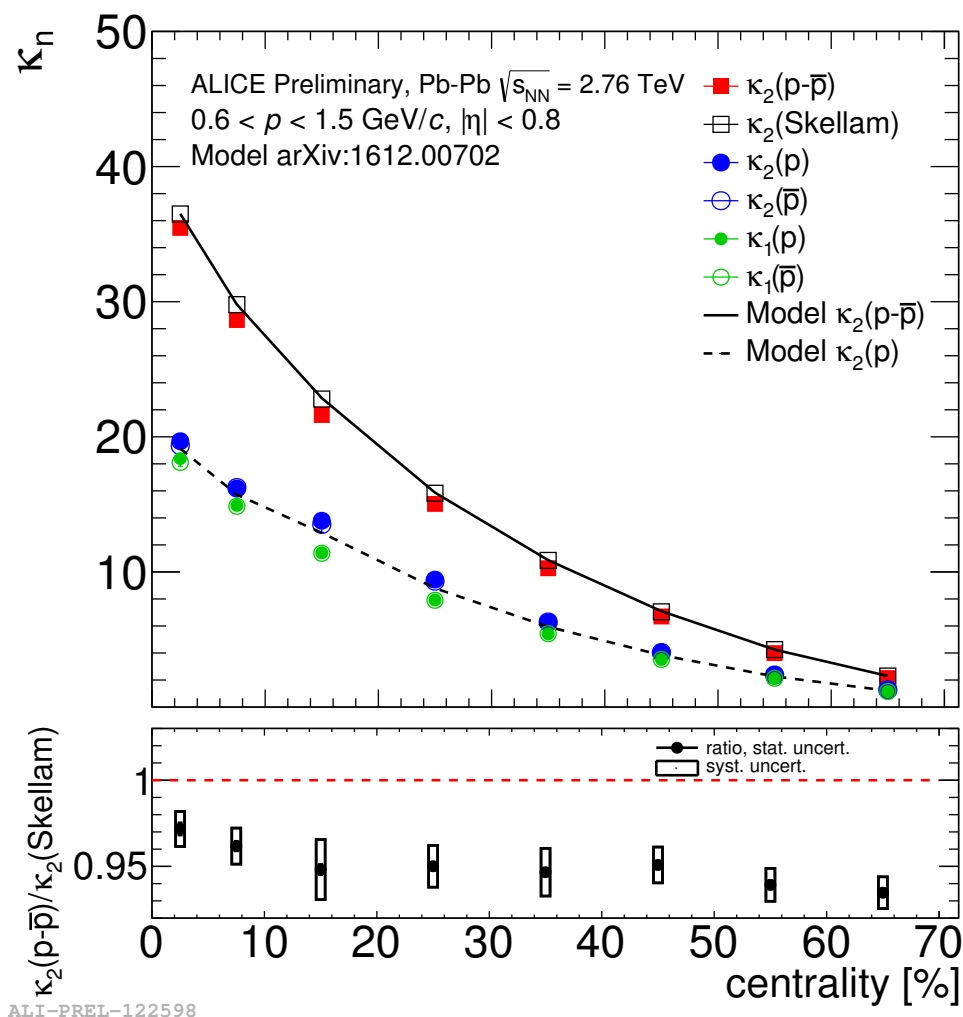
$$= \kappa_2(p) + \kappa_2(\bar{p}) - 2 \underbrace{\left( \langle N_p N_{\bar{p}} \rangle - \langle N_p \rangle \langle N_{\bar{p}} \rangle \right)}_{\text{correlation term}}$$

- $\kappa_2(p-\bar{p})$  shows deviation from Skellam prediction
  - due to correlation term?
  - are protons and anti-protons Poissonian?

$$\kappa_2(\text{Skellam}) = \kappa_1(p) + \kappa_1(\bar{p})$$

# Net-proton fluctuations

A. Rustamov for ALICE, QM2017

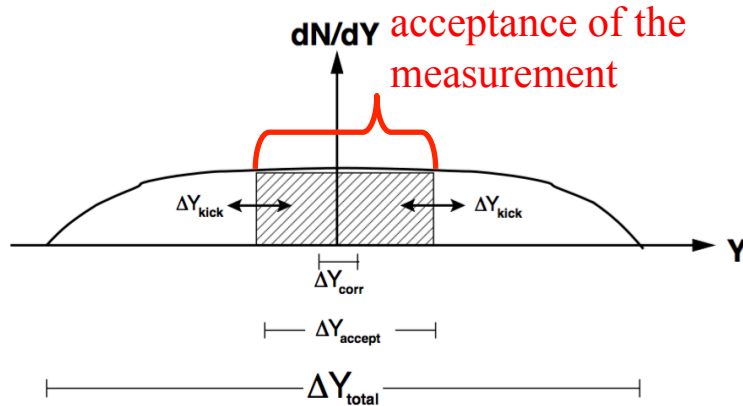


- Modeling the effects of participant fluctuations

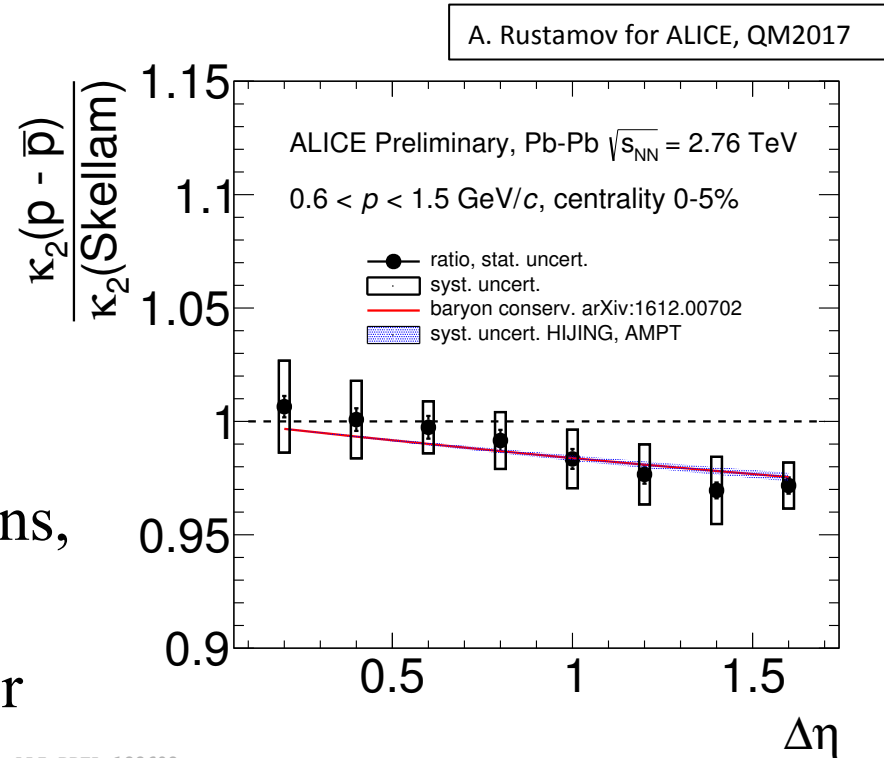
P. Braun-Munzinger et al., NPA 960 (2017)  
 114, arXiv:1612.00702 [nucl-th]

- Inputs to the model:  
 $\kappa_1(p)$ ,  $\kappa_1(\bar{p})$ , centrality determination procedure
- Model gives a consistent picture of  $\kappa_2(p)$ ,  $\kappa_2(\bar{p})$  and  $\kappa_2(p-\bar{p})$  without need of correlations or critical fluctuations

# Global conservation laws



- Small  $\Delta\eta \rightarrow$  Poissonian fluctuations, ratio to Skellam  $\sim 1$
- Large  $\Delta\eta \rightarrow$  global baryon number and strangeness conservation effects, ratio to Skellam  $< 1$
- $\Delta\eta$  dependence consistent with effects of baryon number conservation



