

# Interplay of Bosonic and Fermionic Fluctuations at Finite Densities

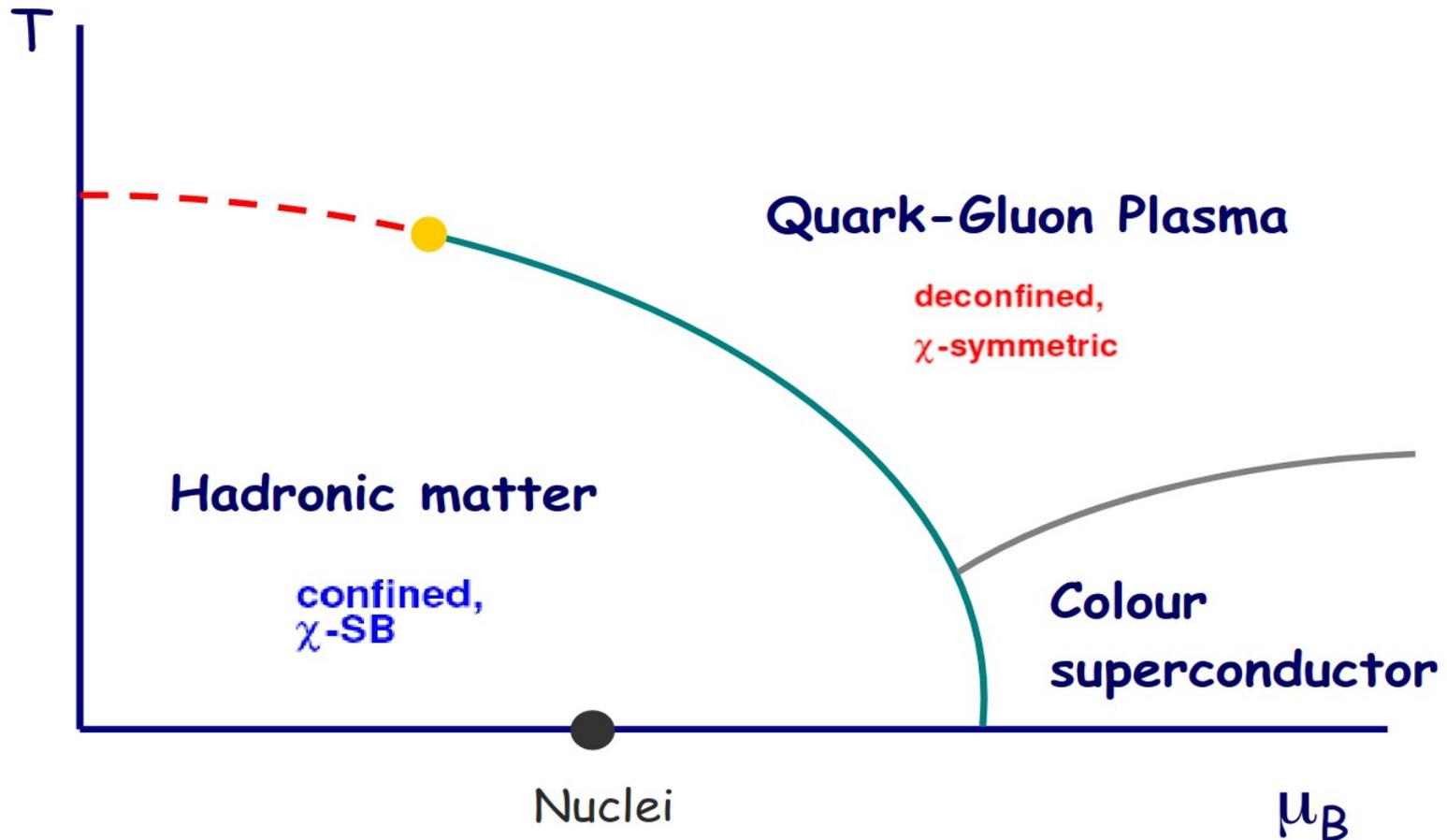
Christopher Busch

based on: Otto,Busch,Schaefer (2022), arXiv:2206.13067

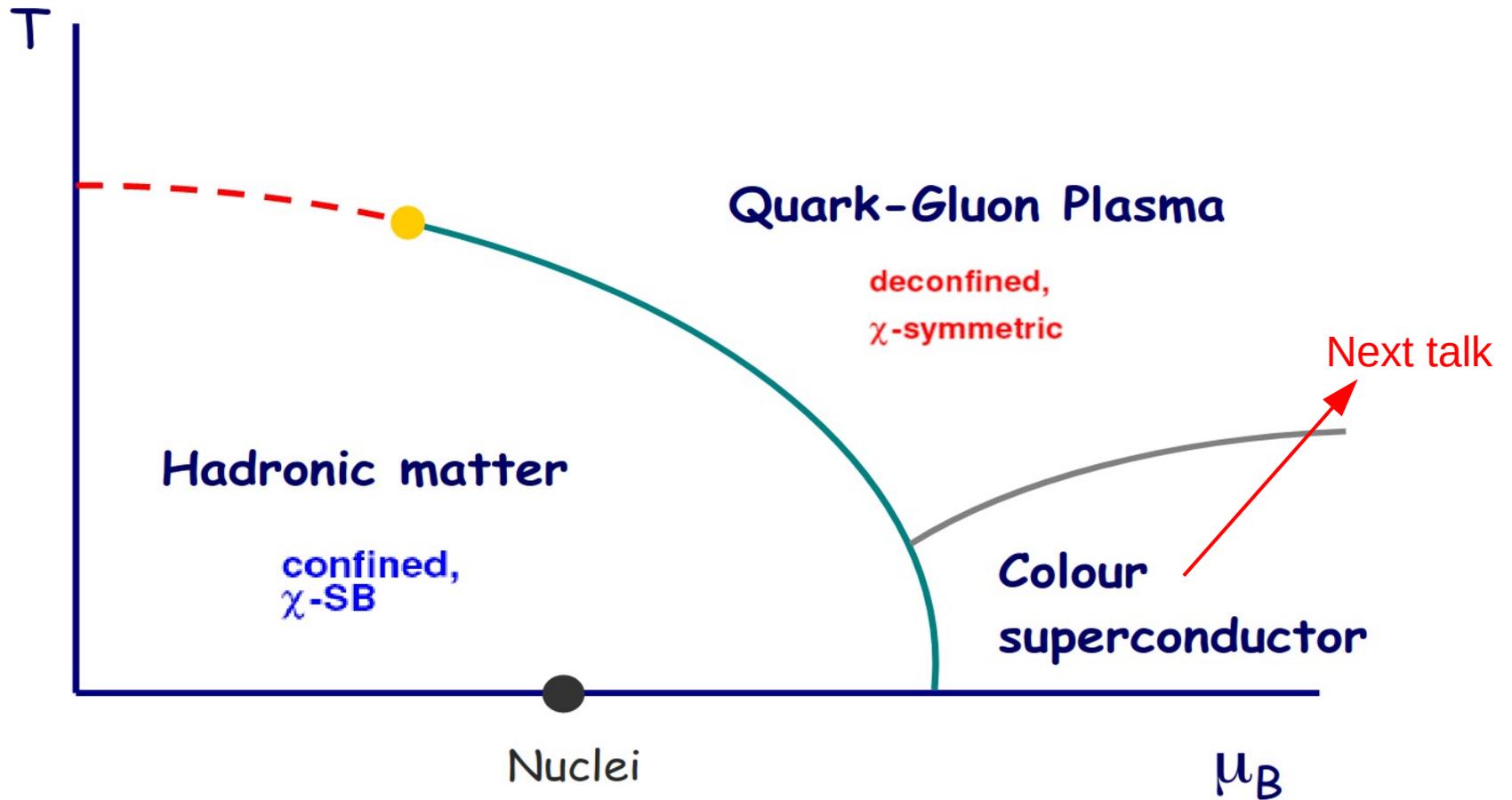
HFHF Theory Retreat – 16.09.2022



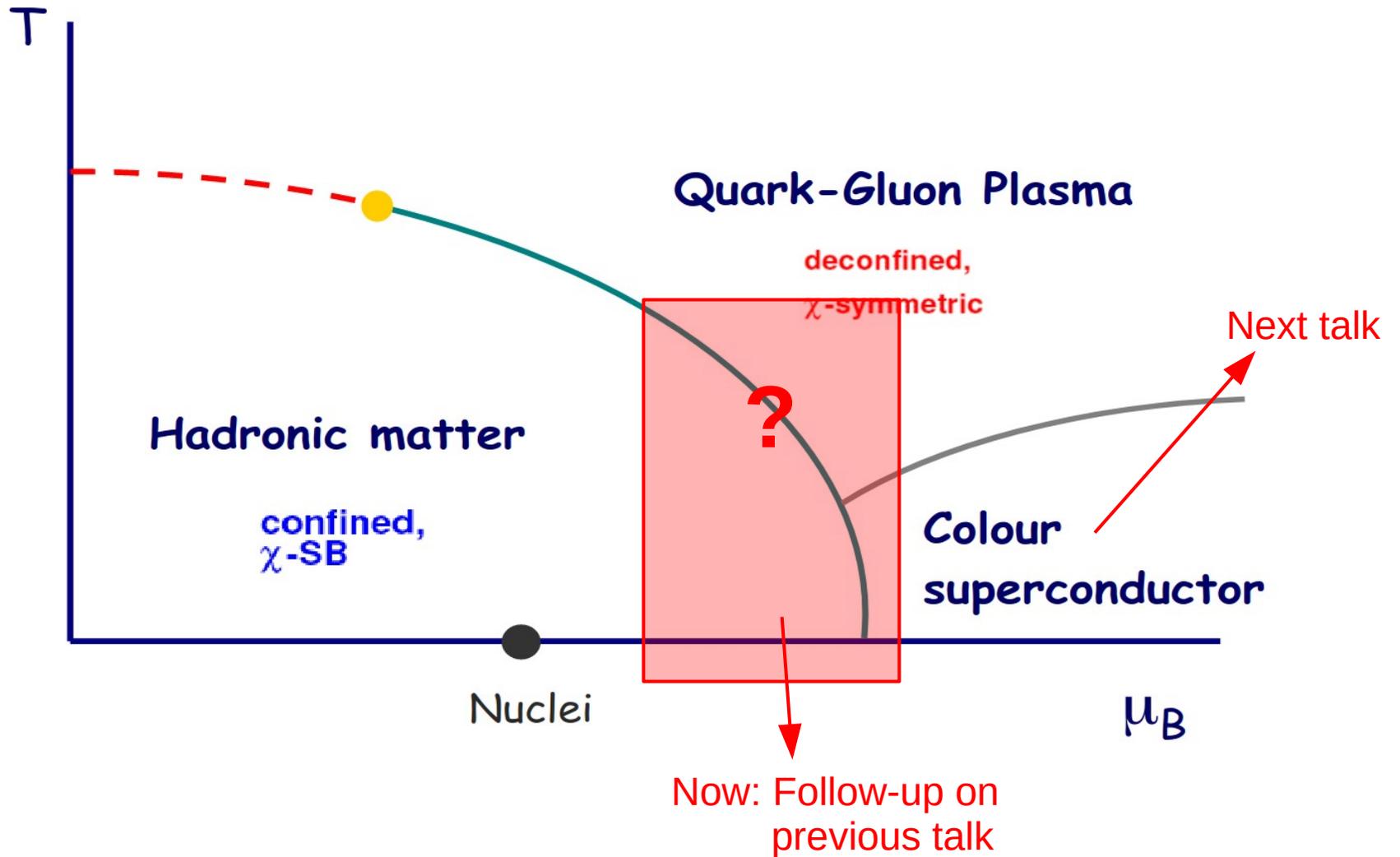
# Objectives



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# Outline

- 1) Scheme Dependence
- 2) Modifications of the LPA Setup
- 3) Results: Regulator Dependence at Large Densities
- 4) Summary & Outlook

# FRG and Truncation Errors

Wetterich equation:  $(\partial_t = k\partial_k)$

$$\partial_t \Gamma_k = \frac{1}{2} \text{STr} \left\{ (\partial_t R_k) \left( \Gamma_k^{(1,1)} + R_k \right)^{-1} \right\}$$

Requirements for regulators:

1.  $\lim_{p^2/k^2 \rightarrow \infty} R_k(p^2) = 0$
2.  $\lim_{p^2/k^2 \rightarrow 0} R_k(p^2) > 0$
3.  $\lim_{k \rightarrow \infty} R_k(p^2) = \infty$

➤ Solving Wetterich eq. for full theory not possible

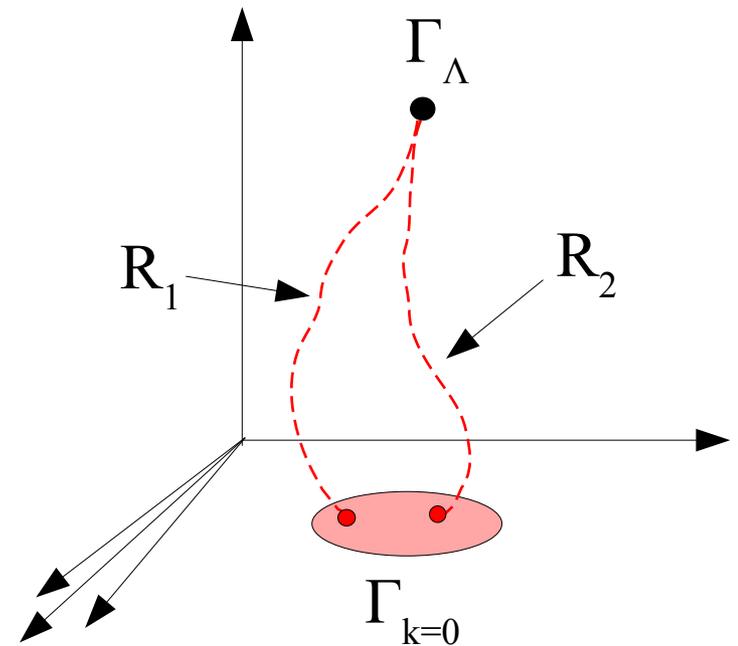
➔ Use of effective models / truncations

➤ Leads to truncation errors and regulator dependence of results

