

# Locating QCD's critical end point

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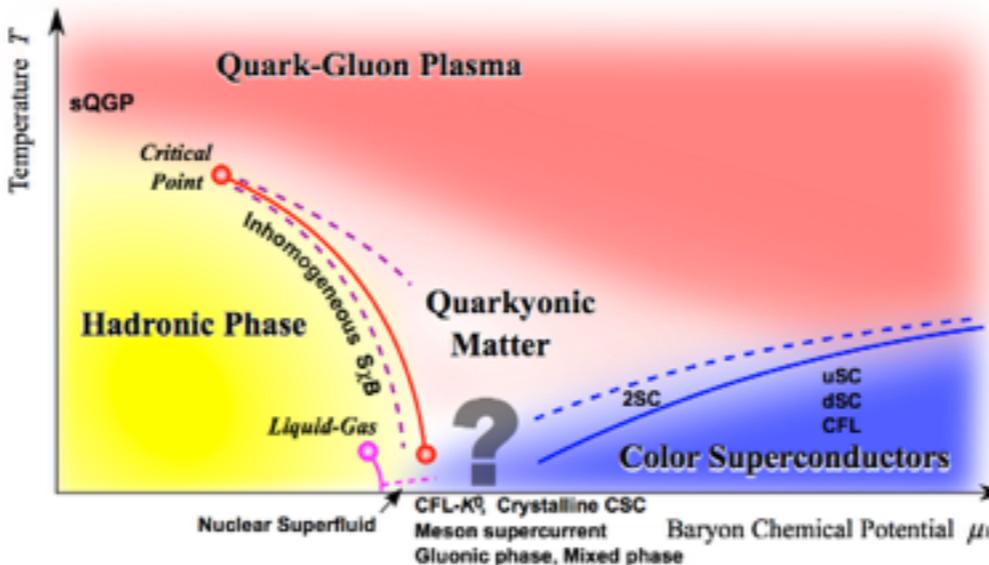
31st of Oct 2016



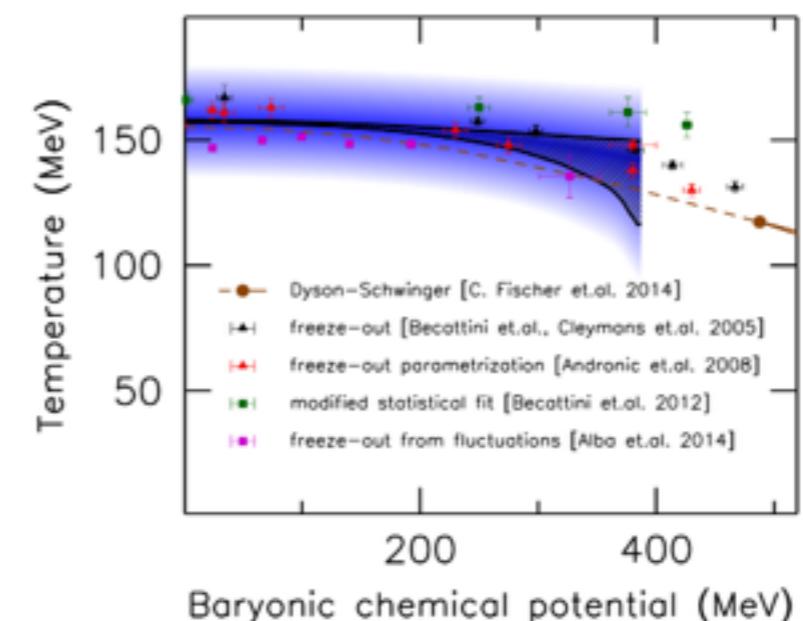
Eichmann, CF, Welzbacher, PRD93 (2016) [1509.02082]  
Eichmann, Sanchis-Alepuz, Williams, Alkofer, CF, PPNP in press [1606.09602]

# Overview

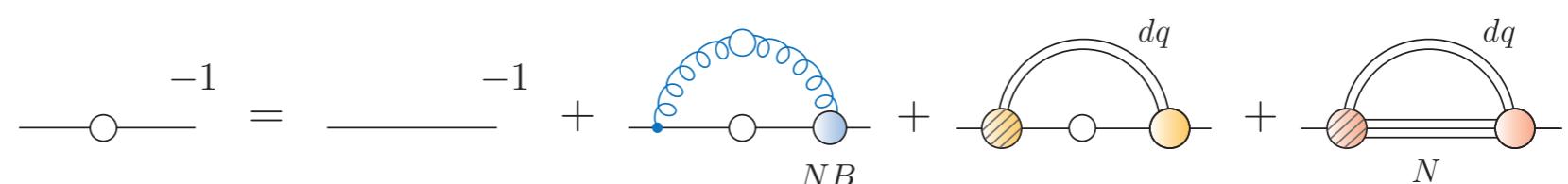
## I. Introduction



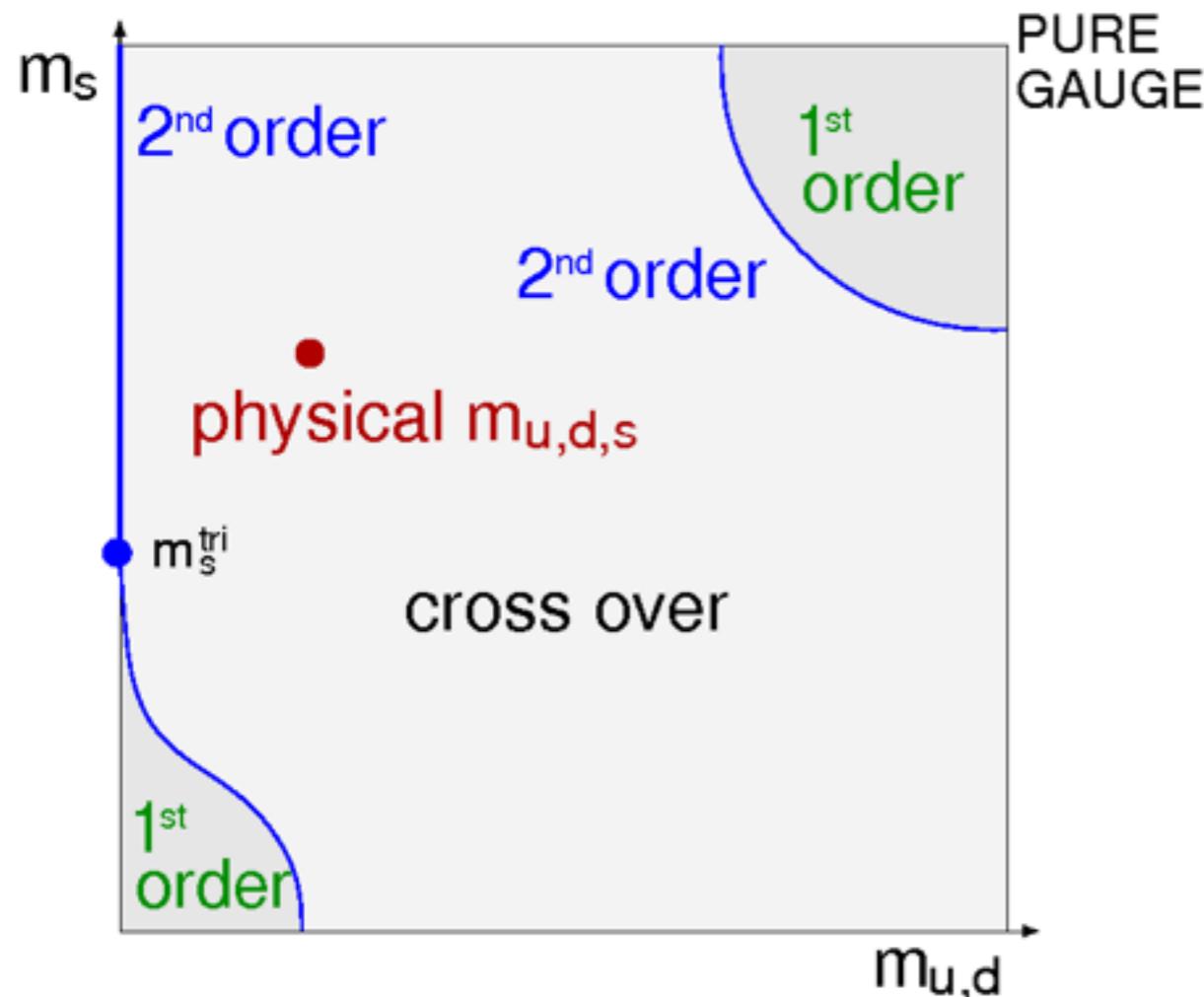
## 2. Gluons, quarks and the CEP



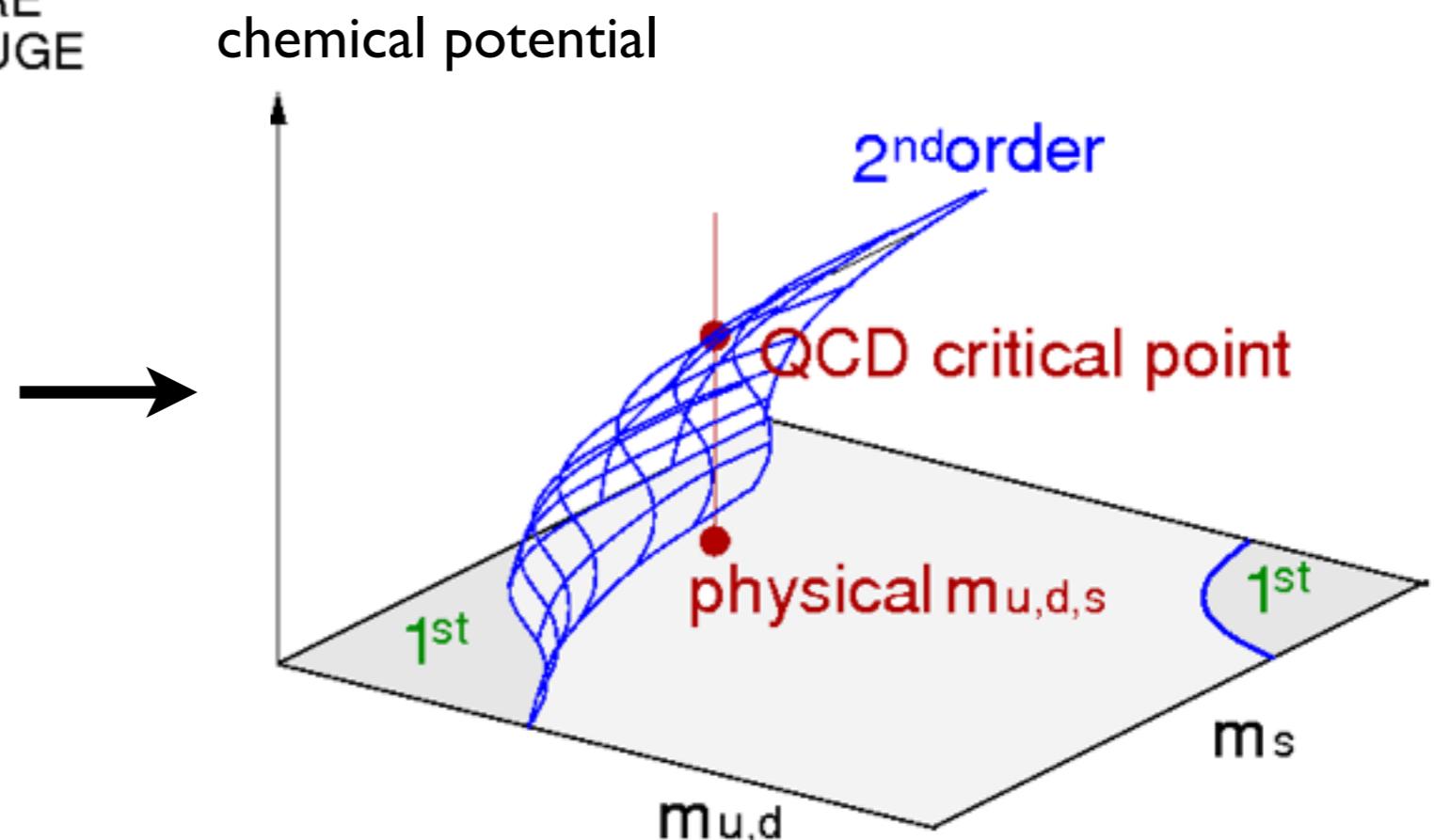
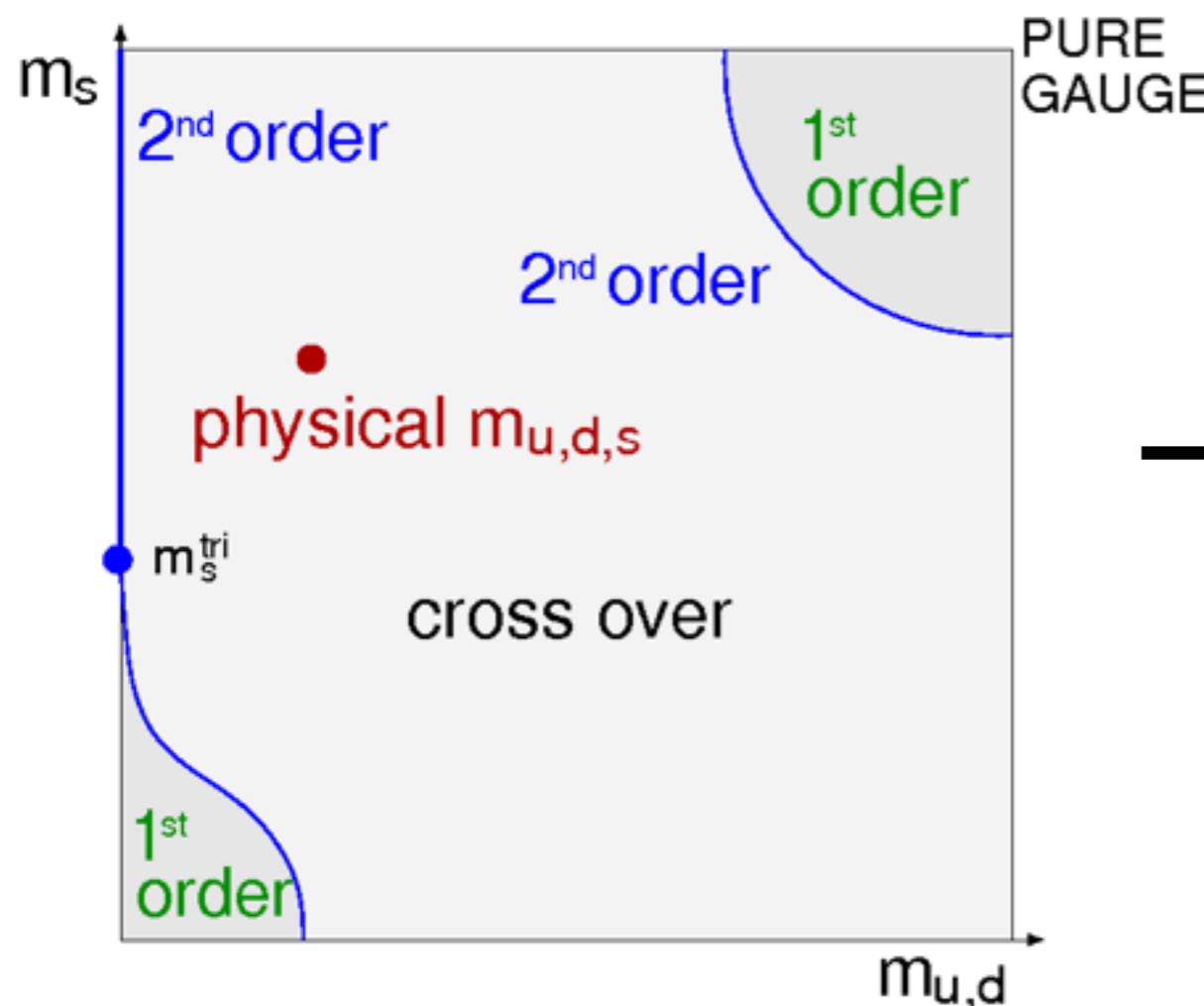
## 3. Baryon effects on the CEP



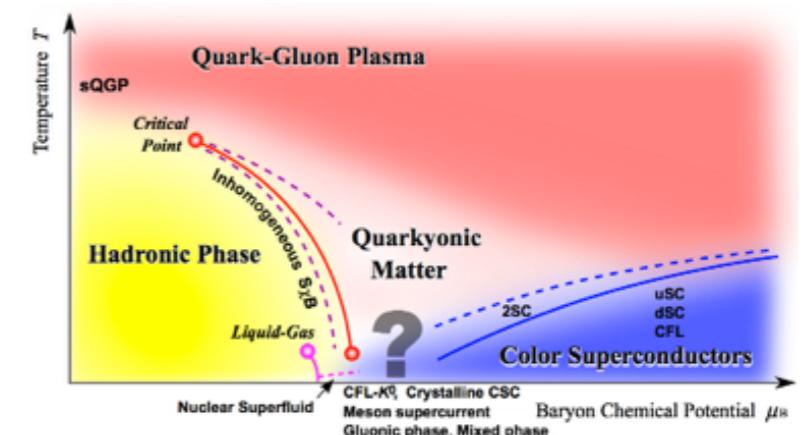
# QCD phase transitions



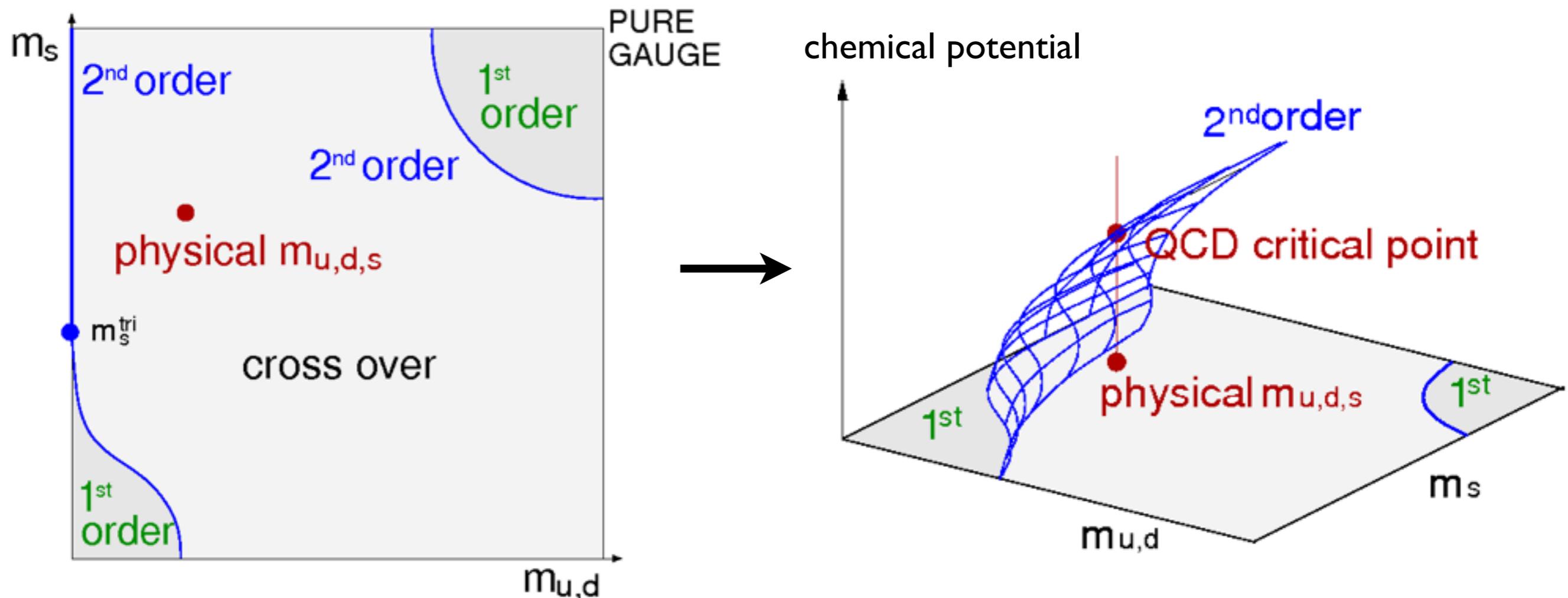
# QCD phase transitions



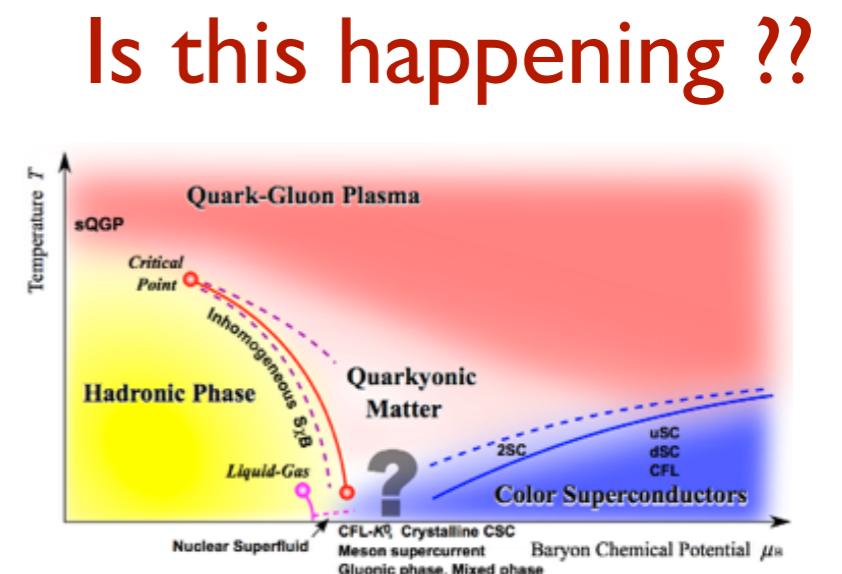
Is this happening ??



