





Dilepton anisotropic flow in hadronic transport

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1. Dilepton as probes

Pairs of single-origin, opposite charge leptons

- No strong interaction \Rightarrow leave the medium undisturbed
- Multi-messenger: spectrometer, chronometer, barometer, thermometer, ...
- Experimental difficulties:
 - rare probes $BR(h \rightarrow l^+ l^-) \sim 10^{-5}$
 - combinatorial background





2 "dilepton" talks + few mentions!

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This work: Ag+Ag at $\sqrt{s_{NN}}$ = 2.55 GeV, 0-40%





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2. the SMASH approach

Simulating Many Strongly-interacting Hadrons

Hadrons

• Hadrons evolved with the relativistic Boltzmann equation (BUU)

 $p^{\mu}\partial_{\mu}f_{i}(x,p) + m_{i}F^{\alpha}\partial^{p}_{\alpha}f_{i}(x,p) = C^{i}_{\text{coll}}$

• Scatterings determined geometrically from "bottom-up" cross sections

$$\pi d_{\text{trans}}^2(a,b) < \sigma_{\text{tot}}(a,b) = \sum_R \sigma_{ab \to R} + \sum_{cd} \sigma_{ab \to cd}$$

• Mass-dependent width for <u>hadronic</u> decays

$$\frac{\operatorname{Prob}(R \operatorname{decays} \operatorname{in} \Delta t)}{\Delta t} = \Gamma_R^{\operatorname{vac}}(m) = \sum_{ab} \Gamma_{R \to ab}(m)$$

[J. Weil et al, PRC 94 (2016) 5, 054905]



https://smash-transport.github.io/



Vacuum properties a priori

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Dileptons

- Electromagnetic coupling is much smaller than strong coupling
- Stable hadrons decay at the end
- Time-integration: perturbative treatment for dilepton emission $w_{R \rightarrow l^+ l^-}(\tau) = \int_{-\infty}^{\tau - t_0} \frac{\mathrm{d}t}{\gamma} \Gamma_{R \rightarrow l^+ l^-}(m_R)$

[U. Heinz and K. S. Lee, Nucl. Phys. A 544, 503 (1992)] [G.-Q. Li and C. M. Ko, Nucl. Phys. A 582, 731 (1995)]

• At every time step the particle radiates a lepton pair, carrying the shining weight



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3. SMASH dileptons

Medium modifications are not fully reproduced by collisional broadening from vacuum interactions ArKCl at 1.76A GeV pp at 3.5 GeV all $- - \eta \rightarrow \gamma e^+ e^-$ ____ 10⁻² $--\cdot \quad \omega \rightarrow \pi^0 e^+ e^$ all $- - n \rightarrow \gamma e^+ e^ \rho \rightarrow e^+ e^-$ SMASH 10^{3} $\mu_B = 330 \,\mathrm{MeV}$ h $--\cdot \quad \omega \rightarrow \pi^0 e^+ e^ \Delta^0 \rightarrow ne^+e^ \rho \rightarrow e^+ e^ \rightarrow e^+ e^-$ ---- Rapp, Wambach 10^{-3} 10² $\Delta^0 \rightarrow ne^+ e^ \cdots \Delta^+ \rightarrow \mathrm{pe}^+ \mathrm{e}^ \phi \rightarrow e^+ e^-$ --- vacuum $N_{\pi^0}^{-1}~dN/dm~[1/GeV]$ $\begin{bmatrix} 10^{1} \\ \eta \\ 0 \end{bmatrix} \frac{10^{1}}{10^{-1}} \frac{10^{-1}}{10^{-2}}$ $\Delta^+ \rightarrow \rm pe^+ \, e^-$ 4 HADES $\rightarrow e^+ e^ \mathcal{A}^{ m med}(m)[m GeV$ $\pi^0 \rightarrow \gamma e^+ e^ 10^{-4}$ T[MeV] ♦ ♦ HADES $\rightarrow \gamma e^+ e^-$ 3 - 120 10⁻⁵ 15010⁻⁶ 2 -18010⁻⁷ 10⁻³ 10⁻⁸ 10-4 10⁻⁹ 0.0 0.0 0.4 0.6 0.8 0.2 1.2 1.0 0.2 0.4 0.6 0.8 1.0 1.2 0.40.60.8 1.0 1.2 $\mathbf{m}_{\mathrm{ee}}\;[\mathrm{GeV}]$ [J. Staudenmaier, H. Elfner, PRC 98 (2018) 3, 054908] $\mathrm{m_{ee}}~[\mathrm{GeV}]$ m[GeV]

