EXPLORING ELASTIC AND RADIATIVE JET QUENCHING IN THE STRONGLY INTERACTING QCD MEDIUM

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PhysRevC.106.014903
arXiv:2308.03105

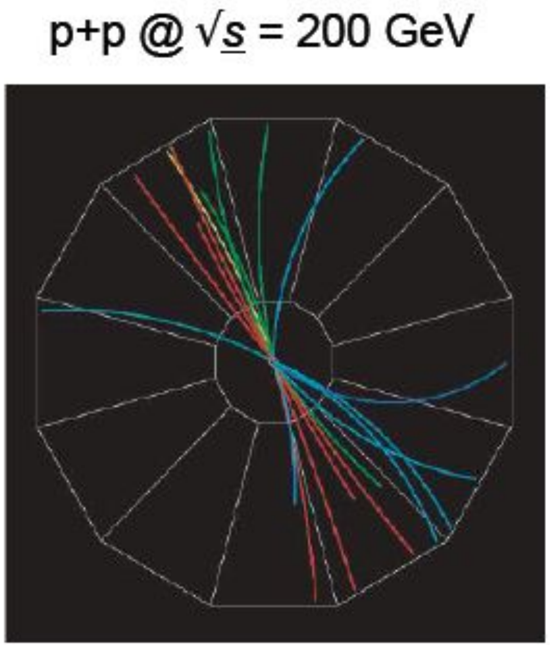
OUTLINE

- Introduction: jets
- Dynamical QuasiParticle Model (DQPM)
- Elastic and inelastic cross sections
- Transport coefficients in kinetic theory
- Summary

INTRODUCTION

What is jet?

A jet is a collimated spray of hadrons generated via successive parton branchings, starting with a highly energetic and highly virtual parton (quark or gluon) produced by the collision



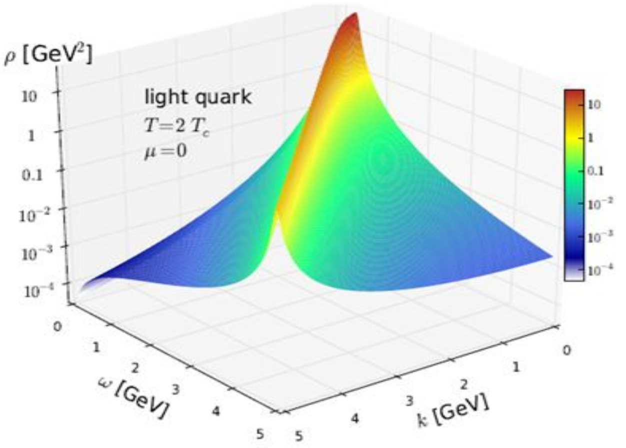
Why do we study jets?

- Early formation time
- Not thermalized in the medium
- Contain the information on the QGP properties

DYNAMICAL QUASIPARTICLE MODEL (DQPM)

- DQPM – effective model for the description of non-perturbative (strongly interacting) QCD based on lQCD EoS
- The QGP phase is described in terms of interacting quasiparticles - massive quarks and gluons - with Lorentzian spectral functions:

$$\rho_j(\omega, \mathbf{p}) = \frac{4\omega\gamma_j}{\left(\omega^2 - \mathbf{p}^2 - M_j^2\right)^2 + 4\gamma_j^2\omega^2}$$



P. Moreau et al., PRC 100, 014911 (2019)

- Field quanta are described in terms of dressed propagators with complex self-energies:

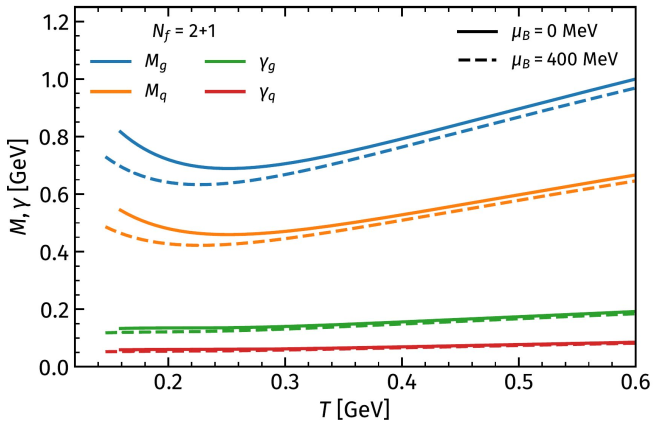
gluon propagator: $\Delta^{-1} = P^2 - \Pi$; quark propagator: $S_q^{-1} = P^2 - \Sigma_q$

gluon self-energy: $\Pi = M_g^2 - 2i\gamma_g\omega$; quark self-energy: $\Sigma_q = M_q^2 - 2i\gamma\omega$

- Real part of the self-energy - thermal masses
- Imaginary part of the self-energy - interaction widths of partons

DQPM INGREDIENTS

Masses and widths of quasiparticles depend on the temperature of the medium and μ_B

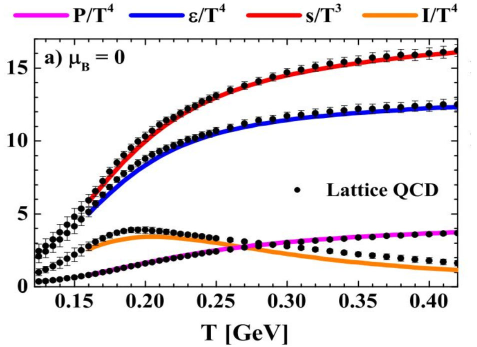


$$m_g^2(T,\mu_B) = C_g \frac{g^2(T,\mu_B)}{6} T^2 \left(1 + \frac{N_f}{2N_c} + \frac{1}{2}\frac{\sum_q \mu_q^2}{T^2\pi^2}\right)$$

$$m_{q(\bar{q})}^2(T,\mu_B) = C_q \frac{g^2(T,\mu_B)}{4} T^2 \left(1 + \frac{\mu_q^2}{T^2\pi^2}\right)$$