# QCD phase transitions



- Lattice-QCD
  - present: extrapolation
  - future: exact methods ?
- DSE/FRG
  - not exact, but allow for '10%-physics'



# Search for the CEP

- Taylor expansion ( $N_f=2$ ):

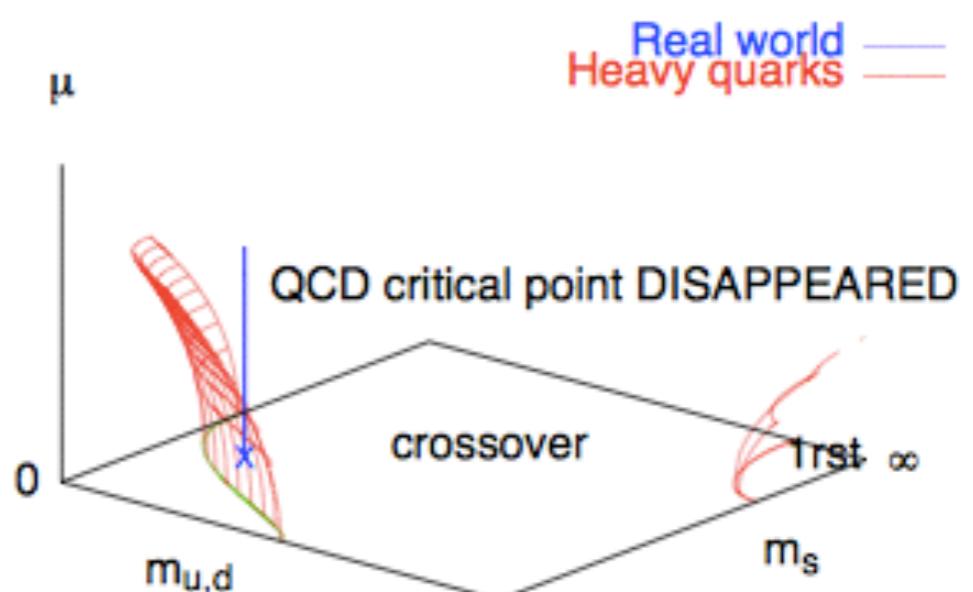
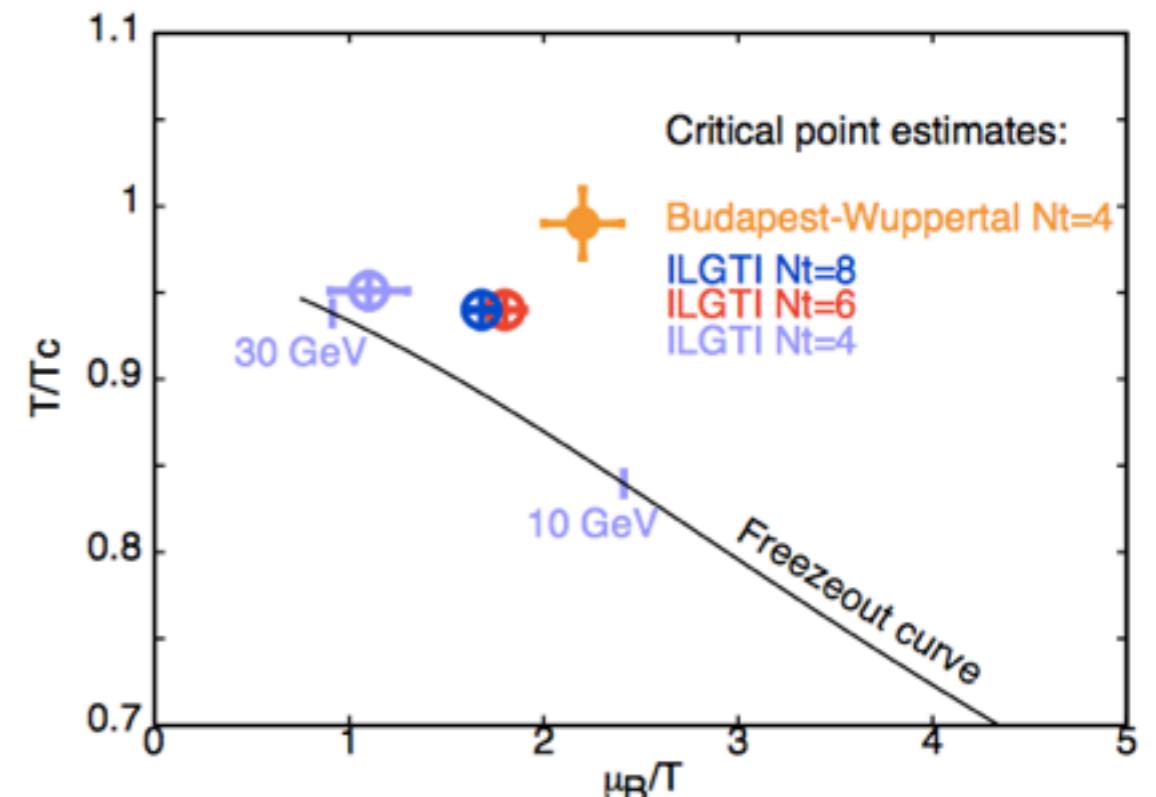
Datta, Gavai and Gupta, NPA 904–905 (2013) 883c  
Gavai, Gupta, PRD 71 (2005) 114014

- Reweighting ( $N_f=2+1$ ):

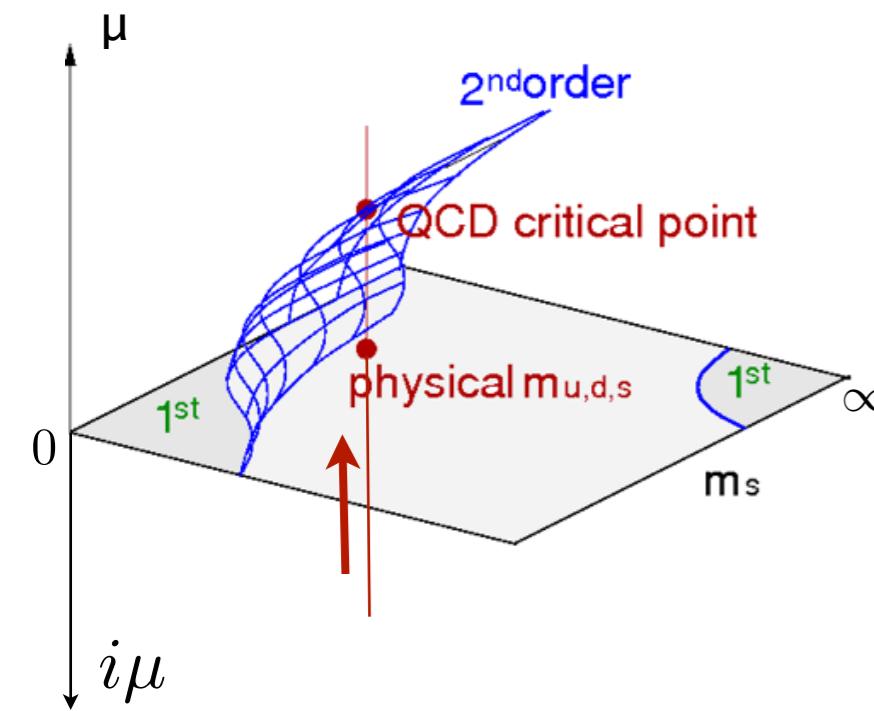
Fodor, Katz, JHEP 0404 (2004) 050

- Analytic continuation ( $N_f=3$ ):

de Forcrand, Philipsen, JHEP 0811 (2008) 012;  
NPB 642 (2002) 290

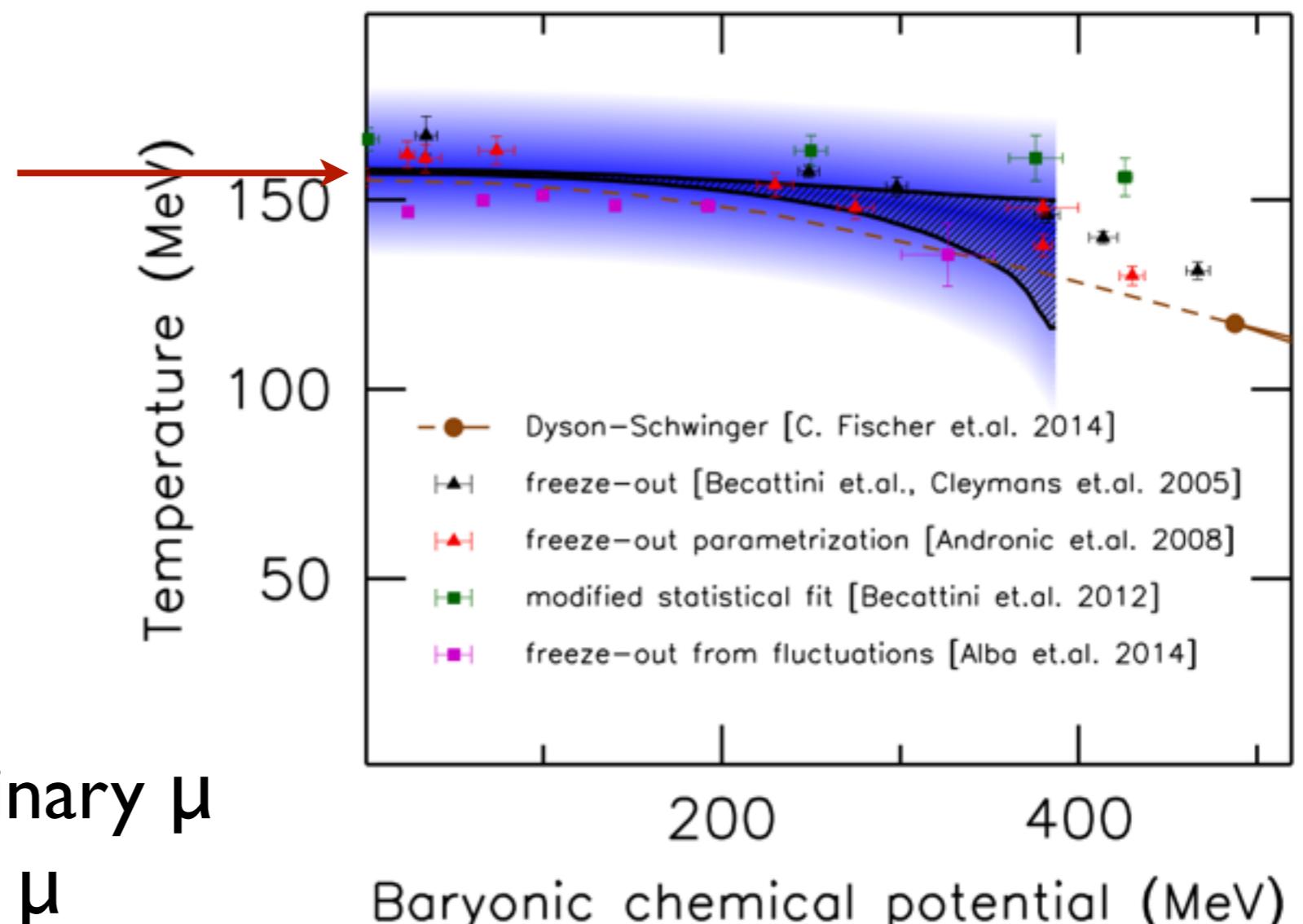


# Chiral transition line from analytic continuation



Lattice method:

- Calc. boundary at imaginary  $\mu$  and extrapolate to real  $\mu$
- Control systematics



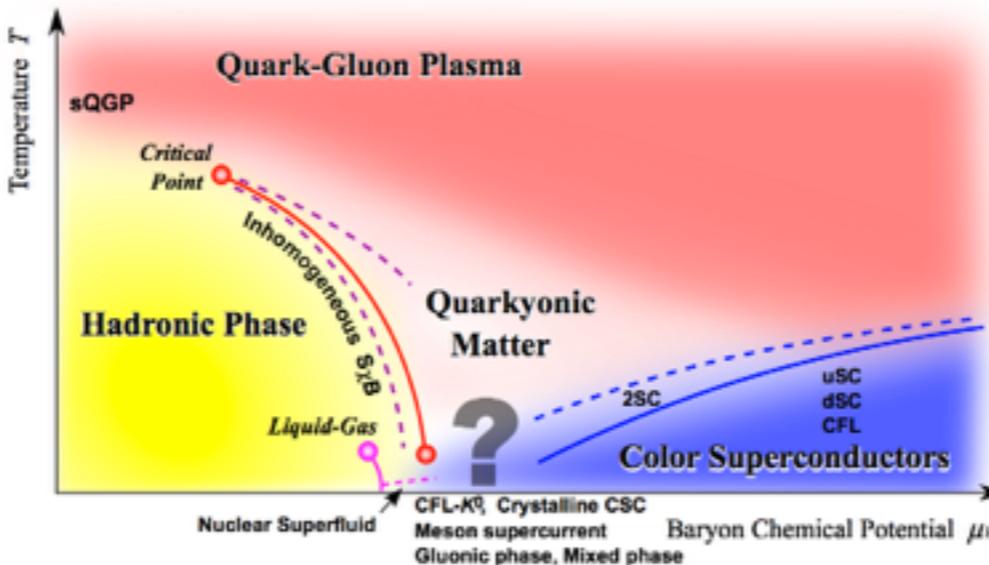
Bellwied, Borsanyi, Fodor, Günther,  
Katz, Ratti and Szabo, PLB B 751 (2015) 559

Results:

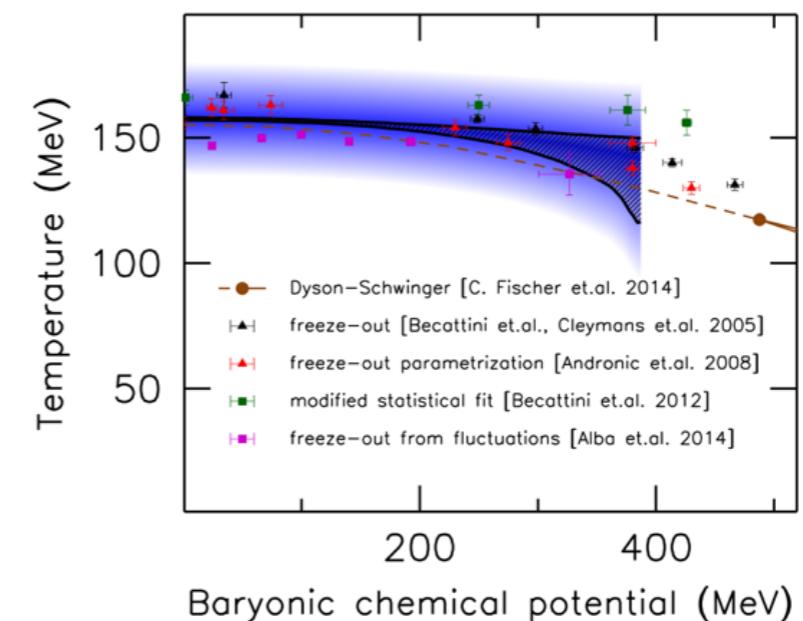
- Larger curvature than previous results  
(but: different definitions and error budget)

# Overview

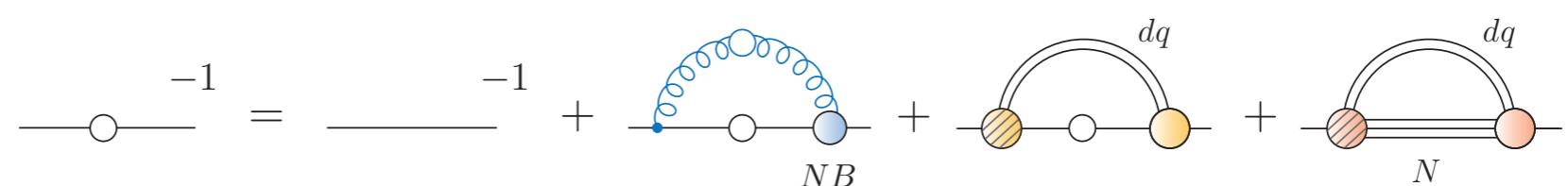
## 1. Introduction



## 2. Gluons, quarks and the CEP



## 3. Baryon effects on the CEP



# QCD order parameters from propagators

$$-\frac{1}{\text{---} \bullet \text{---}} = -\frac{1}{\text{---} \bullet \text{---}} - \text{---} \bullet \text{---}$$

**Chiral order parameter:**

$$\langle \bar{\Psi} \Psi \rangle = Z_2 N_c \text{Tr}_D \frac{1}{T} \sum_{\omega} \int \frac{d^3 p}{(2\pi)^3} S(\vec{p}, \omega)$$

**Deconfinement:**

- dressed Polyakov loop

$$\Sigma = - \int_0^{2\pi} \frac{d\varphi}{2\pi} e^{-i\varphi} \langle \bar{\Psi} \Psi \rangle_{\varphi}$$

Synatschke, Wipf, Wozar, PRD 75, 114003 (2007)  
 Bilgici, Bruckmann, Gattringer, Hagen, PRD 77 094007 (2008)  
 CF, PRL 103 052003 (2009)

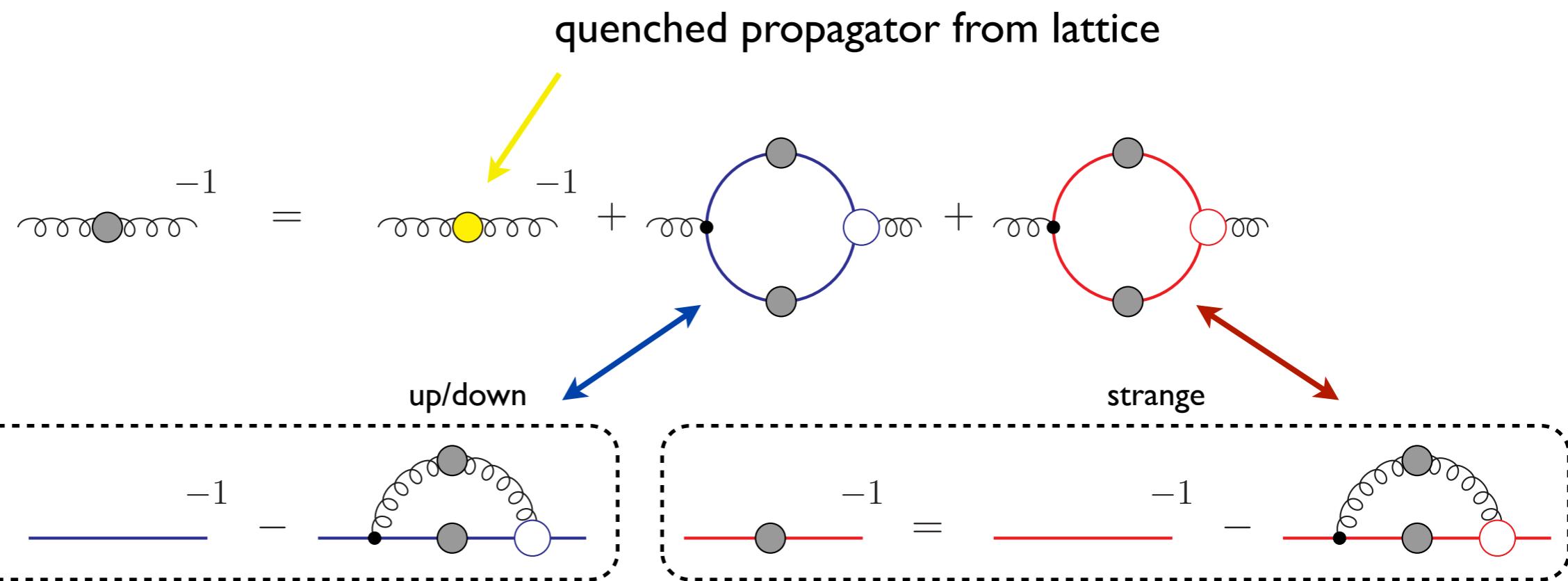
- Polyakov loop potential

$$L = \frac{1}{N_c} \text{Tr} e^{ig\beta A_0}$$

$$\frac{\delta (\Gamma - S)}{\delta A_0} = \frac{1}{2} \text{---} \bullet \text{---} + \text{---} \circlearrowleft \text{---} - \text{---} \bullet \text{---} - \frac{1}{6} \text{---} \bullet \text{---} + \text{---} \circlearrowright \text{---}$$

Braun, Gies, Pawłowski, PLB 684, 262 (2010)  
 Braun, Haas, Marhauser, Pawłowski, PRL 106 (2011)  
 Fister, Pawłowski, PRD 88 045010 (2013)  
 CF, Fister, Luecker, Pawłowski, PLB 732 (2013)

