SUBA-Jet

A New Model for Jet Energy Loss in Heavy Ion Collisions

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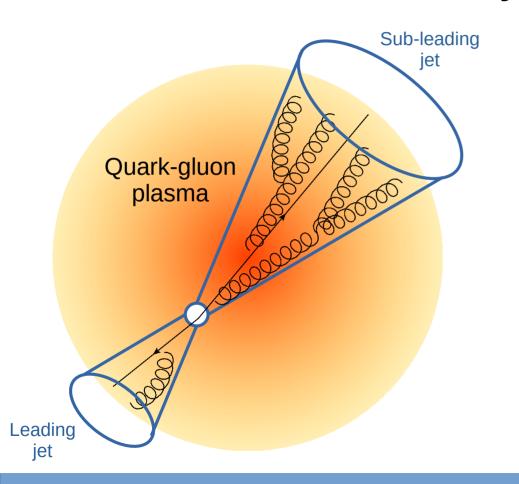
with Iurii Karpenko, Joerg Aichelin, Pol-Bernard Gossiaux, Martin Rohrmoser, and Klaus Werner



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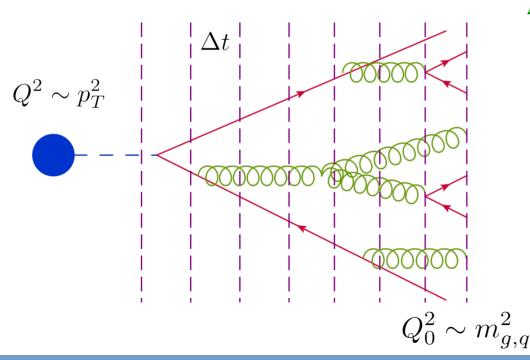
Jets in Heavy Ion Collisions



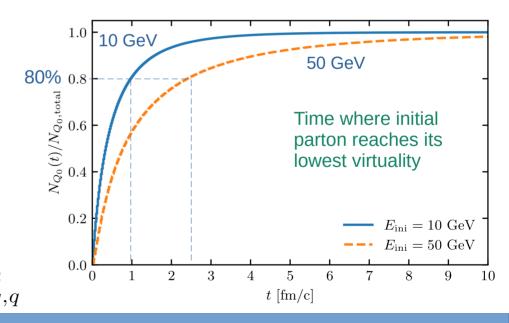
- Jets can be produced alongside the QGP in heavy ion collisions
- Interactions between jet partons and the QGP medium leads to modifications of jet properties
 - → Jet Energy Loss / Quenching
- SUBA-Jet:
 Monte Carlo for jet energy loss in heavy ion collisions

Vacuum Parton Shower

- Monte Carlo of a vacuum parton shower originally developed by Martin Rohrmoser
- Evolution according to the DGLAP equations from high virtuality $Q_{max} \sim p_T$ to low virtuality Q_0
- Time evolution split into time steps, mean life time



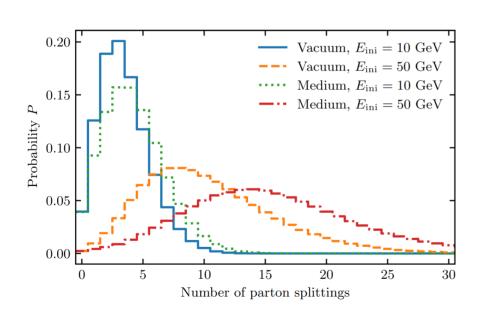
$$\Delta t = \tau = \frac{E}{Q^2}$$

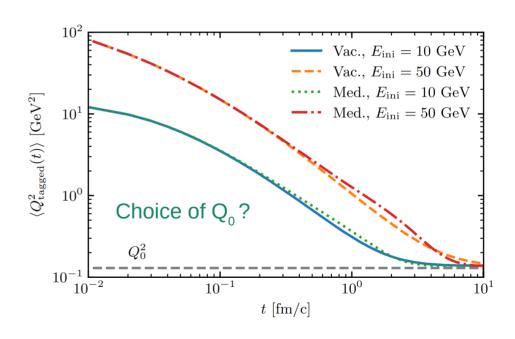


"Vacuum" Parton Shower in Medium

 Medium interactions for high Q regime resulting in virtuality increase, similar to YaJEM (T. Renk, 2008)