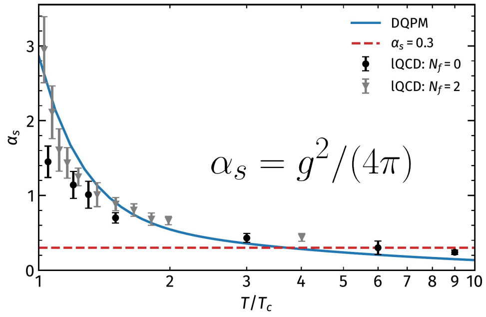
$$\gamma_j(T,\mu_B) = \frac{1}{3} C_j \frac{g^2(T,\mu_B)T}{8\pi} \ln\left(\frac{2c_m}{g^2(T,\mu_B)} + 1\right)$$

Input: entropy density vs T for μ_B=0



$$g^2(s/s_{SB}) = d((s/s_{SB})^e - 1)^f$$

$$s_{SB}^{QCD} = 19/9\pi^2 T^3$$



O. Kaczmarek, F. Zantow, Phys. Rev. D 71, 114510

DQPM: SUMMARY

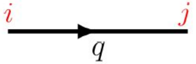
There are four effects that make the DQPM different from the “pure” pQCD:

- non-perturbative origin of the strong coupling which depends on (T, μ_B);
- finite masses of the intermediate parton propagators (screening masses);
- finite masses of the medium partons;
- finite widths of partons.

PARTONIC ELASTIC INTERACTIONS

DQPM partonic interactions are described in terms of leading order diagrams:

quark propagator:

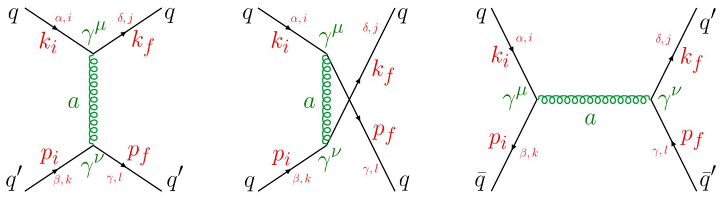


$$= i\delta_{ij}\frac{\not{q} + M_q}{q^2 - M_q^2 + 2i\gamma_q q_0}$$

gluon propagator:

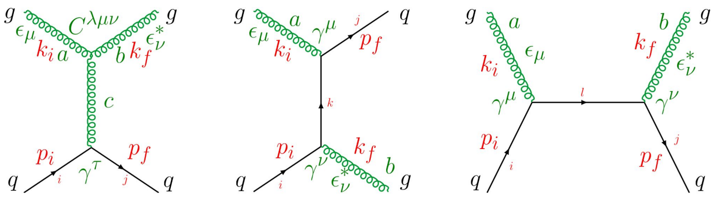
$$= -i\delta_{ab}\frac{g^{\mu\nu} - q^\mu q^\nu / M_g^2}{q^2 - M_g^2 + 2i\gamma_g q_0}$$

qq' → qq' scattering



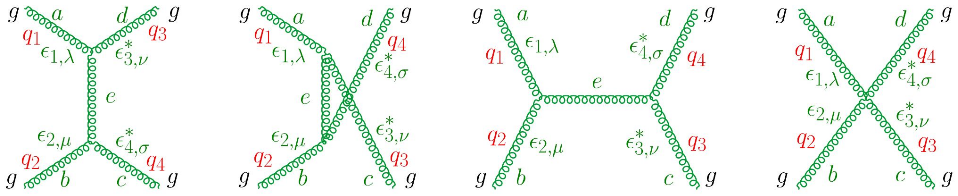
$t-$channel $u-$channel $s-$channel

qg → qg scattering



$t-$channel $u-$channel $s-$channel

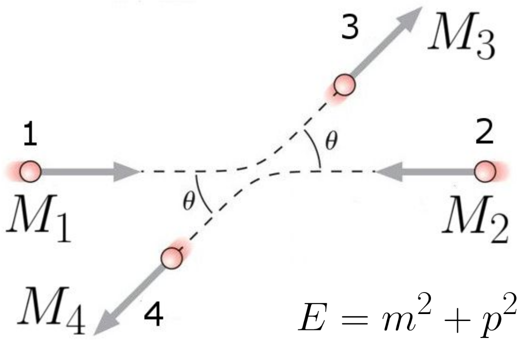
gg → gg scattering



$t-$channel $u-$channel $s-$channel $4-$point

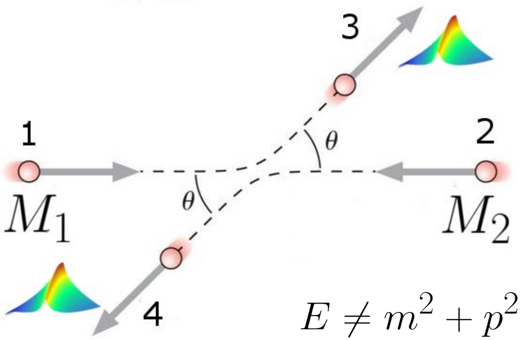
DQPM PARTONIC CROSS SECTIONS

On-shell: final masses = pole masses



$$d\sigma^{\rm on} = \frac{d^3p_3}{(2\pi)^3 2E_3}\frac{d^3p_4}{(2\pi)^3 2E_4}(2\pi)^4\delta^{(4)}\left(p_1+p_2-p_3-p_4\right)\frac{|\bar{\mathcal{M}}|^2}{F}$$