[R. Rapp, J. Wambach, Eur.Phys.J.A 6 (1999) 415-420]
 [H. van Hees, R. Rapp, Nucl.Phys.A 806 (2008) 339-387]
 [RH, J. Staudenmaier, H. Elfner, PRC 107 (2022) 2, 025208]

Caveat

3. SMASH dileptons





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Dilepton flow in hadronic transport

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4. Anisotropic flow 101

Azimuthal Fourier decomposition of yields:

$$\frac{dN}{d\phi} \propto 1 + 2\sum_{n=1}^{\infty} v_n \cos\left[n(\phi - \Psi_R)\right]$$
Anisotropic flow coefficient

$$v_n \approx v_n \{ \text{RP} \} = \langle \cos[n(\phi - \Psi_R)] \rangle$$

Reaction plane $\, \Psi_R : \,$ spanned by beam axis and impact parameter







In particular: v₂ follows from initial geometry

5. Comparisons to HADES (preliminary)



- Out-of-plane flow in Dalitz $\pi^{\scriptscriptstyle 0}$ region
- Overall consistency to experiment
- "Zero" flow for $M_{ee} > m_{\pi^{\circ}}$
- Error bars (both experimental and from SMASH) are still very large!
- Centrality dependence follows expected trend (less central, more flow)





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5. Comparisons to HADES (preliminary)



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6. Checks within SMASH



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- Dilepton from resonances have flow!
- Cancellations lead to "zero" flow in higher masses: in-plane flow of Δ vs. out-of-plane flow of η

6. Checks within SMASH



• Self-consistency: low-mass SMASH dileptons follow charged pion flow



- Dilepton from resonances have flow!
- Cancellations lead to "zero" flow in higher masses: in-plane flow of Δ vs. out-of-plane flow of η

Still: error bars are very large

7. Can we improve this?

Scalar product method: correlate global hadron flow with differential dilepton flow

 $\operatorname{Cov}(\mathbf{q}^h, \mathbf{q}^{ll}(X)) = \left\langle \mathbf{q}^h \mathbf{q}^{ll}(X)^* \right\rangle - \left\langle \mathbf{q}^h \right\rangle \left\langle \mathbf{q}^{ll}(X)^* \right\rangle$

Eve flow vect

$$v_{2}^{ll}(X) = \frac{\left\langle |\mathbf{q}_{n}^{h}||\mathbf{q}_{n}^{ll}(X)|\cos[n(\Psi_{n}^{h} - \Psi_{n}^{ll})]\right\rangle_{\text{ev}}}{\left\langle \sqrt{\langle |\mathbf{q}_{n}^{h}|^{2}\rangle_{\text{ev}}}} \right|^{2\pi} \frac{1}{N^{h}} \sum_{j}^{N^{h}} e^{in\phi_{j}} \text{ final state}$$

$$\frac{\mathbf{q}_{n}^{ll}(\text{bin}) = \frac{1}{N^{ll}} \int_{\text{bin}} dt dy p_{T} dp_{T} \int_{0}^{2\pi} d\phi \frac{dN^{ll}}{d^{4}p} e^{in\phi_{j}} \xrightarrow{\text{SMASH}} \frac{1}{\sum w_{j}} \sum_{j \in \text{bin}}^{N^{ll}} w_{j} e^{in\phi_{j}} \text{ full evolution}$$

$$\frac{v_{2}^{ll}(X) = \frac{\left\langle |\mathbf{q}_{n}^{h}||\mathbf{q}_{n}^{ll}(X)|\cos[n(\Psi_{n}^{h} - \Psi_{n}^{ll})]\right\rangle_{\text{ev}}}{\sqrt{\langle |\mathbf{q}_{n}^{h}|^{2}\rangle_{\text{ev}}}} \xrightarrow{\text{Proposed for LHC energies:}} \begin{bmatrix} \text{Jean-François Paquet et al. PRC 93, 044906 (2016)} \\ \text{[Gojko Vujanovic et al. PRC 101, 044904 (2020)]} \end{bmatrix}$$



low resolution

6(2016)]))]