➔ Choice of regulator becomes relevant

↳ Optimization criteria for regulators

- Principle of minimum sensitivity
- “Gap Criterion” [Litim(2000)]
- “Shortest Path” [Pawlowski(2007)]



# Flow Equation

(for dim. reduced regulators)

➤ Quark-Meson model in Local Potential Approximation (LPA):

$$\Gamma_k = \int \left\{ \bar{\psi} [\not{\partial} - \mu\gamma_0 + g(\sigma + i\gamma_5 \vec{\tau}\vec{\pi})] \psi + \frac{1}{2} (\partial_\nu \phi)^2 + \mathbf{U}_k(\phi^2) - c\sigma \right\}$$

➤ General form of the (LPA) flow equation:

$$\tilde{m}_j := m_j/k$$

$$\partial_t \mathbf{U}_k = \frac{k^4}{4\pi^2} \left\{ l_0^B(\tilde{m}_\sigma^2) + 3l_0^B(\tilde{m}_\pi^2) - 4N_f N_c l_0^F(\tilde{m}_\psi^2) \right\}$$

- Splits into bosonic and fermionic loop contributions
- Threshold functions  $l_0^{F/B}$  comprise distribution functions
- Fermions at T=0:

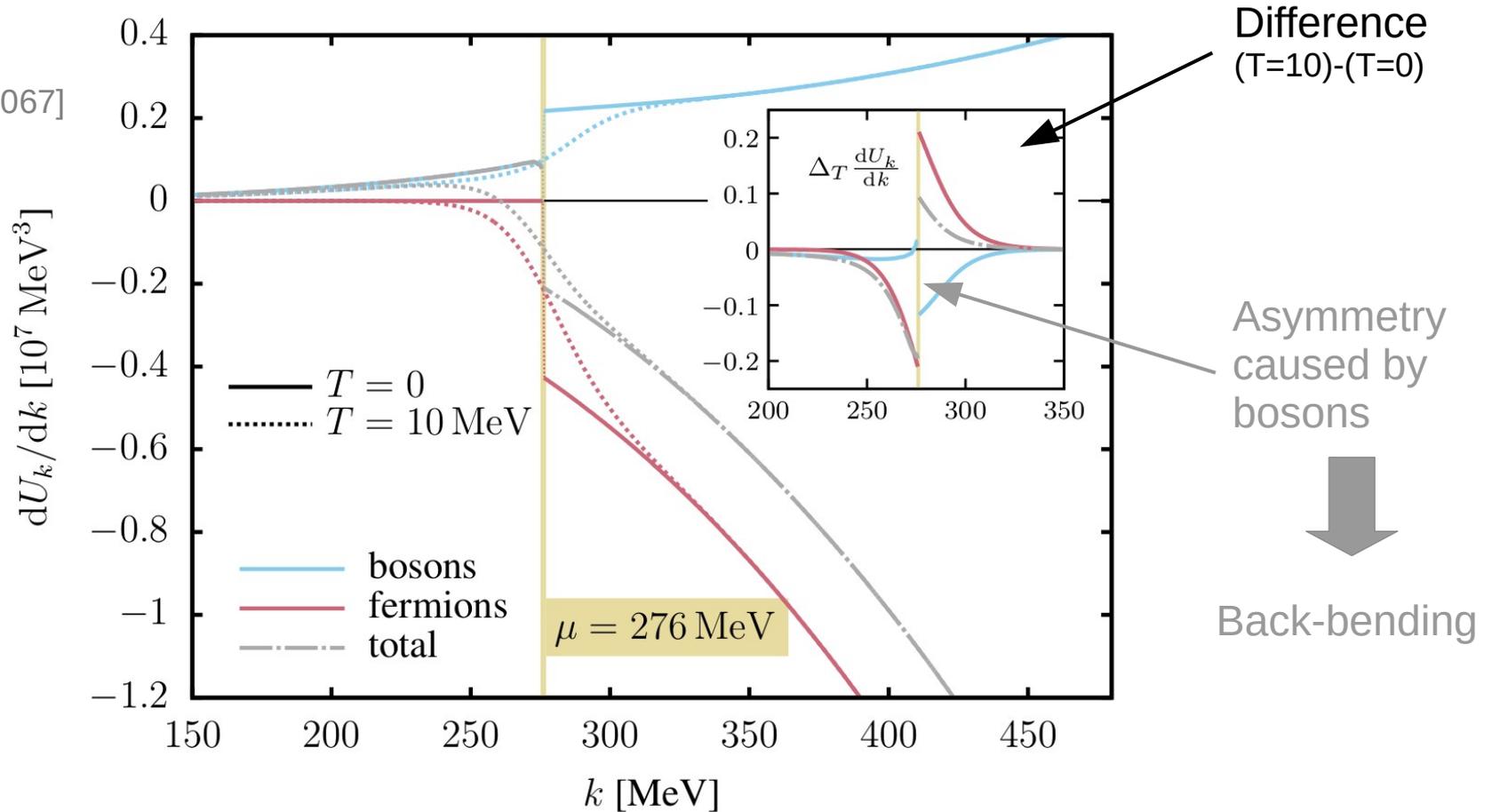
$$E_\psi(x) = \sqrt{x(1+r_\psi(x))^2 + \tilde{m}_\psi^2}$$

$$l_0^F(\tilde{m}_\psi^2)|_{T=0} = \int_0^\infty \frac{dx}{2} x^{3/2} \partial_t (1+r_\psi(x))^2 \frac{1}{2E_\psi(x)} \left\{ 1 - \Theta(\mu - E_\psi(x)) \right\}$$