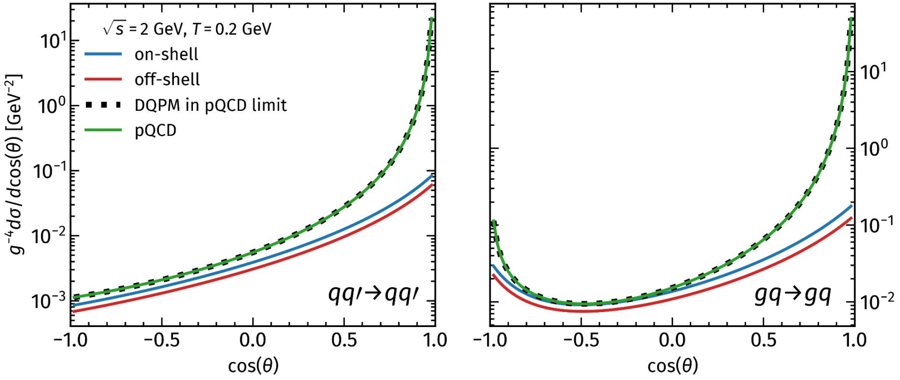
Off-shell: integration over final masses



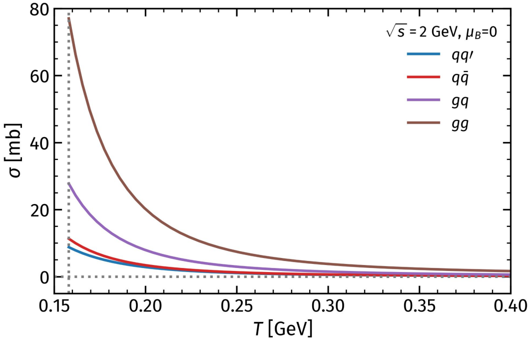
$$F d\sigma^{\rm off} = \frac{d^4p_3}{(2\pi)^4}\ \frac{d^4p_4}{(2\pi)^4}\ \tilde{\rho}_3(\omega_3,\mathbf{p}_3)\ \theta(\omega_3)\ \tilde{\rho}_4(\omega_4,\mathbf{p}_4)\ \theta(\omega_4)$$
$$\times\,(2\pi)^4\delta^{(4)}\left(p_1+p_2-p_3-p_4\right)|\bar{\mathcal{M}}|^2$$

DQPM PARTONIC CROSS SECTIONS

DQPM angular dependence for differential cross sections (scaled by g^4)



DQPM total cross sections



✓ DQPM reproduces pQCD cross sections for masses and widths → 0

✓ DQPM angular distribution is more “isotropic” then pQCD

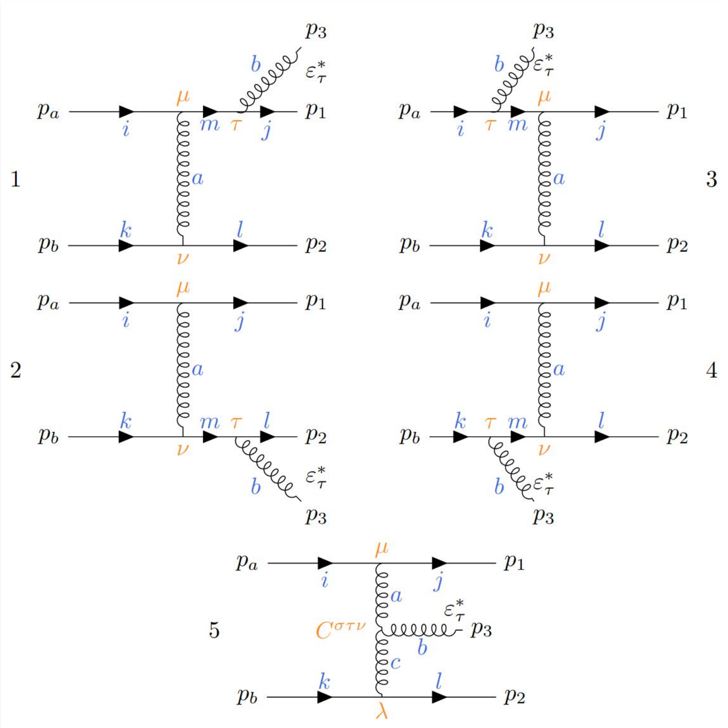
✓ the off-shell effects are small for energetic partons and for high T

✓ strong T-dependence

PARTONIC INELASTIC INTERACTIONS: Q+Q $\longrightarrow$ Q+Q+G

pQCD result: F. A. Berends et al., Phys. Lett., B103, 124 (1981)