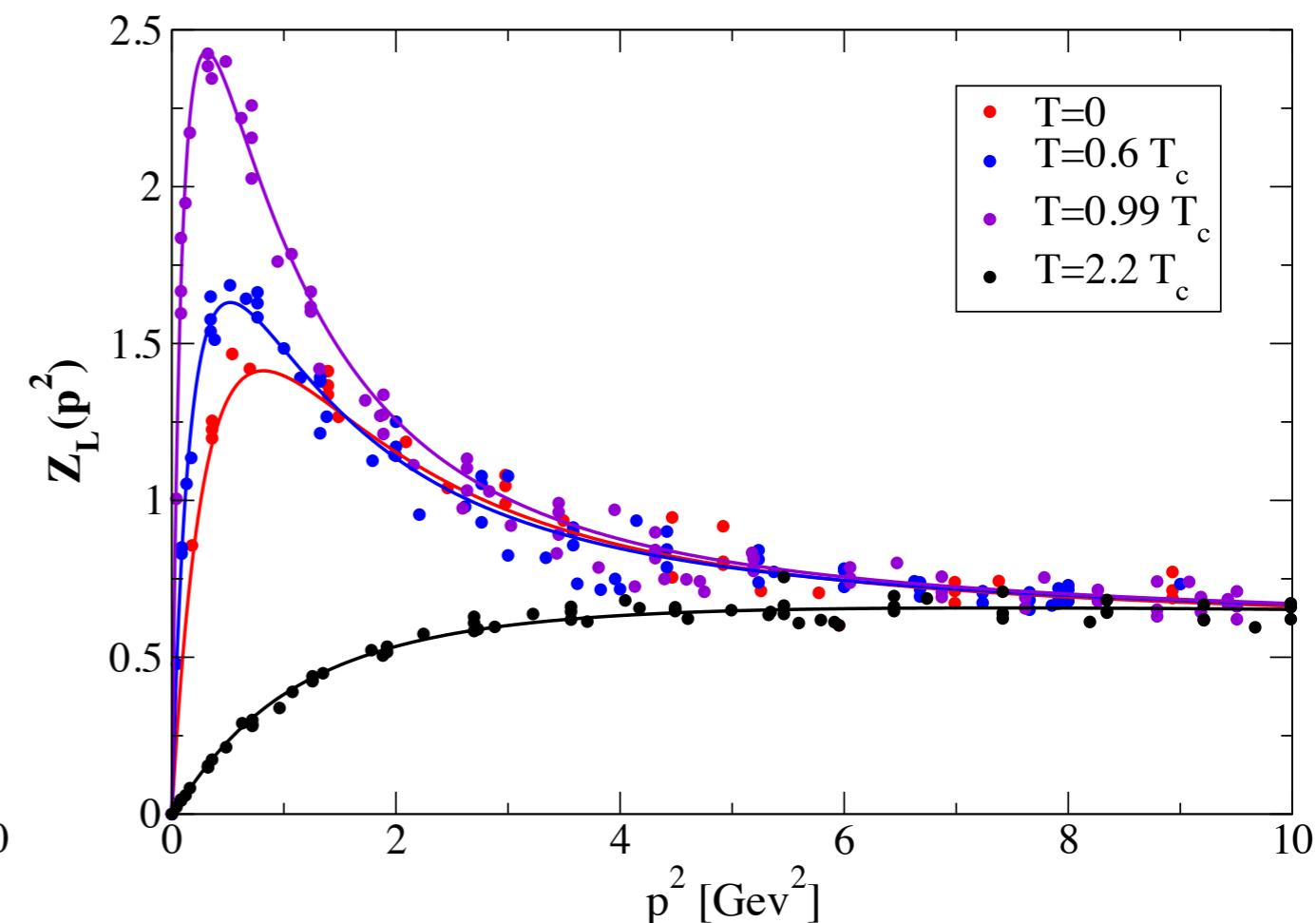
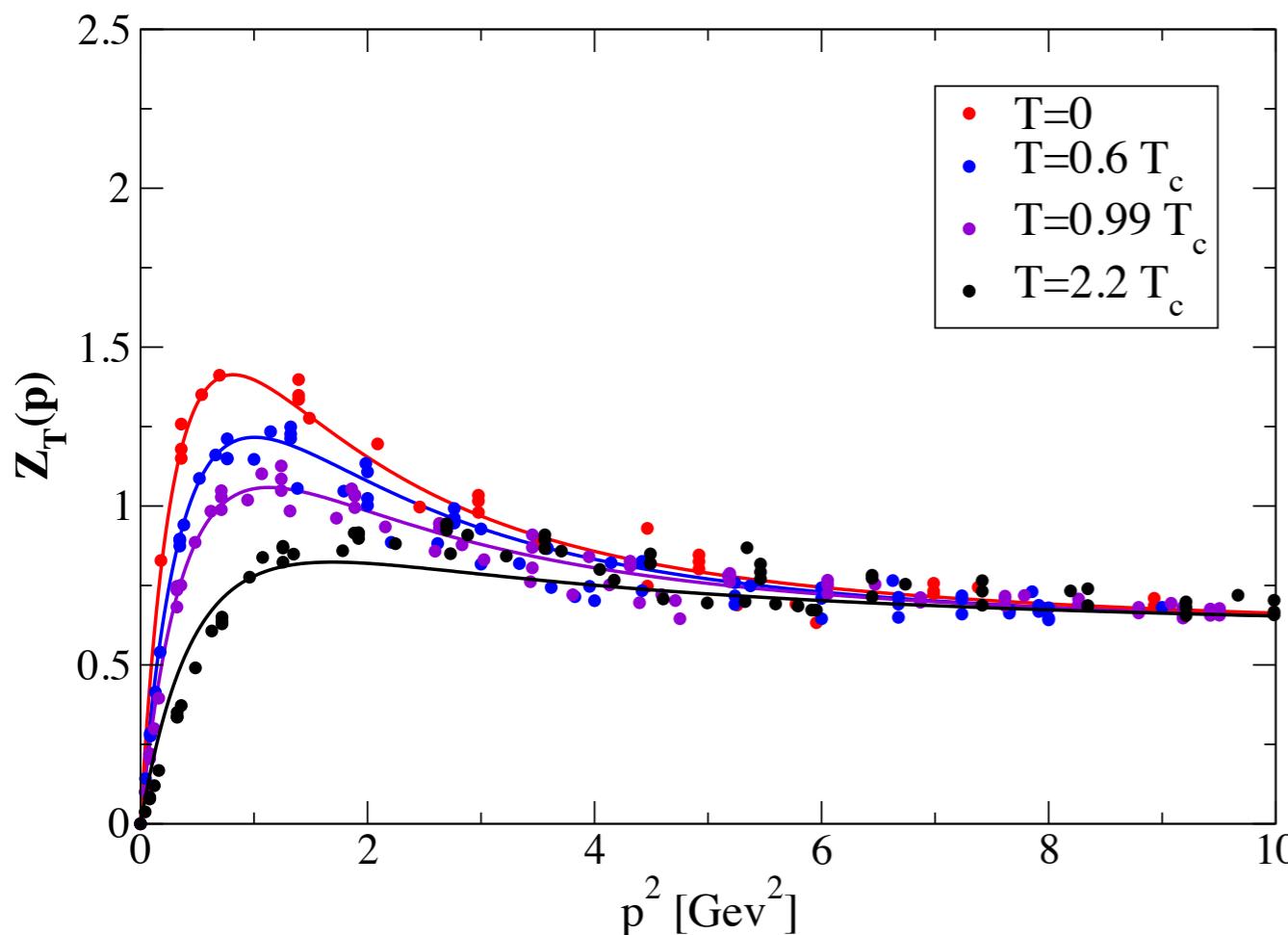
# $N_f=2+1$ -QCD with DSEs



- quenched: without quark-loop
- $N_f=2$ : isospin symmetry
- $N_f=2+1$ : solve coupled system of 2+3+3 equations
- Vertex: ansatz built along STI and known UV/IR behaviour

# Glue at finite temperature ( $T \neq 0$ )

T-dependent gluon propagator from quenched lattice simulations:



- Crucial difference between magnetic and electric gluon
- Maximum of electric gluon near  $T_c$

Cucchieri, Maas, Mendes, PRD 75 (2007)

CF Maas, Mueller, EPJC 68 (2010)

Cucchieri, Mendes, PoS FACESQCD 007 (2010)

Aouane, Bornyakov, Ilgenfritz, Mitrjushkin, Muller-Preussker and Sternbeck, PRD 85 (2012) 034501

Silva, Oliveira, Bicudo, Cardoso, PRD 89 (2014) 074503

FRG: Fister, Pawłowski, arXiv:1112.5440

# Approximation for Quark-Gluon interaction

- T, μ, m-dependent vertex:

$$\Gamma_\nu(q, k, p) = \tilde{Z}_3 \left( \delta_{4\nu} \gamma_4 \frac{C(k) + C(p)}{2} + \delta_{j\nu} \gamma_j \frac{A(k) + A(p)}{2} \right) \times \\ \times \left( \frac{d_1}{d_2 + q^2} + \frac{q^2}{\Lambda^2 + q^2} \left( \frac{\beta_0 \alpha(\mu) \ln[q^2/\Lambda^2 + 1]}{4\pi} \right)^{2\delta} \right)$$

Abelian WTI

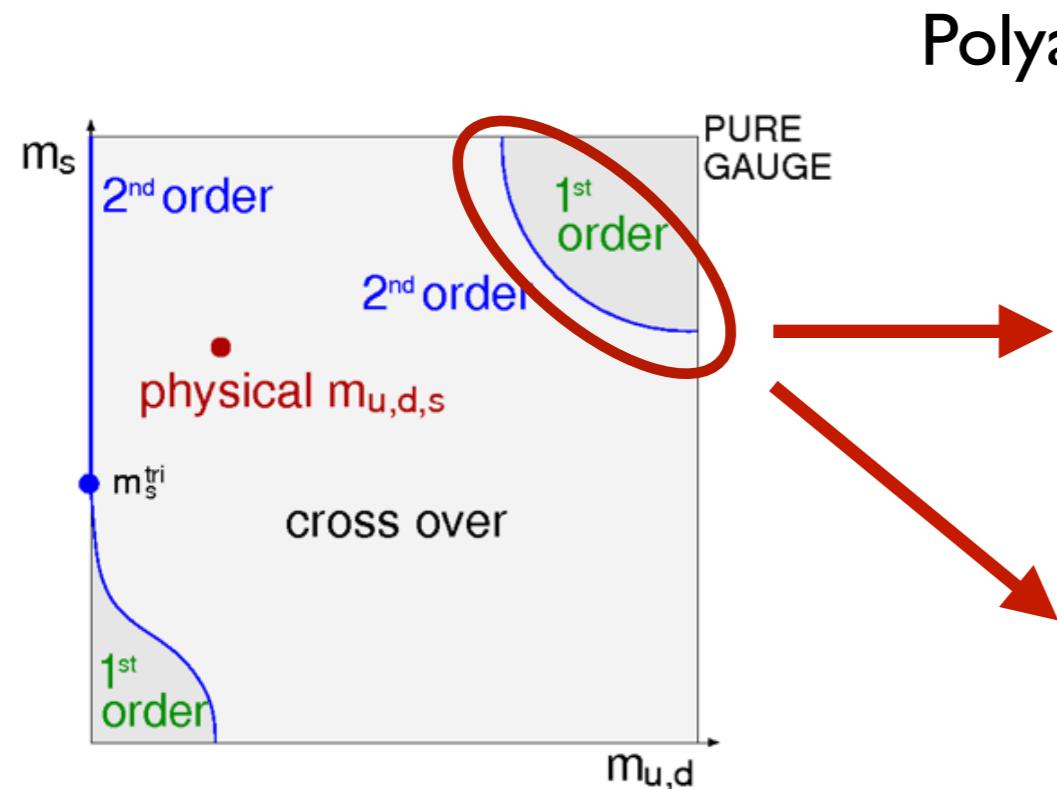
perturbation theory

Infrared ansatz:

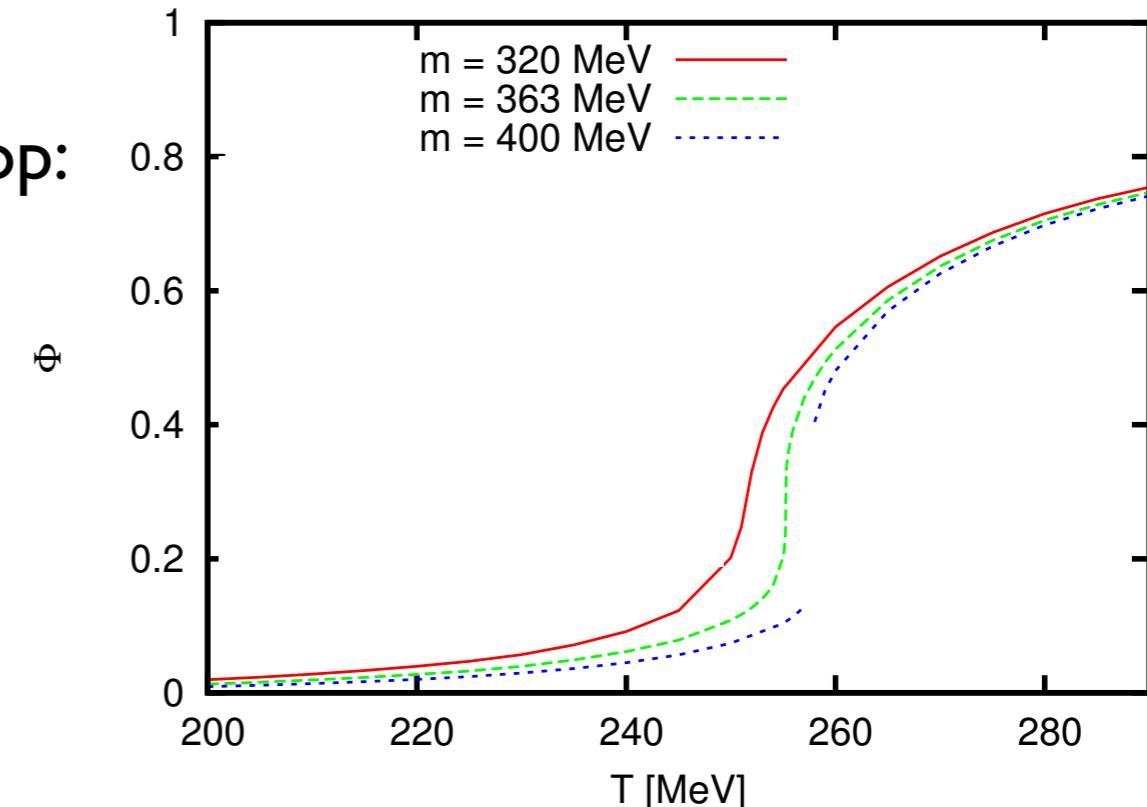
- d2 fixed to match gluon input
- d1 fixed via quark condensate (see later)
- correct UV (quant.) and IR-behavior (qual.)

explicit solutions at T=0: Mitter, Pawłowski and Strodthoff, PRD 91 (2015) 054035  
Williams, CF, Heupel, PRD PRD 93 (2016) 034026

# Critical line/surface for heavy quarks

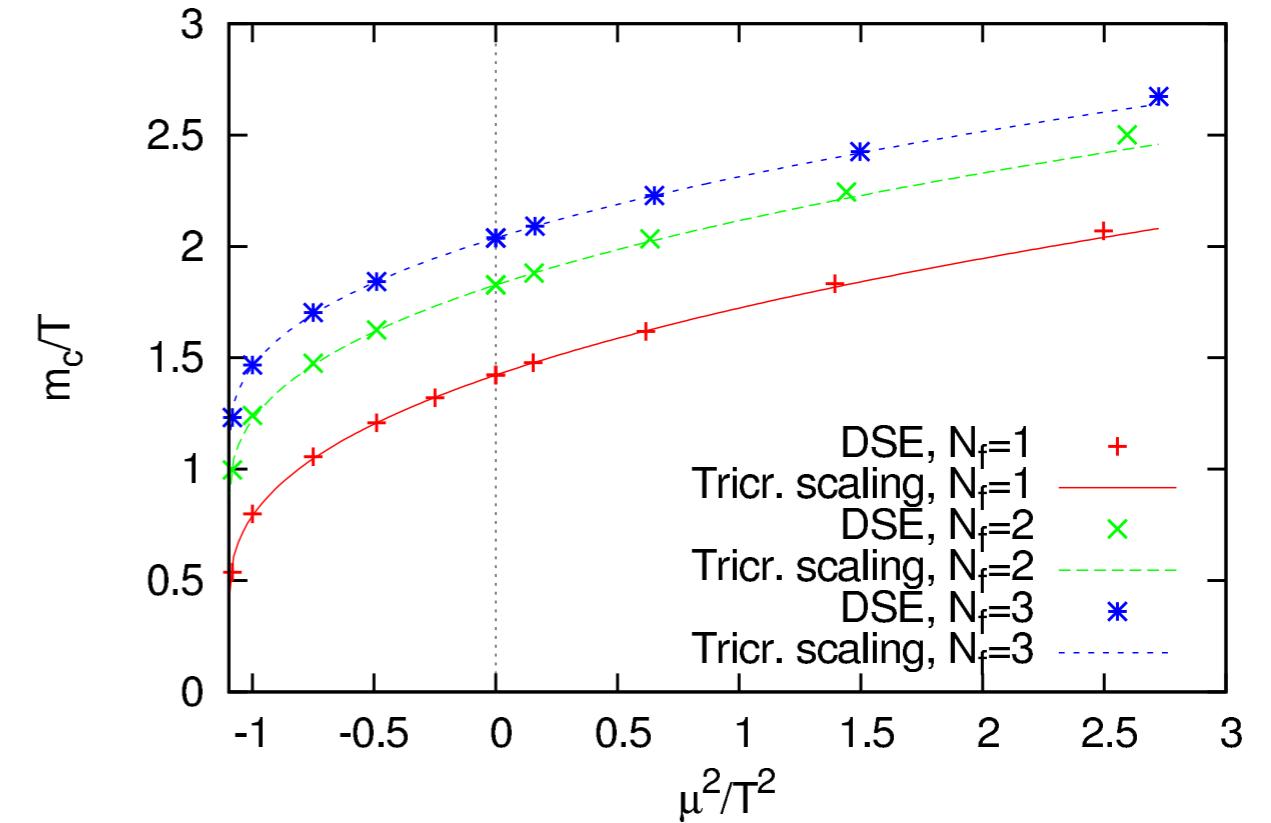


Polyakov Loop:



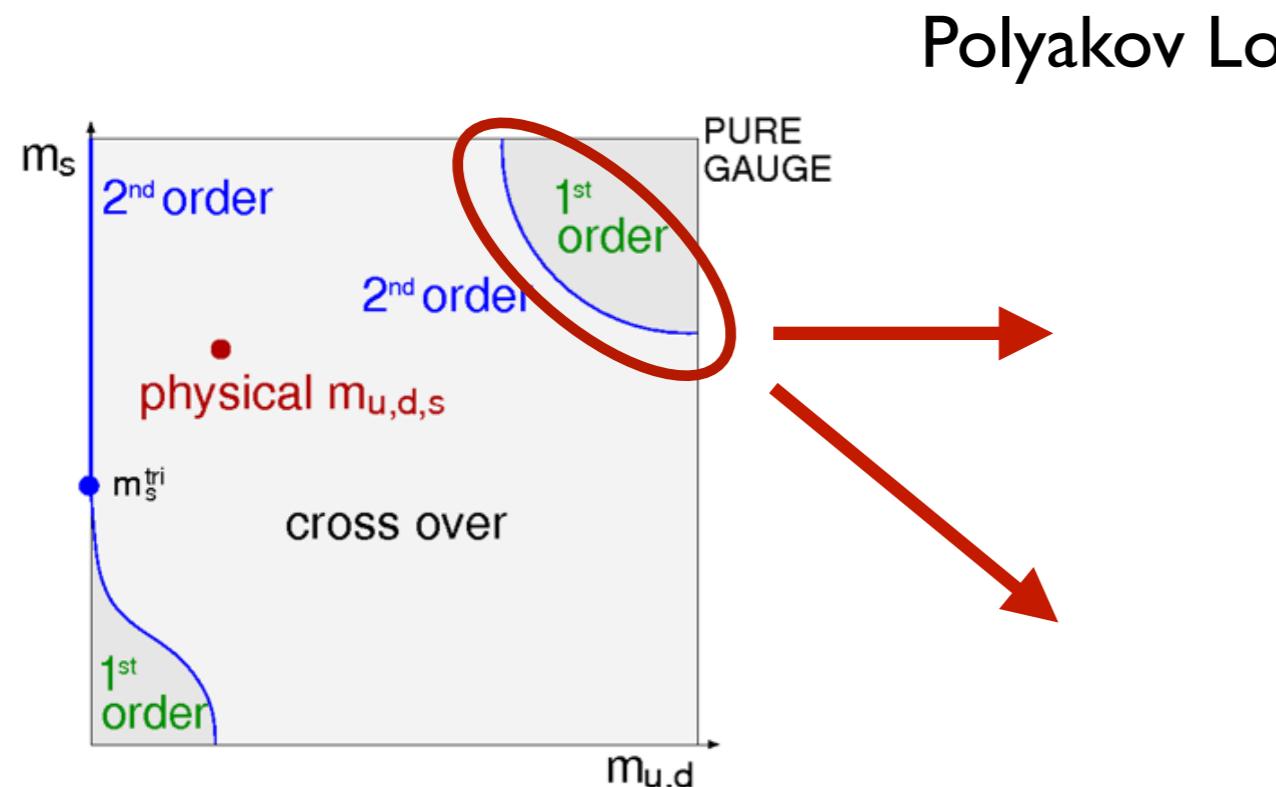
- Deconfinement transition in agreement with lattice QCD
- Correct tricritical scaling