Fermions decouple regulator dependent

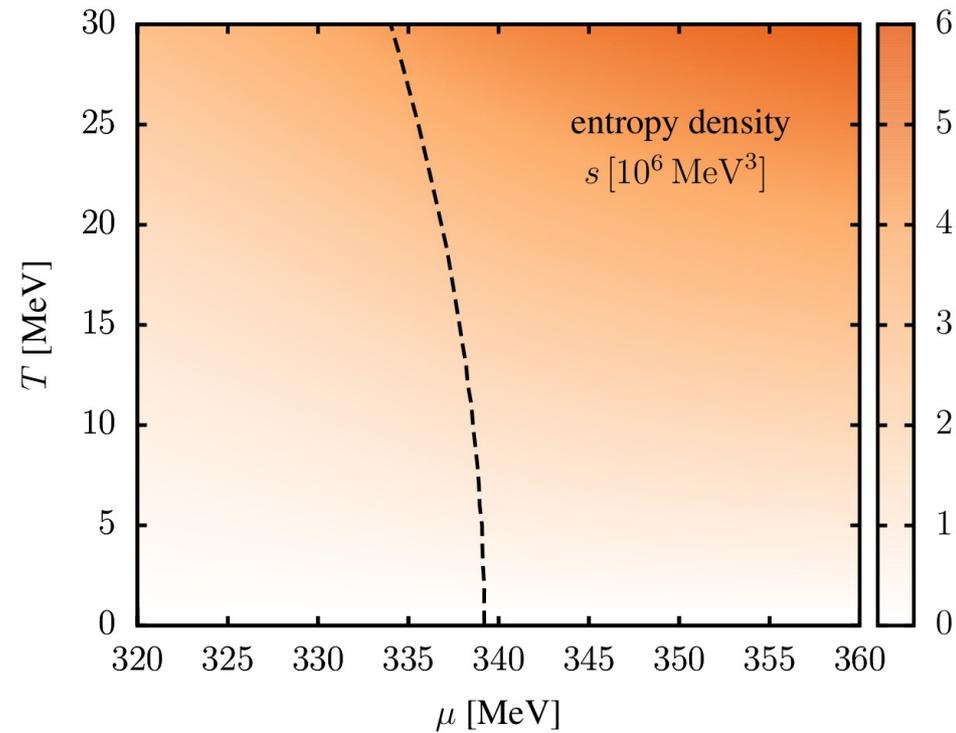
# Flat Regulator: High Density Flow

[arXiv:2206.13067]



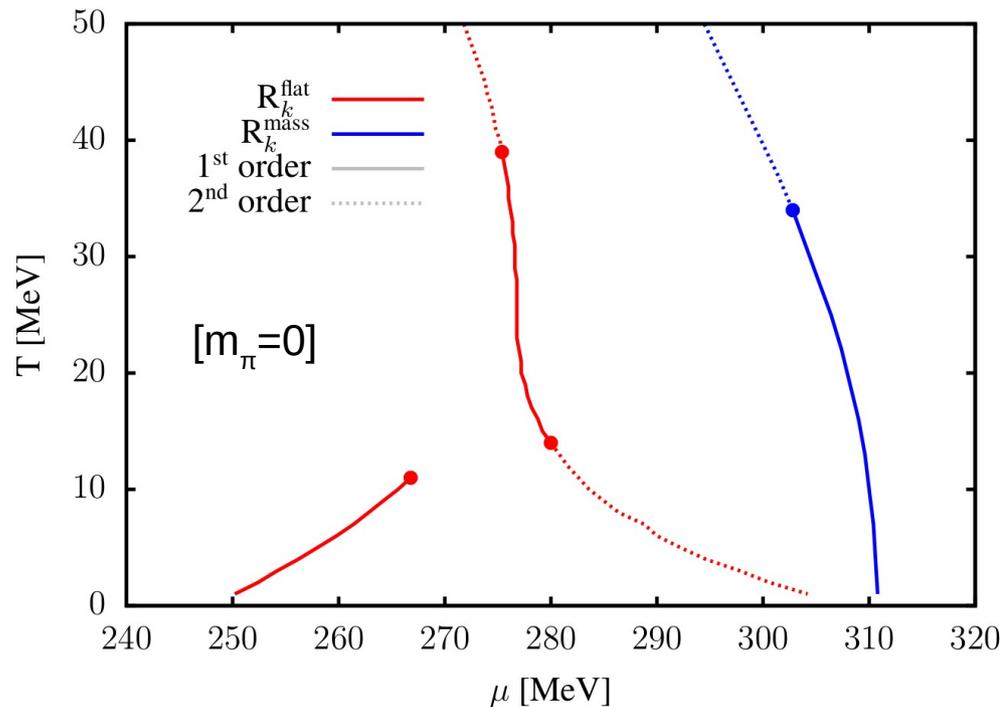
$$R_k^{\text{flat}} = (k^2 - \vec{p}^2) \Theta(k^2 - \vec{p}^2) \longrightarrow l_0^F \propto 1 - \Theta(\mu - E_\psi)$$

# Endpoint with Mass-like Regulator



# Endpoint with Mass-like Regulator

➤ Missing endpoint? Check chiral limit:



➤ Next with physical pion mass:

- Want to vary  $m_\sigma$
- Not possible in current setup!

➤ Flat regulator:

- Phase structure similar to previous findings: “triangular region”

[Schaefer, Wambach(2005)]

➤ Mass-like regulator:

- No splitting of the phase transition line
- First order transition and corresponding endpoint

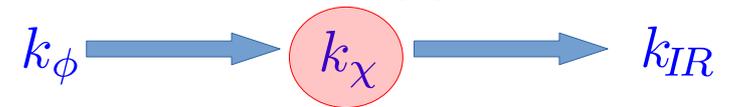
↳ Existence of CEP depends on masses

# How to Improve the LPA Setup?

- Limitations in the previous setup:
  - ➔ Range of possible sigma masses strongly restricted
  - ➔ Only low temperatures not affected by cutoff effects
  - ➔ Vacuum calculations not feasible for mass-like regulators

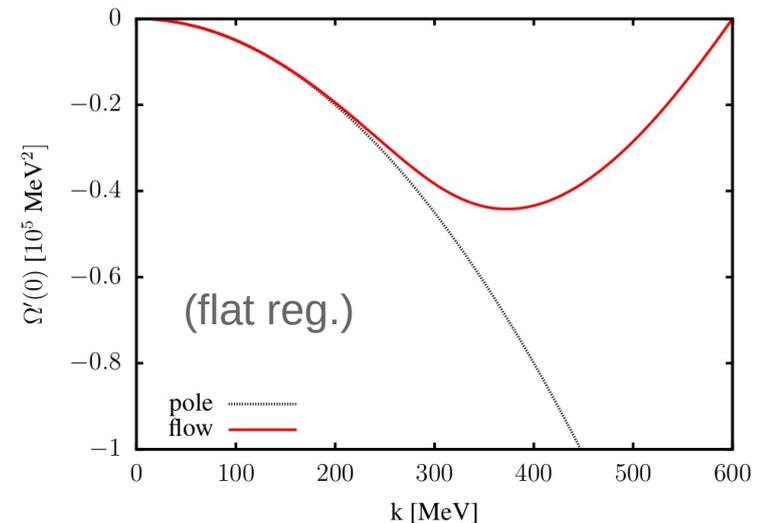
...caused by:

Parameter finding procedure:



- $k_\chi$  and  $k_\phi$  as only free parameters
- small UV-cutoff

Proximity to pion pole:

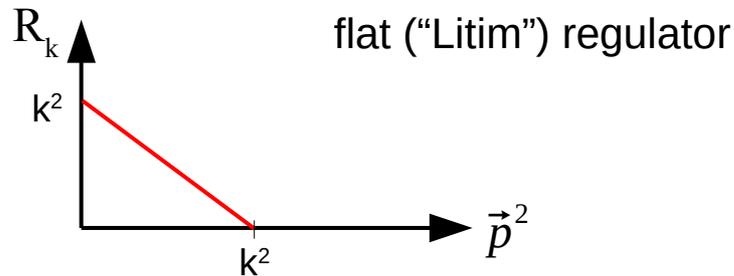


# How to Improve the LPA Setup?

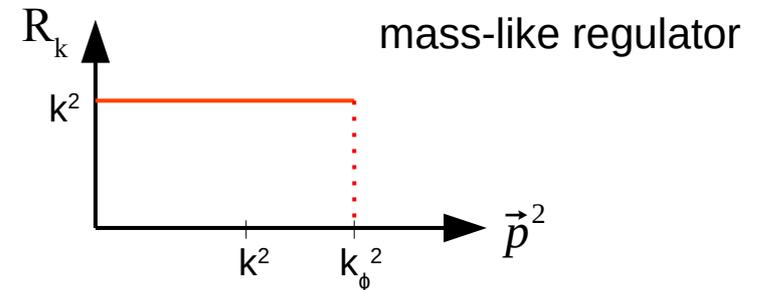
- Idea: Use two different types regulators for fermions and bosons
  - Proximity to pion pole determined by bosonic regulator
    - ➔ Choose flat regulator here to stabilize calculations and allow “usual” way of parameter fixing
  - Fermion decoupling sensitive to regulator choice
    - ➔ Vary fermionic regulator:  
Check how bosons “react” to different flows

# Regulator Choices: Varying the Cutoff

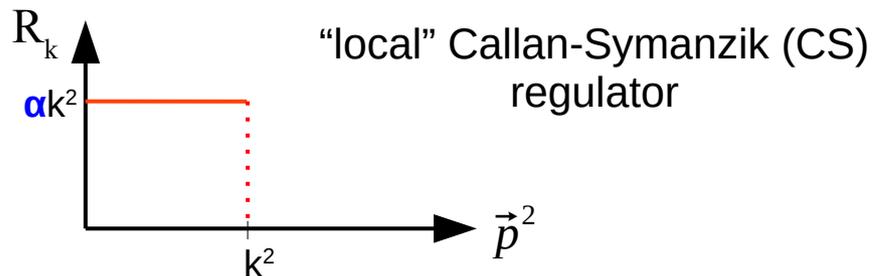
- Different types of (3d-) regulators, “bosonic form”:



$$R_k^{\text{flat}} = (k^2 - \vec{p}^2) \Theta(k^2 - \vec{p}^2)$$

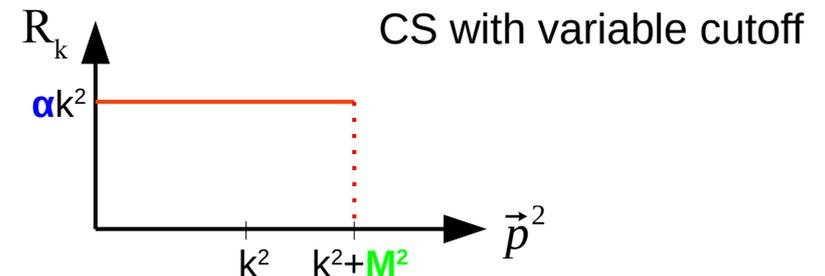


$$R_k^{\text{mass}} = k^2 \Theta(k_\phi^2 - \vec{p}^2)$$



$$R_k^{\text{CS,loc}} = \alpha k^2 \Theta(k^2 - \vec{p}^2)$$

allows fit to  $R_k^{\text{flat}}$ -flow



$$R_k^{\text{CS,var}} = \alpha k^2 \Theta(k^2 + M^2 - \vec{p}^2)$$

see also: “Renormalized spectral flows”,  
(Braun et al., 2022)