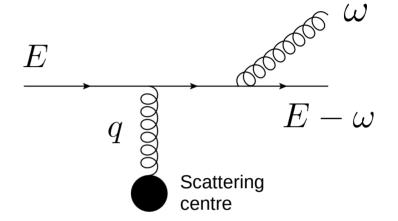
$$\frac{\mathrm{d}Q^2}{\mathrm{d}t} = \hat{q}(T)$$





Medium-Induced Single Radiation

- Inelastic collision:
 Single gluon emission from single medium scattering
- Original result from Gunion-Bertsch (1982)
 Generalised to massive case by Aichelin, Gossiaux, Gousset (2014)

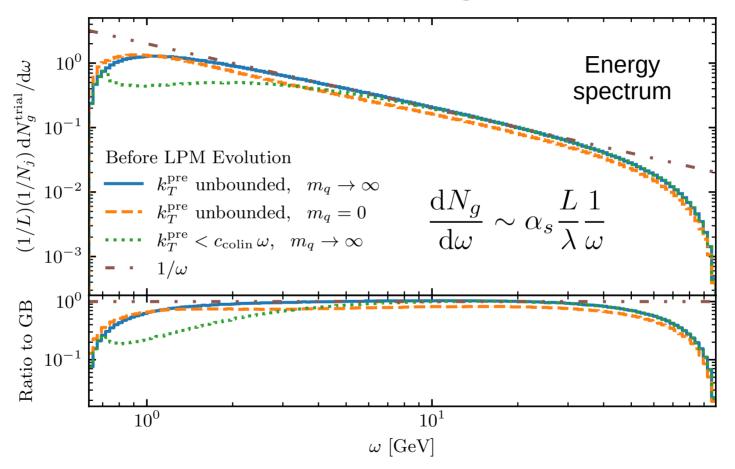


- Initial Gunion-Bertsch seed: i.e. radiation of a preformed gluon from a single scattering (Each parton can generate a number of preformed gluons)
- Gunion-Bertsch cross-section from scalar QCD

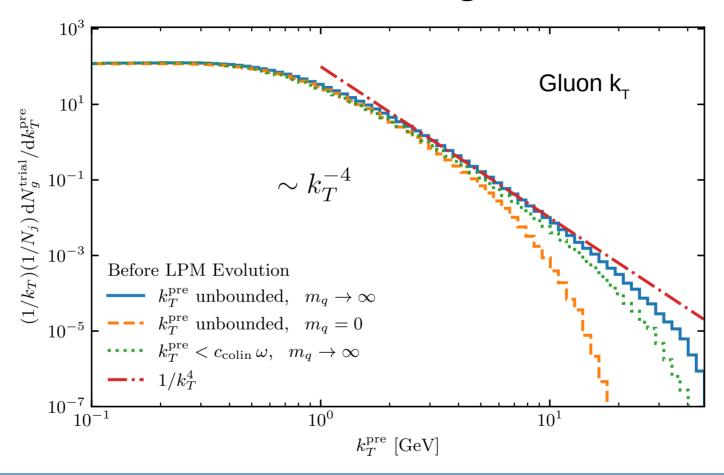
$$\frac{\mathrm{d}\sigma^{Qq \to Qqg}}{\mathrm{d}x \,\mathrm{d}^2 k_T \,\mathrm{d}^2 l_t} = \frac{\mathrm{d}\sigma_{\mathrm{el}}}{\mathrm{d}^2 l_t} P_g(x, k_T, l_T) \theta(\Delta)$$

$$\frac{\mathrm{d}\sigma_{\mathrm{el}}}{\mathrm{d}^2 l_t} \sim \frac{8\alpha_s^2}{9(l_T^2 + \mu^2)^2}$$

Medium-Induced Single Radiation



Medium-Induced Single Radiation



Coherency and the LPM Effect

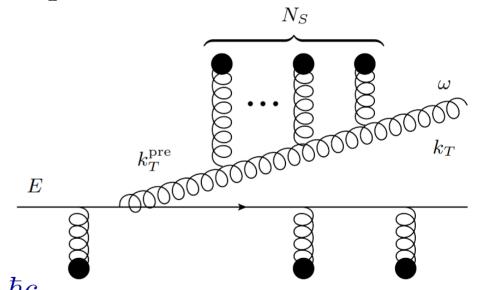
The formation of the radiated gluon is a quantum mechanical process