Fromm, Langelage, Lottini, Philipsen, JHEP 1201 (2012) 042

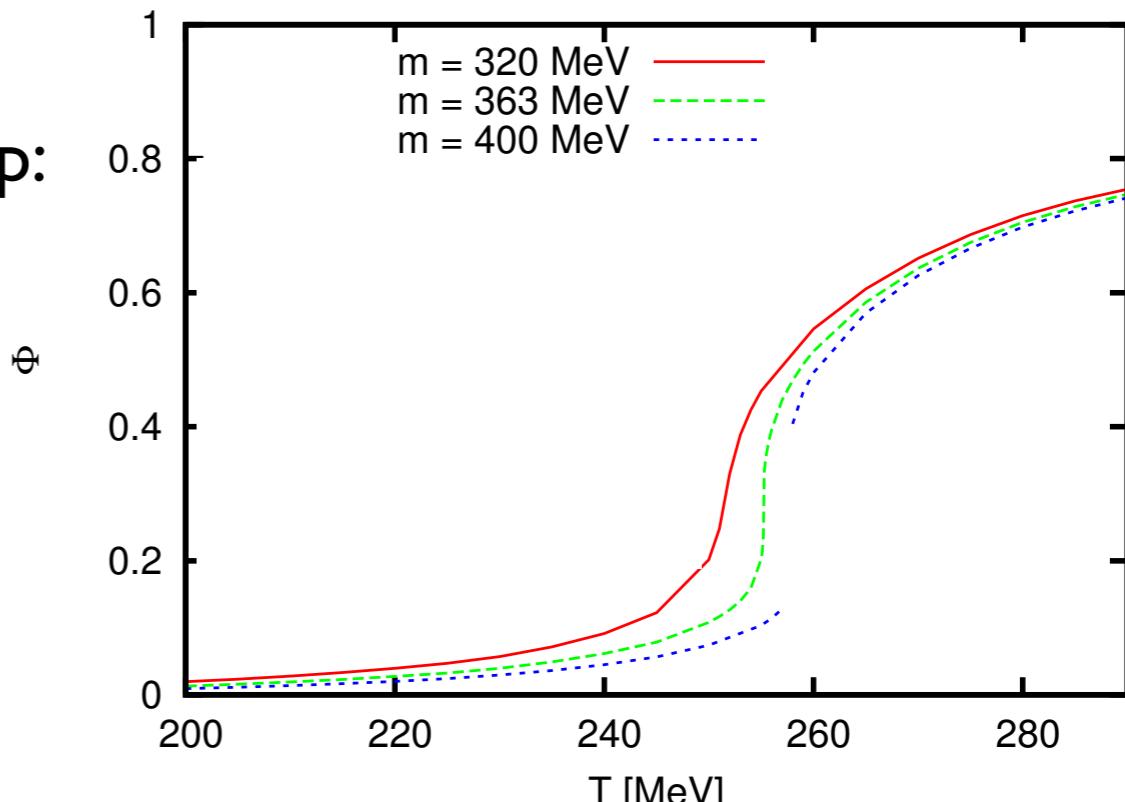


CF, Luecker, Pawłowski, PRD 91 (2015) 1

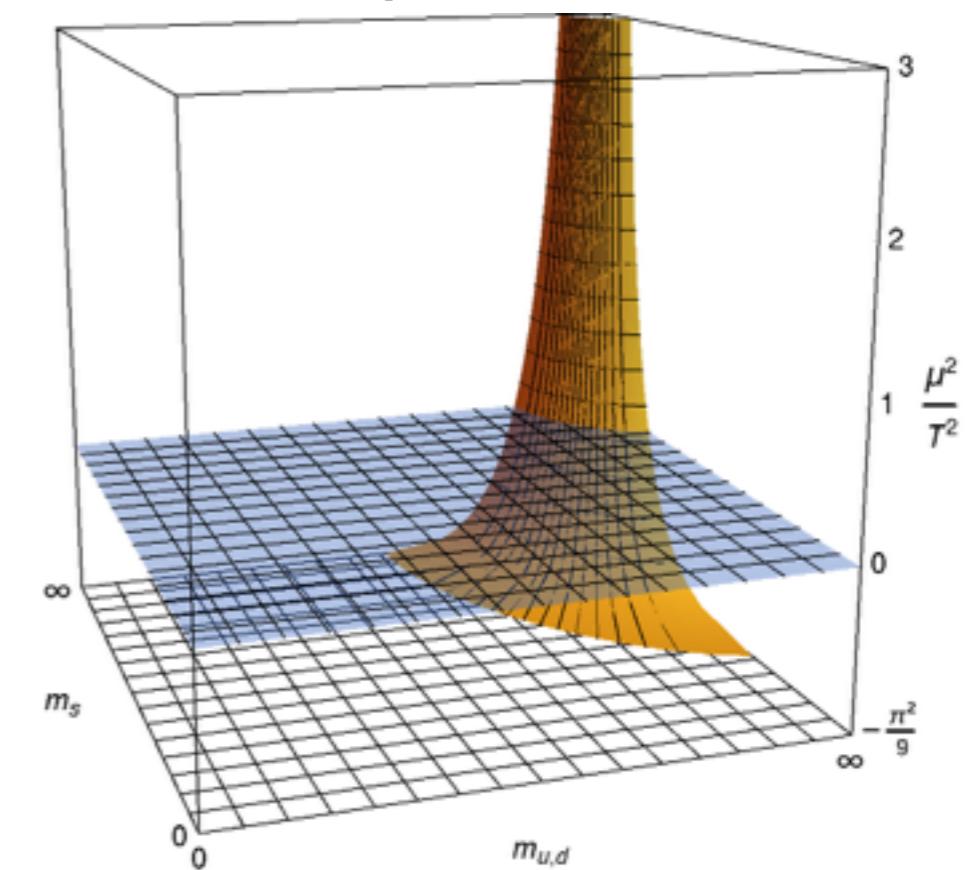
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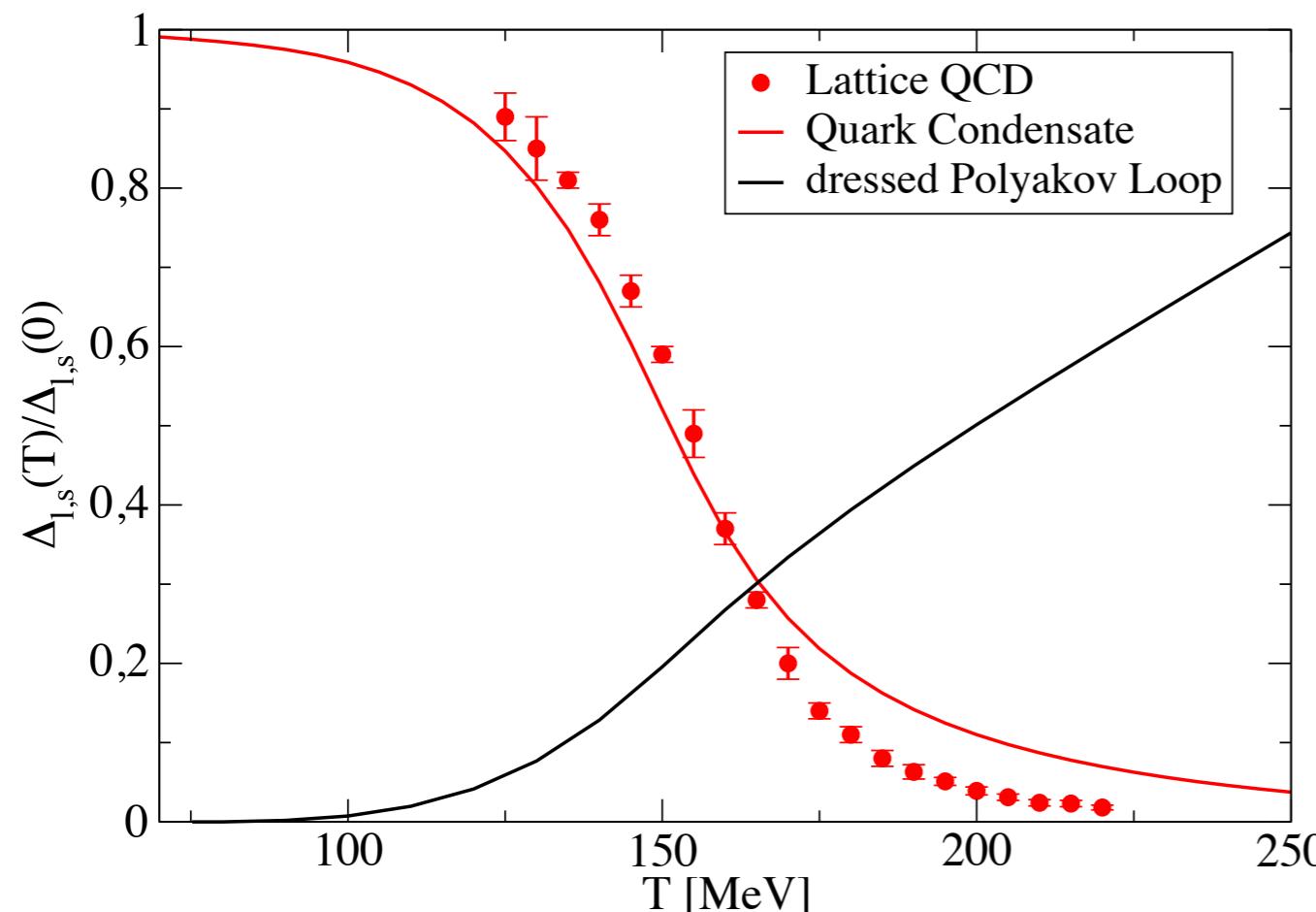
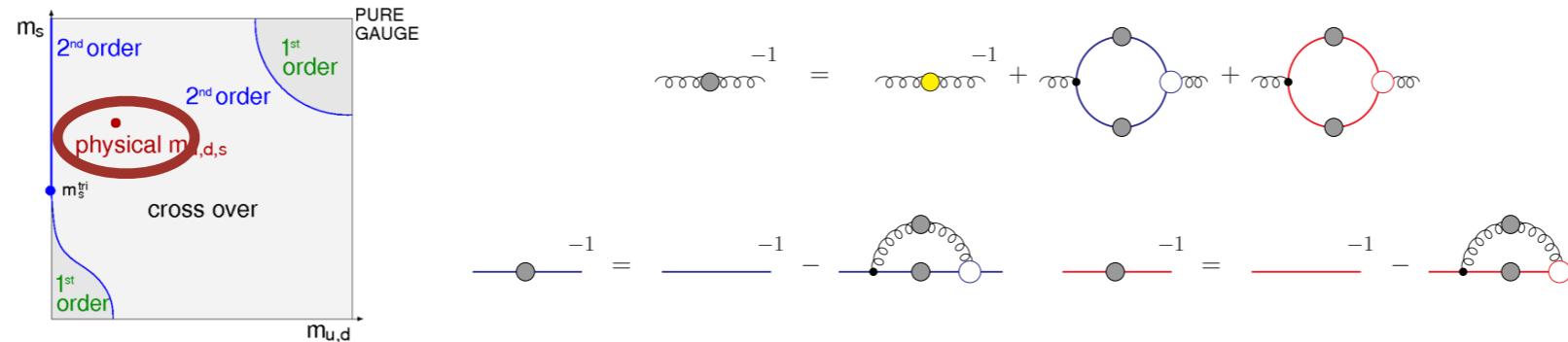
- Deconfinement transition in agreement with lattice QCD
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Fromm, Langelage, Lottini, Philipsen, JHEP 1201 (2012) 042

CF, Luecker, Pawłowski, PRD 91 (2015) 1

# $N_f=2+1$ , $\mu=0$ , physical point

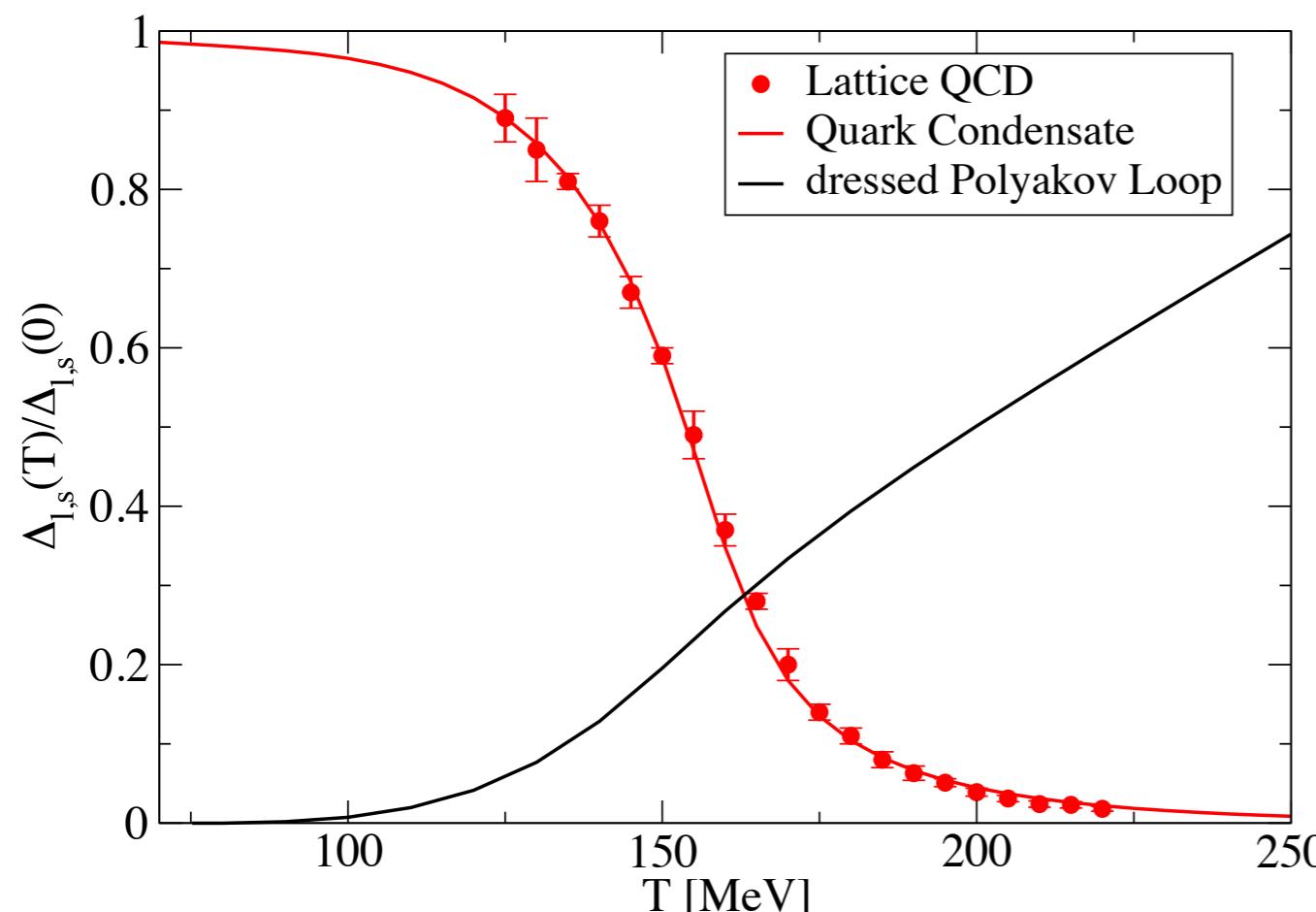
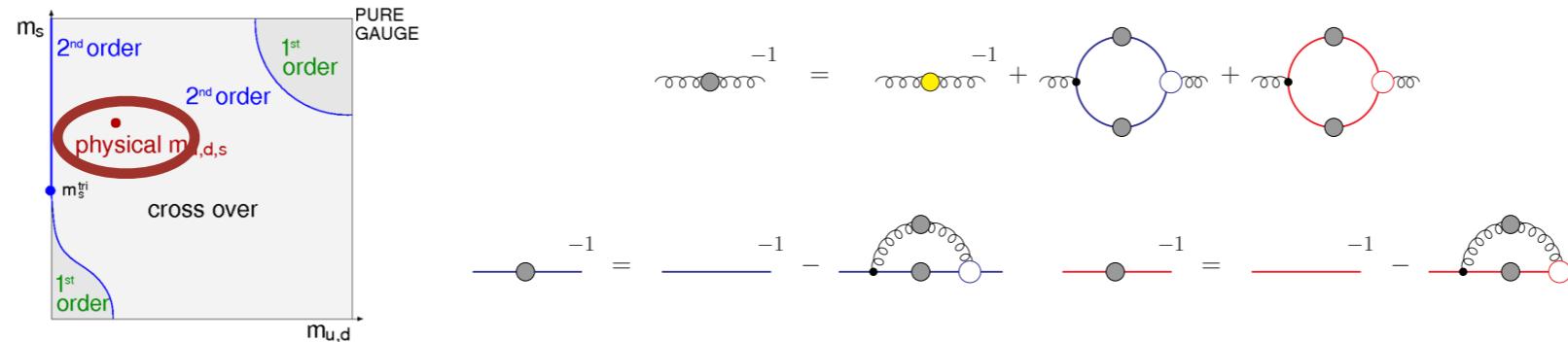


Lattice: Borsanyi et al. [Wuppertal-Budapest], JHEP 1009(2010) 073

DSE: CF, Luecker, PLB 718 (2013) 1036,

CF, Luecker, Welzbacher, PRD 90 (2014) 034022

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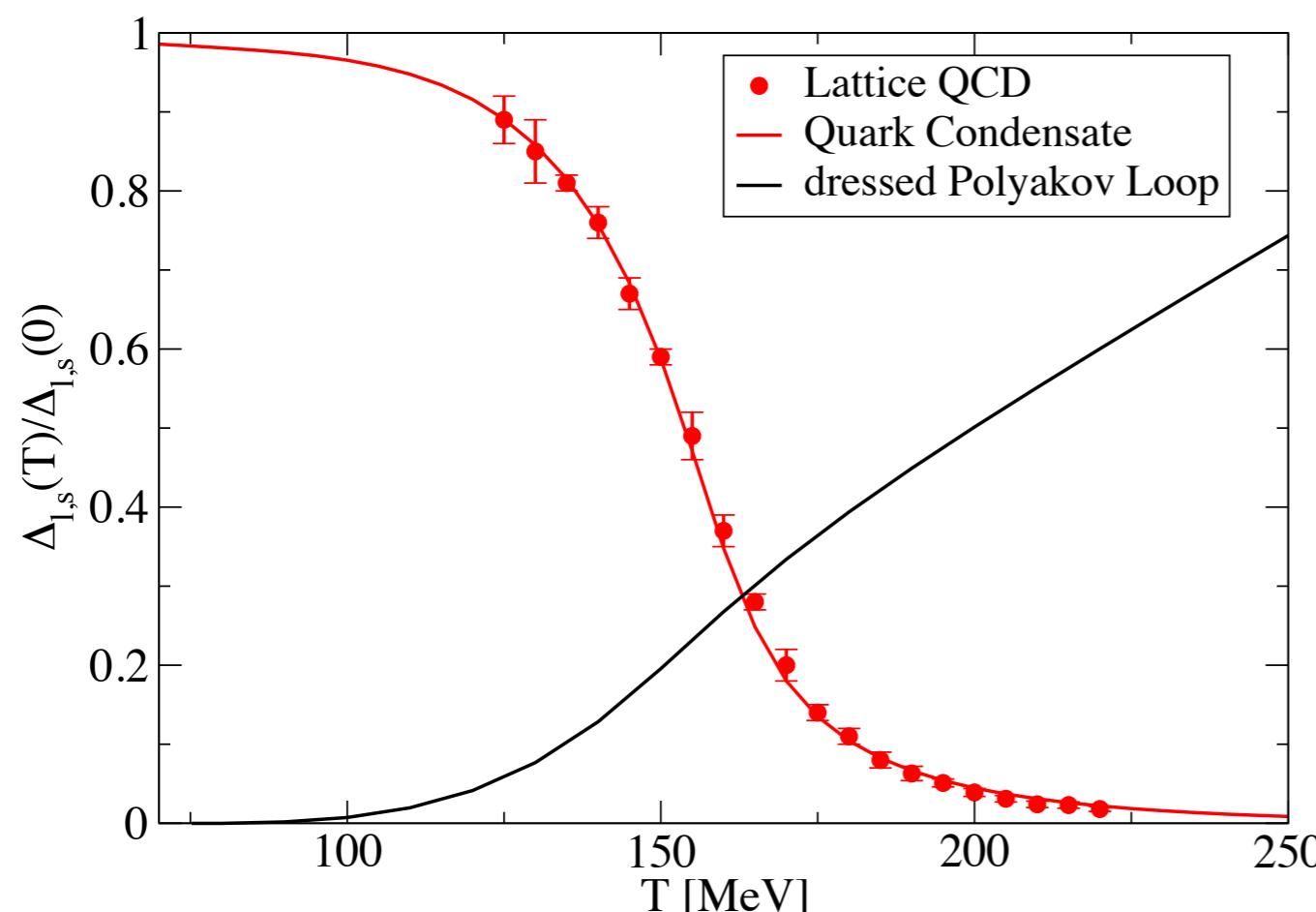
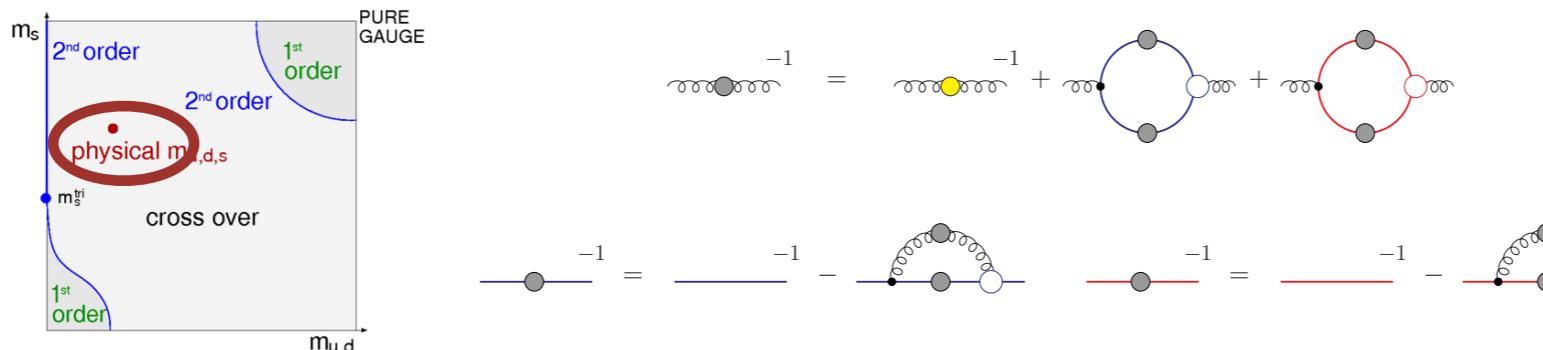