# Regulator Choices

$$R_k^{\text{CS,var}} = \alpha k^2 \Theta(k^2 + M^2 - \vec{p}^2)$$

## ➤ Parameter M:

- Momentum boundary M, enables testing for relevant momenta
- “Interpolation” between other regulators:

$$R_k^{\text{CS,var}} \approx \begin{cases} R_k^{\text{CS,loc}} & , k \gg M \\ R_k^{\text{mass}} & , k \ll M \end{cases}$$

## ➤ Connection between fermionic and bosonic regulators:

$$R_k(\vec{p}^2) = \vec{p}^2 r_\phi(\vec{p}^2/k^2)$$

$$R_k^\psi(\vec{p}^2) = i\vec{p} r_\psi(\vec{p}^2/k^2) \quad \text{with} \quad r_\psi(x) := \sqrt{1 + r_\phi(x)} - 1$$

# Mismatch in Contributions

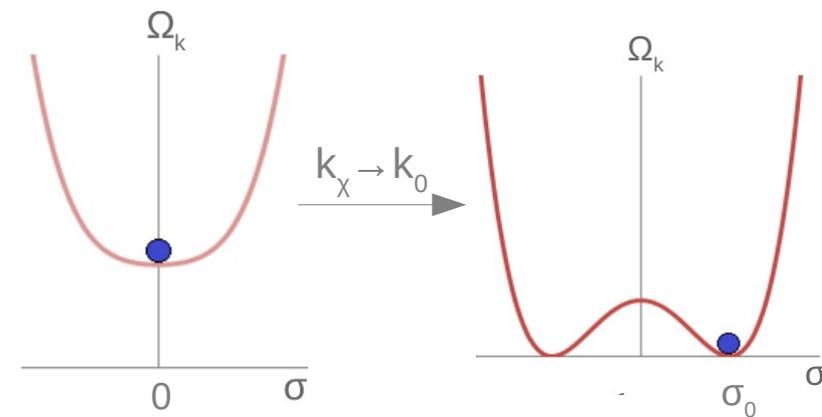
- Large differences between flows from different regulators
  - ↳ Mismatch: Different effective scales for bosons and fermions
  - ↳ partly cured by rescaling / matching vacuum flows

- Chiral sym. breaking scale  $k_\chi$ :

Fermi Regulator	$\Lambda$	$\lambda$	$v^2$	$k_\chi$
flat	0.9	20.9	0.355	0.605
mass	0.6	6.5	0.351	0.206
CS,loc	0.9	24.4	0.305	0.626
CS,var (M=0.3)	0.9	17.5	0.377	0.588
CS,var (M=0.5)	0.9	9.1	0.496	0.454

(in GeV units)

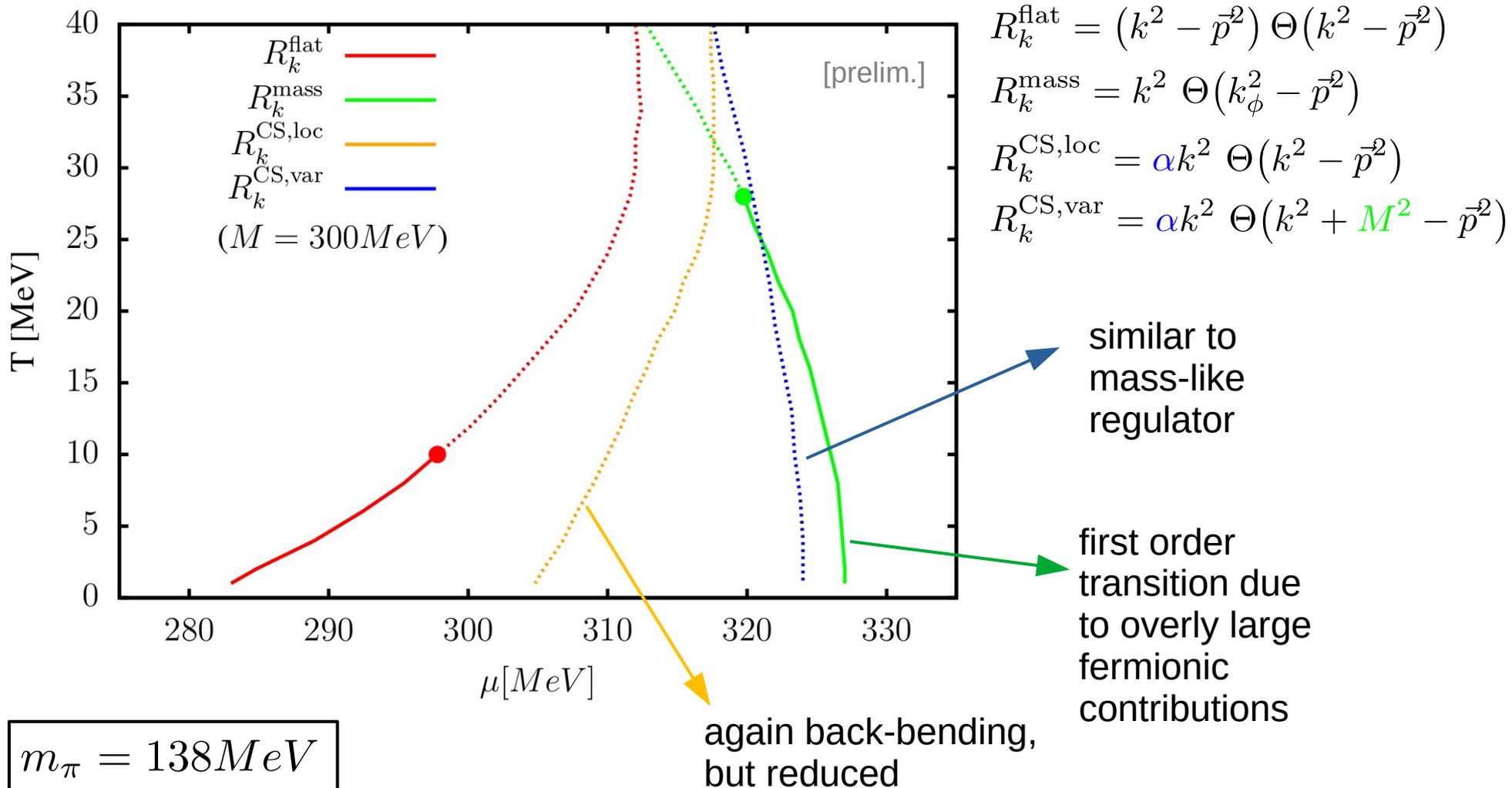
➔ Fermionic contributions strongly overestimated



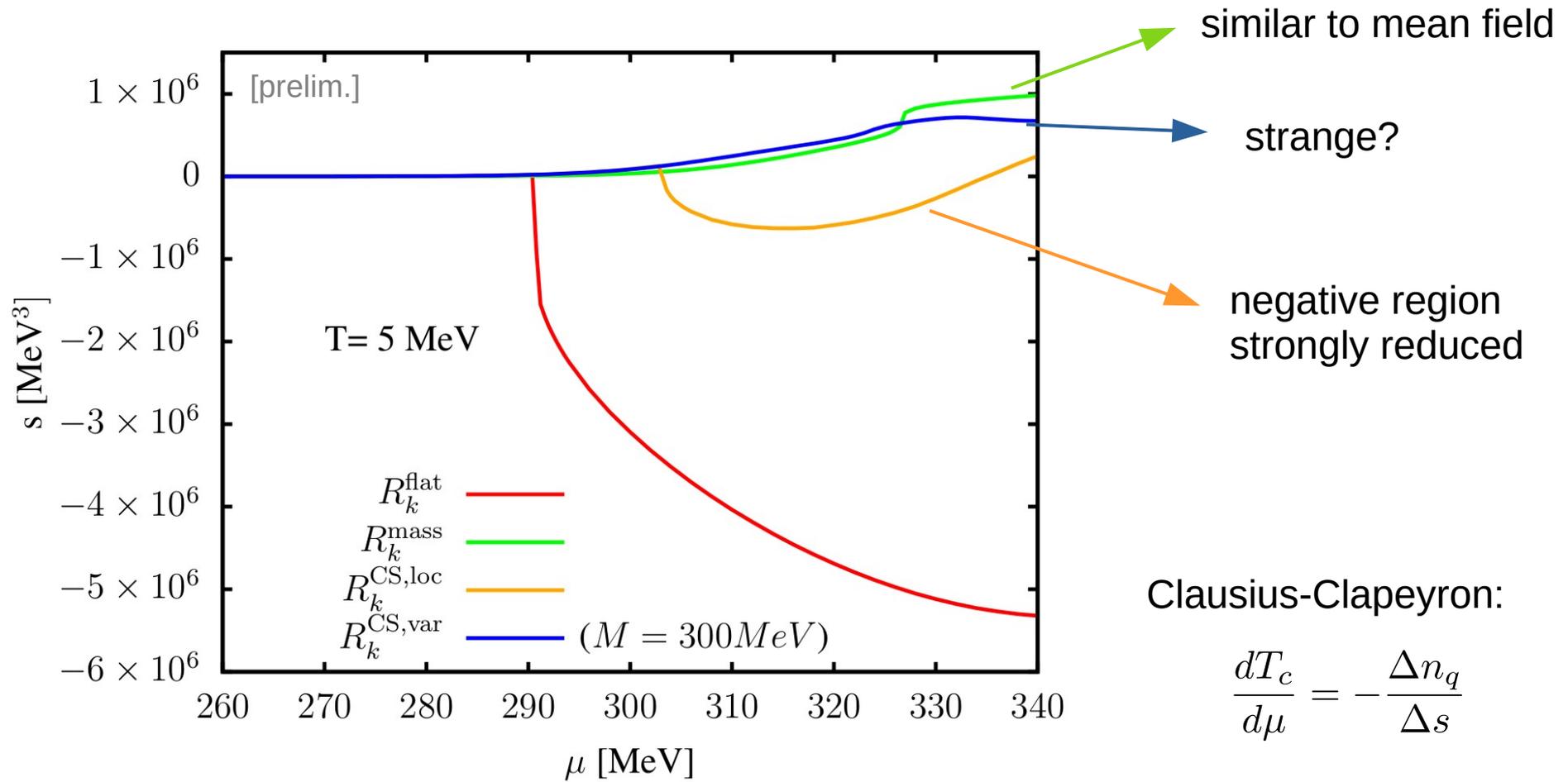
# Numerical Results

- In the following:
  - For bosons the flat regulator is used for all results
  - Comparison of different fermionic regulators