P. Braun-Munzinger et al., NPA 960 (2017) 114, arXiv:1612.00702 [nucl-th]

# Global conservation laws

- Contribution from global baryon number conservation calculated as

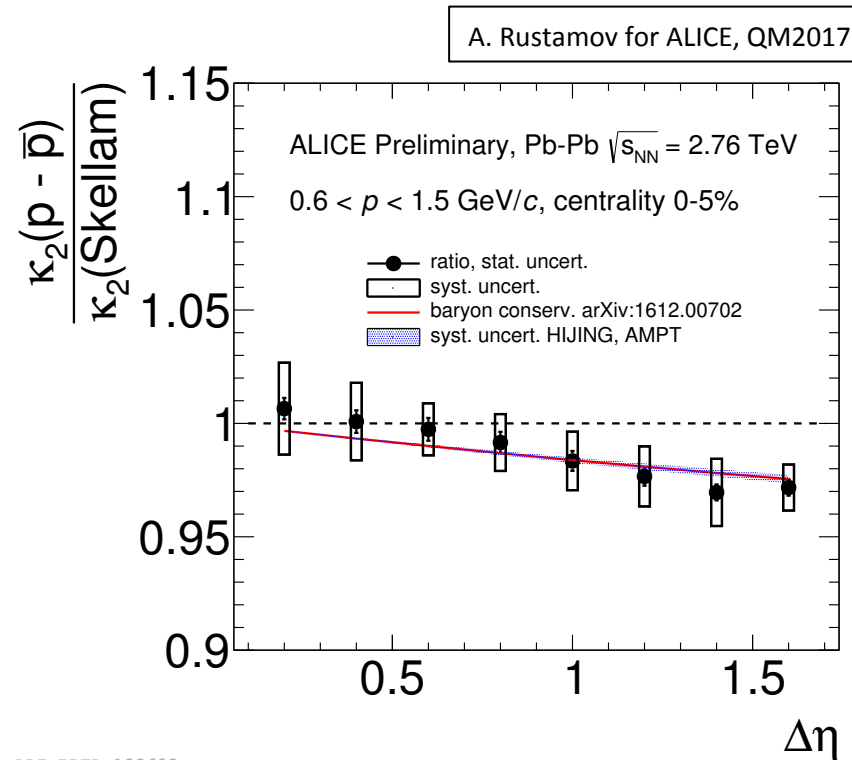
$$\frac{\kappa_2(p - \bar{p})}{\kappa_2(\text{Skellam})} = 1 - \frac{\langle N_p^{meas} \rangle}{\langle N_B^{4\pi} \rangle} = 1 - \alpha$$

- Inputs for  $\langle N_B^{acc} \rangle$  from

P. Braun-Munzinger et al., PLB 747 (2015) 292,  
arXiv:1412.8614 [hep-ph]

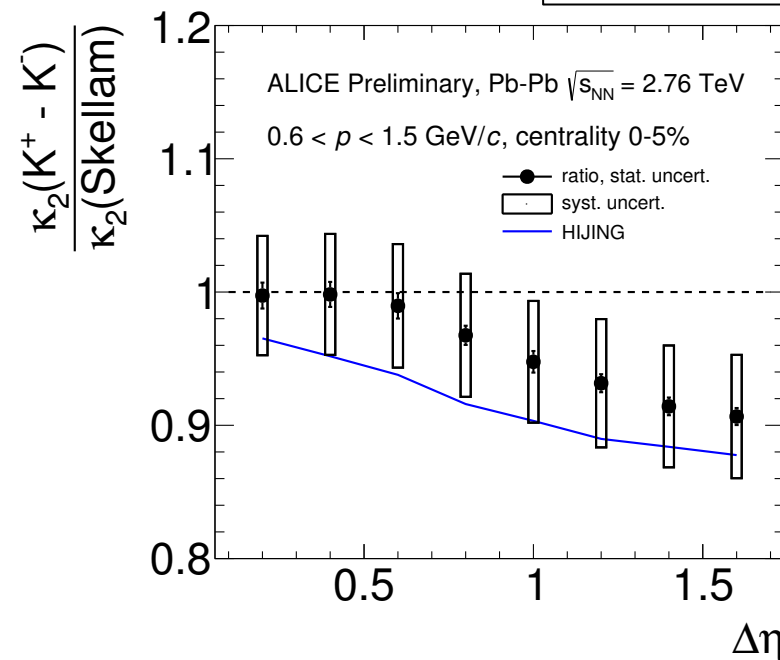
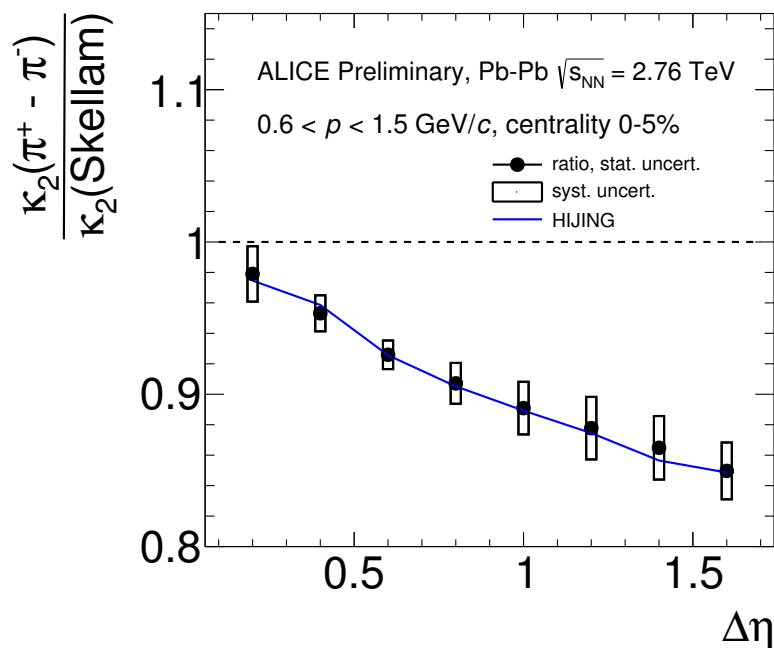
Extrapolation from  $\langle N_B^{acc} \rangle$  to  $\langle N_B^{4\pi} \rangle$  using AMPT and HIJING

- Deviation from Skellam baseline accounted for by global baryon number conservation



# Net-pion and net-kaon fluctuations

A. Rustamov for ALICE, QM2017



ALI-PREL-122614

ALI-PREL-122618

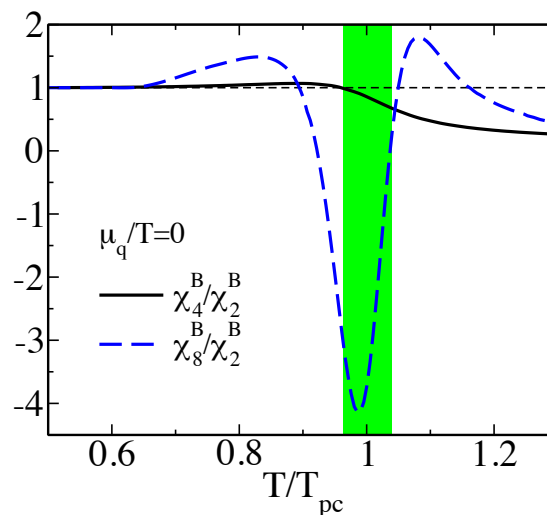
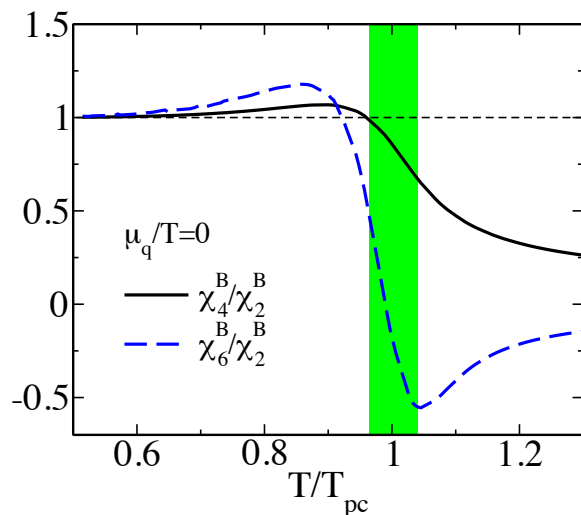
- Pions show good agreement with HIJING
- Production of pions and kaons from resonance decays contributes significantly to the measurement
- Skellam distribution is not a proper baseline for net-pions and net-kaons