Lattice: Borsanyi et al. [Wuppertal-Budapest], JHEP 1009(2010) 073

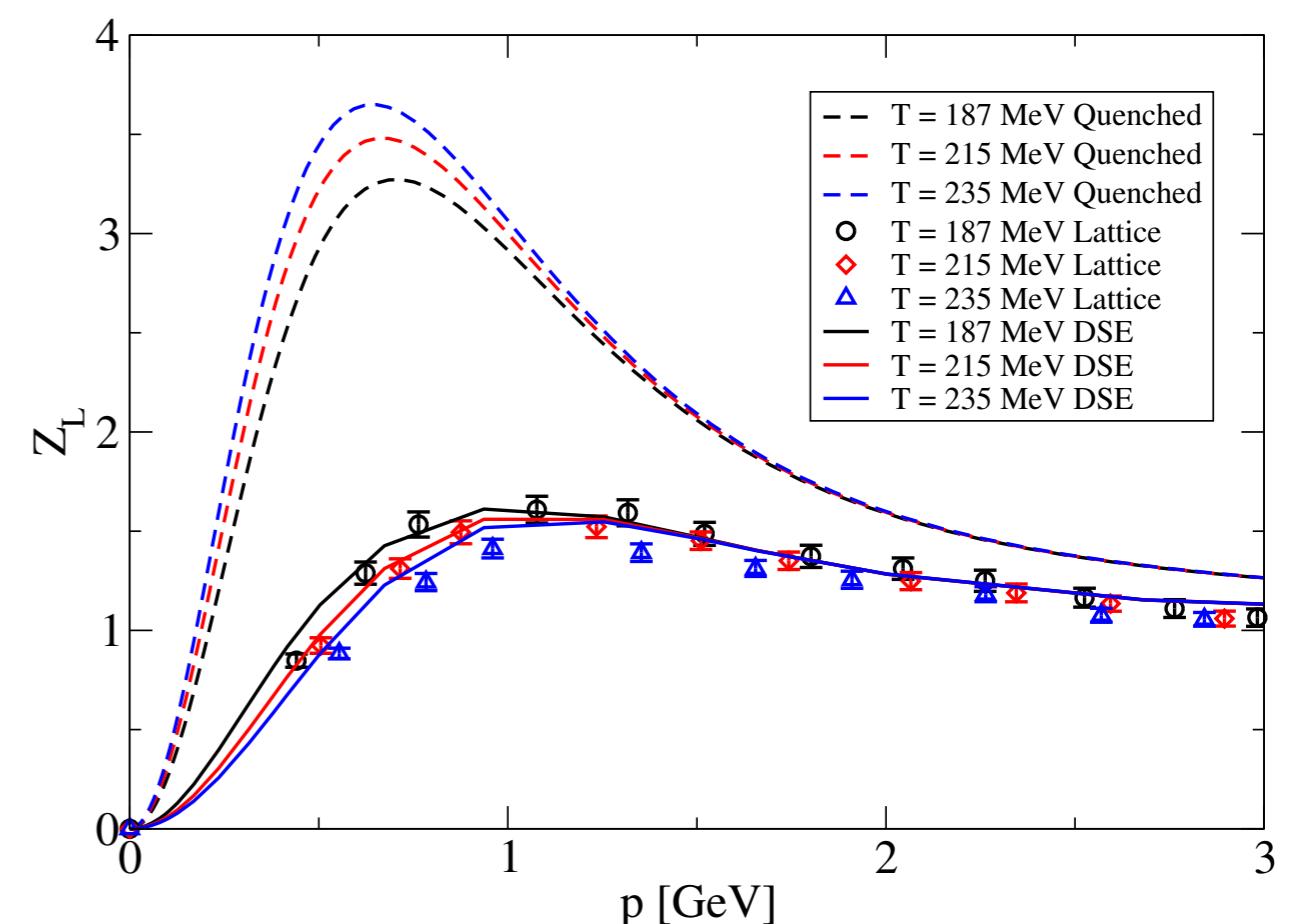
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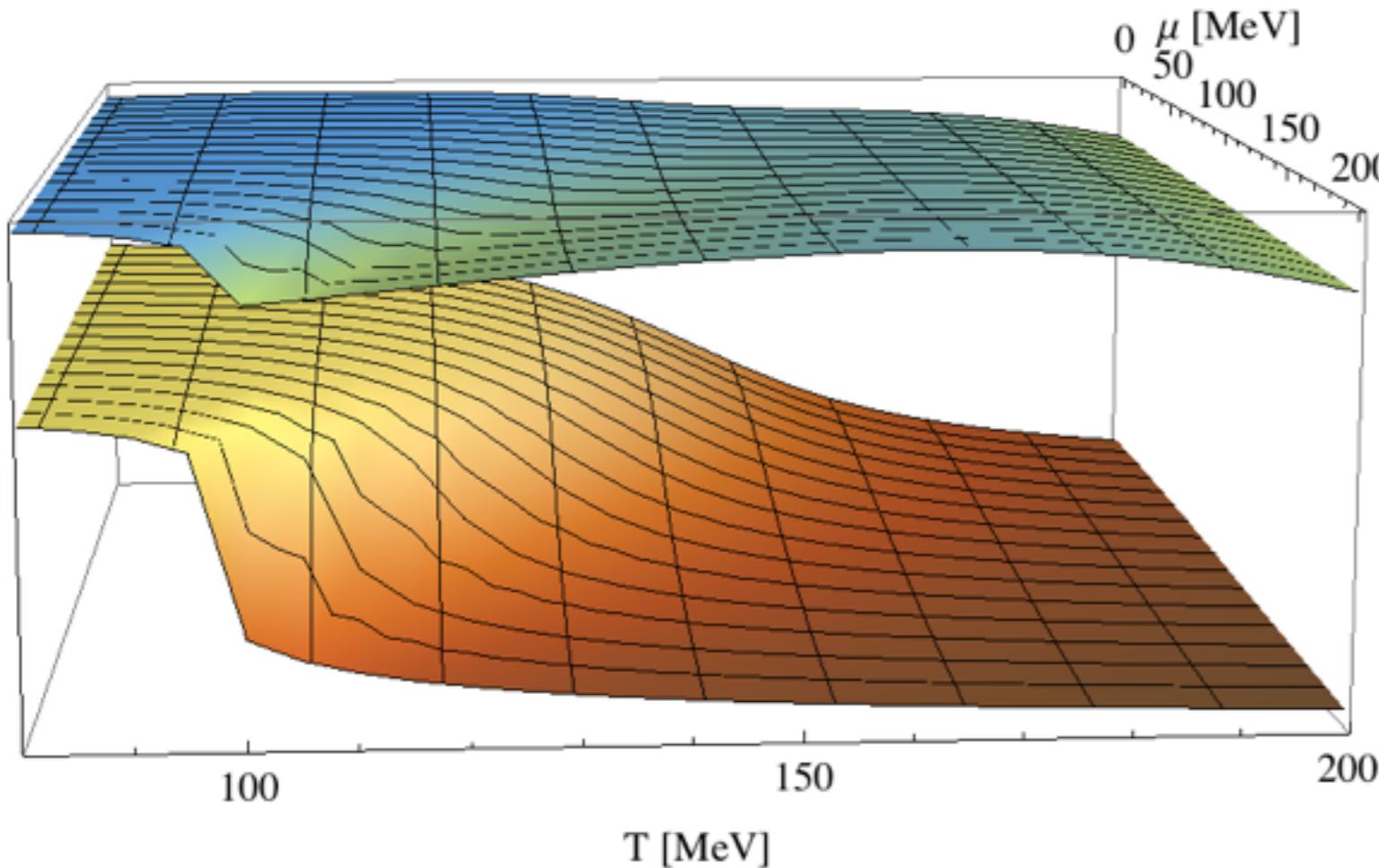
Lattice: Borsanyi et al. [Wuppertal-Budapest], JHEP 1009(2010) 073  
 DSE: CF, Luecker, PLB 718 (2013) 1036,  
 CF, Luecker, Welzbacher, PRD 90 (2014) 034022



Lattice: Aouane, et al. PRD 87 (2013), [arXiv:1212.1102]  
 DSE: CF, Luecker, PLB 718 (2013) 1036, [arXiv:1206.5191]

● quantitative agreement: DSE prediction verified by lattice

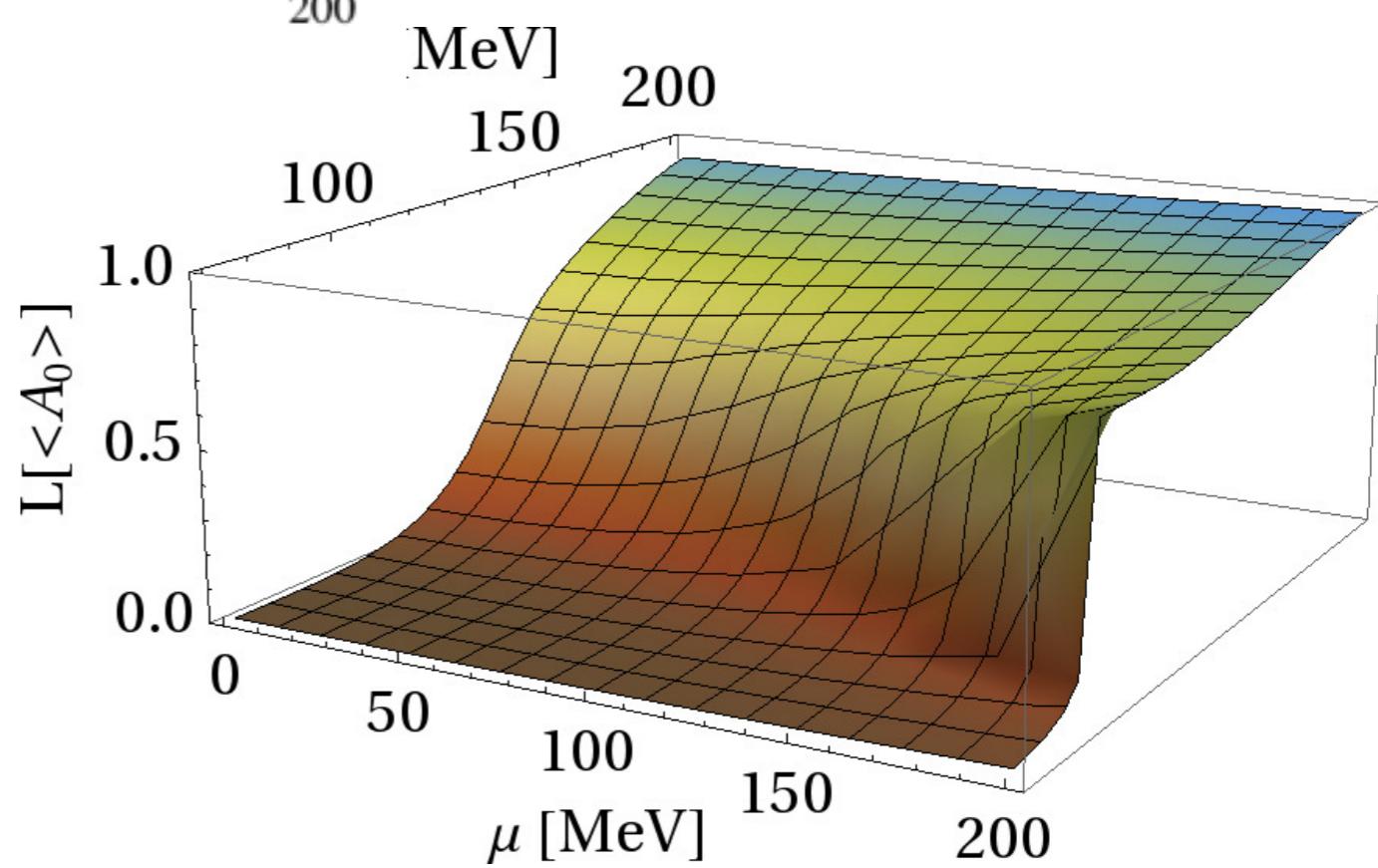
# Nf=2+1: Condensate and dressed Polyakov Loop



Quark condensate

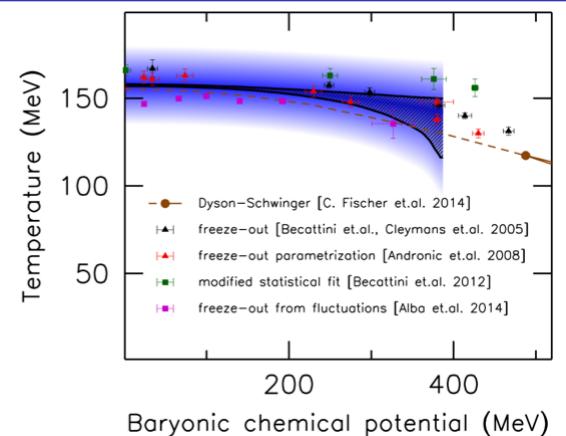
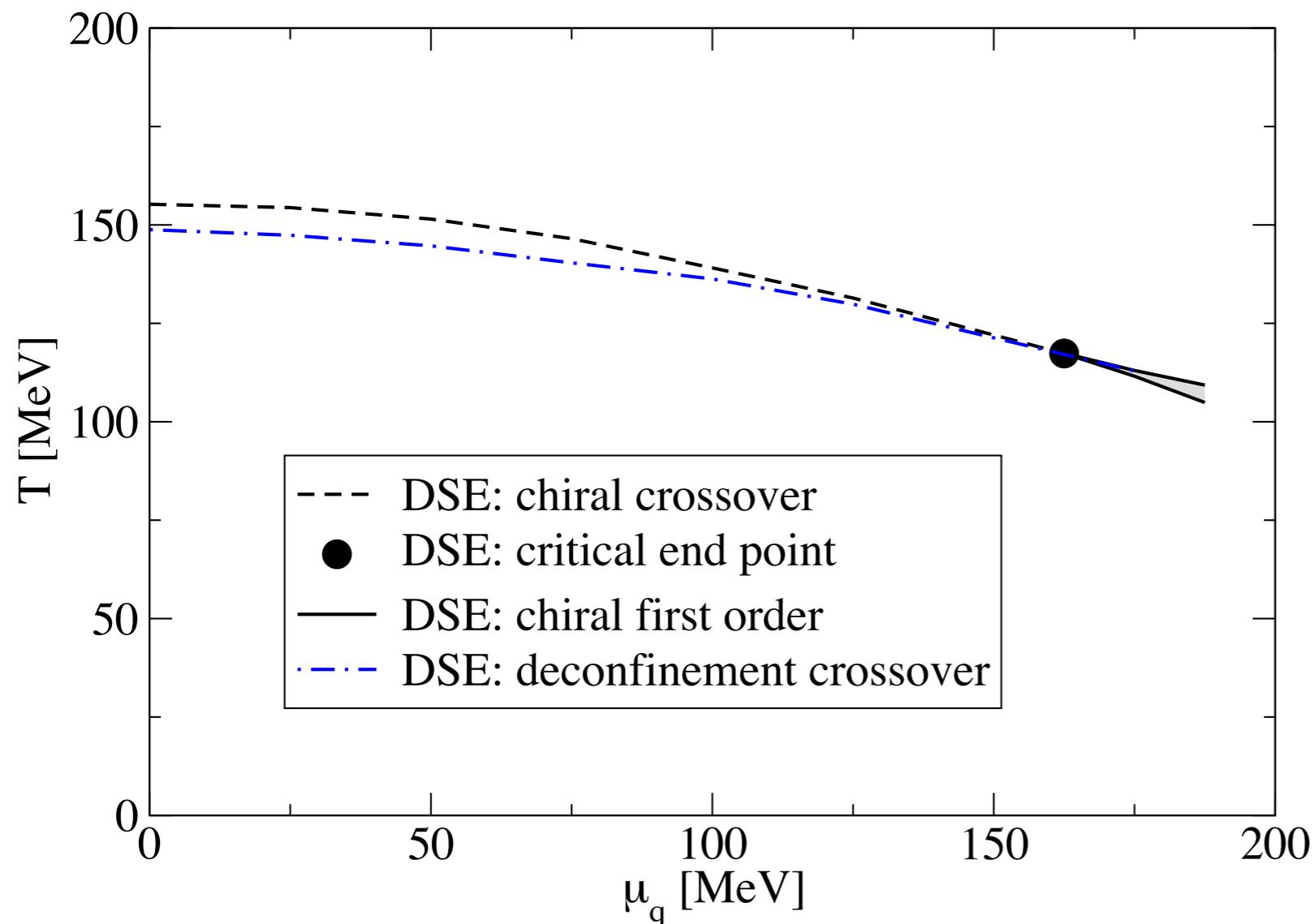
Polyakov-Loop

$$L = \frac{1}{N_c} \text{tr } e^{ig \int A_0}$$

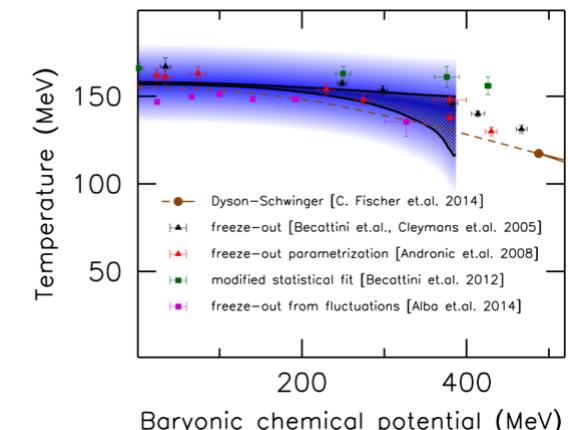
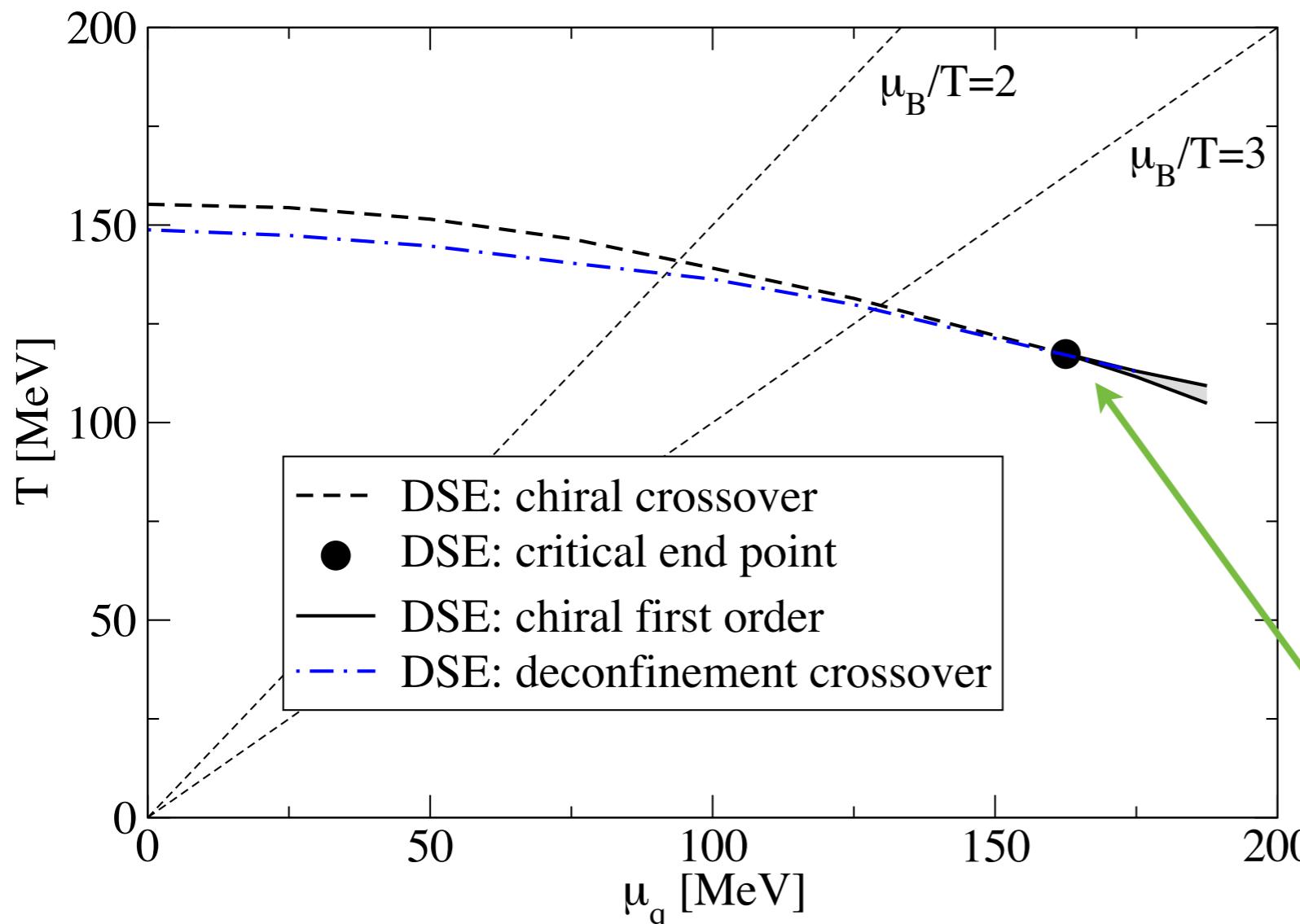


CF, Fischer, Luecker, Pawłowski, PLB 732 (2014) 273

# $N_f=2+1$ : phase diagram



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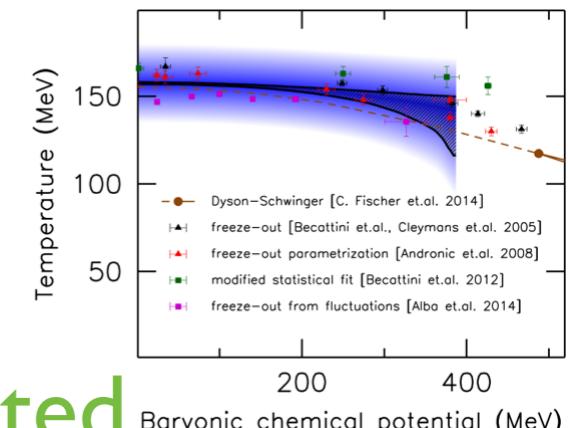
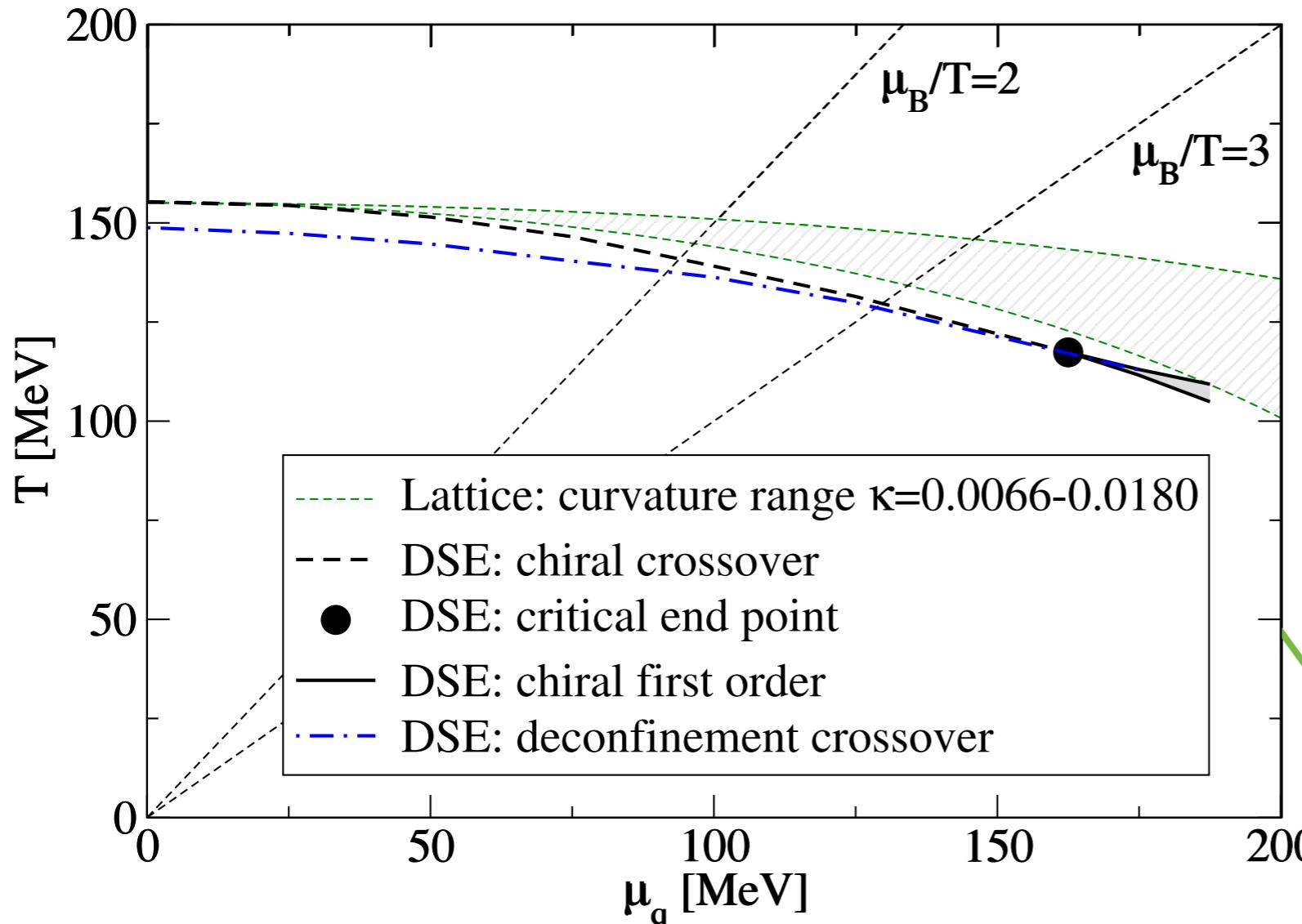