Formation time:
$$t_f \sim \sqrt{\frac{\omega}{\hat{q}}}$$

- Coherence effects: Landau-Pomeranchuk-Migdal (LPM) effect
- Have to take into account multiple scatterings with the medium during the formation time

$$N_s = rac{t_f}{\lambda}$$

$$\omega = {
m gluon\ energy}$$
 $\hat{q} = {
m medium\ modifications}$



L= path length of medium

Implementation of the LPM Effect

- At each timestep:
 - Elastic scattering with prob. $\Gamma_{\rm el}\Delta t$

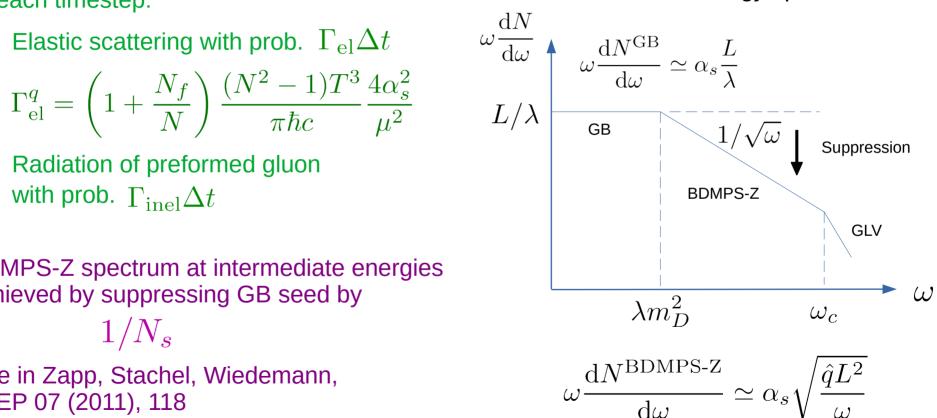
$$\Gamma_{\rm el}^q = \left(1 + \frac{N_f}{N}\right) \frac{(N^2 - 1)T^3}{\pi \hbar c} \frac{4\alpha_s^2}{\mu^2}$$

- Radiation of preformed gluon with prob. $\Gamma_{\rm inel}\Delta t$
- BDMPS-Z spectrum at intermediate energies achieved by suppressing GB seed by

$$1/N_s$$

Like in Zapp, Stachel, Wiedemann, JHEP 07 (2011), 118

Radiation energy spectrum:

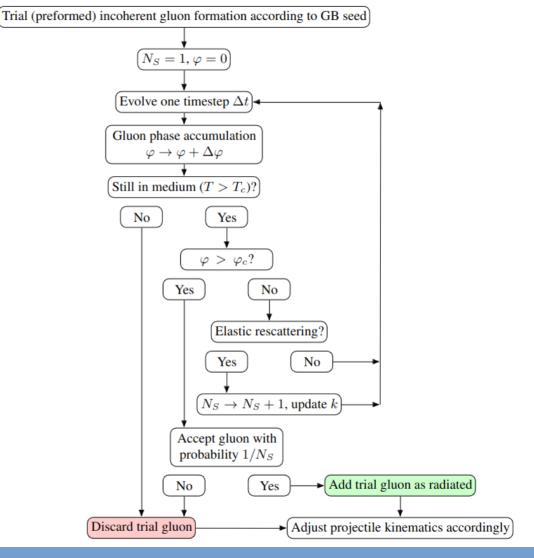


The Algorithm

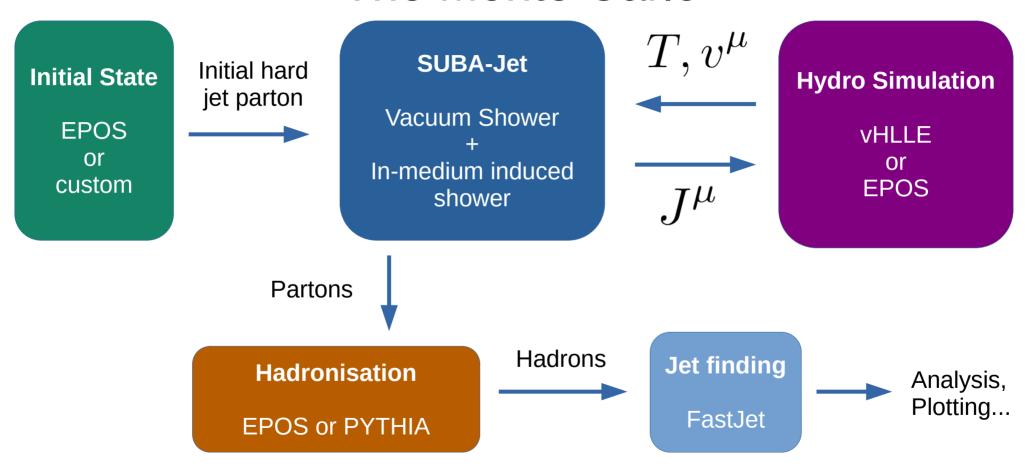
Flow diagram:

Monte Carlo algorithm for the coherent mediuminduced gluon radiation in our model

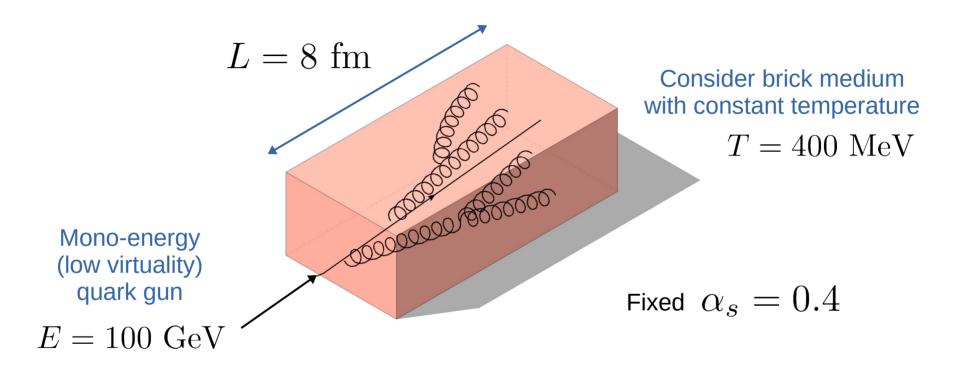
Various parameters and settings can be changed and tuned to compare distributions



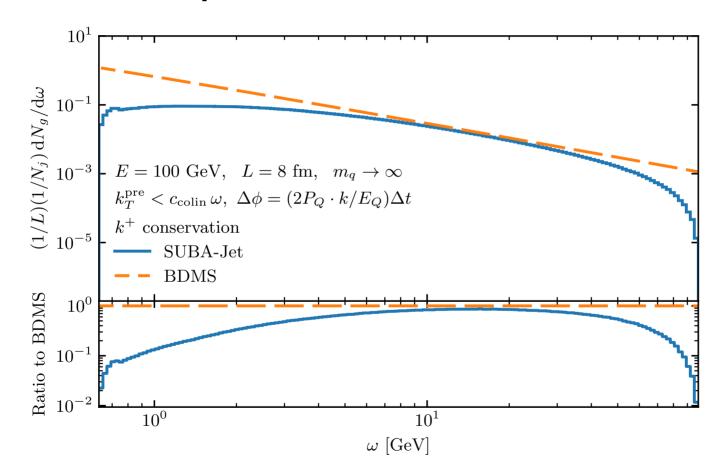
The Monte Carlo



First Results



We want to reproduce theoretical expectation and check effect of model parameters