# Higher moments from the LHC



# Higher moments

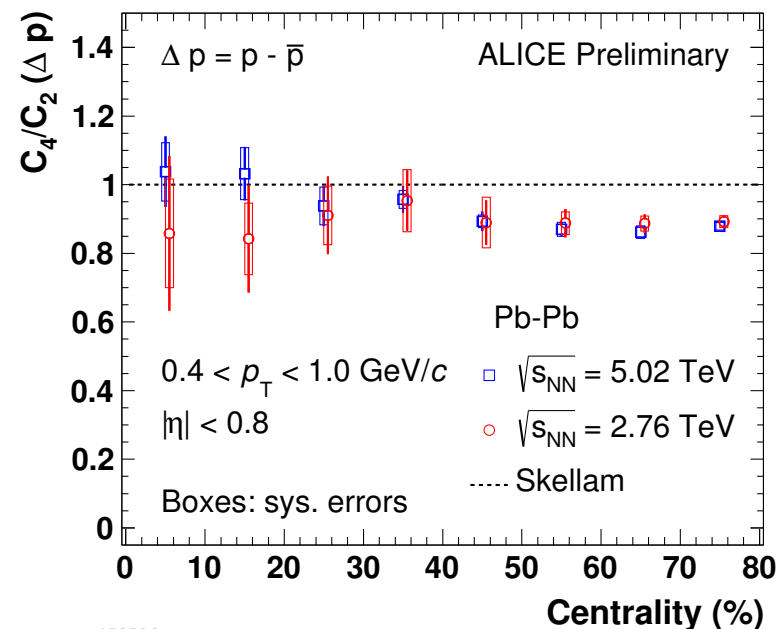
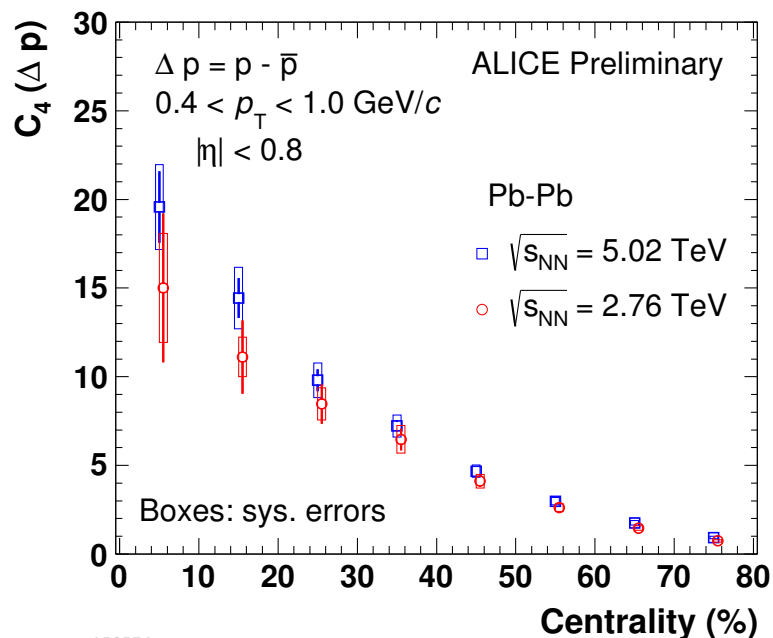
- Deviations from unity and signs of criticality are greatly enhanced for the higher moments (4<sup>th</sup>, 6<sup>th</sup>, 8<sup>th</sup>,...)



Friman, B., et al. Eur. Phys. J. C 71 (2011) 1694, arXiv:1103.3511 [hep-ph]

- But huge statistics are needed and experimental effects must be carefully controlled

# First higher moments from ALICE



ALI-PREL-159574

ALI-PREL-159586

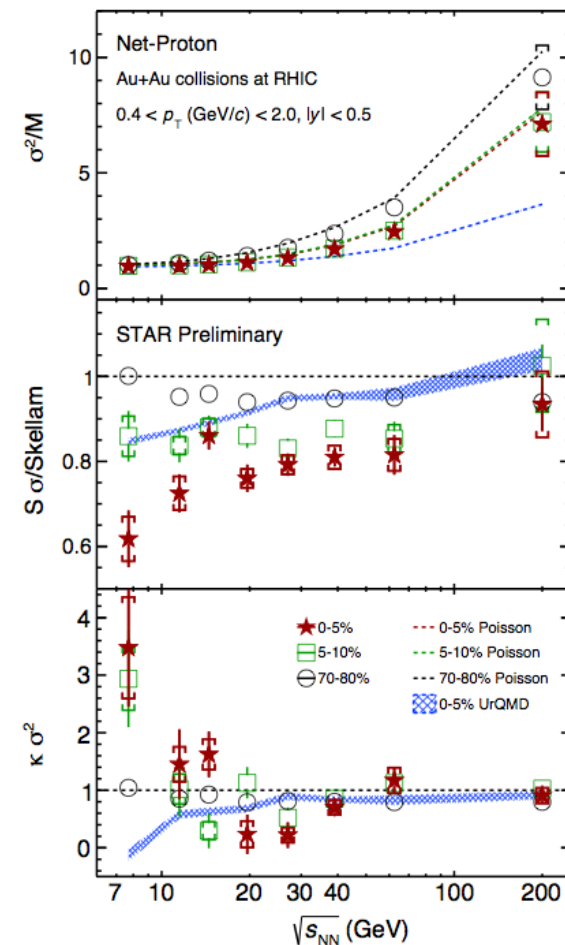
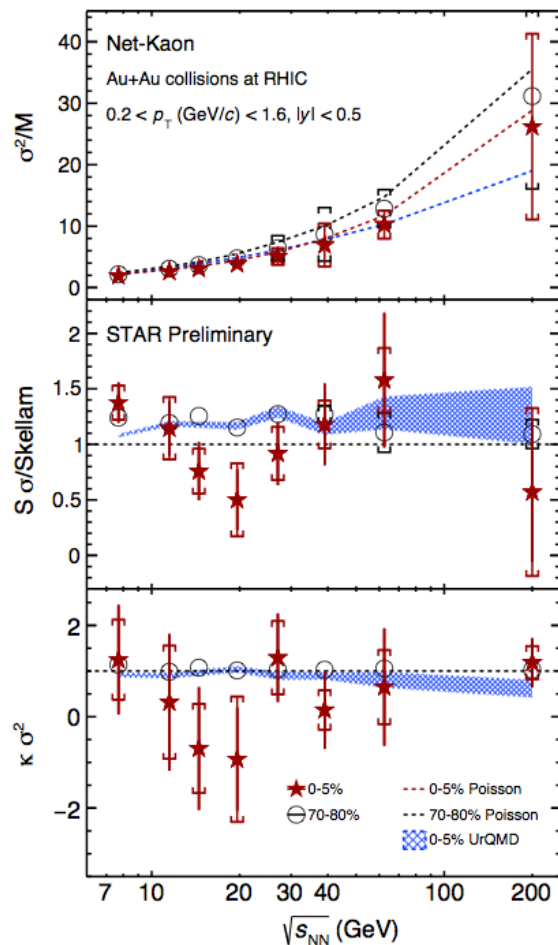
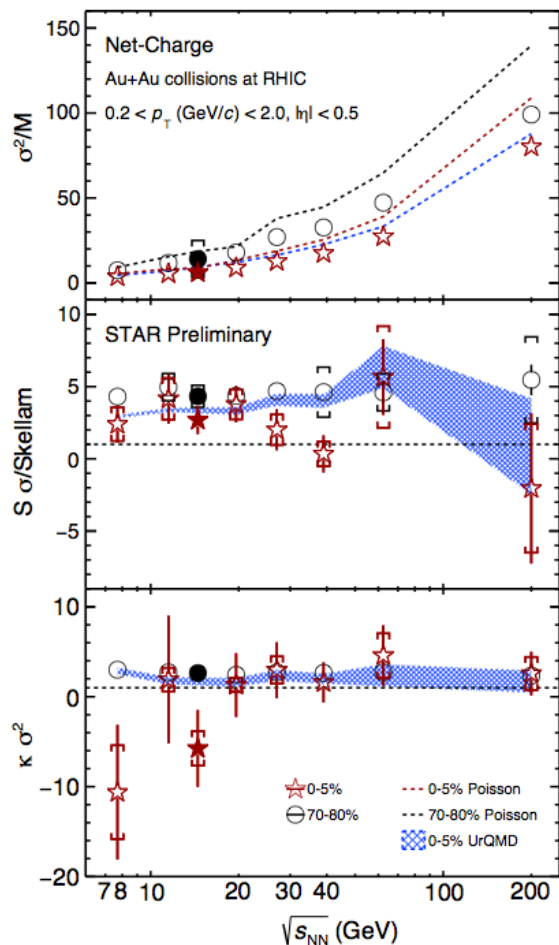
- Consistent results between  $\sqrt{s_{NN}} = 2.76 \text{ TeV}$  and  $5.02 \text{ TeV}$  within statistical and systematic uncertainties
- In central events, consistency with Skellam baseline ( $C_4/C_2 = 1$ )
- Higher statistics and improved understanding of systematics are needed to obtain the precision needed for LQCD comparisons