***t*-channel**



$$\Pi_{\mu\nu}(k) = \left[-i\frac{g_{\mu\nu} - (k_\mu k_\nu)/M_g^2}{k^2 - M_g^2 + 2i\gamma_g\omega_k}\right] \quad \text{(gluon propagator)},$$

$$\Lambda(k) = \left[i\frac{\not{k} + M_q}{k^2 - M_q^2 + 2i\gamma_q\omega_k}\right] \quad \text{(quark propagator)},$$

$$V^{\nu,a}_{ik} = (-ig\gamma^\nu T^a_{ik}) \quad \text{(vertex)},$$

$$i\mathcal{M}_1 = \bar{u}^l(p_2)V^{\nu,a}_{lk}u^k(p_b)\Pi_{\mu\nu}(p_b - p_2)\bar{u}^j(p_1)\varepsilon^*_\tau(p_3)V^{\tau,b}_{jm}\Lambda(p_1 + p_3)V^{\mu,a}_{mi}u^i(p_a)$$

$$i\mathcal{M}_2 = \bar{u}^j(p_1)V^{\mu,a}_{ji}u^i(p_a)\Pi_{\mu\nu}(p_a - p_1)\bar{u}^l(p_2)\varepsilon^*_\tau(p_3)V^{\tau,b}_{lm}\Lambda(p_2 + p_3)V^{\nu,a}_{mk}u^k(p_b)$$

$$i\mathcal{M}_3 = \bar{u}^l(p_2)V^{\nu,a}_{lk}u^k(p_b)\Pi_{\mu\nu}(p_b - p_2)\bar{u}^j(p_1)V^{\mu,a}_{jm}\Lambda(p_a - p_3)\varepsilon^*_\tau(p_3)V^{\tau,b}_{mi}u^i(p_a)$$

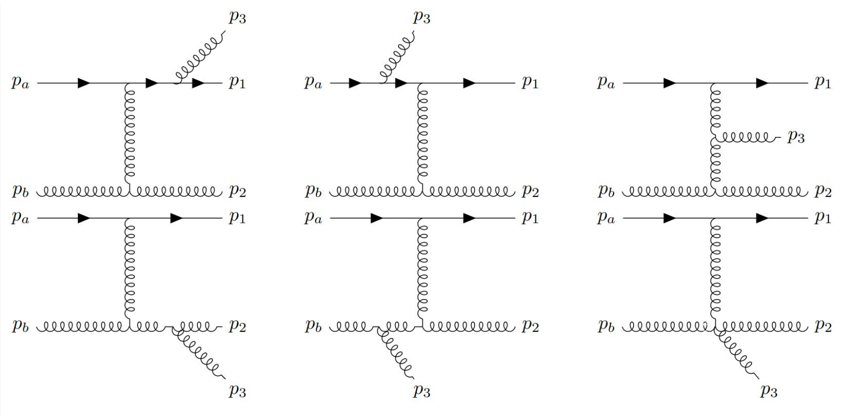
$$i\mathcal{M}_4 = \bar{u}^j(p_1)V^{\mu,a}_{ji}u^i(p_a)\Pi_{\mu\nu}(p_a - p_1)\bar{u}^l(p_2)V^{\nu,a}_{lm}\Lambda(p_b - p_3)\varepsilon^*_\tau(p_3)V^{\tau,b}_{mk}u^k(p_b)$$

$$i\mathcal{M}_5 = \bar{u}^j(p_1)V^{\mu,a}_{ji}u^i(p_a)\ \bar{u}^l(p_2)V^{\lambda,c}_{lk}u^k(p_b)\Pi_{\mu\nu}(p_a - p_1)$$
$$\times\ \Pi_{\lambda\sigma}(p_b - p_2)\varepsilon^*_\tau(p_3)\left(-gf^{abc}C^{\sigma\tau\nu}(p_b - p_2, -p_3, p_2 - p_b + p_3)\right)$$

✓ emitted gluon is massive!

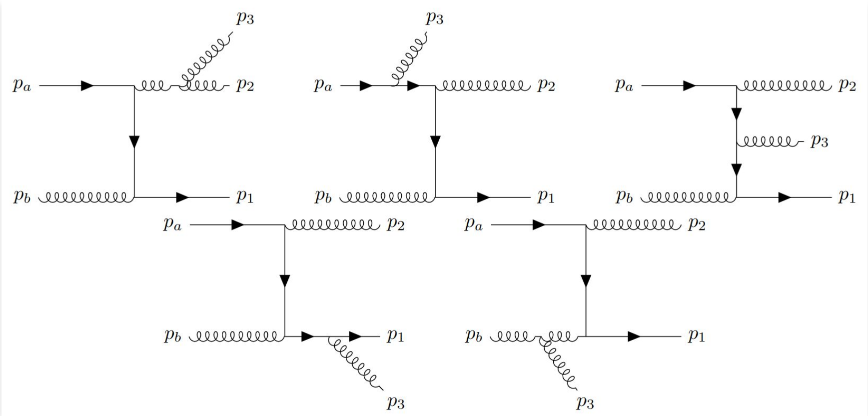
PARTONIC INELASTIC INTERACTIONS: Q+G → Q+G+G