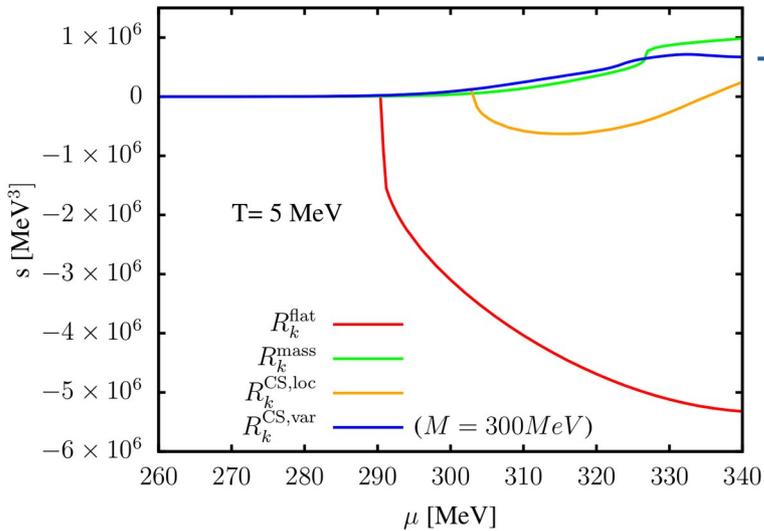
# Phase Transition with Different Regulators



# Entropy Density

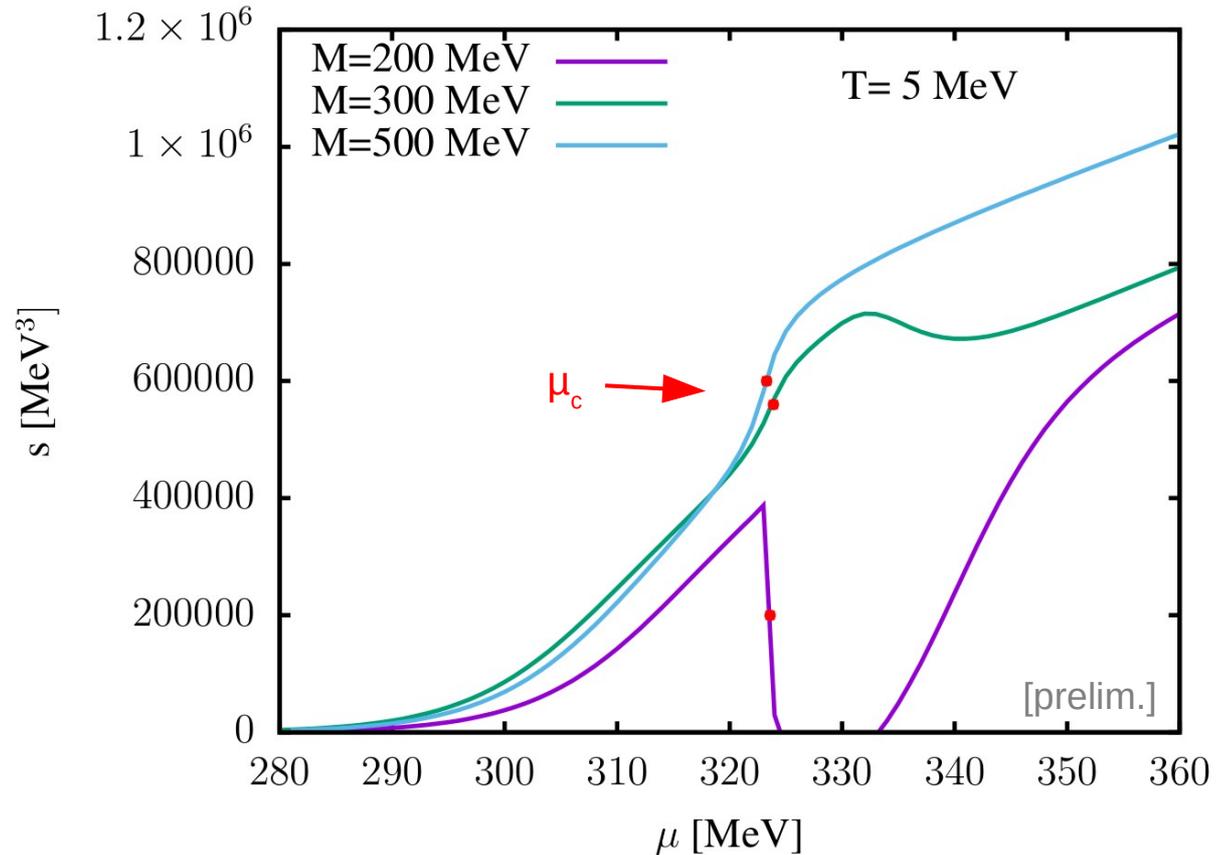


# Entropy Density



Strange behavior for  $R_k^{\text{CS,var}}$  ?

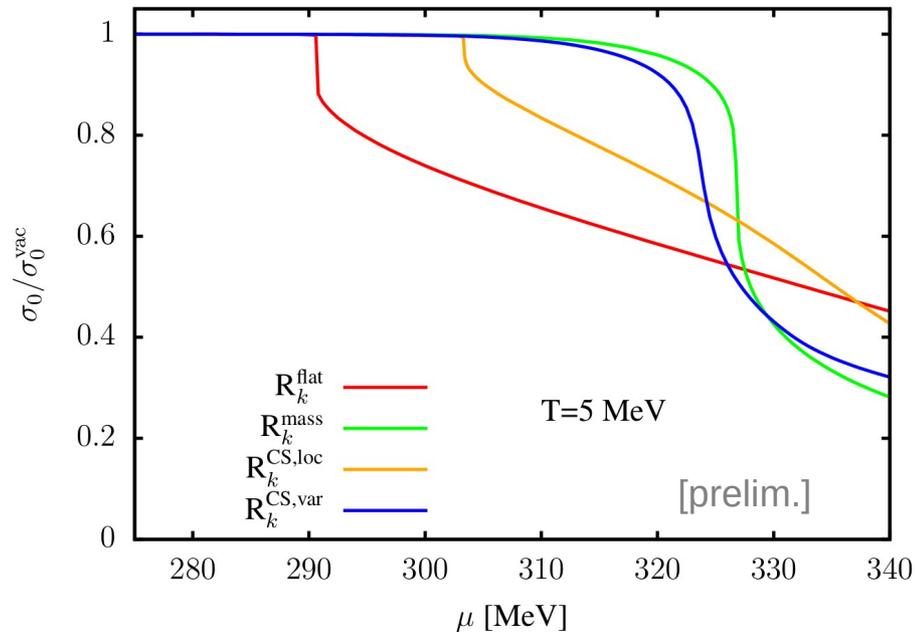
Compare different M-values:



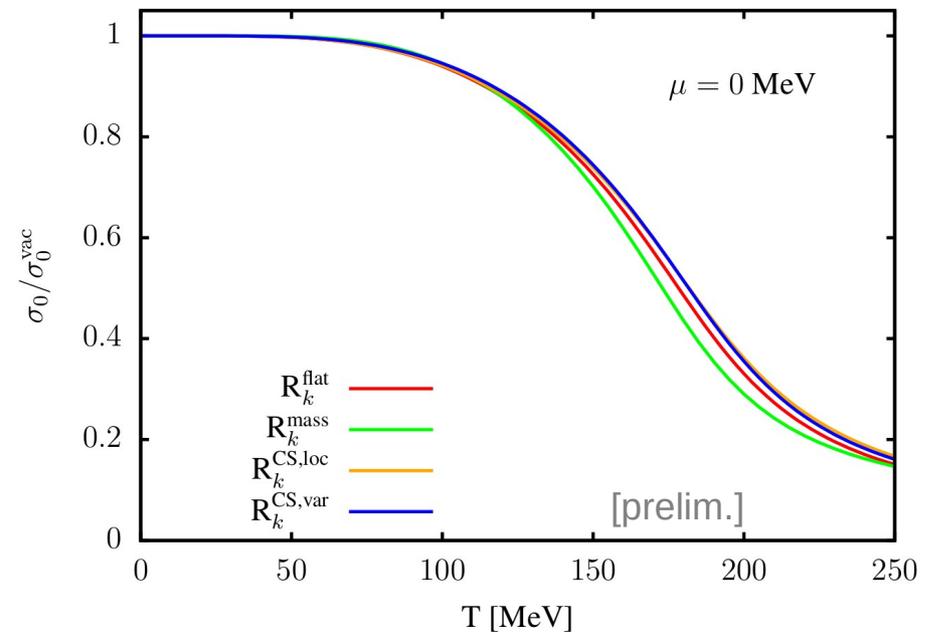
➤ Negative entropy for  $M \lesssim 300 \text{ MeV}$