N. Behera for ALICE, QM2018

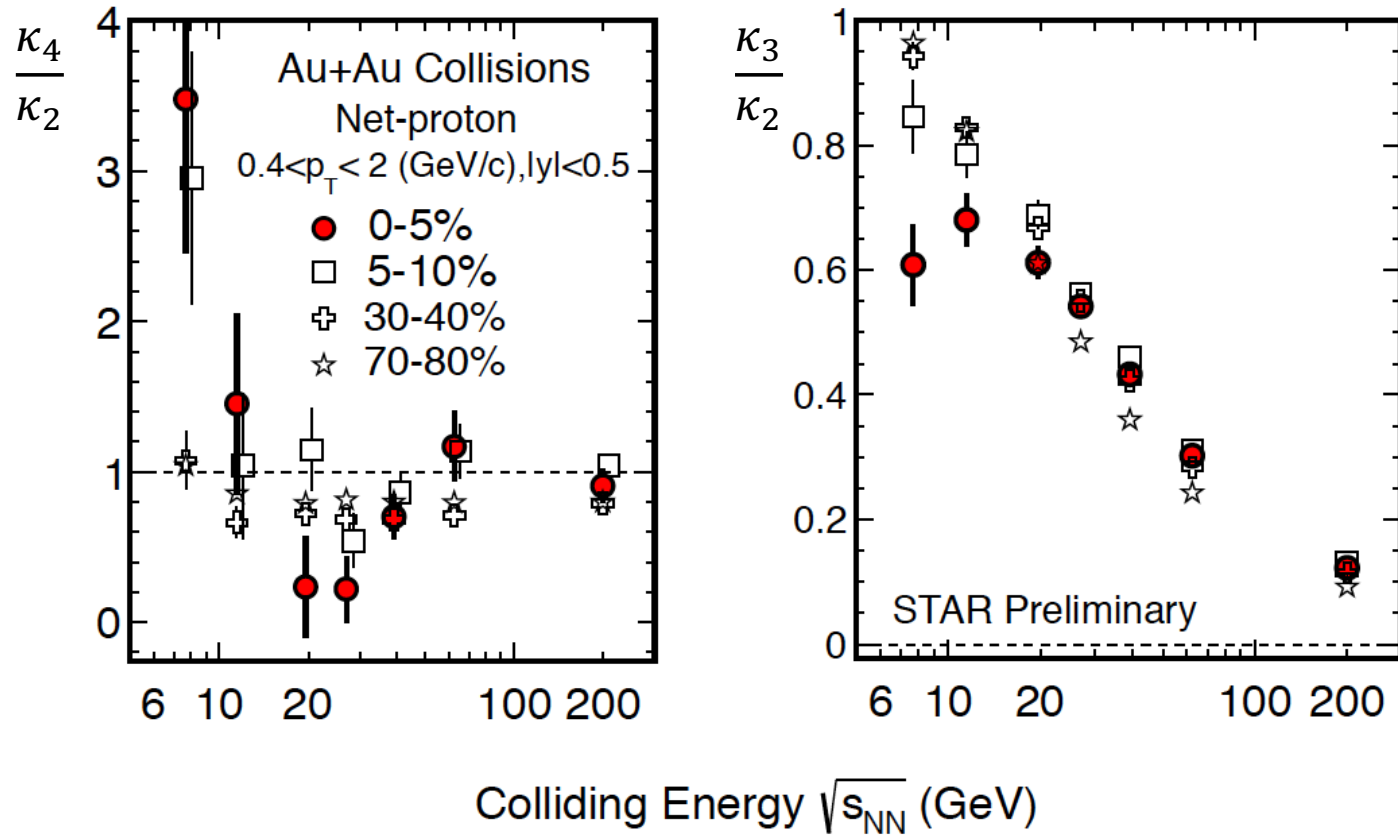
# Higher moments at RHIC

# STAR results: net-charge, net-K, net-p

J. Thäder for STAR, QM2015



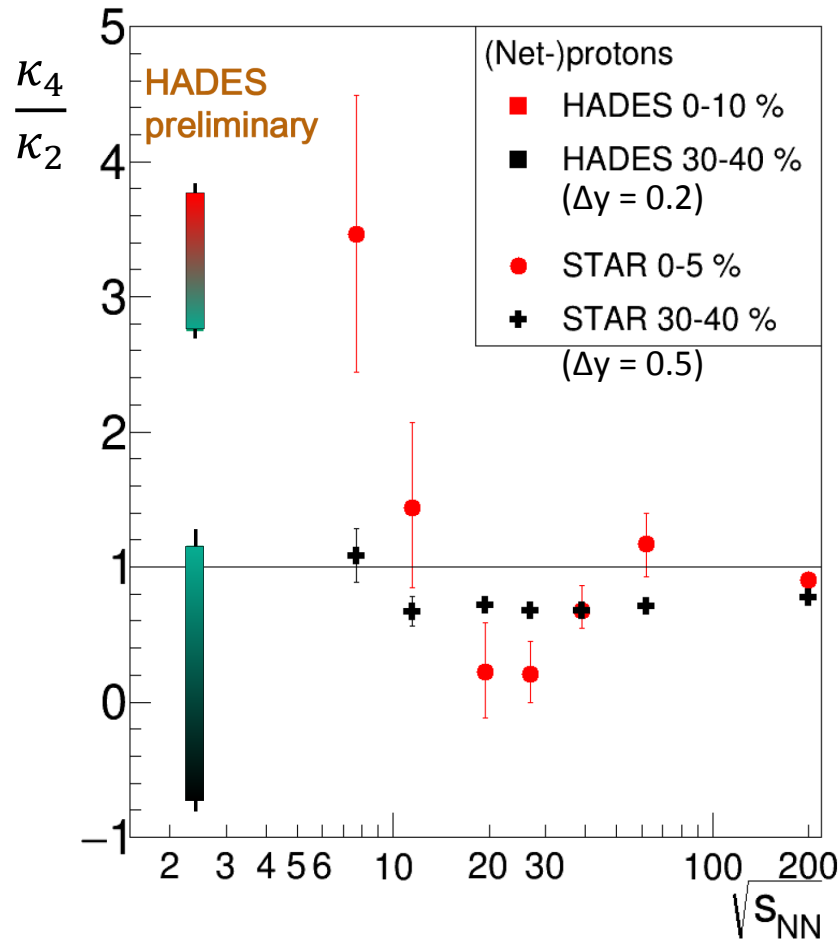
# STAR results: net-protons in the BES



- non-monotonic behavior observed below  $\sqrt{s_{NN}} = 39$  GeV

X. Luo, PoS CPOD2014 (2015) 019  
STAR, PRL 112 (2014) 032302

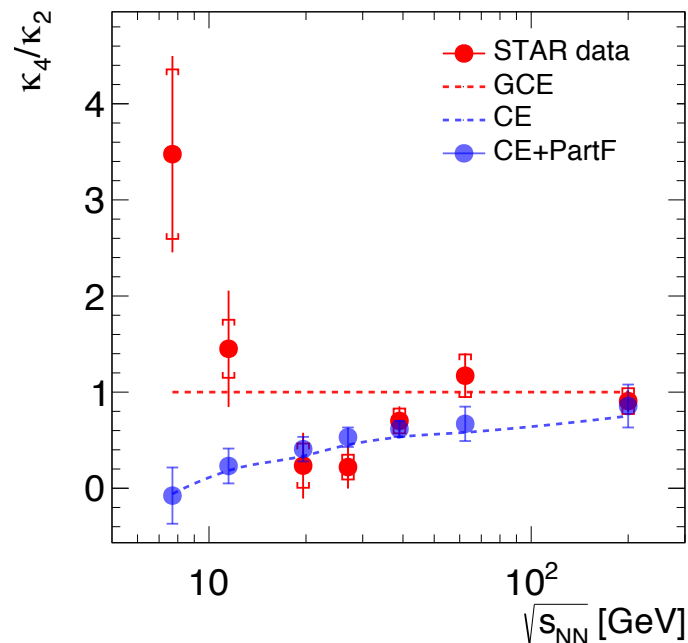
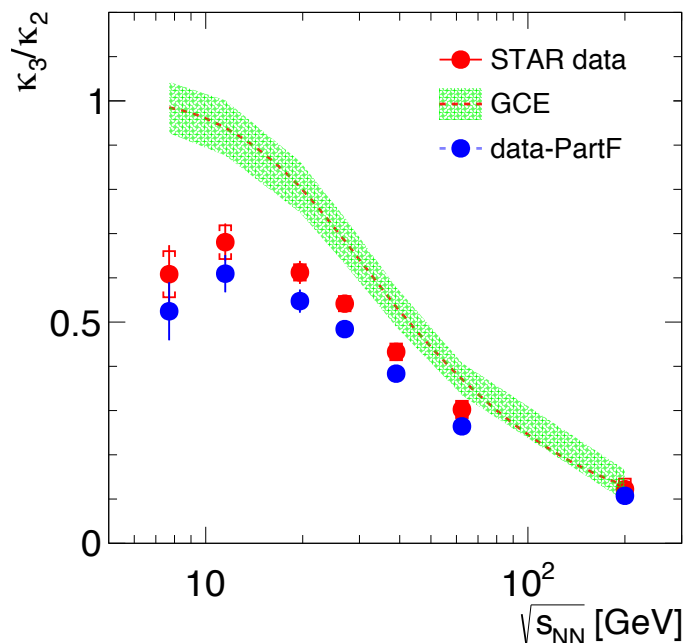
# STAR + HADES: net-protons vs $\sqrt{s_{NN}}$



T. Galatyuk, CPOD 2018

- different correction methods:
  - unfolding + volume fluctuation correction
  - E-by-E correction of factorial moments + vol. fluct. corr.
 → large differences in results (still under investigation)

# Effects of conservation laws + vol. fluct.



- At RHIC, proton and anti-proton multiplicities not equal

$$\frac{\kappa_3(n_p - n_{\bar{p}})}{\kappa_2(n_p - n_{\bar{p}})} = \frac{\langle n_p - n_{\bar{p}} \rangle_{CE}}{\langle n_p + n_{\bar{p}} \rangle_{CE}} (1 - 2\alpha)$$

$$\frac{\kappa_4(n_B - n_{\bar{B}})}{\kappa_2(n_B - n_{\bar{B}})} = 1 - 6\alpha(1 - \alpha) \left[ 1 - \frac{2}{\langle N_B + N_{\bar{B}} \rangle_{CE}} \left( \langle N_B \rangle_{GCE} \langle N_{\bar{B}} \rangle_{GCE} - \langle N_B \rangle_{CE} \langle N_{\bar{B}} \rangle_{CE} \right) \right]$$