***t*-channel**

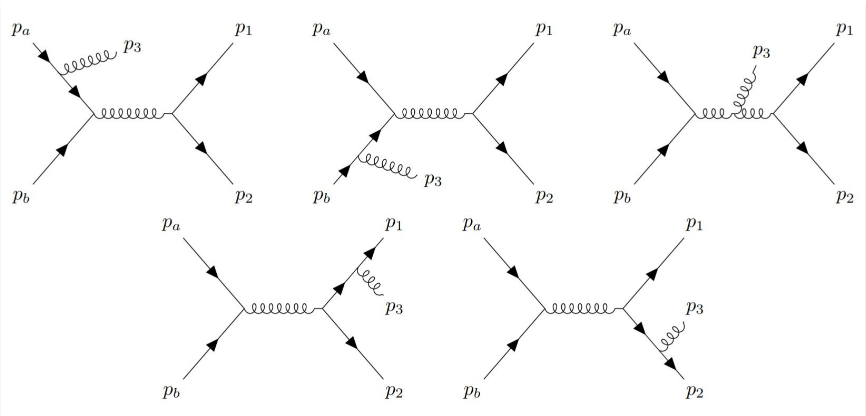


→ most dominant

***u*-channel**

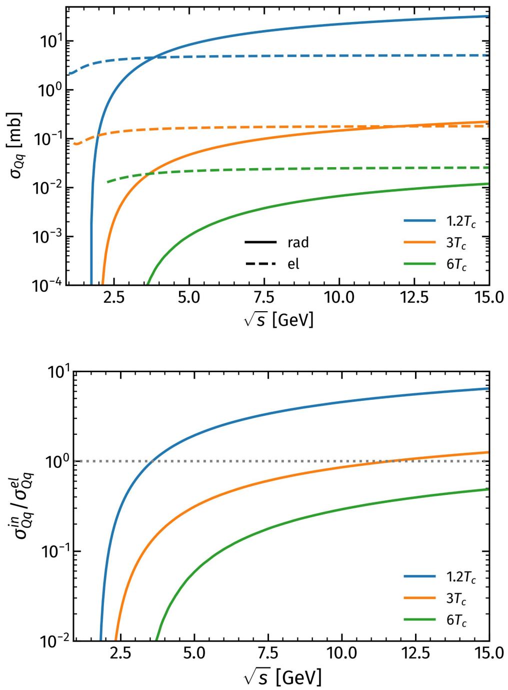


***s*-channel**

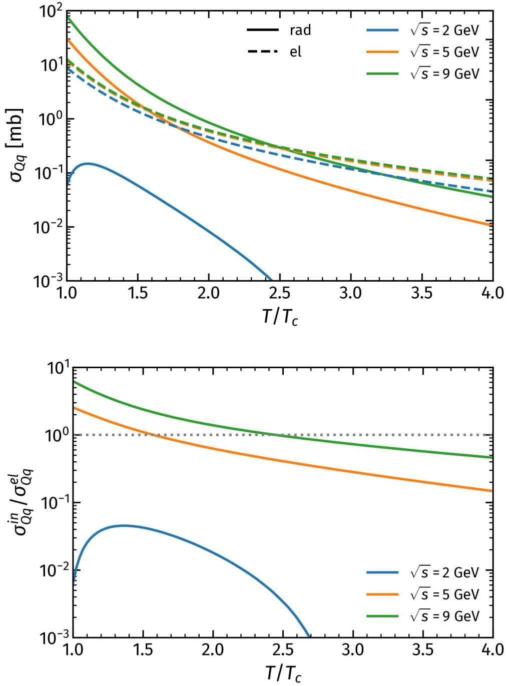


PARTONIC CROSS SECTIONS: ELASTIC VS INELASTIC

Energy dependence



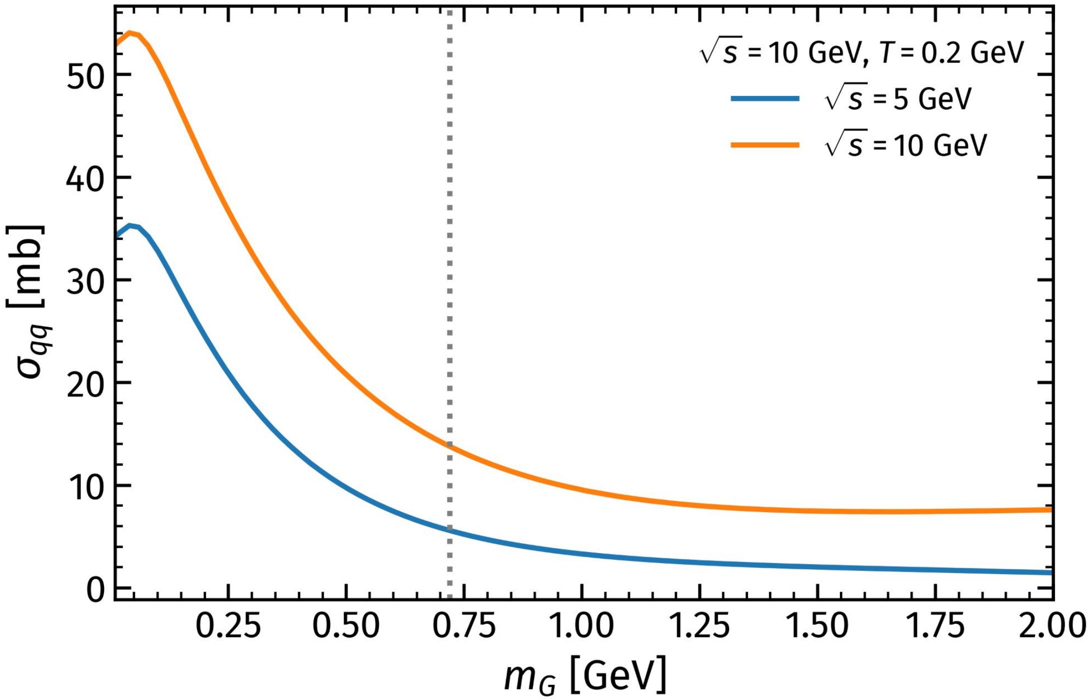
Temperature dependence



- ✓ suppression of radiative cross section for small energies
- ✓ enhancement of radiative cross section for small temperatures

PARTONIC CROSS SECTIONS: EMITTED GLUON MASS

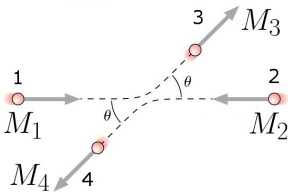
Emitted gluon mass dependence



TRANSPORT COEFFICIENTS IN KINETIC THEORY

On-shell:

- integration over momentums
- masses = pole masses

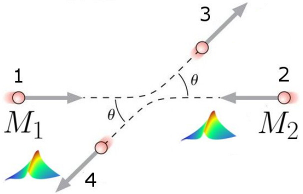


$$E^2 = m^2 + p^2$$

$$\langle \mathcal{O} \rangle^{\text{on}} = \frac{1}{2E_i} \sum_{j=q,\bar{q},g} d_j f_j \int \frac{d^3 p_j}{(2\pi)^3 2E_j} \times \int \frac{d^3 p_1}{(2\pi)^3 2E_1} \int \frac{d^3 p_2}{(2\pi)^3 2E_2} \times (1 \pm f_1)(1 \pm f_2) \mathcal{O} |\overline{\mathcal{M}}|^2 (2\pi)^4 \delta^{(4)}(p_i + p_j - p_1 - p_2)$$

Off-shell:

- integration over momentums
- + two additional integrations over medium partons energy



$$\frac{1}{2E} \to \int \frac{d\omega}{(2\pi)} \rho(\omega, \mathbf{p}) \theta(\omega)$$

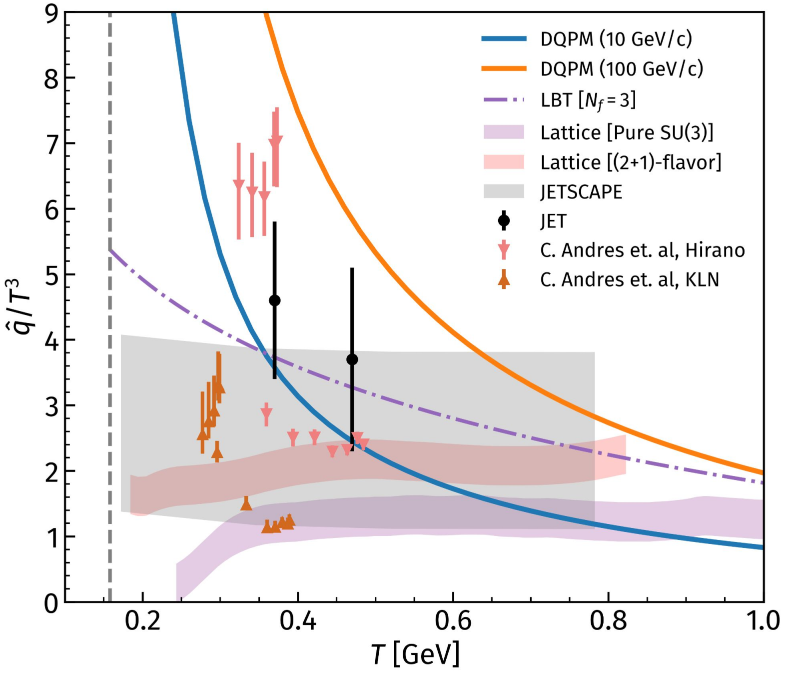
$$\langle \mathcal{O} \rangle^{\text{off}} = \frac{1}{2E_i} \sum_{j=q,\bar{q},g} d_j f_j \int \frac{d^4 p_j}{(2\pi)^4} \rho(\omega_j, \mathbf{p}_j) \theta(\omega_j) \times \int \frac{d^3 p_1}{(2\pi)^3 2E_1} \int \frac{d^4 p_2}{(2\pi)^4} \rho(\omega_2, \mathbf{p}_2) \theta(\omega_2) \times (1 \pm f_1)(1 \pm f_2) \mathcal{O} |\overline{\mathcal{M}}|^2 (2\pi)^4 \delta^{(4)}(p_i + p_j - p_1 - p_2)$$

$$\mathcal{O} = |\vec{p_T} - \vec{p_T}'|^2 \to \langle O \rangle = \hat{q}$$

$$\mathcal{O} = (E - E') \to \langle O \rangle = dE/dx$$

RESULTS: Q-HAT FROM ELASTIC PROCESSES

The DQPM q-hat(T) for elastic scattering of a jet quark vs other models



DQPM: I.Grishmanovskii, T.Song, O.Soloveva, C.Greiner, E.Bratkovskaya, *PRC 106, 014903*