 $v_2(a) > v_2(b) > 0$ $v_2(c) = 0$ $v_2(d) < 0$

Usual geometric interpretation is lost!

8. Invariant mass



- Peak from Dalitz decay of $\pi^0 \rightarrow \gamma e^+ e^-$
- Small *non-zero* signal in resonance region
- Reverse centrality dependence!

9. Rapidity



- v₂(y) behavior seems to mimic multiplicity
- v₂{SP} is a correlation between time-integrated dileptons with final-state hadrons
- <u>Guess</u>: central events have larger contributions of π⁰/Δ

[Jan Staudenmaier, PhD thesis]

10. Transverse momentum

- Low mass, low p_T: bump structure with opposite centrality dependence
- At p_{τ} ~0.3 GeV: switch to intuitive behavior
- 👍 Resonance region: no unexpected structures



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11. Time dependence



11. Time dependence



- Bump appears only in the low mass region, and when the system is dilute
- Non-flow: resonances survive the medium but decay into $\pi^{\scriptscriptstyle 0}$

ightarrow They inherit the flow, but have a smaller $m p_{ au}$



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13. Dilepton "tagging"

$$\begin{aligned} v_2^{ll}(X) &= \frac{\left< |\mathbf{q}_n^h| |\mathbf{q}_n^{ll}(X)| \cos[n(\Psi_n^h - \Psi_n^{ll})] \right>_{\mathrm{ev}}}{\sqrt{\left< |\mathbf{q}_n^h|^2 \right>_{\mathrm{ev}}}} \\ \end{aligned} \\ \end{aligned} \\ \begin{aligned} & \text{We can choose which hadrons} \\ & \text{are correlated to dileptons!} \end{aligned}$$

13. Dilepton "tagging"



$$v_2^{ll}(X) = \frac{\left\langle |\mathbf{q}_n^h| |\mathbf{q}_n^{ll}(X)| \cos[n(\Psi_n^h - \Psi_n^{ll})] \right\rangle_{\text{ev}}}{\sqrt{\left\langle |\mathbf{q}_n^h|^2 \right\rangle_{\text{ev}}}}$$

We can choose which hadrons are correlated to dileptons!

- When a ∆ decays into dileptons, it also produces a nucleon
- Protons are then a natural probe for $\Delta \rightarrow Ne^+e^-$ processes



13. Dilepton "tagging"



$$\psi_2^{ll}(X) = \frac{\left\langle |\mathbf{q}_n^h| |\mathbf{q}_n^{ll}(X)| \cos[n(\Psi_n^h - \Psi_n^{ll})] \right\rangle_{\text{ev}}}{\sqrt{\left\langle |\mathbf{q}_n^h|^2 \right\rangle_{\text{ev}}}}$$

We can choose which hadrons are correlated to dileptons!

- When a ∆ decays into dileptons, it also produces a nucleon
- Protons are then a natural probe for $\Delta \rightarrow Ne^+e^-$ processes
- Charged pions ~ neutral pion. Using them isolates the low mass region and the bump
- Intuitive centrality dependence restored!



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Conclusions

- Scalar Product method has smaller fluctuations than reaction plane, currently used by HADES
- Geometric interpretation of flow is lost and replaced with correlations, which don't always have a clear origin
- Hadronic transport allows for demystification of this signal
- Different tags are accessible in experiment and can distinguish sources of correlations!

Outlook

- Include nuclear potentials for realistic calculation
- Compare to coarse-grained signal and to the new dynamic initialization of our hybrid model



Thank you for the attention!