P. Braun-Munzinger, A. Rustamov,  
J. Stachel, arXiv:1807.08927 [nucl-th]

- Above  $\sqrt{s_{NN}} = 11.5$  GeV: deviation from unity can be described by global baryon number conservation

# Net- $\Lambda$ moments

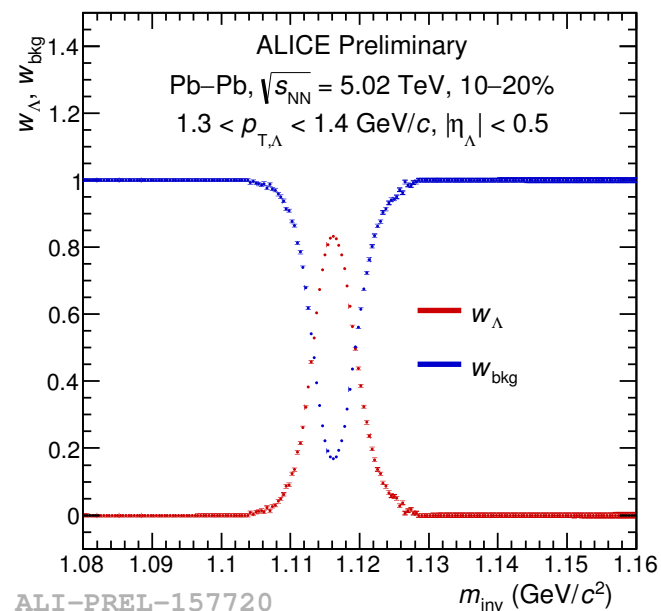
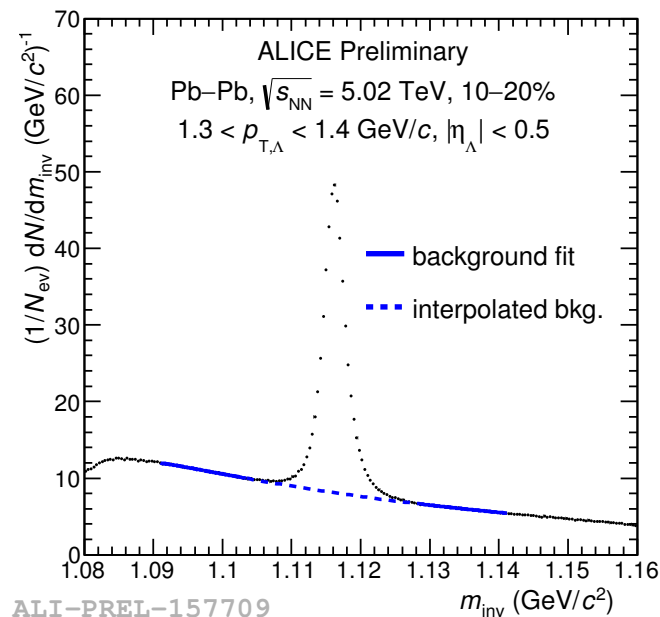


# From net- $\pi$ , $K$ , $p$ to net- $\Lambda$ moments

- Explore correlated fluctuations of baryon number and strangeness
- Establish baseline for future measurements of higher moments in the strangeness sector
- Improve understanding of net-baryon fluctuations
  - different contributions from resonances, etc, than in net-proton measurement
- $\Lambda$ s can be “added” to net-proton or net-kaon results to get closer to net-baryon and net-strangeness fluctuations

# Identity method for invariant mass

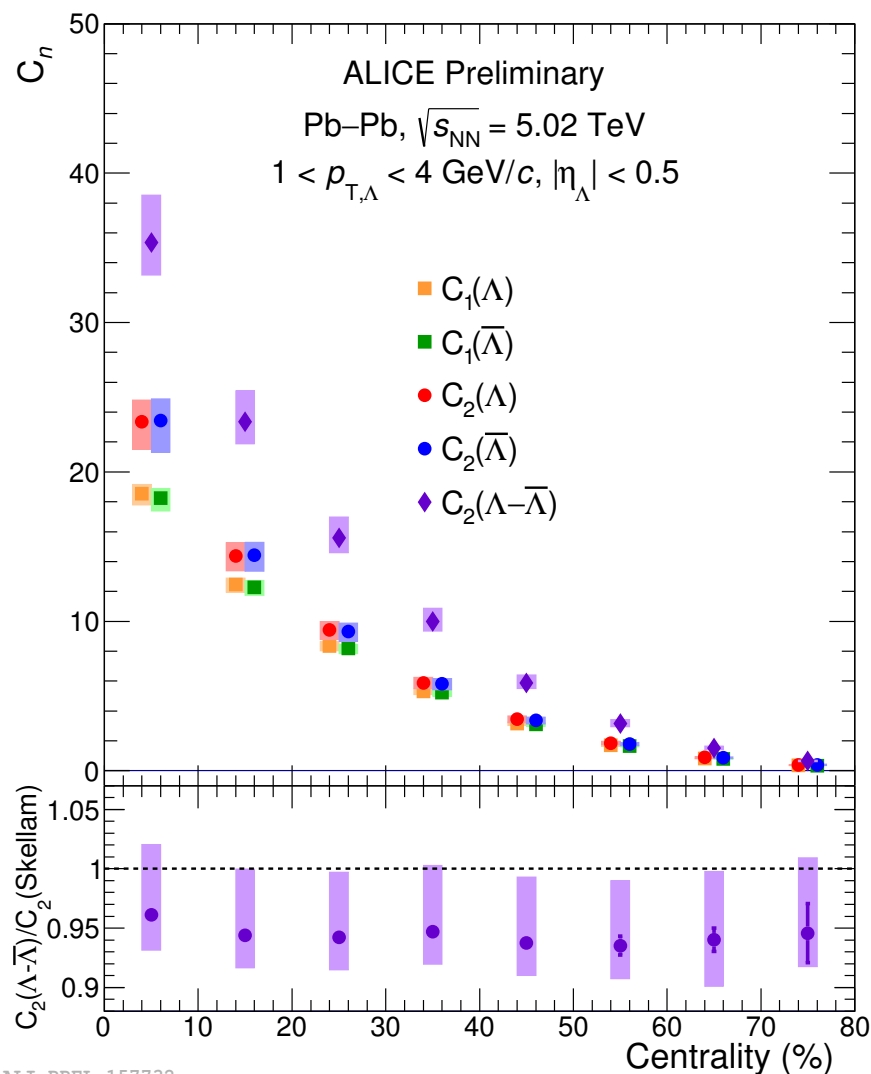
- For any value of  $m_{\text{inv}}$ , the probability that a  $\pi p$  pair comes from the decay of a  $\Lambda$  baryon is known
- Apply Identity Method for four “species”:  
 $\Lambda$ ,  $\bar{\Lambda}$ , combinatoric  $\pi p$ , combinatoric  $\pi^+ \bar{p}$