CEP at large  $\mu$

CF, Luecker, PLB 718 (2013) 1036,  
CF, Fister, Luecker, Pawłowski, PLB 732 (2014) 273  
CF, Luecker, Welzbacher, PRD 90 (2014) 034022

● combined evidence of FRG and DSE: no CEP at  $\mu_B/T < 2$

# $N_f=2+1$ : phase diagram



Extrapolated  
curvature from lattice

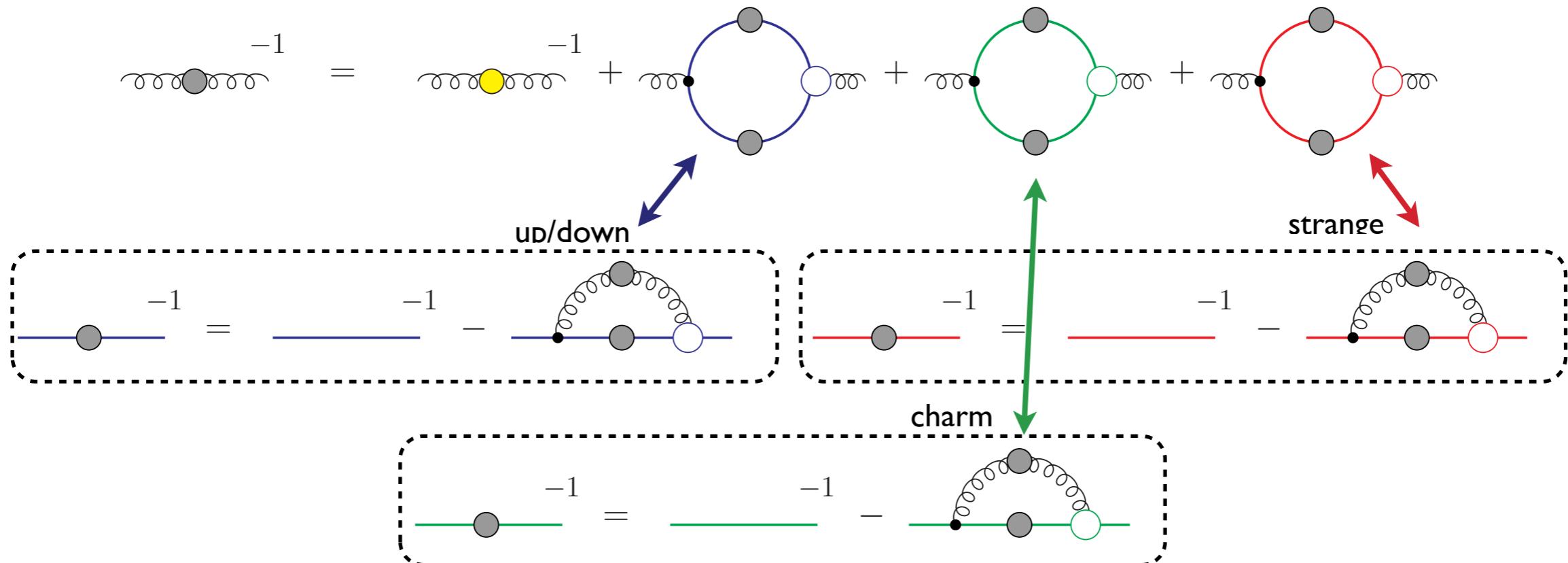
Kaczmarek et al. PRD 83 (2011) 014504,  
Endrodi, Fodor, Katz, Szabo, JHEP 1104 (2011) 001  
Cea, Cosmai, Papa, PRD 89 (2014), PRD 93 (2016)  
Bonati et al., PRD 92 (2015) 054503  
Bellwied et al. PLB 751 (2015) 559

CEP at large  $\mu$

CF, Luecker, PLB 718 (2013) 1036,  
CF, Fister, Luecker, Pawłowski, PLB 732 (2014) 273  
CF, Luecker, Welzbacher, PRD 90 (2014) 034022

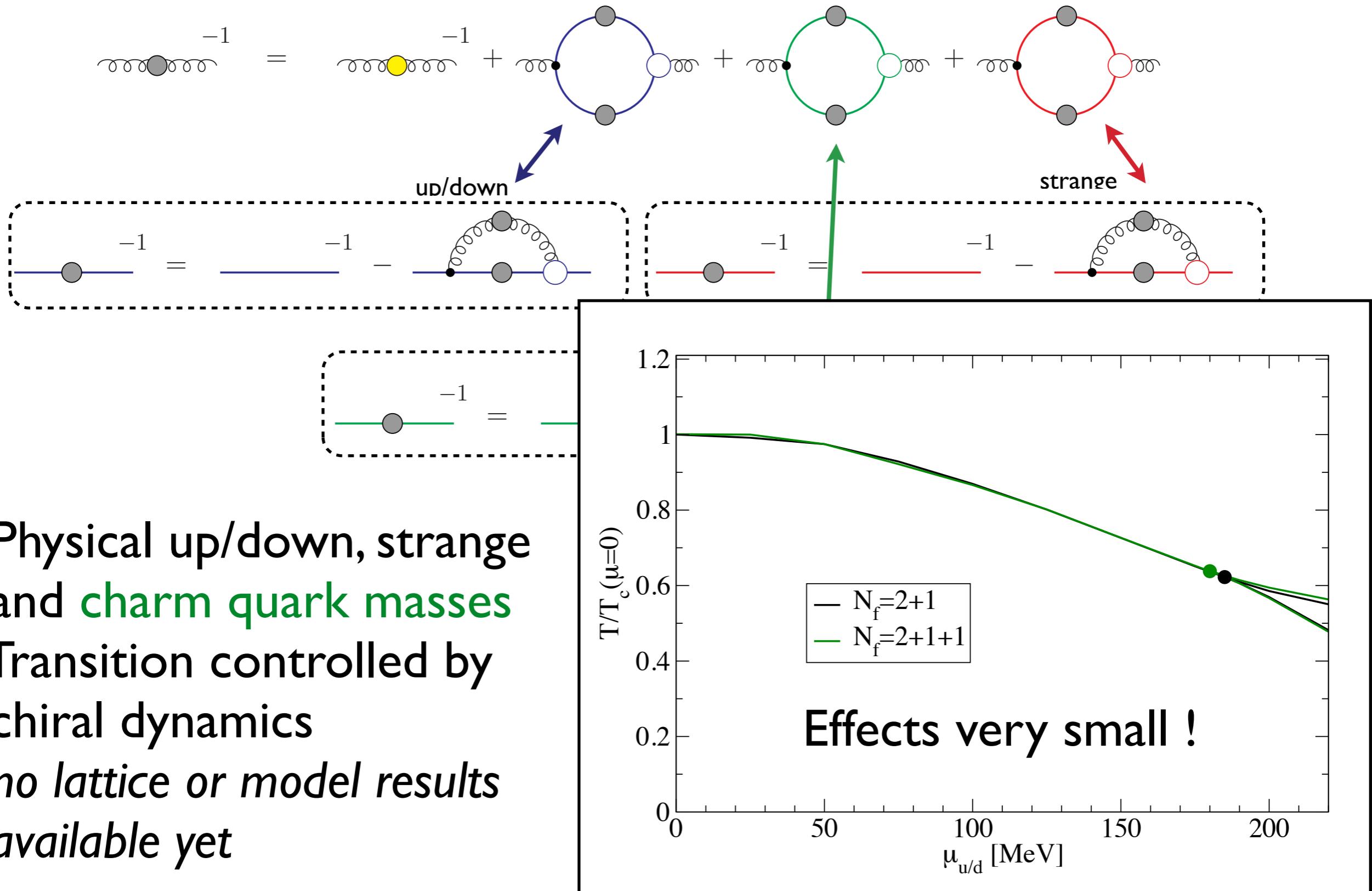
● combined evidence of FRG and DSE: no CEP at  $\mu_B/T < 2$

# $N_f=2+1+1$ : effects of charm



- Physical up/down, strange and **charm quark masses**
- Transition controlled by chiral dynamics
- *no lattice or model results available yet*

# $N_f=2+1+1$ : effects of charm



CF, Luecker, Welzbacher, PRD 90 (2014) 034022

# Location of CEP in freeze-out landscape

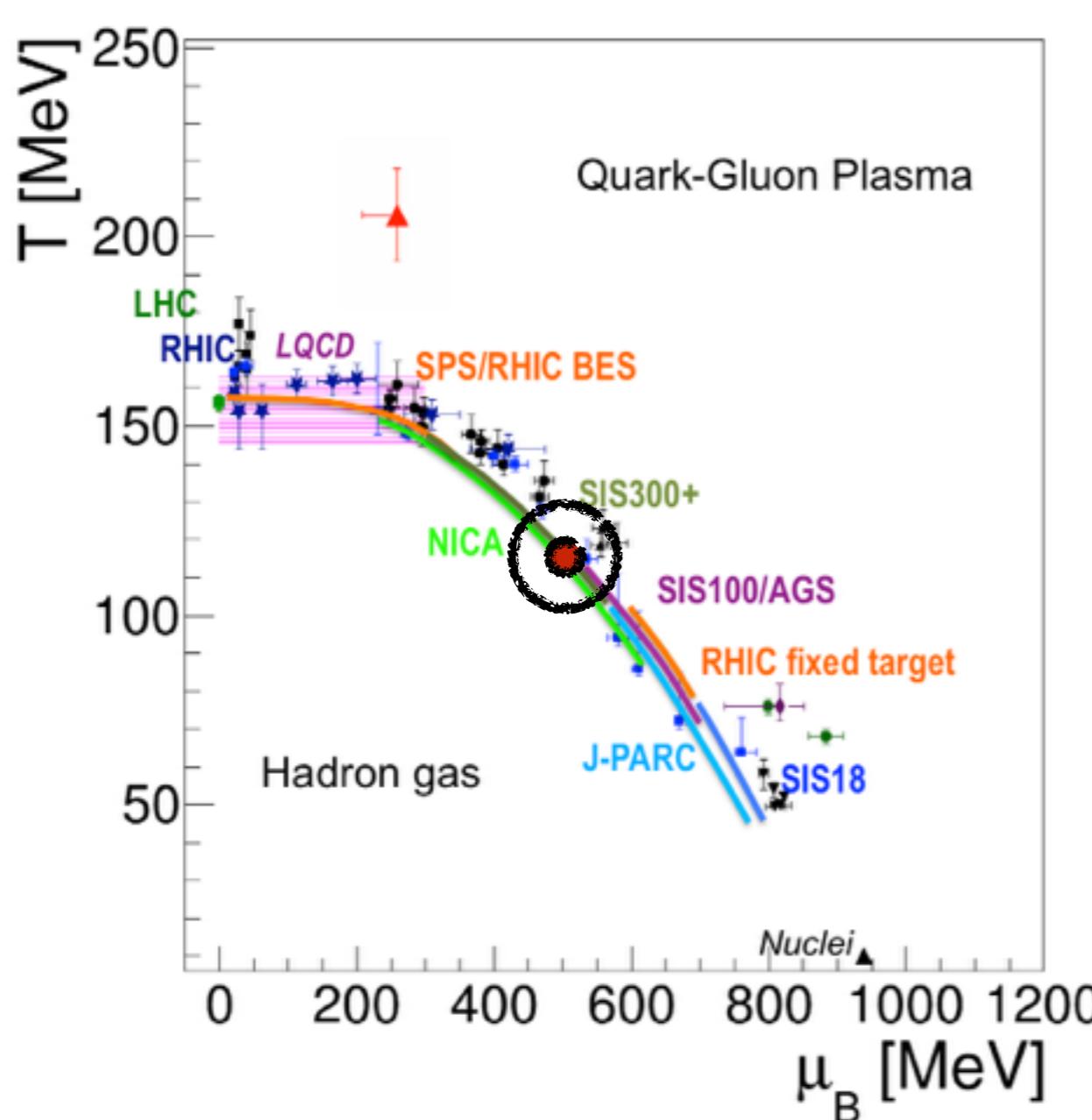


Figure taken from talk of T. Galatyuk, Erice 2016

## Caveats:

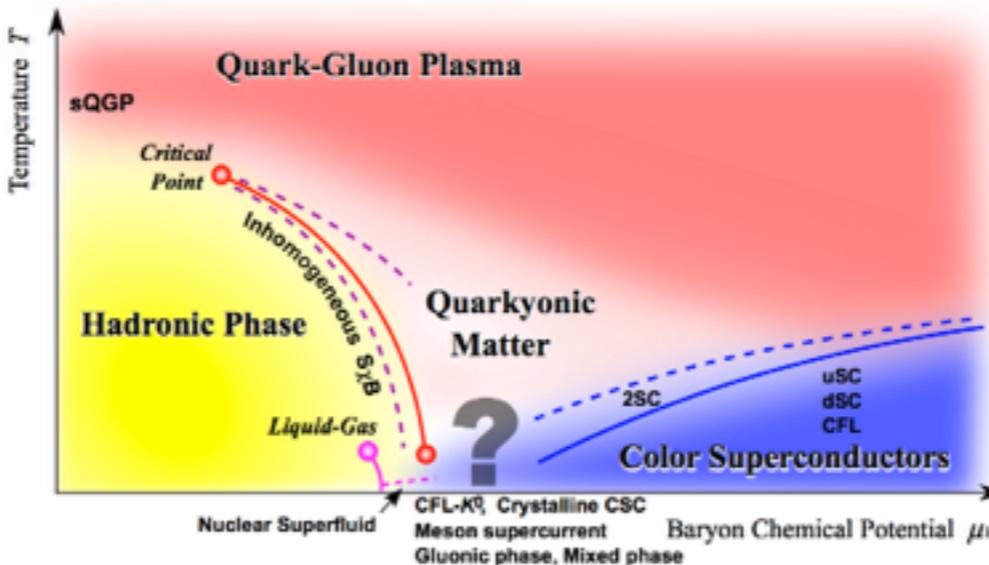
- inhomogeneous phases
- effects of baryons ?
- ...

Müller, Buballa and Wambach, PLB 727 (2013) 240

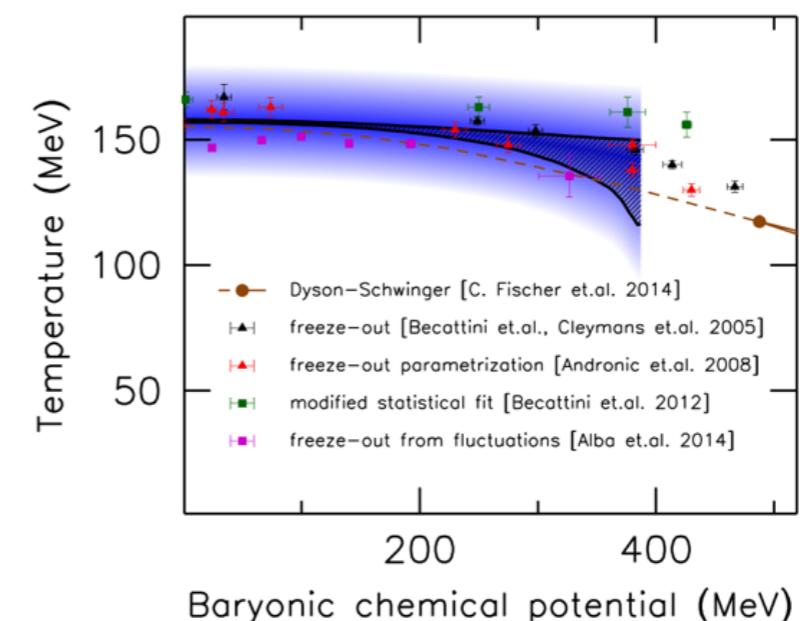
$N_c=2$ : Brauner, Fukushima and Hidaka, PRD 80 (2009) 74035  
Strodthoff, Schaefer and Smekal, PRD 85 (2012) 074007

# Overview

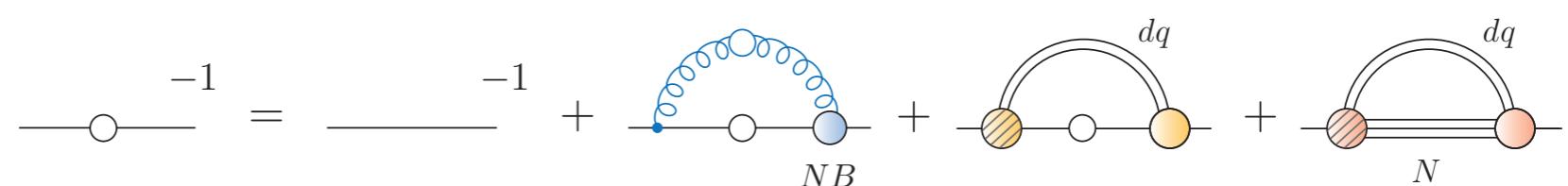
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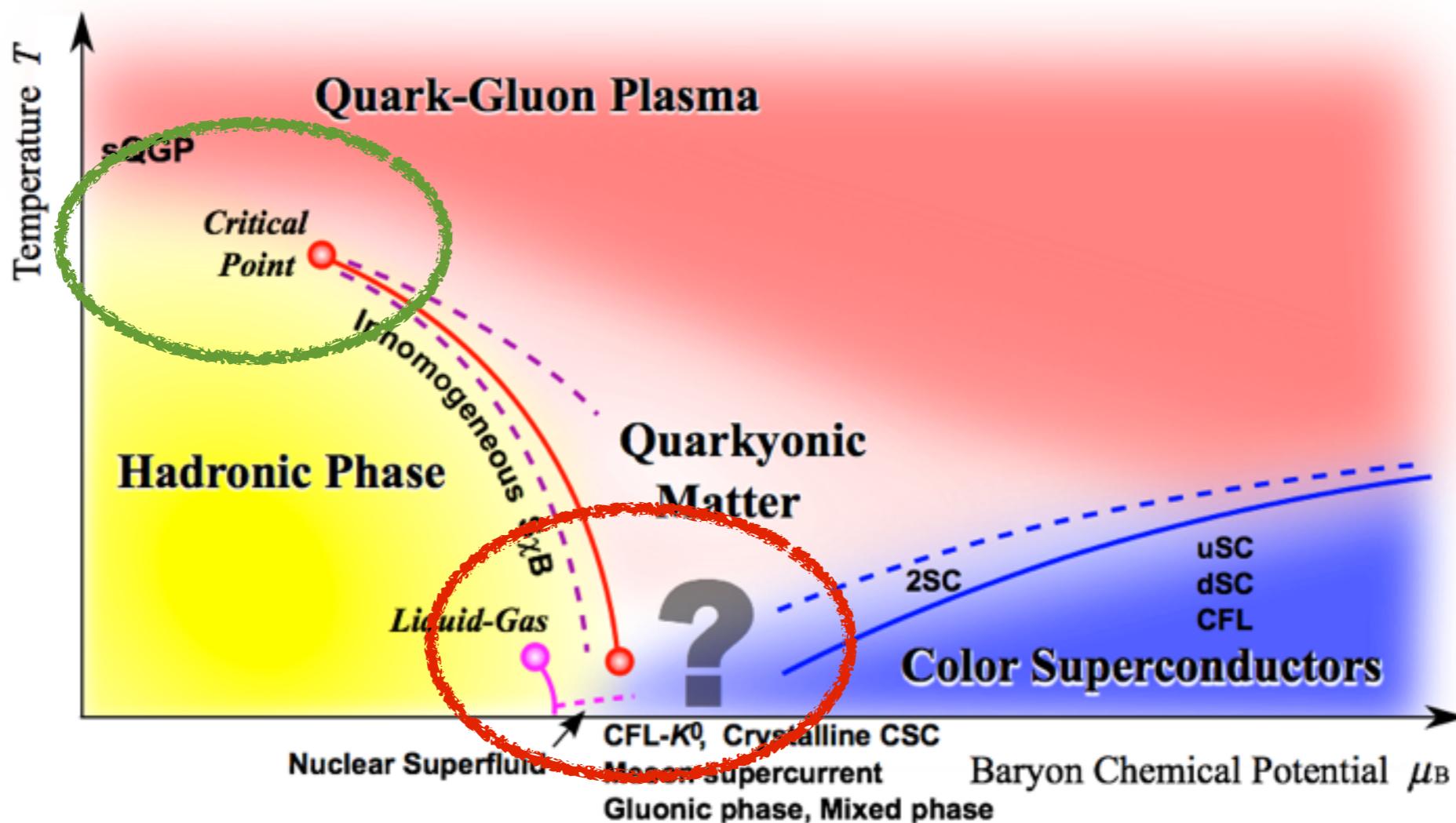