JET: K. M. Burke et al., *PRC 90, 014909 (2014)*

lQCD: A. Kumar et al., PRD.106.034505

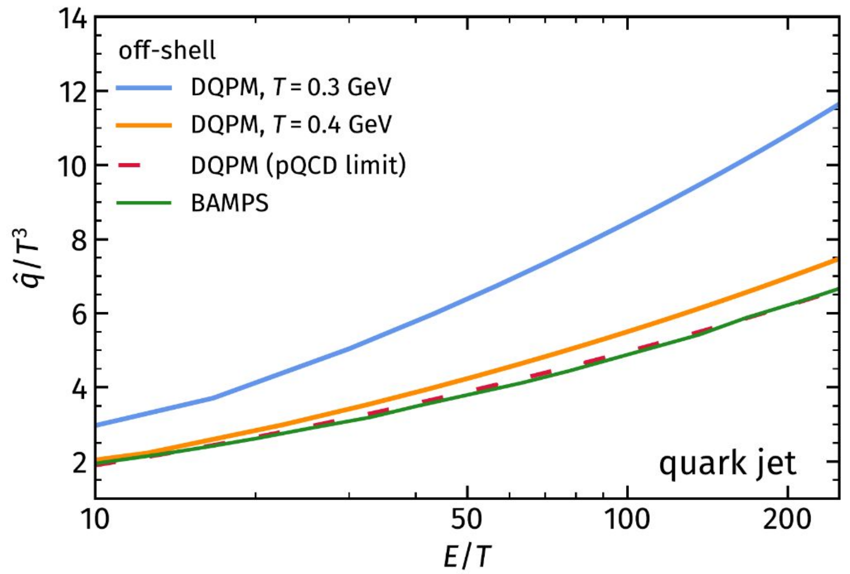
LBT: Y. He et al., *PRC 91 (2015)*

JETSCAPE: S. Cao *et al. PRC 104, 024905 (2021)*

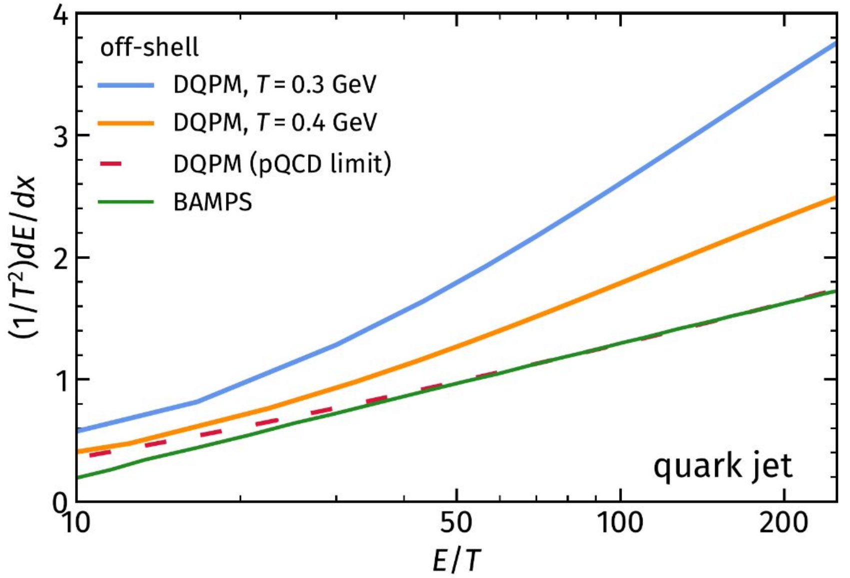
BDMPS: C.Andres *et al., Eur.Phys.J.C 76 (2016) 9, 475*

RESULTS: Q-HAT AND ENERGY LOSS

Energy dependence of the scaled q-hat

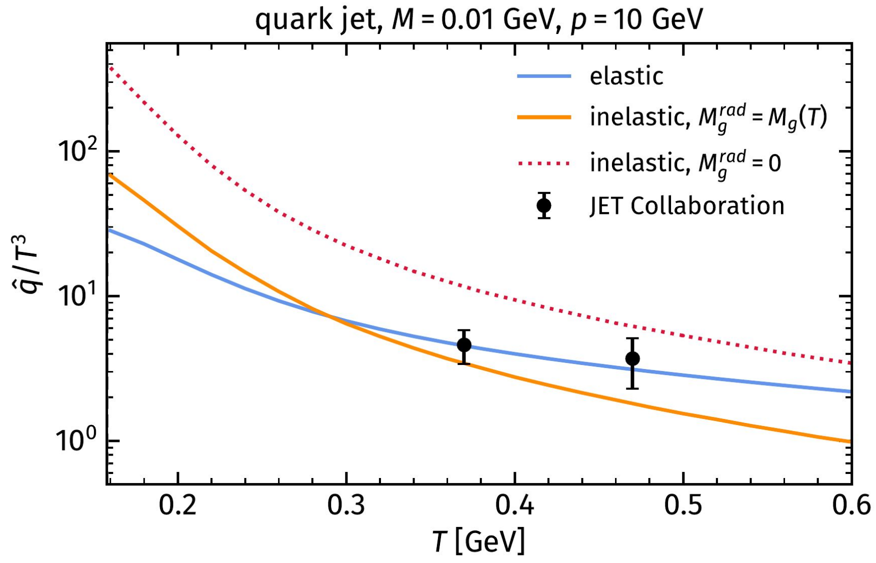


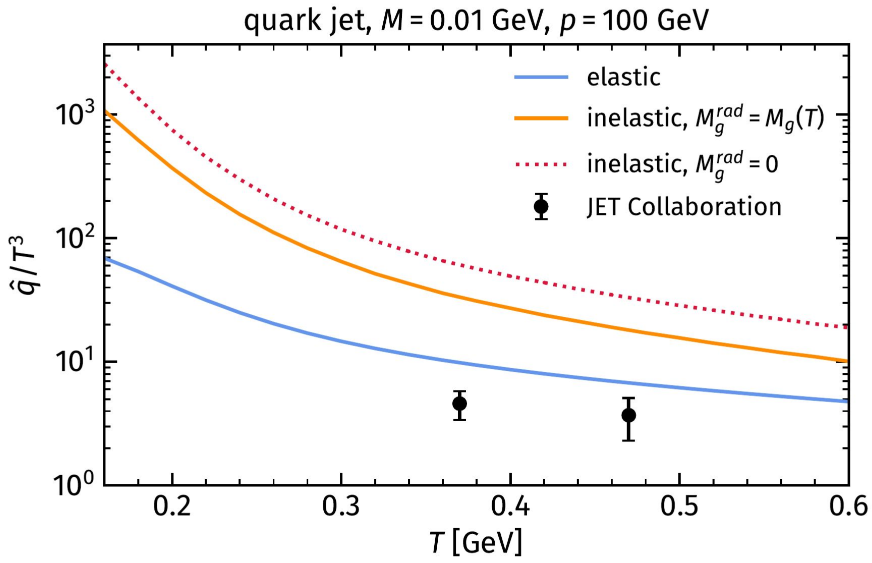
Energy dependence of the scaled energy loss dE/dx