A. Ohlson for ALICE, QM2018

# Net- $\Lambda$ fluctuations

A. Ohlson for ALICE, QM2018



ALI-PREL-157732

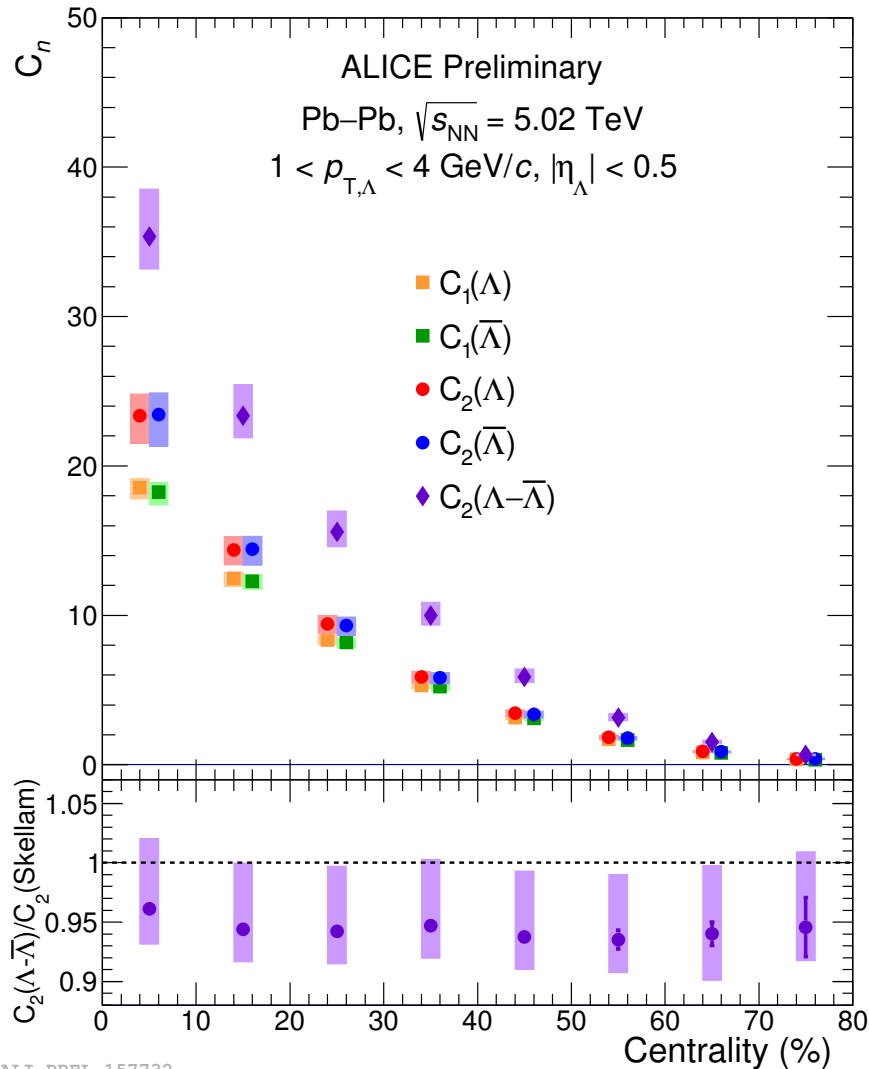
$$C_1(\Lambda) = \langle N_\Lambda \rangle \quad C_2(\Lambda) = \left\langle \left( N_\Lambda - \langle N_\Lambda \rangle \right)^2 \right\rangle$$

$$C_2(\Lambda - \bar{\Lambda}) = \left\langle \left( N_\Lambda - N_{\bar{\Lambda}} - \langle N_\Lambda - N_{\bar{\Lambda}} \rangle \right)^2 \right\rangle$$

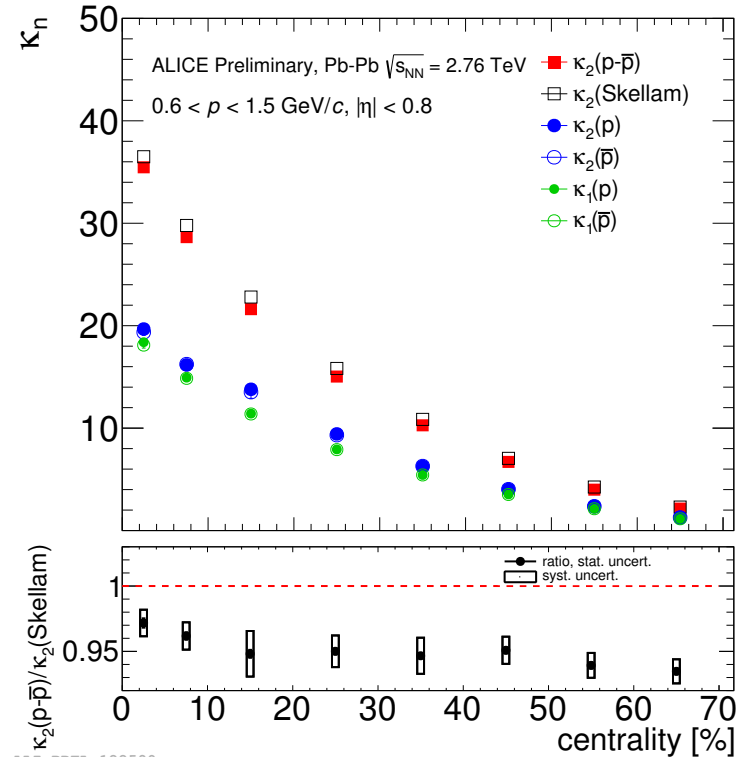
$$= C_2(\Lambda) + C_2(\bar{\Lambda}) - 2 \left( \langle N_\Lambda N_{\bar{\Lambda}} \rangle - \langle N_\Lambda \rangle \langle N_{\bar{\Lambda}} \rangle \right)$$

- Small deviations from Skellam baseline  $\rightarrow$  correlation term? non-Poissonian  $\Lambda$  or  $\bar{\Lambda}$  distributions? critical fluctuations? effects of volume fluctuations and global conservation laws?