## 3. Baryon effects on the CEP



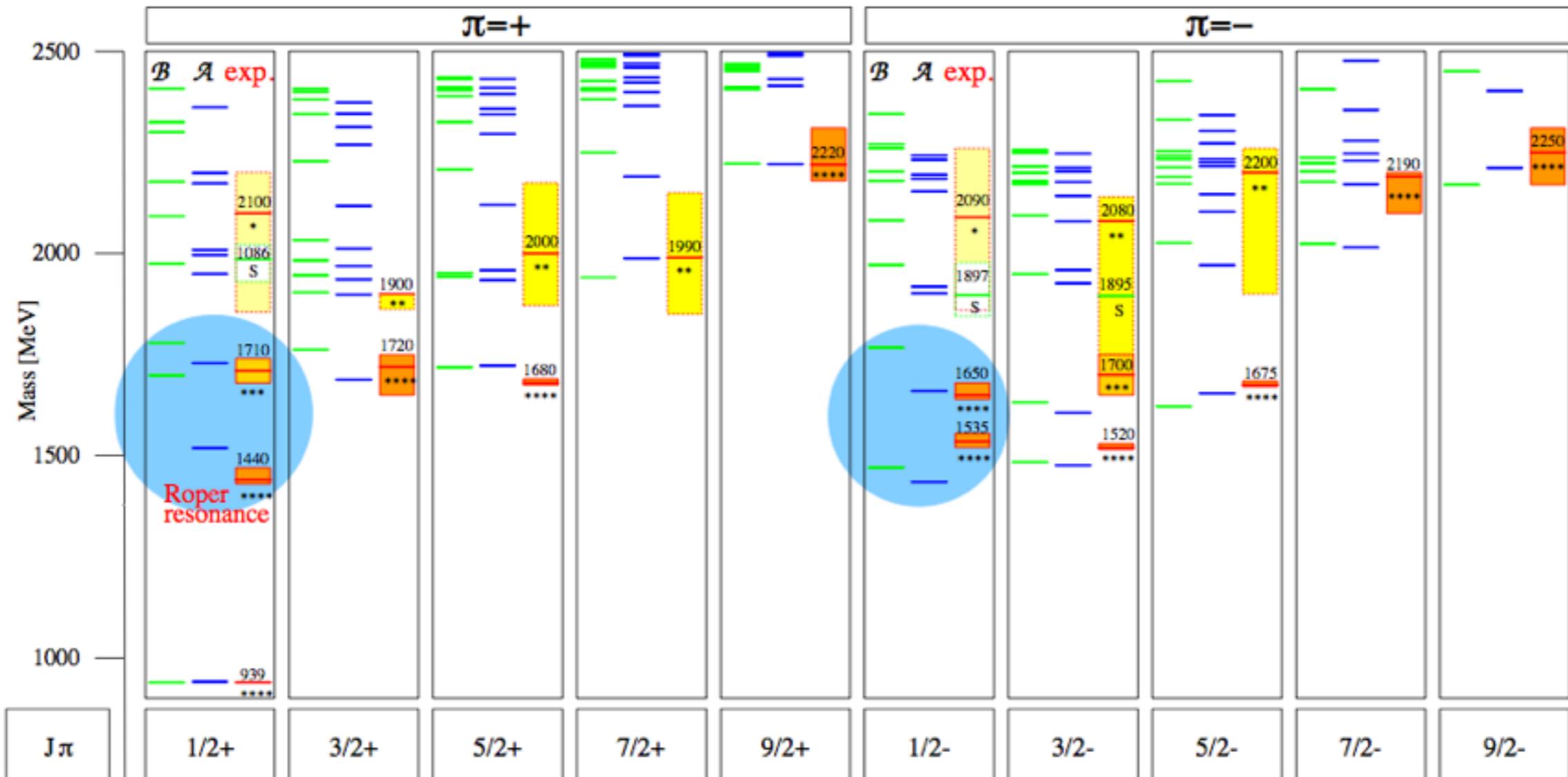
# QCD phase transitions I

Fukushima, Hatsuda, Rept. Prog. Phys. 74 (2011) 014001



- Low temperatures, large chemical potential:  
**baryons are important degrees of freedom**
- How do baryons affect the quark condensate ??

# Baryons: Quark model



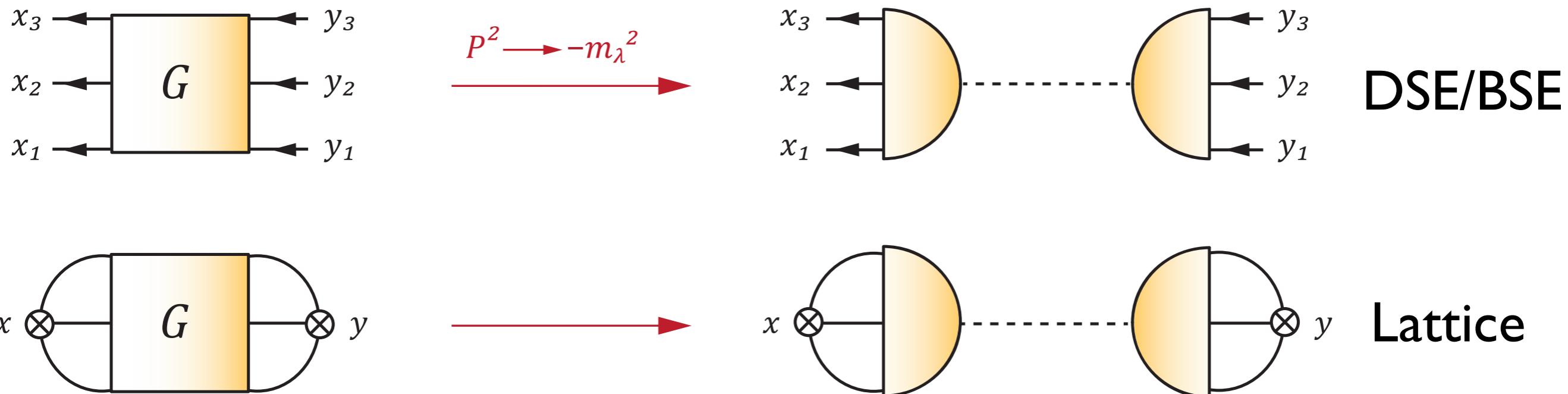
Loring, Metsch, Petry, EPJA 10 (2001) 395

- ‘missing resonances’ - three-body vs. quark-diquark
- level ordering:

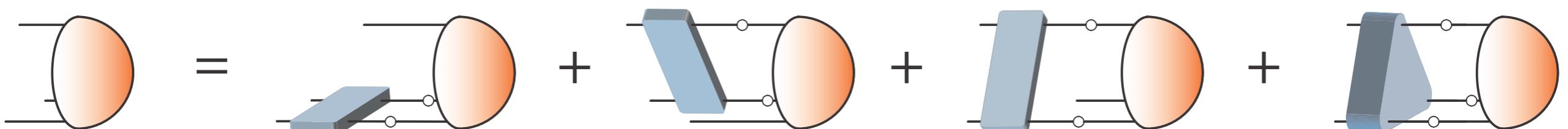
$$N \frac{1}{2}^\pm \quad \text{vs.} \quad \Lambda \frac{1}{2}^\pm$$

→ Eichmann, A2

# Extracting spectra from correlators



BSE for baryons (derived from equation of motion for  $G$ )



- exact equation for baryon ‘wave function’

# Baryon effects onto quark

$$\text{---}^{-1} = \text{---}^{-1} + \text{---} \quad NB \quad + \quad \text{---} \quad B$$

The diagram illustrates the decomposition of a quark's self-energy into two components. On the left, a horizontal line with a circle at its right end represents the full quark self-energy. An equals sign follows, and another horizontal line with a circle at its right end represents the non-baryonic part (NB). A plus sign indicates the sum of the non-baryonic part and the baryonic part (B). The baryonic part is shown as a loop containing a quark loop and a baryon loop.

Off-shell baryon  
in vertex

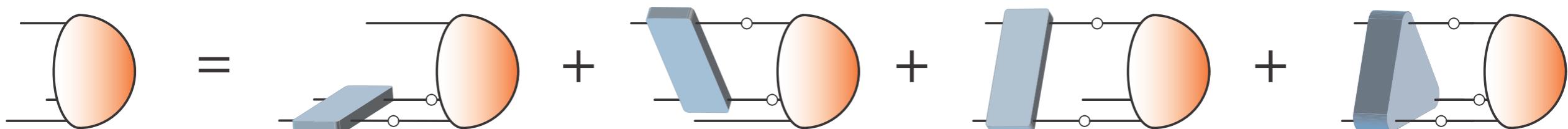
$$\text{---}^{-1} = \text{---}^{-1} + \text{---} \quad NB \quad + \quad \text{---} \quad dq \quad + \quad \text{---} \quad dq$$

This diagram shows the quark self-energy decomposition similar to the one above, but it includes a quark loop (dq) and a gluon loop (dq) contribution. The gluon loop is shaded with diagonal lines. A gluon exchange between two quarks is labeled with the letter  $N$ .

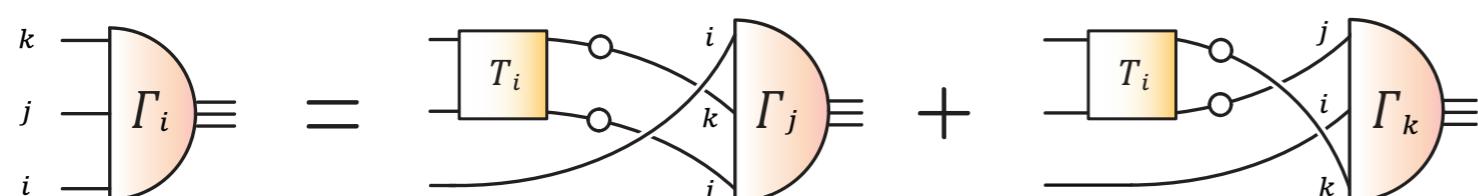
- Dependence on  $T$  and  $\mu$  via -propagators  
-wave functions
- Exploratory calculation: use wave functions from  $T=\mu=0$

# Vacuum: Baryons from BSEs

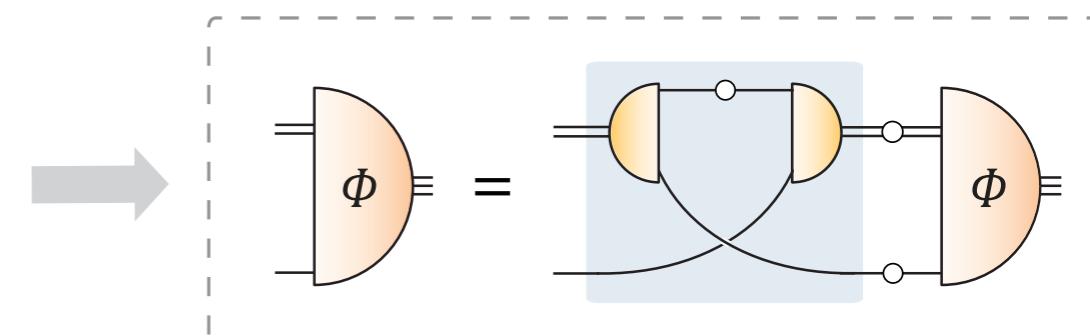
BSE for baryons (derived from equation of motion for G)



Faddeev equation (no three-body forces)

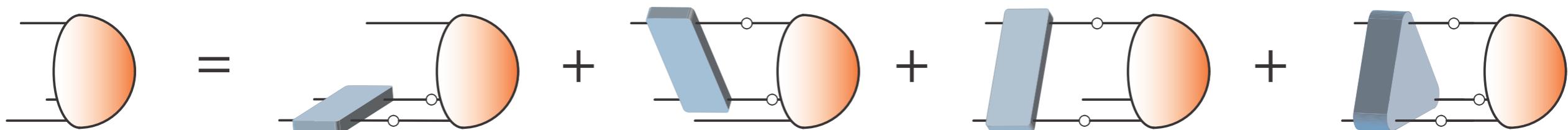


Diquark-quark

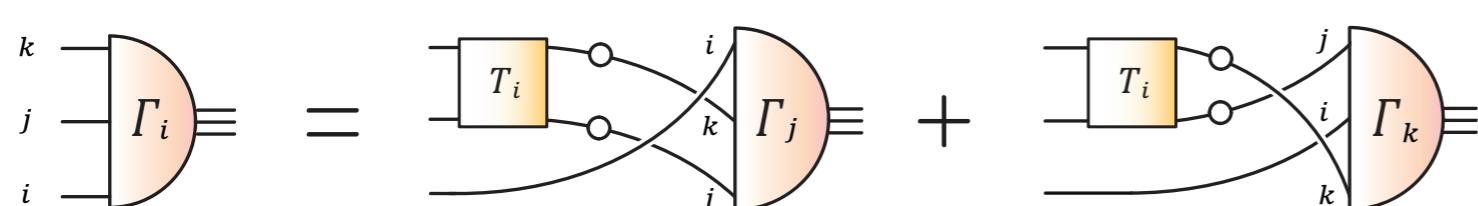


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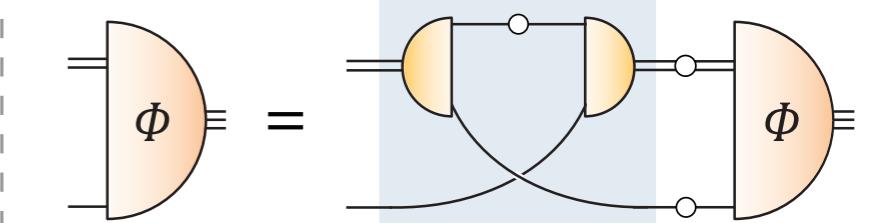
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Diquark-quark



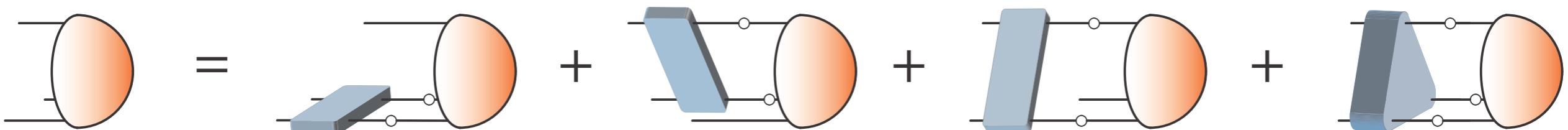
$$\text{---}^{-1} = \text{---}^{-1} + \text{---}$$

$$= \text{---}$$

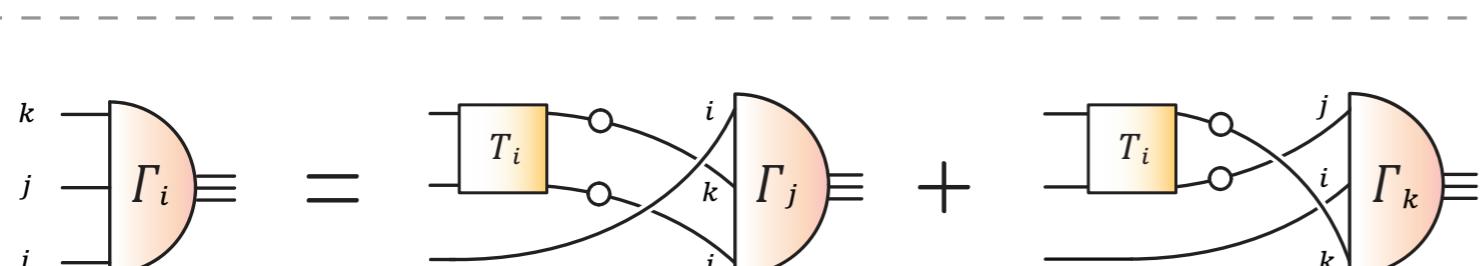
$$= \text{---}^{-1} = \text{---} + \text{---}$$

# Vacuum: Baryons from BSEs

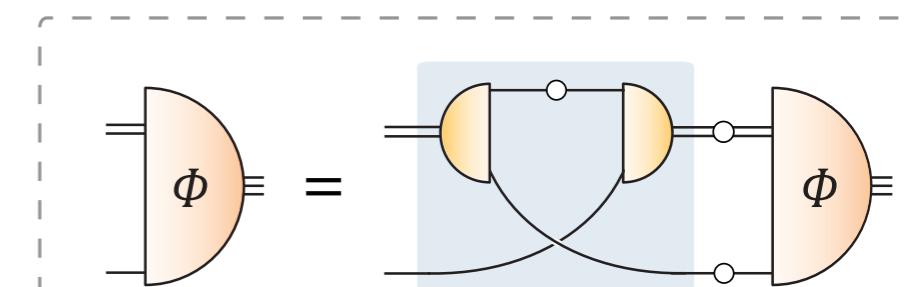
BSE for baryons (derived from equation of motion for G)



Faddeev equation (no three-body forces)



Diquark-quark



$$\text{---} \circ \text{---}^{-1} = \text{---} \text{---}^{-1} + \text{---} \text{---}$$

$$\text{---} \circ \text{---} = \text{---} \text{---} \text{---} \text{---}$$

$$\text{---} \circ \text{---}^{-1} = \text{---} \text{---} \text{---} \text{---} + \text{---} \text{---} \text{---} \text{---}$$

- Input: Non-perturbative quark, quark-gluon interaction (RL)

$$\text{---} \circ \text{---}^{-1} = \text{---} \text{---}^{-1} + \text{---} \text{---}$$

$$\alpha(k^2) = \pi \eta^7 \left( \frac{k^2}{\Lambda^2} \right) e^{-\eta^2 \left( \frac{k^2}{\Lambda^2} \right)} + \alpha_{UV}(k^2)$$

# Vacuum: DSE/Faddeev landscape

	Quark-diquark			Three-quark		
	Contact interaction	QCD-based model	DSE (RL)	RL	bRL	bRL + 3q
$N, \Delta$ masses	✓	✓	✓	✓	✓	...
$N, \Delta$ em. FFs	✓	✓	✓	✓		
$N \rightarrow \Delta \gamma$	✓	✓	✓	...		
Roper	✓	✓		...		
$N \rightarrow N^* \gamma$	✓	✓		...		
$N^*(1535), \dots$	...	...		...	...	
$N \rightarrow N^* \gamma$	...	...		...	...	

Roberts et al

Oettel, Alkofer  
Roberts, Bloch  
Segovia et al.

Eichmann, Alkofer  
Nicmorus, Krassnigg

Eichmann, Alkofer  
Sanchis-Alepuz, CF

Sanchis-Alepuz, CF  
Williams

Eichmann,  $N^*$ -Workshop, Trento 2015

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$N^*(1535), \dots$	...	...	...			...
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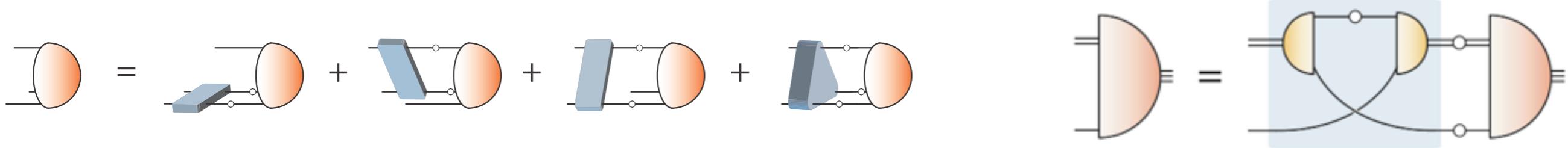
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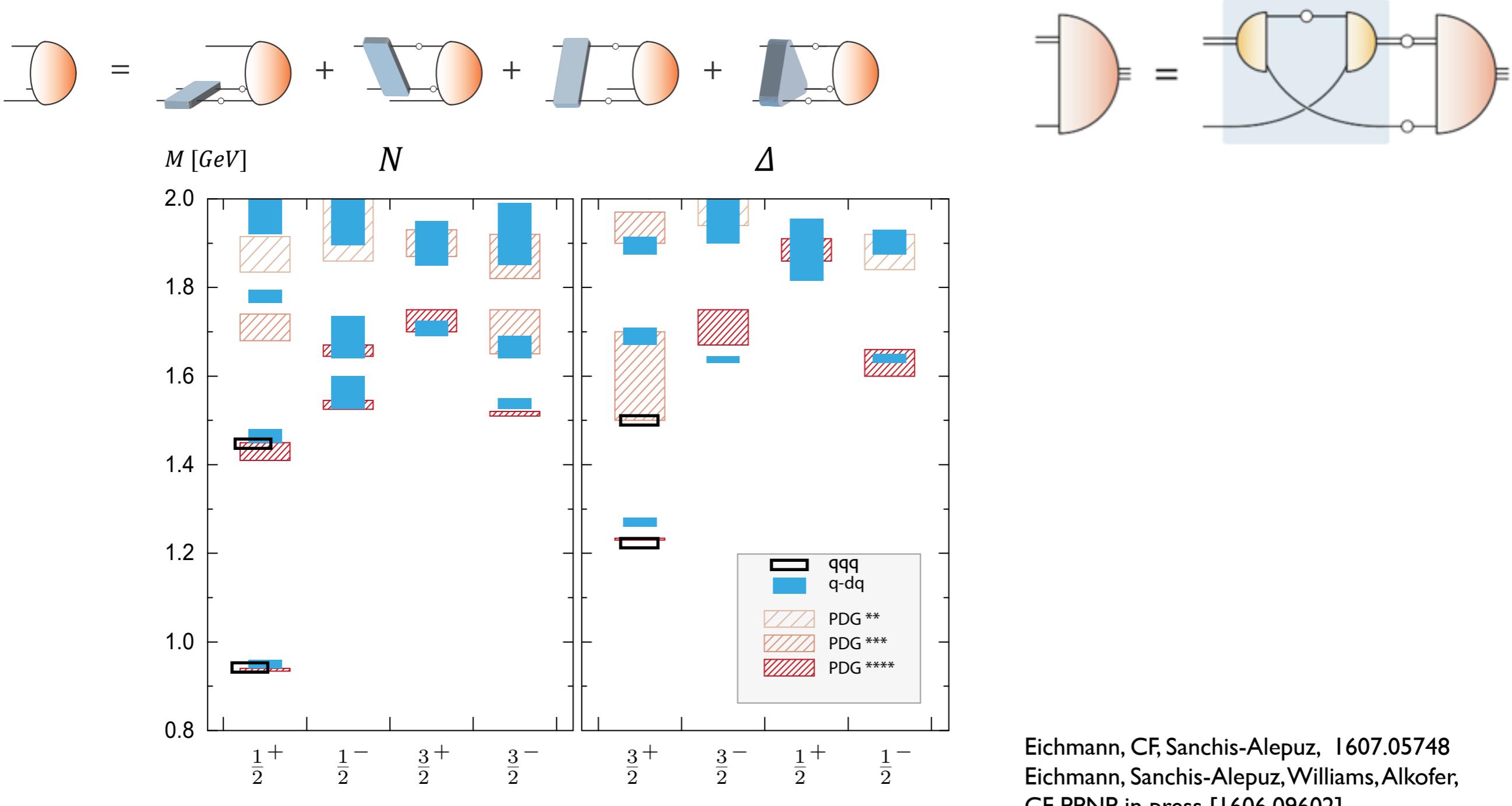
Sanchis-Alepuz, CF  
Williams