✓ all models predict logarithmic growth of q-hat and dE/dx with jet energy

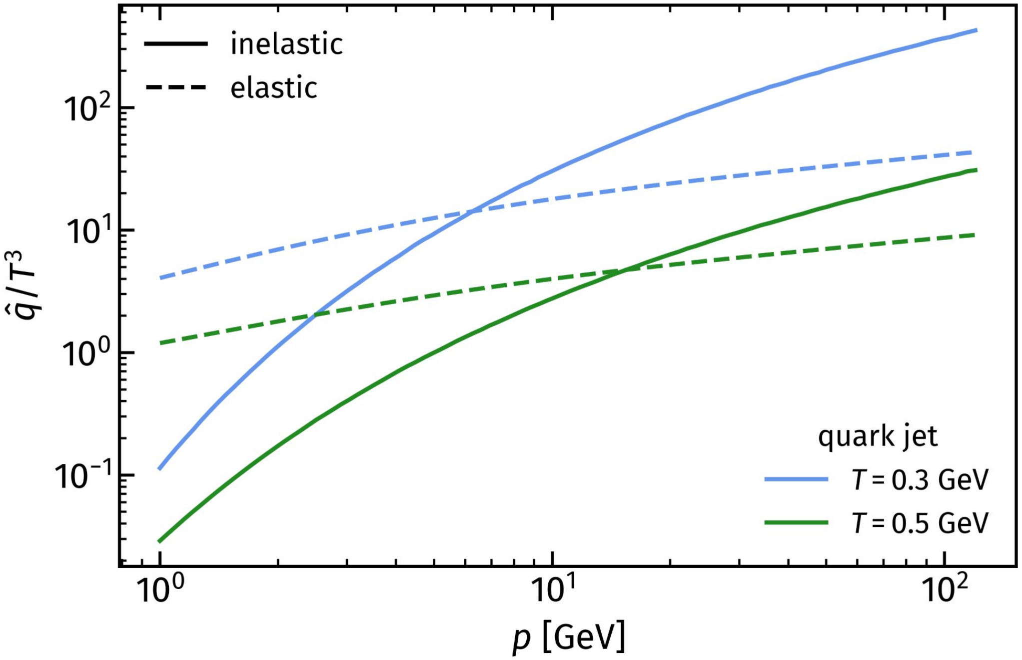
RESULTS: Q-HAT FROM ELASTIC + INELASTIC PROCESSES



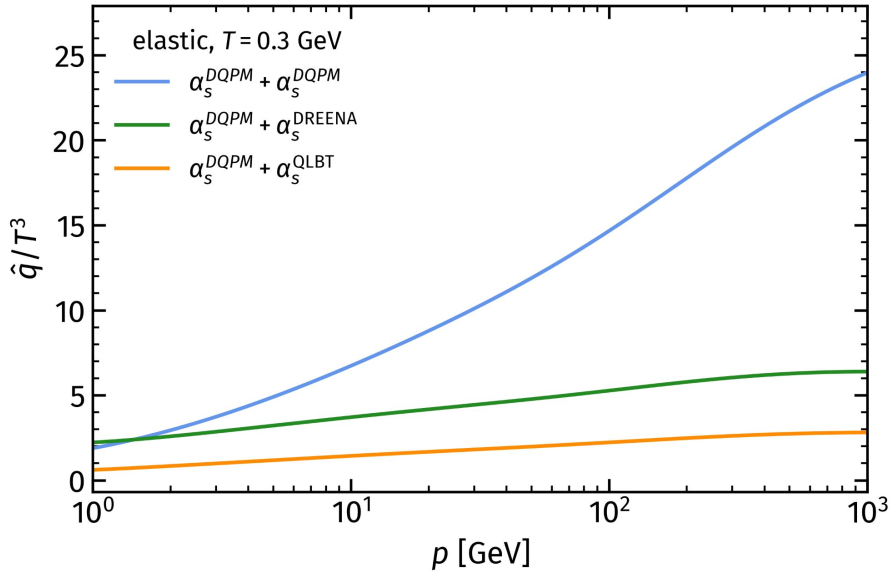


- ✓ inelastic q-hat is suppressed for low jet momentum, but is significant for high momentum
- ✓ emitted gluon mass is important

RESULTS: Q-HAT FROM ELASTIC + INELASTIC PROCESSES



RESULTS: RUNNING COUPLING DEPENDENCE (PRELIMINARY)



$$\alpha_s^{\text{jet}} \neq \alpha_s^{\text{QGP}}$$

$$\alpha_s^{\text{DQPM}} \to \alpha_s(ET) \text{ or } \alpha_s(T)$$

1. $\alpha(ET) = \dfrac{4\pi}{(11 - \frac{2}{3}n_f)} \dfrac{1}{\ln\left(\frac{ET}{\Lambda^2}\right)}$

2. $g^2(E) = \dfrac{48\pi^2}{(11N_c - 2N_f)\ln\left[(AE/T_c + B)^2\right]},$

$$\alpha_s = g^2/(4\pi)$$

1. B. Karmakar, D. Zigic, I. Salom et al., arXiv:2305.11318
2. F.Liu, X.-Y. Wu, S. Cao et al., arXiv:2304.08787

OUTLOOK

Summary:

- Elastic and inelastic cross sections are calculated within DQPM
- Transport coefficients (q-hat and dE/dx) are evaluated for the propagation of the jet parton through the strongly interacting QGP based on the DQPM
- DQPM predicts stronger energy loss than pQCD models
- DQPM reproduces the pQCD limits for zero masses and widths of medium partons

Future:

- Implementing inelastic $2 \to 3$ cross sections into full transport simulation (PHSD)