# Comparison to net-protons



ALI-PREL-157732

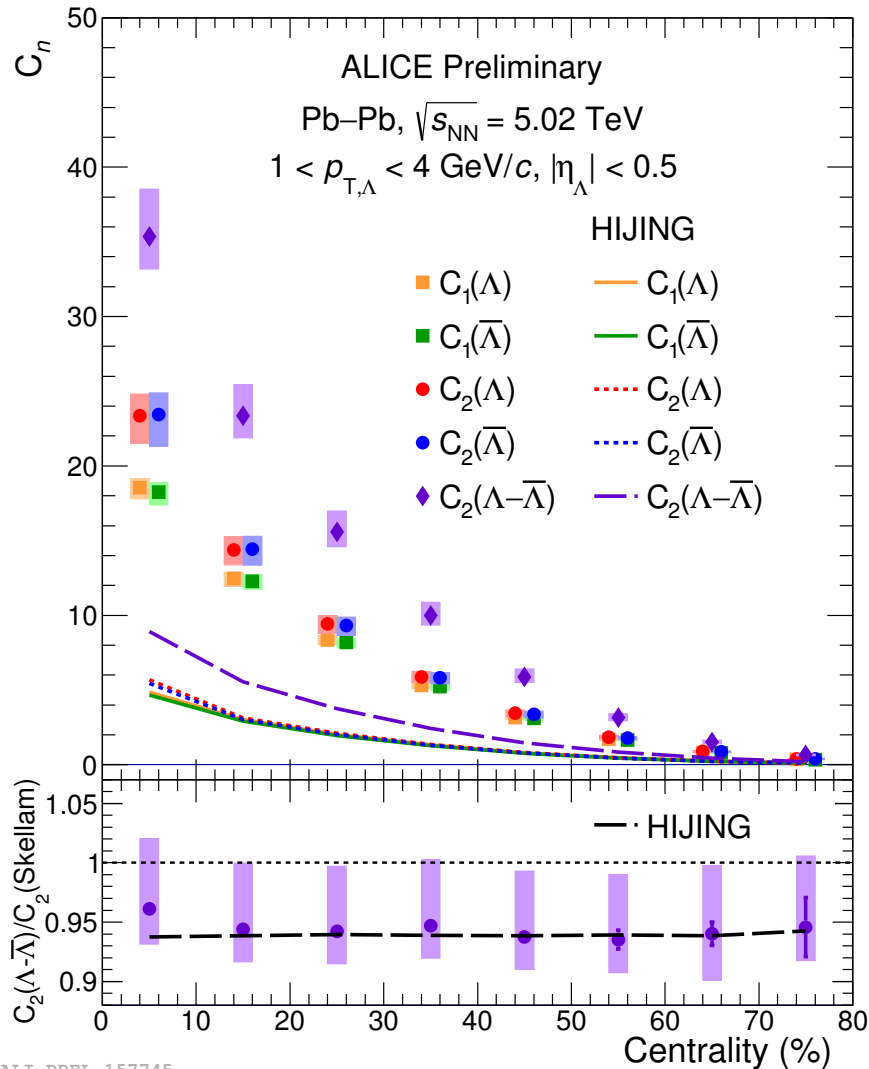


ALI-PREL-122590

- Qualitatively similar results for net-protons
  - different kinematic range,
  - different contributions from resonance decays

# Comparison to HIJING

A. Ohlson for ALICE, QM2018

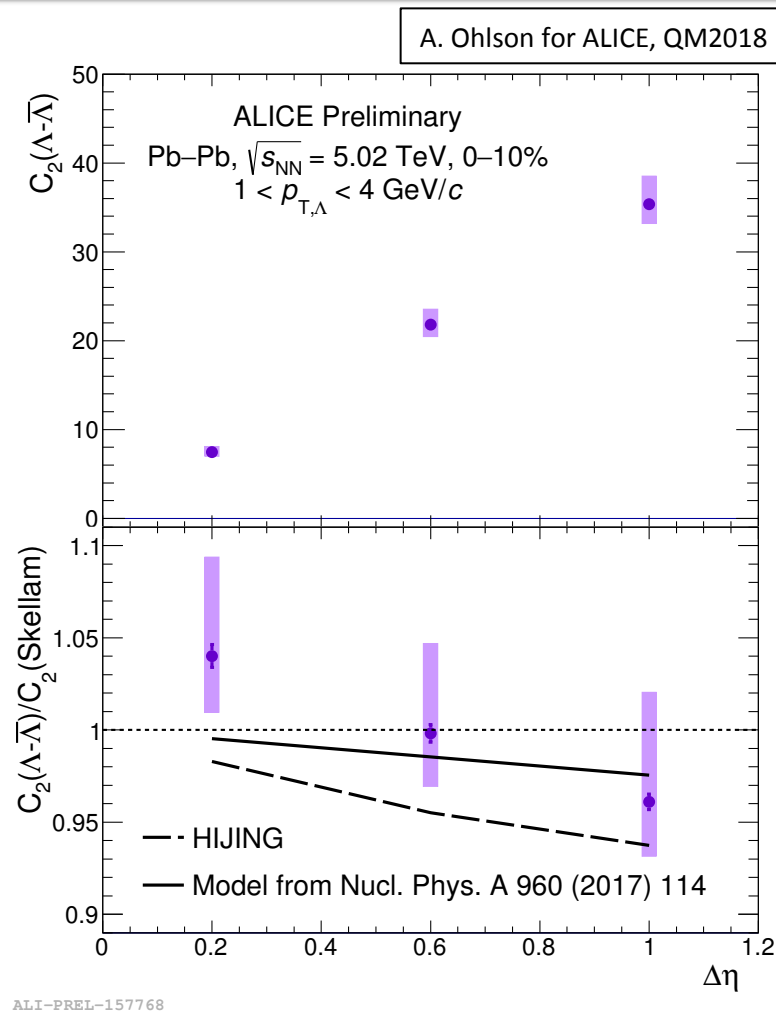


ALI-PREL-157745

- HIJING does not describe strangeness production well
  - underestimates  $C_1$  and  $C_2$  by factor  $\sim 4$
- However,  $C_2(\Lambda-\bar{\Lambda})/C_2(\text{Skellam})$  ratio agrees with data
  - coincidence? or due to description of fluctuations and resonance contributions in HIJING?

# $\Delta\eta$ dependence of net- $\Lambda$ fluctuations

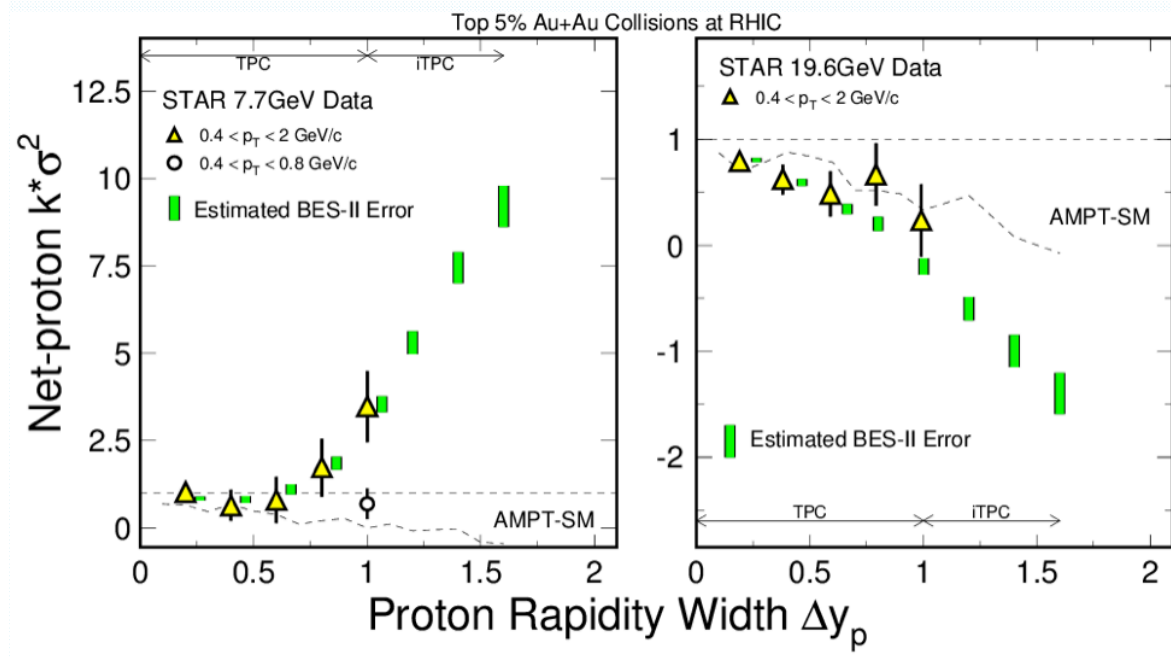
- Small  $\Delta\eta \rightarrow$  Poissonian fluctuations, ratio to Skellam  $\sim 1$
- Large  $\Delta\eta \rightarrow$  global baryon number and strangeness conservation effects, ratio to Skellam  $< 1$
- Systematic uncertainties are highly correlated point-to-point
- $\Delta\eta$  dependence consistent with effects of baryon number conservation  $\rightarrow$  strangeness conservation should also be considered
- consistency also with HIJING



# Outlook

- Runs 3+4 at the LHC will allow us to measure the fourth and sixth moments of the net-proton distribution with unprecedented precision
- BES-II + detector upgrades at RHIC will allow us to probe fluctuations across a wide range of the phase diagram

LHC Yellow Report: arXiv:1812.06772 [hep-ph]