Eichmann,  $N^*$ -Workshop, Trento 2015

# Vacuum: Light baryon spectrum



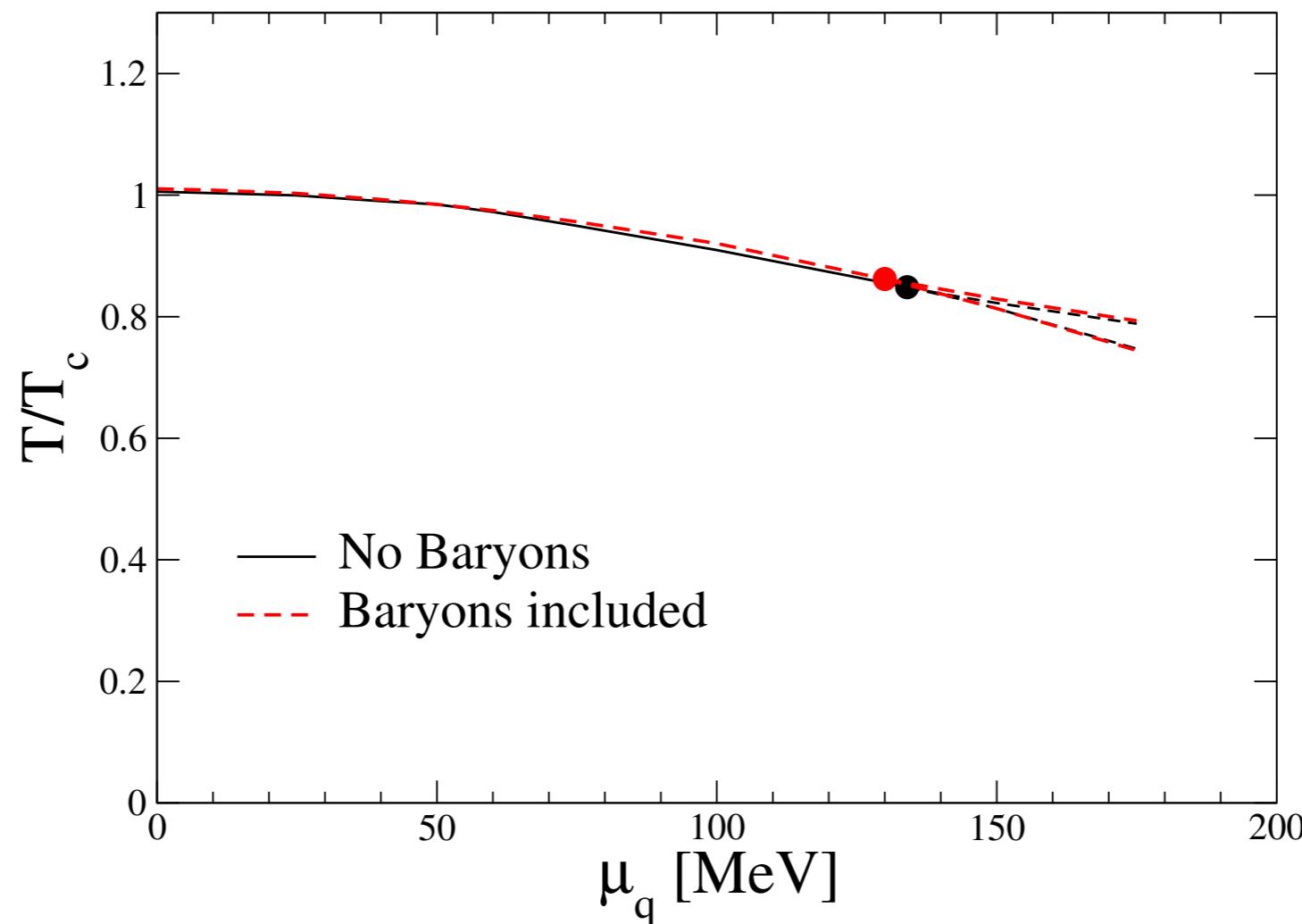
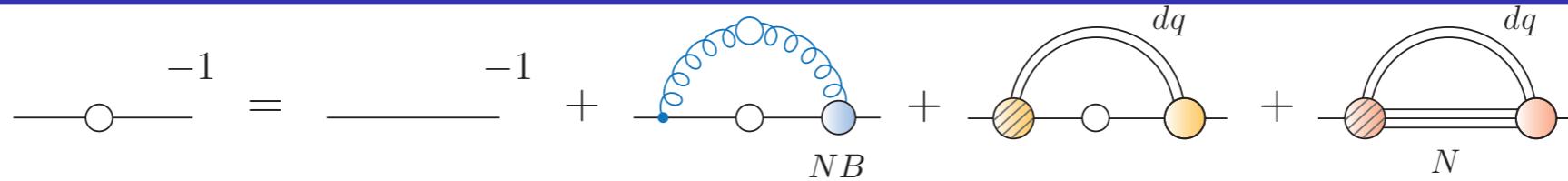
# Vacuum: Light baryon spectrum



Eichmann, CF, Sanchis-Alepuz, 1607.05748  
Eichmann, Sanchis-Alepuz, Williams, Alkofer,  
CF, PPNP in press [1606.09602]

- Three-body and diquark-quark approach agree qualitatively
- Spectrum in one-to-one agreement with experiment
- Correct level ordering (wo. coupled channel effects) !

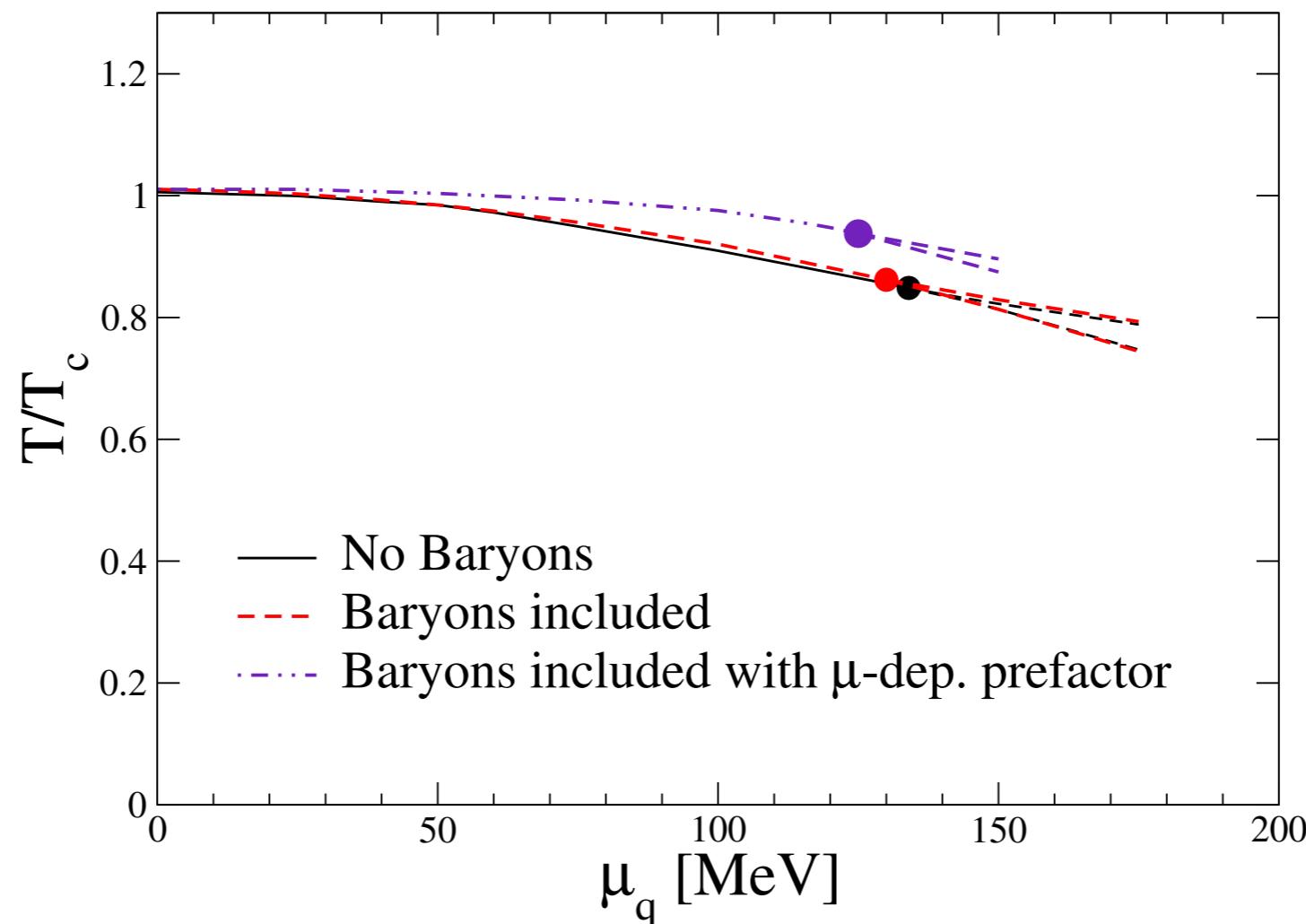
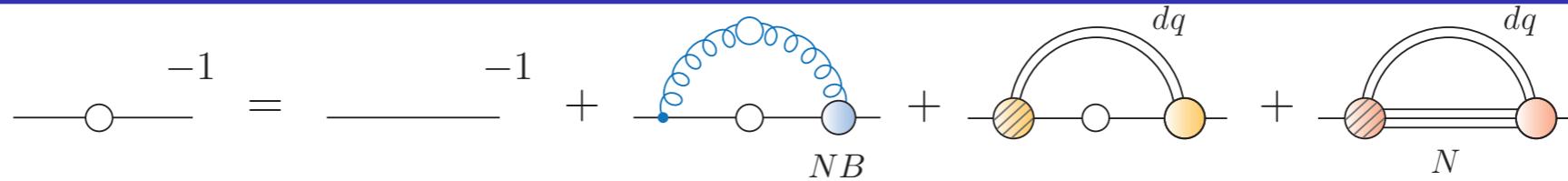
# Baryon effects on the CEP - results ( $N_f=2$ )



- Zero chemical potential: no effects after rescaling
- CEP: almost no effects

Eichmann, CF, Welzbacher, PRD93 (2016) [1509.02082]

# Baryon effects on the CEP - results ( $N_f=2$ )



- Zero chemical potential: no effects after rescaling
- CEP: almost no effects
- But: strong  $\mu$ -dependence of baryon wave function may change situation...

Eichmann, CF, Welzbacher, PRD93 (2016) [1509.02082]

# Summary

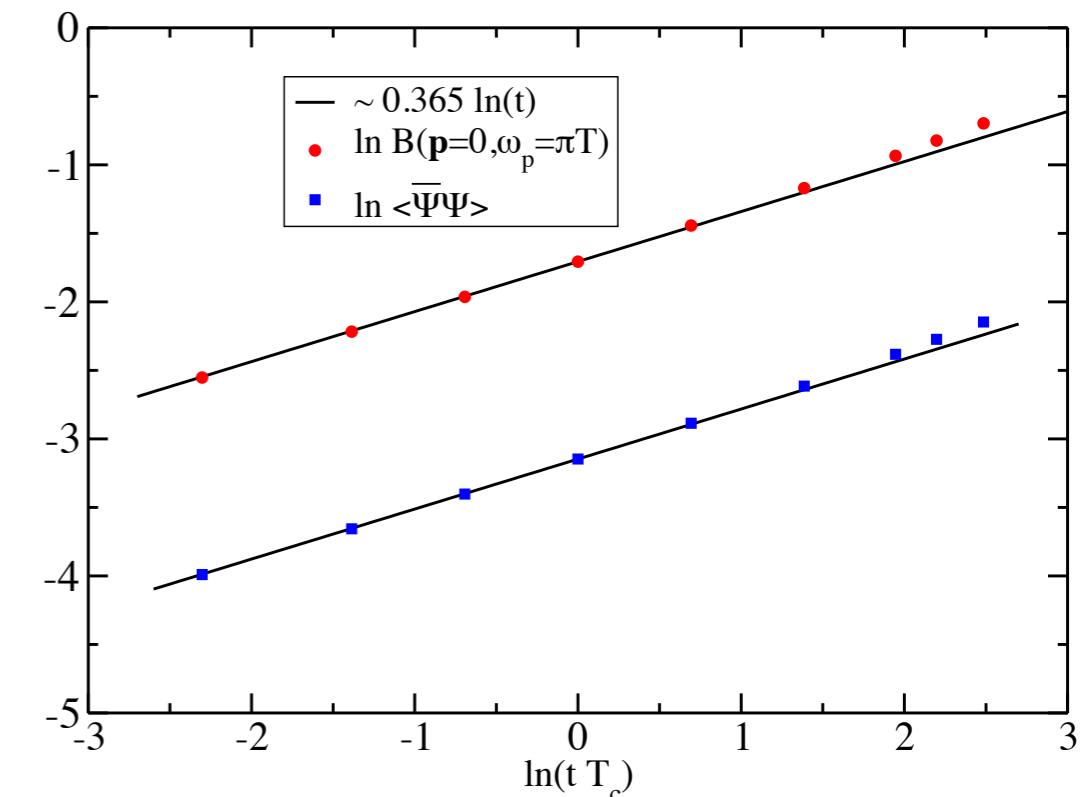
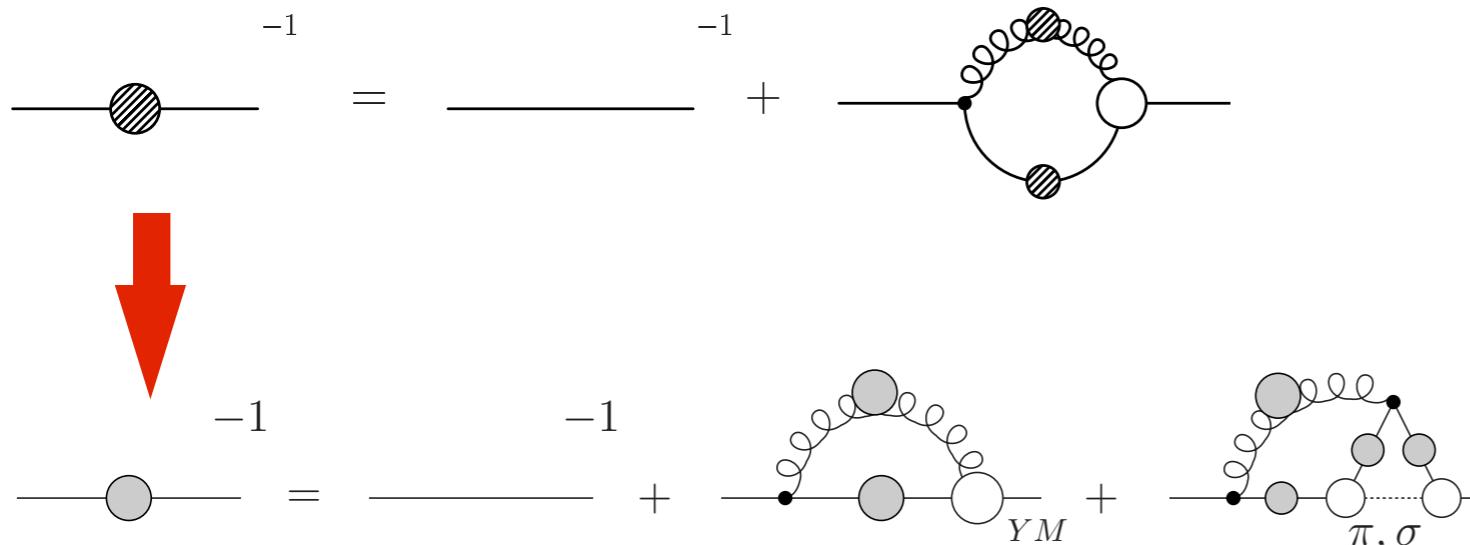
QCD with finite chemical potential:

- back-reaction of quarks onto gluons important
- $N_f=2+l$  and  $N_f=2+l+l$ : CEP at  $\mu_c/T_c > 3$
- charm quark does not influence CEP
- Baryon effects may or may not be significant for CEP...

Work in progress: - mesons and baryons at finite T and  $\mu$   
- volume effects of CEP from DSEs

# Backup

# Critical scaling from DSEs



- Need to take meson part of vertex explicitly into account
- $T=0$ : meson cloud corrections of order of 10-20 %  
CF, Williams,  
78 (2008) 074006
- $T=T_c$ : meson corrections are dominant !
- Critical scaling:  $\langle \bar{\Psi} \Psi \rangle(t) \sim B(t) \sim t^{\nu/2}$

$$f_{\pi,s}^2 \sim t^\nu \quad (t = (T_c - T)/T_c)$$

CF and Mueller, PRD 84 (2011) 054013

# Approximation for Quark-Gluon interaction

- T,  $\mu$ , m-dependent vertex:

$$\Gamma_\nu(q, k, p) = \tilde{Z}_3 \left( \delta_{4\nu} \gamma_4 \frac{C(k) + C(p)}{2} + \delta_{j\nu} \gamma_j \frac{A(k) + A(p)}{2} \right) \times \\ \times \left( \frac{d_1}{d_2 + q^2} + \frac{q^2}{\Lambda^2 + q^2} \left( \frac{\beta_0 \alpha(\mu) \ln[q^2/\Lambda^2 + 1]}{4\pi} \right)^{2\delta} \right)$$

Abelian WTI

perturbation theory

Infrared ansatz:

- d2 fixed to match gluon input
- d1 fixed via quark condensate (see later)
- correct UV (quant.) and IR-behavior (qual.)

CF, Pawłowski, PRD 80 (2009) 025023  
Mitter, Pawłowski and Strodthoff, PRD 91 (2015) 054035  
Williams, Fischer, Heupel, PRD 93 (2016) 034026

# DSE/Faddeev landscape ( $T=\mu=0$ )

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Oettel, Alkofer  
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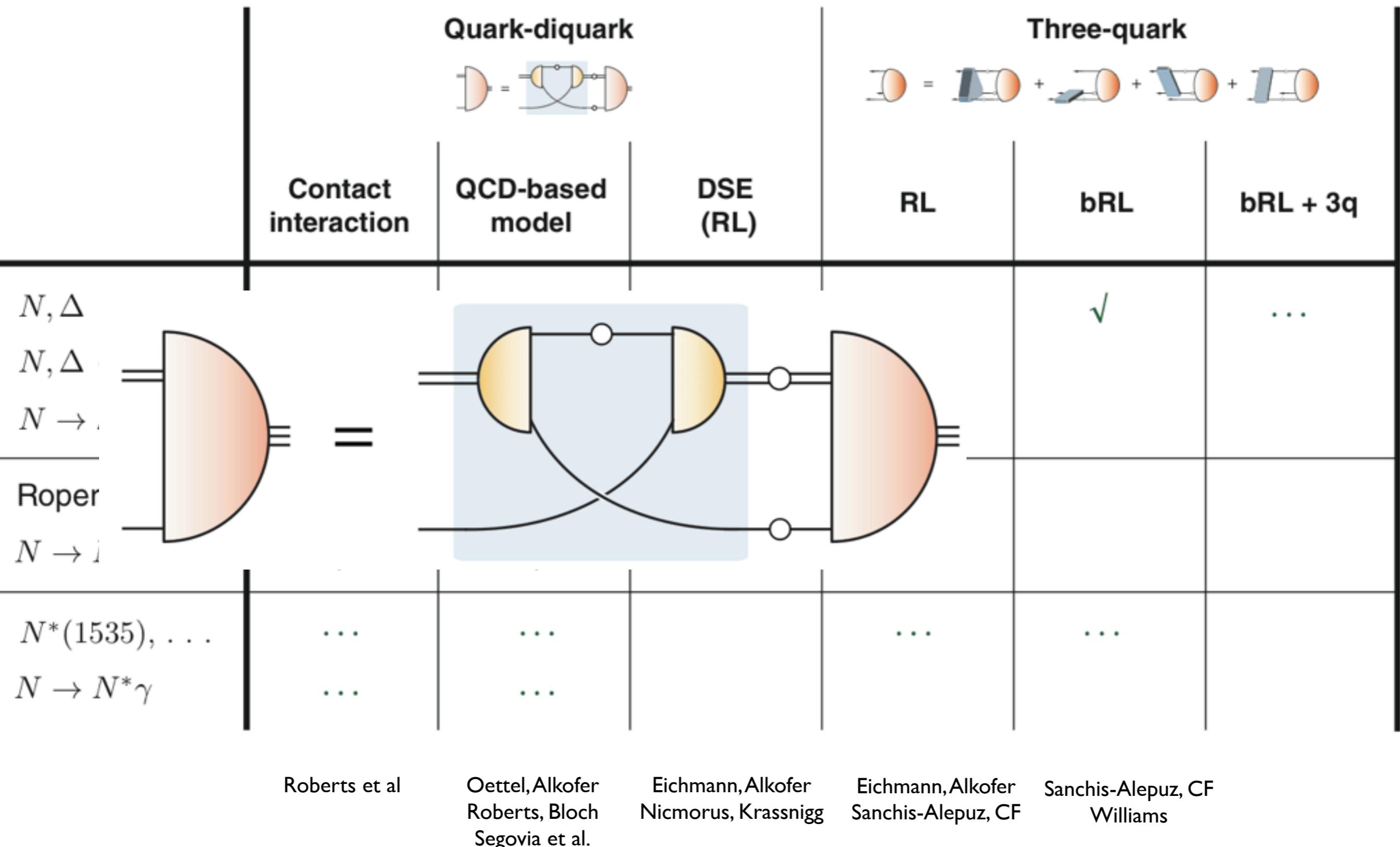
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Eichmann,  $N^*$ -Workshop, Trento 2015

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