➤ “Dip” if  $M \lesssim \mu$  as remnant away from phase transition

# $\sigma$ -Condensate: Comparison with $\mu=0$ -axis

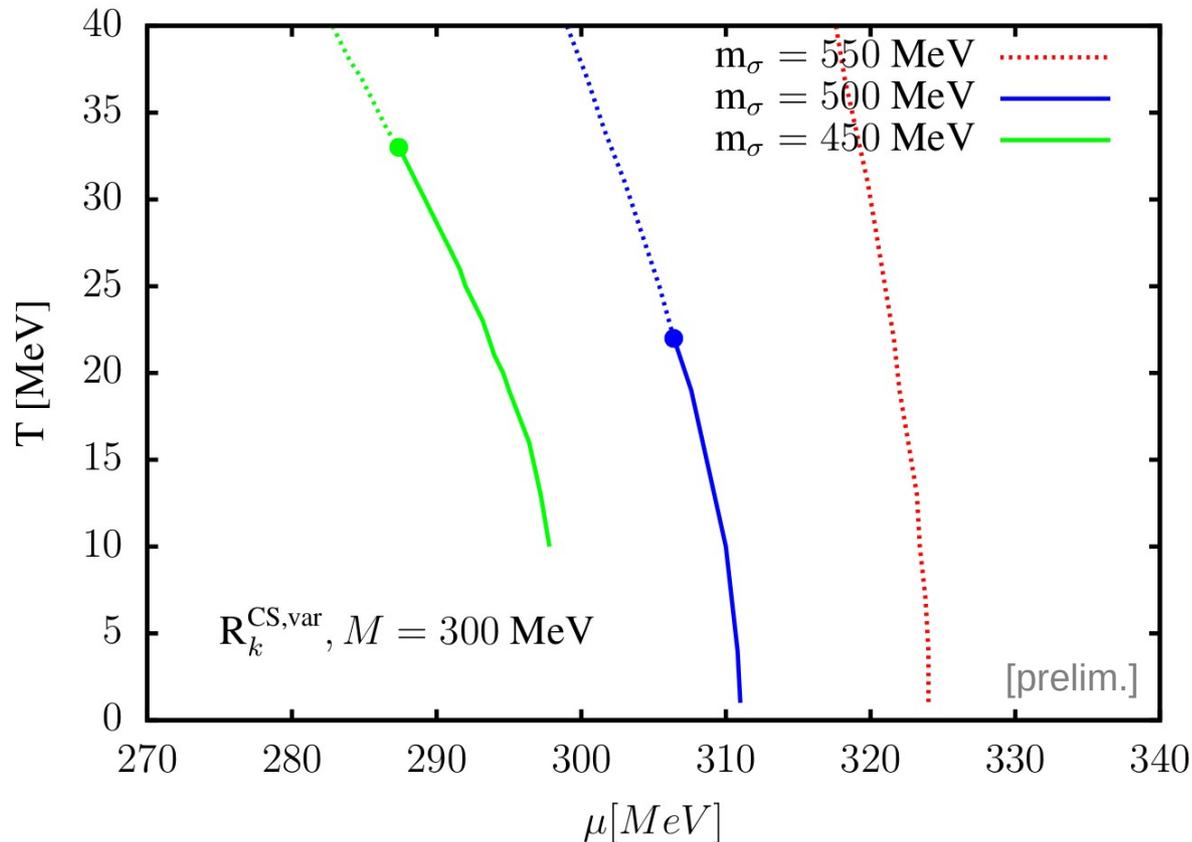


- Large distinctions for different  $R_k$
- Back-bending: Small transition and residual condensate
- others more similar to MF result, mass-like much steeper



- Regulator effects much smaller, largest in crossover region
- Critical Temperature varies only by  $\sim 5$  MeV

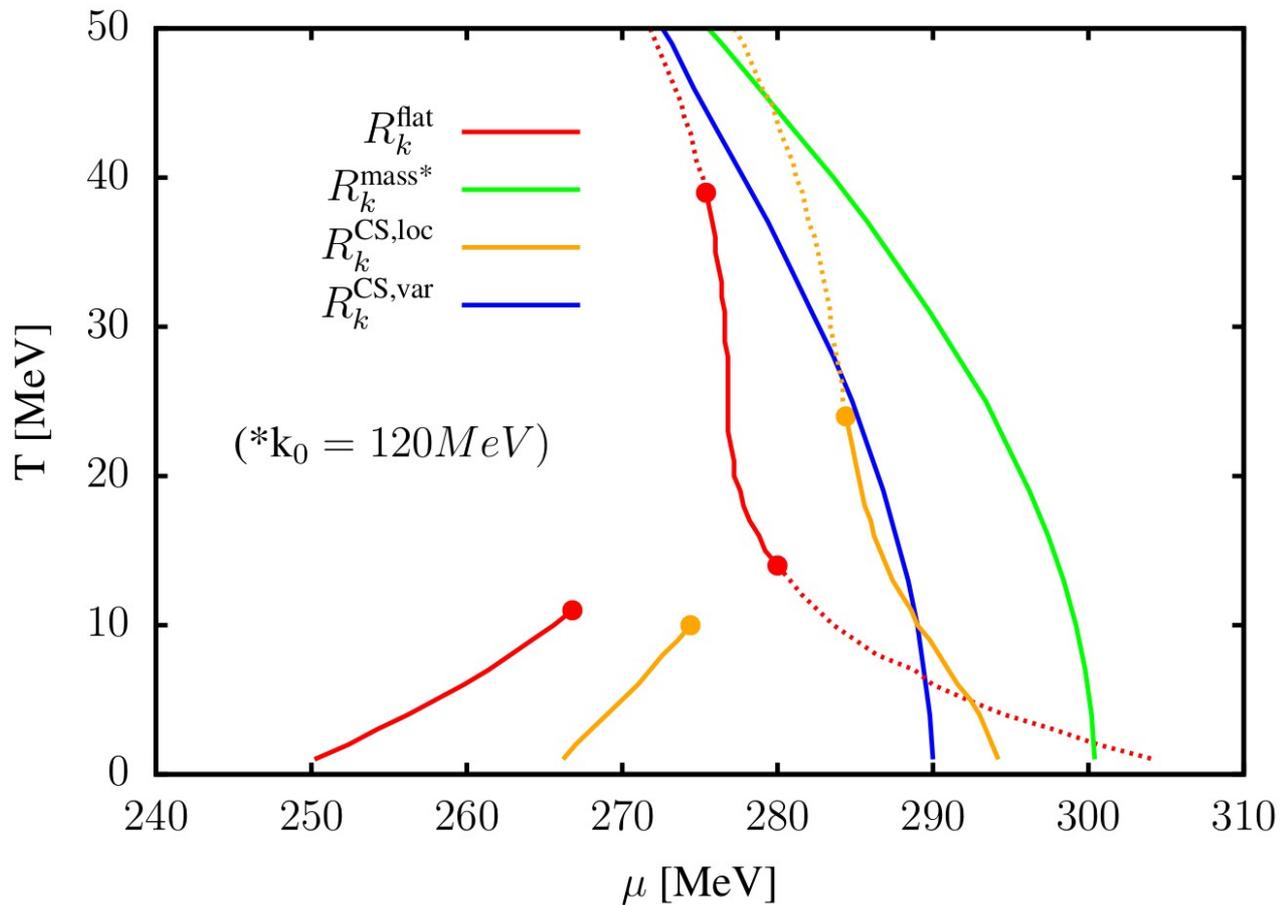
# CEP for Different $\sigma$ -masses



- No CEP when  $\sigma$ -mass similar to previous calculations
- CEP appears & crit. temperature increases when  $m_\sigma$  is lowered

$$R_k^{\text{CS,var}} = \alpha k^2 \Theta(k^2 + M^2 - \vec{p}^2)$$

# Chiral Limit



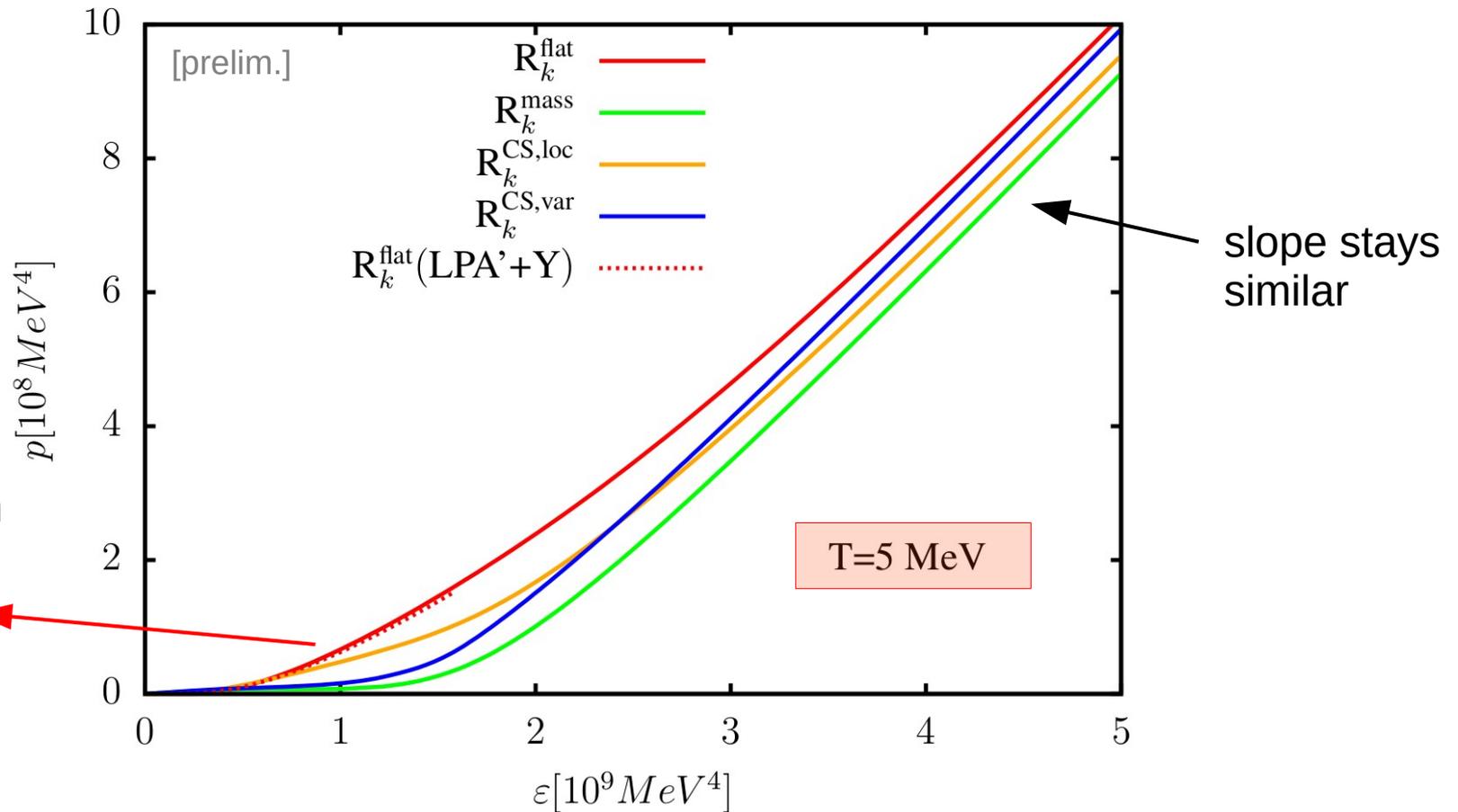
$$R_k^{\text{flat}} = (k^2 - \vec{p}^2) \Theta(k^2 - \vec{p}^2)$$

$$R_k^{\text{mass}} = k^2 \Theta(k_\phi^2 - \vec{p}^2)$$

$$R_k^{\text{CS,loc}} = \alpha k^2 \Theta(k^2 - \vec{p}^2)$$

$$R_k^{\text{CS,var}} = \alpha k^2 \Theta(k^2 + M^2 - \vec{p}^2)$$

# Equation of State



# Summary

- Investigated truncation effects in the FRG framework
- Mass-like regulator: **No** back-bending, entropy remains **positive** but limited usability
- Tested Setup with fixed bosonic and different fermionic regulators
  - UV-cutoff of CS-type regulators crucial for results
  - artifacts found as soon as momenta  $\vec{p}^2 < \mu^2$  are cut
  - only minor regulator effects at vanishing chemical potential
  - Nice testing ground but surely not “the final answer”