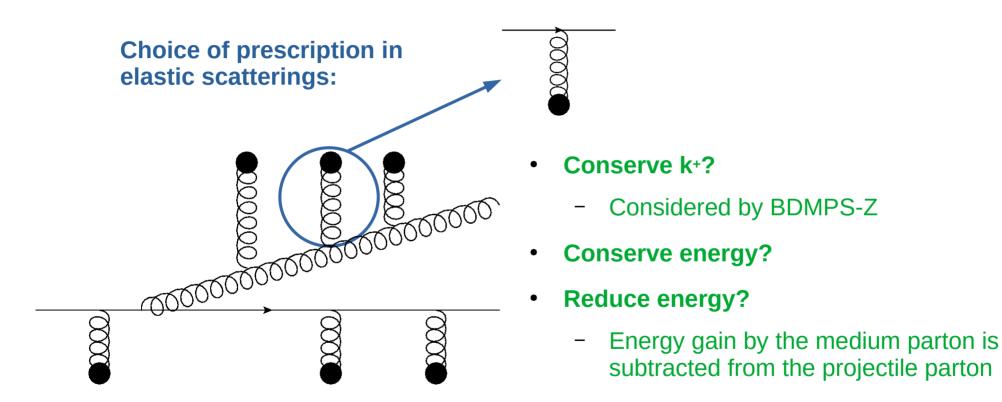


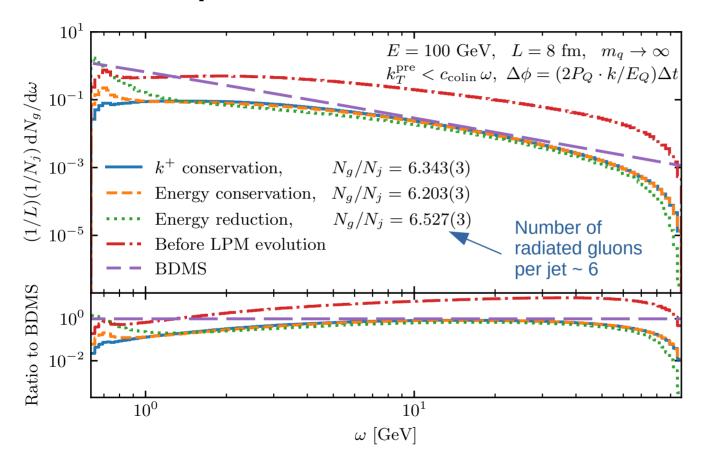
Energy spectrum

$$rac{\mathrm{d}N}{\mathrm{d}\omega}$$
 vs ω

Reproduces BDMPS-Z for intermediate energies

$$\frac{\mathrm{d}N}{\mathrm{d}\omega} \sim \frac{1}{\omega^{3/2}}$$

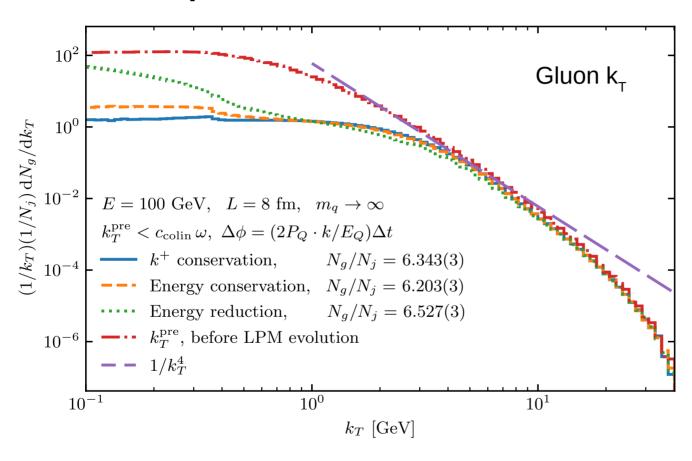




Same BDMS behaviour at intermediate energies

Difference at small energies

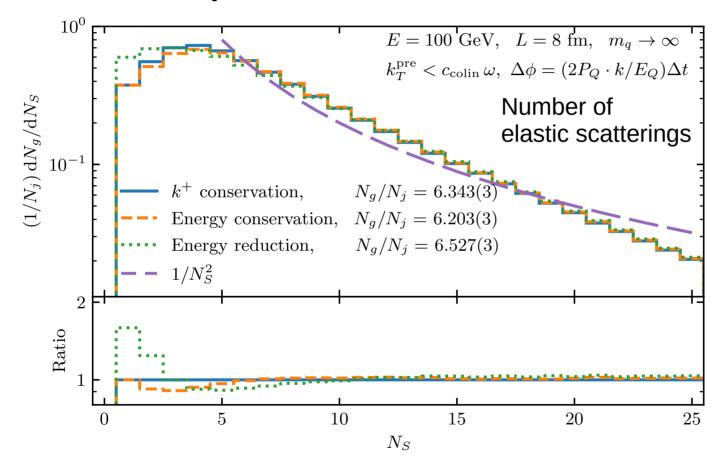
Effect of the phase accumulation on the energy distribution