<https://drupal.star.bnl.gov/STAR/starnotes/public/sn0619>

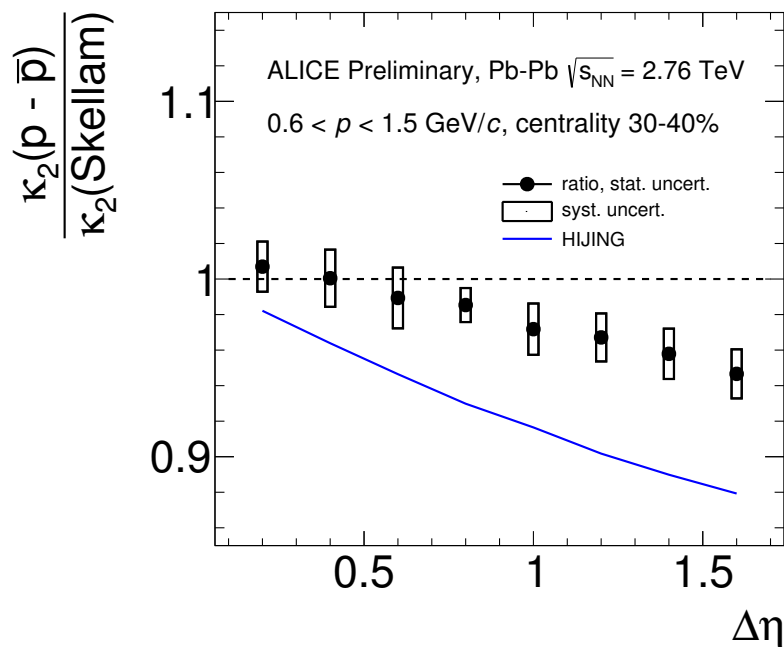
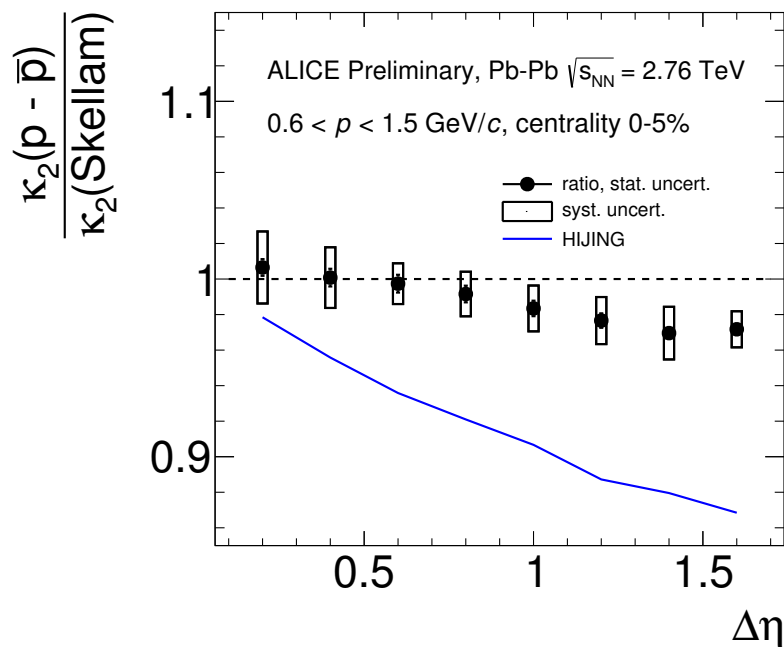
# Conclusions

- Event-by-event fluctuations of identified particles
  - yield information on properties of the QGP medium
  - test lattice QCD predictions at  $\mu_B = 0$
  - allow us to look for effects of criticality
- Effects of detector inefficiency and particle misidentification being brought under control
- Effects of volume fluctuations and global baryon number conservation are assessed
- Net-proton and net- $\Lambda$  fluctuations at LHC energies: no deviations from Skellam baseline observed after accounting for baryon number conservation, agreement with LQCD predictions
- Net-proton fluctuations at RHIC energies: can be described above  $\sqrt{s_{NN}} = 11.5$  GeV by baryon number conservation



backup

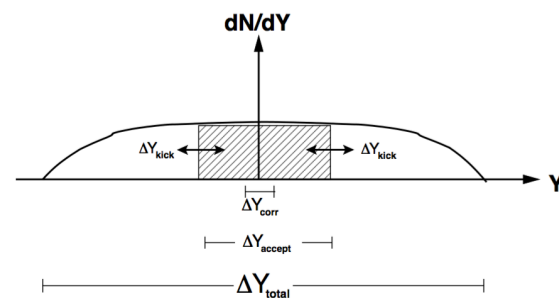
# Centrality dependence



ALI-PREL-122606

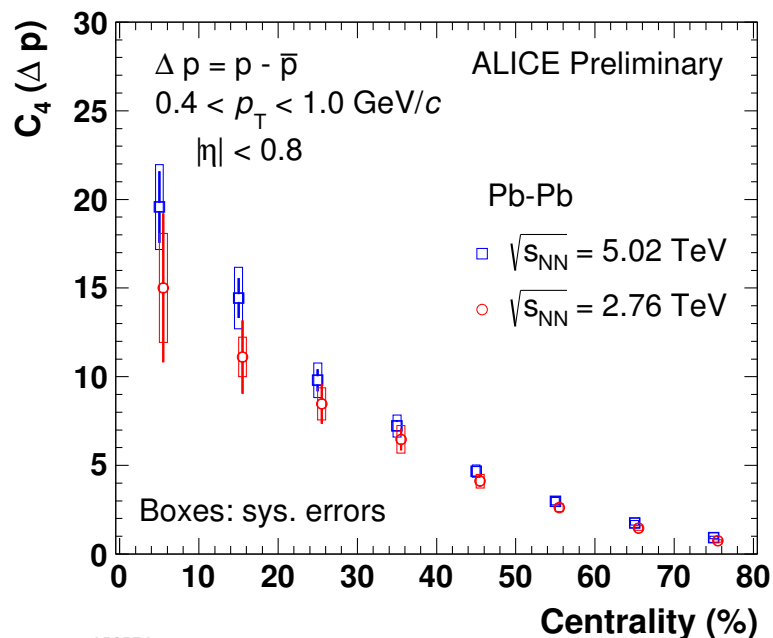
ALI-PREL-122610

- Deviations from Skellam can be attributed global baryon number conservation, more significant in more peripheral collisions
- Disagreement with HIJING

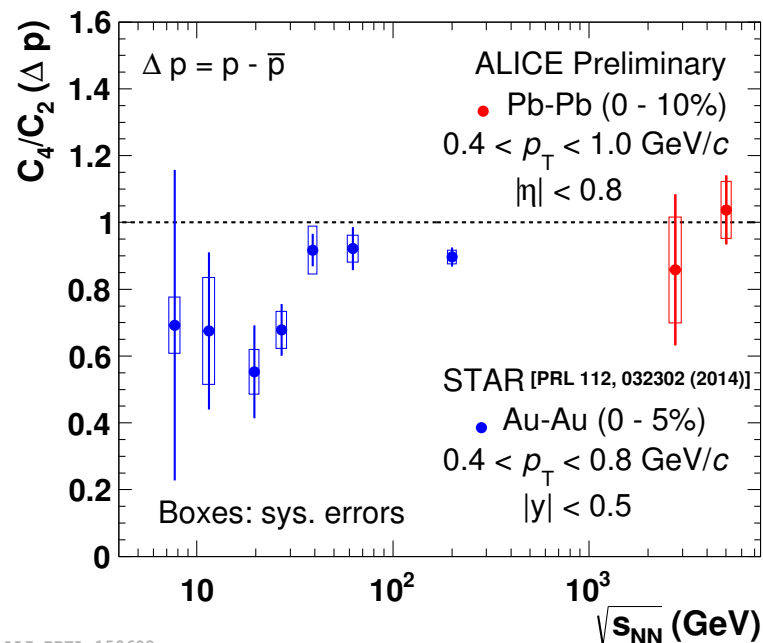


A. Rustamov for ALICE, QM2017

# First higher moments from ALICE!



ALI-PREL-159574



ALI-PREL-159602

- Measured with traditional (cut-based) PID method
- Consistent results between  $\sqrt{s_{NN}} = 2.76$  TeV and 5.02 TeV within statistical and systematic uncertainties
- In central events, consistency with Skellam baseline ( $C_4/C_2 = 1$ ) at LHC energies

N. Behera for ALICE, QM2018