# Outlook

- Effects on neutron star equation of state (lower  $T$ ) and mass-radius relations?
- Find solutions/regulators for more advanced truncations, e.g. LPA' or higher order derivative expansions
- Additional channels (e.g. via dynamical hadronization)
- Upgrade momentum dependencies

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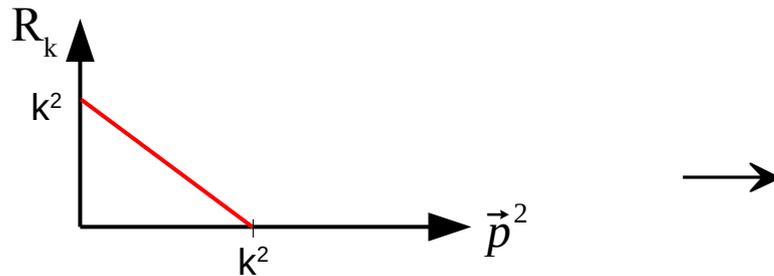
GRAZIE MILLE !

# Backup

# Throwback to the Previous Talk

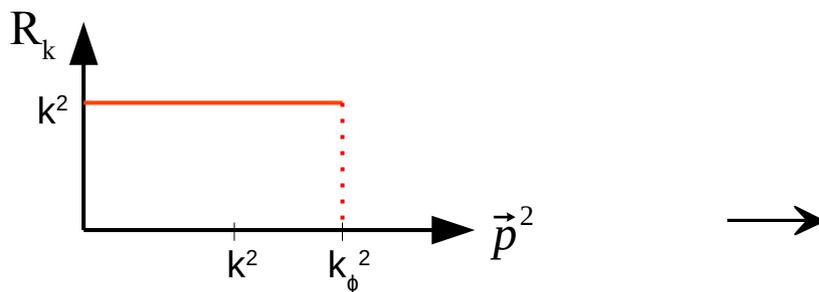
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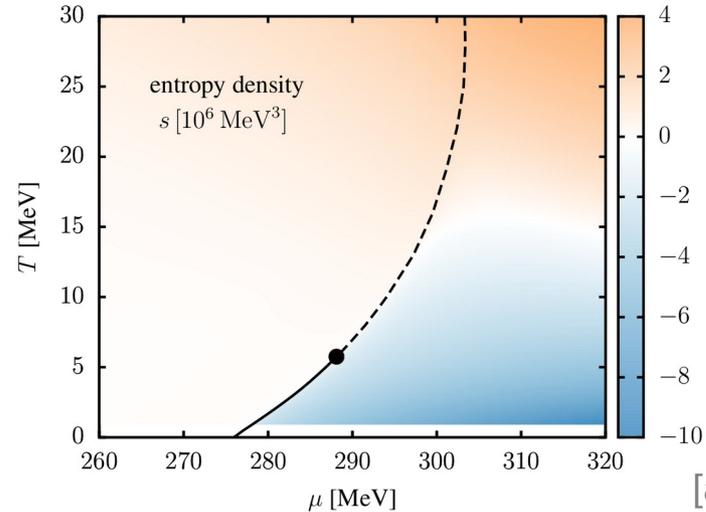
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dim.reduced regulators

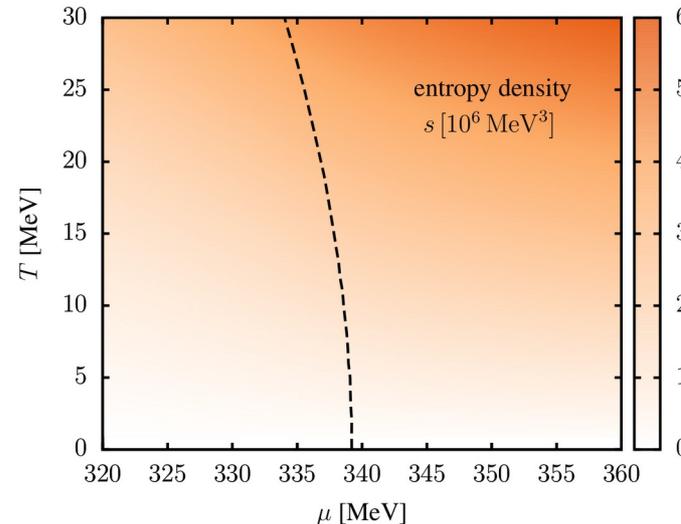


$$R_k^{\text{mass}} = k^2 \Theta(k_\phi^2 - \vec{p}^2)$$

compositeness scale



[arXiv:2206.13067]



# Why are we using 3d-regulators?

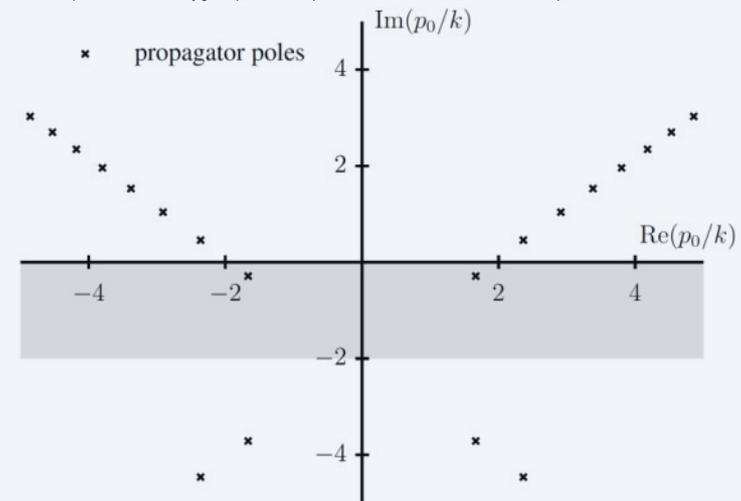
- Downside: Dimensionally reduced regulators break  $O(4)$  symmetry
- Upsides:
  - Matsubara summation can be performed analytically
  - Easiest way to prevent regulators from breaking Silver Blaze symmetry

## 4d fermionic regulators:

- Necessary for Silver Blaze:
- Silver Blaze violation still possible due to regulator induced complex poles, e.g. with exponential reg.:
- Discussion and Solutions:  
[Pawlowski, Strodhoff (2015)]

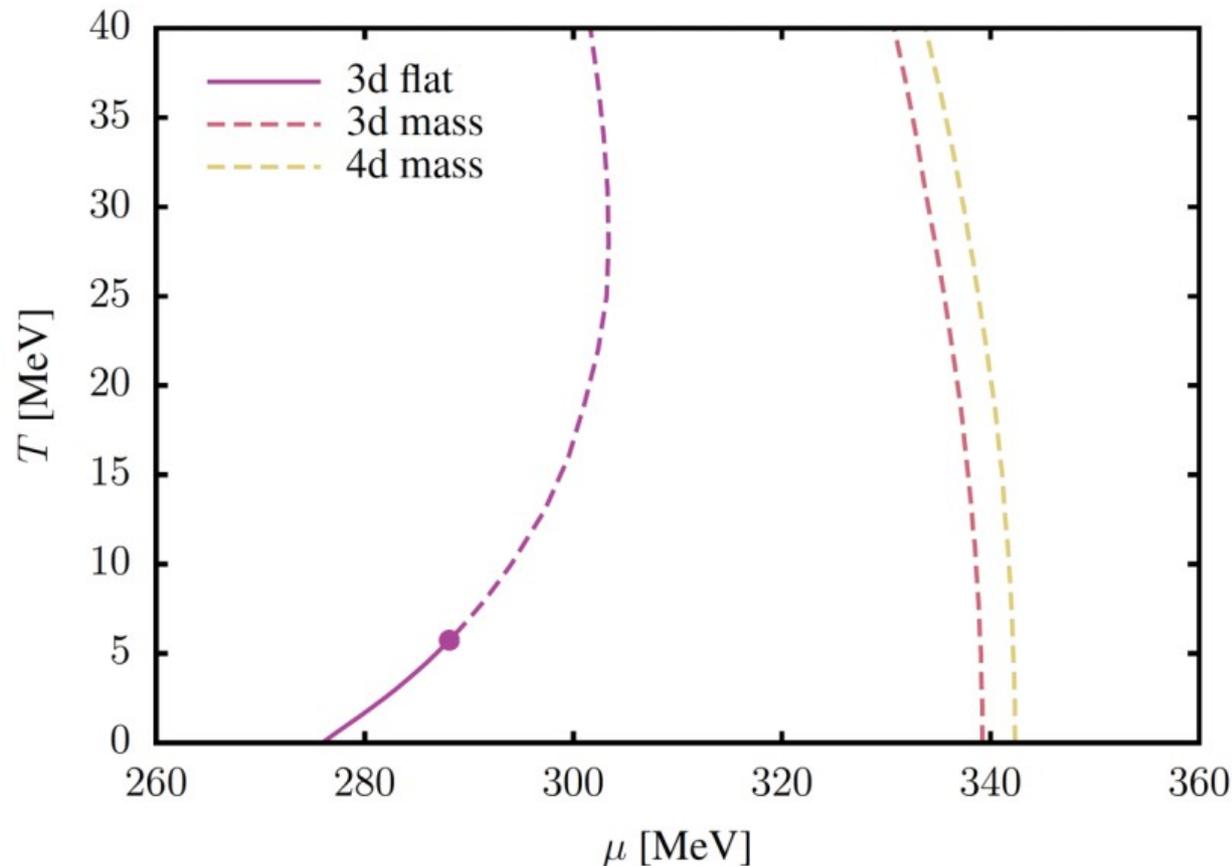
$$R_k^F = i\not{p} r^F (p^2/k^2)$$

$$R_k^F(p, \mu) = R_k^F(\tilde{p}, 0) \text{ with } \tilde{p} = (p_0 + i\mu, \vec{p})$$