Large difference at small $k_{\scriptscriptstyle T}$

The large k_T tail goes as $\sim k_T^{-4}$ as expected

$$\frac{\mathrm{d}N}{\mathrm{d}k_T} \sim \frac{1}{k_T^3}$$



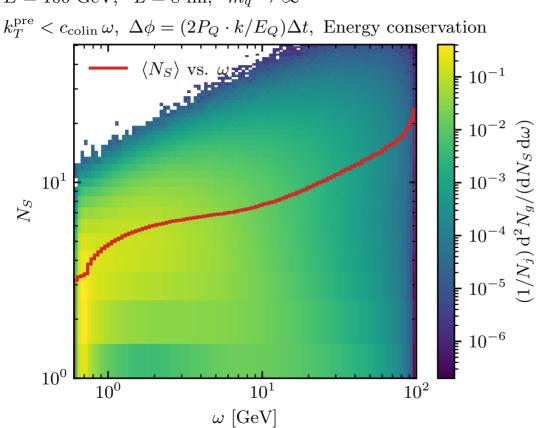
The number of elastic scatterings during LPM evolution is expected to go as

$$\frac{\mathrm{d}N}{\mathrm{d}N_S} \sim \frac{1}{N_S^2}$$

The energy reduction case is larger at $N_s = 1$

→ Larger probability of emission

$$E = 100 \text{ GeV}, L = 8 \text{ fm}, m_q \to \infty$$

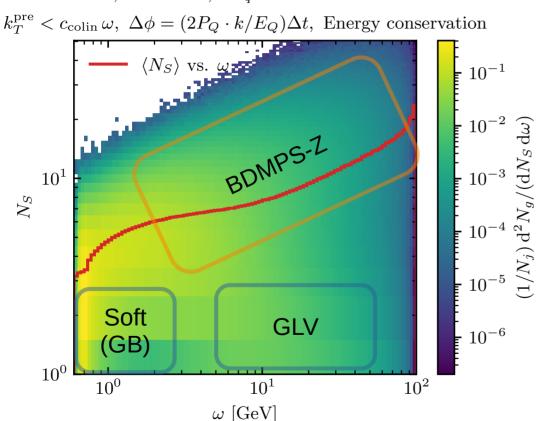


Double differential plot in N_s and ω

Red line: $\langle N_s \rangle$ vs. ω

$$N_S \sim t_f \sim \sqrt{\omega}$$

$$E = 100 \text{ GeV}, \quad L = 8 \text{ fm}, \quad m_q \to \infty$$



Double differential plot in N_s and ω

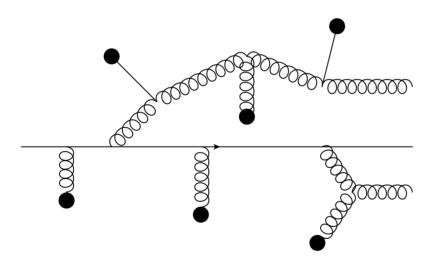
Red line: $\langle N_s \rangle$ vs. ω

$$N_S \sim t_f \sim \sqrt{\omega}$$

Convolution of different distributions

The Role of the Phase Accumulation

Choice of phase accumulation of the preformed (trial) gluons:



What is used in JEWEL:

$$\Delta \varphi = \frac{k_T^2}{\omega} \Delta t$$

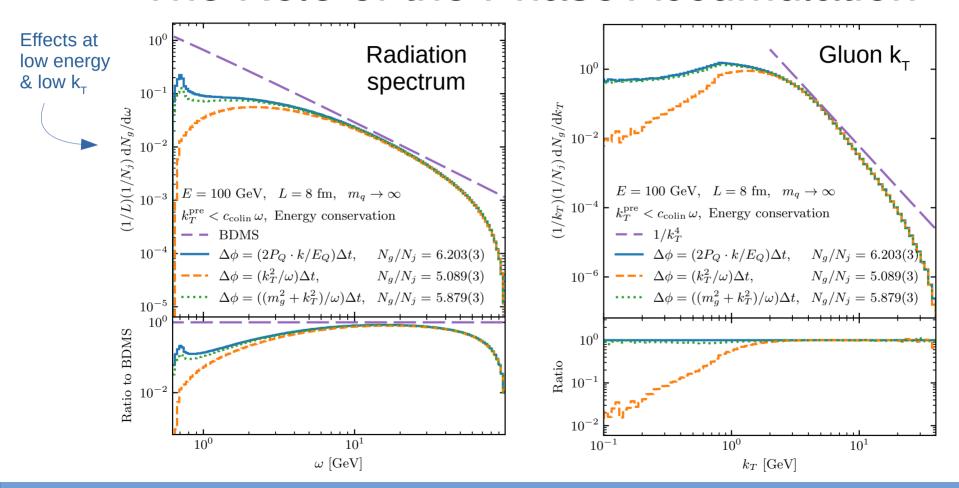
Including thermal gluon mass:

$$\Delta \varphi = \frac{m_g^2 + k_T^2}{\omega} \Delta t$$

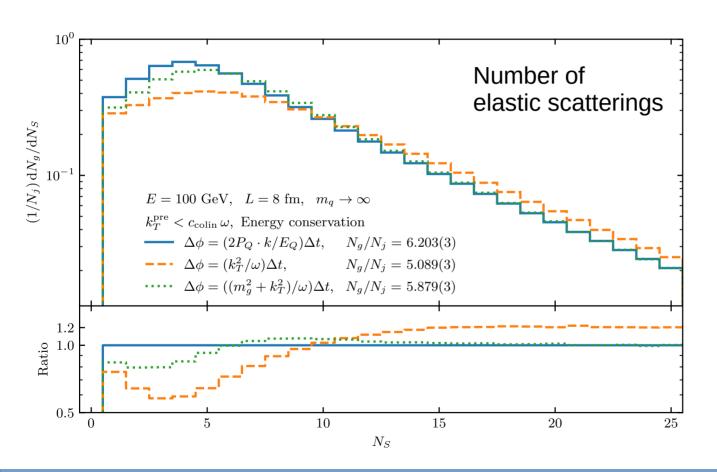
More general formula:

$$\Delta \varphi = \frac{2P_Q \cdot k}{E_Q} \Delta t$$

The Role of the Phase Accumulation



The Role of the Phase Accumulation



When neglecting the gluon mass in the phase accumulation, a larger path length is required to have a comparable overall mumber off radiations

Looking Forward: Towards More Realism

Next step:

- Interface with vHLLE to get hydro evolution of the medium
- Running strong coupling in elastic scatterings
- Start with high virtuality partons
- Sampling of initial parton p_⊤

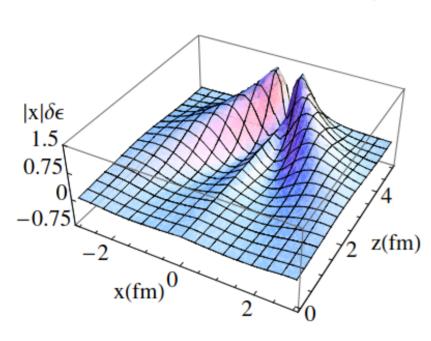
$$\frac{\mathrm{d}\sigma}{\mathrm{d}p_T} \sim p_T^{-6.5}$$

Run with hadronisation and jet finding

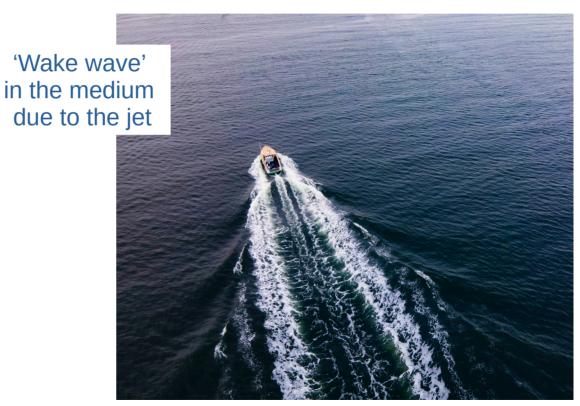


Looking Forward: Effect on the Medium

The jet also affects the medium



G.-Y. Qin, A. Majumder, H. Song, U. Heinz 0903.2255 [nucl-th]



Summary

- We have presented a new model for jet energy loss in heavy ion collisions
- Implementation in a Monte Carlo framework
- 1st step done:
 - Reproduction of the BDMS radiation energy spectrum
 - Shown effects of different model assumptions
- 2nd step: First results with hydro evolution interface to vHLLE
- **3**rd **step**: Implementation within the new EPOS4
 - EPOS4+JETS Initial state, hydro, and hadronisation from EPOS4

Thank you for your attention!