# Why are we using 3d-regulators?

- Mass-like regulators allow us to circumvent this problems
- Only small quantitative differences found:

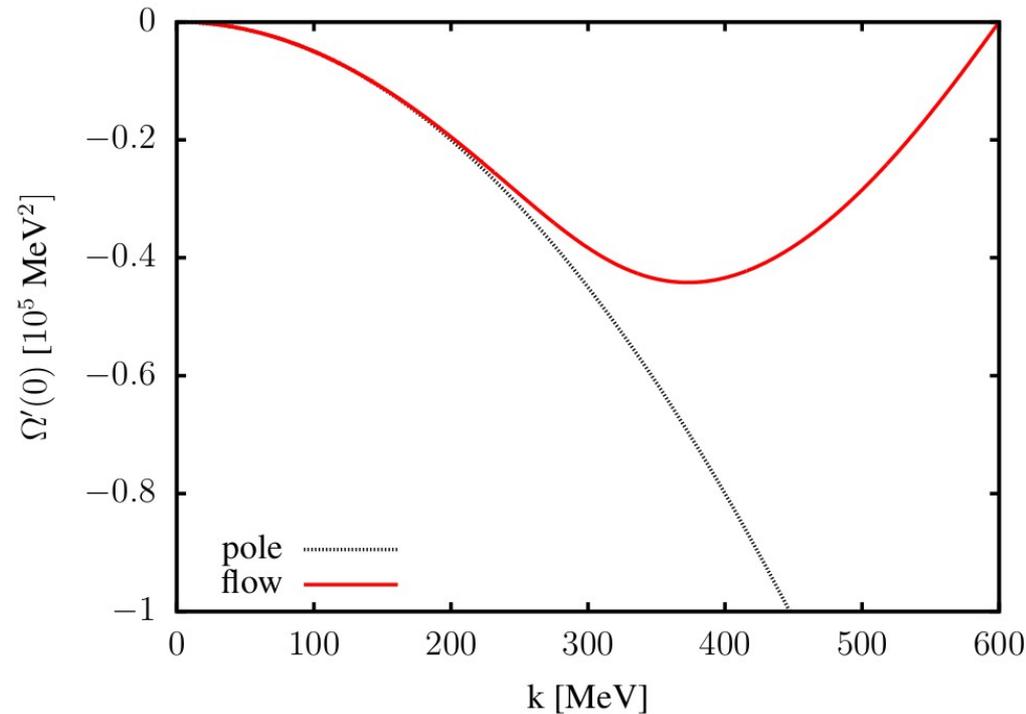


→ Minor impact of dimensional reduction on results

→ Especially: No connection between back-bending and 3d-regulators

# Pole Proximity: Estimation

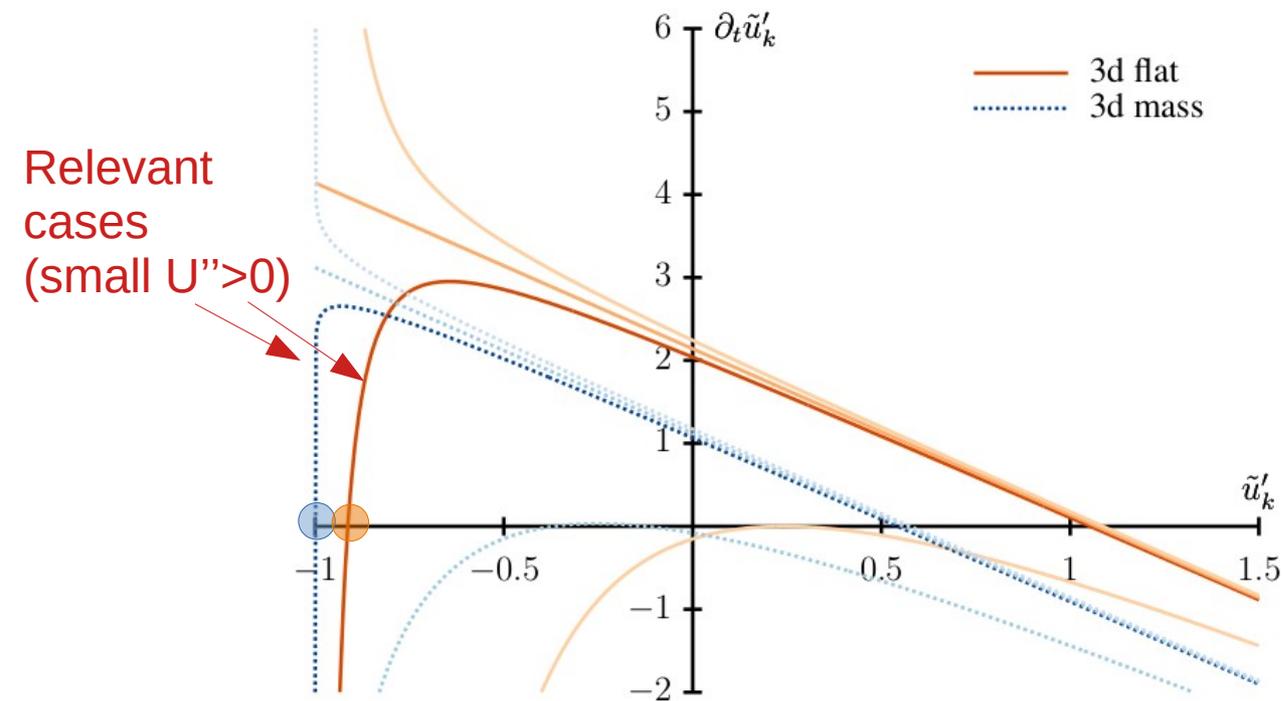
- Reminder: Vacuum flow runs close to the pole at  $E_\pi=0$ :



How can we estimate the distance to this pole for a given regulator?

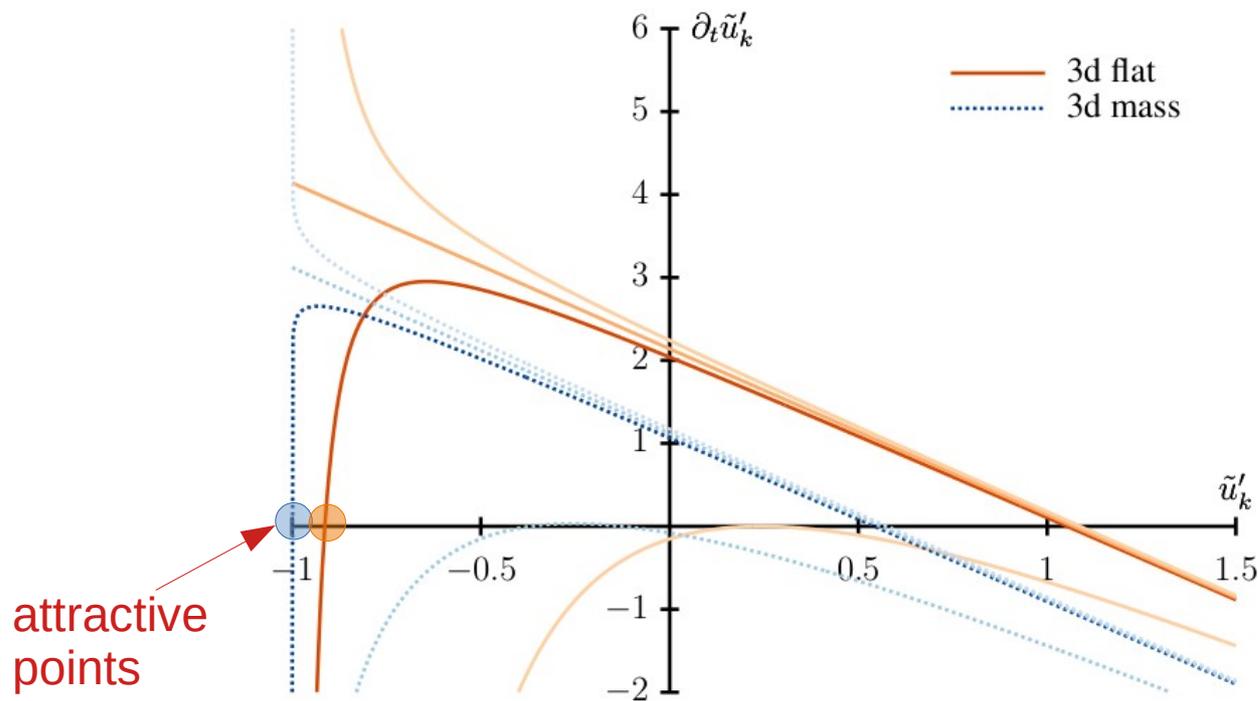
# Pole Proximity: Estimation

- For convenience define  $\tilde{u}'_k = 2U'_k(0)/k^2 \longrightarrow$  Pole at  $\tilde{u}'_k = -1$
- Examine it's flow equation for different (fixed) values  $U''$  :



# Pole Proximity: Estimation

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- Examine it's flow equation for different (fixed) values  $U''$  :



➔ Relevant cases have attractive stationary point near pole

↳ no fixed point (depends on  $U''$ )

↳ Explains behavior observed in the flow

↳ Setting  $\partial_t \tilde{u}'_k = 0$  and expanding in  $\delta u = \tilde{u}'_k + 1$  allows to approximate distance between pole and attractive point