

Towards a proper modeling of anisotropic magnetized compact objects

S. López Pérez¹

In collaboration with:

D. Alvear Terrero², A. Pérez Martínez²,
G. Quintero Angulo¹, D. Manreza Paret¹,
V. Hernández Mederos²

¹Havana University, ²ICIMAF

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Outline

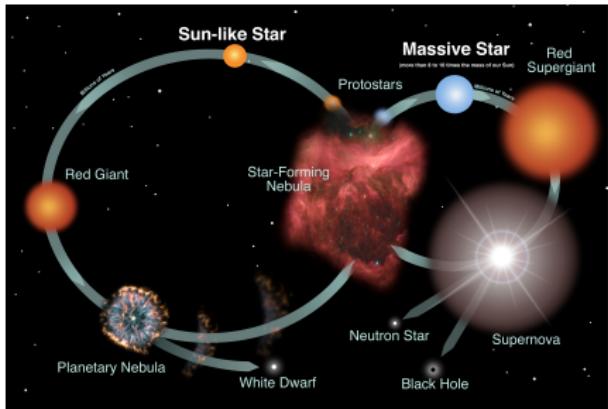
- Motivation
- Precedents
- Anistropic model for Compact Objects
- Preliminary results for White Dwarfs
- Summary and Future Work

Motivation

Compact Objects

- Final stage in the evolution of ordinary stars.
- Formed when thermonuclear fusion cannot compensate the gravitational collapse.

White Dwarfs	Neutron Stars
$M \sim M_{\odot}$	$M \sim 2M_{\odot}$
$R \sim 10^3 km$	$R \sim 10 km$
$B < 10^{13} G$	$B \sim 10^{18} G$
$M \sim R^{-3}$	$M \sim R^{-3}$
$\rho < 10^{11} \frac{g}{cm^3}$	$\rho < 10^{15} \frac{g}{cm^3}$



Motivation

Constant magnetic field effects on a fermion gas:

- Anisotropic energy-momentum tensor

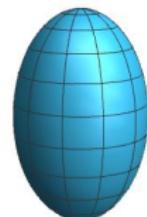
$$T_{\mu\nu} = \text{diag}(E, P_\perp, P_\perp, P_\parallel)$$

- Anisotropic equation of state.

$$E = \Omega + \mu N$$

$$P_\parallel = -\Omega$$

$$P_\perp = -\Omega - BM$$



Spherical symmetry is broken.

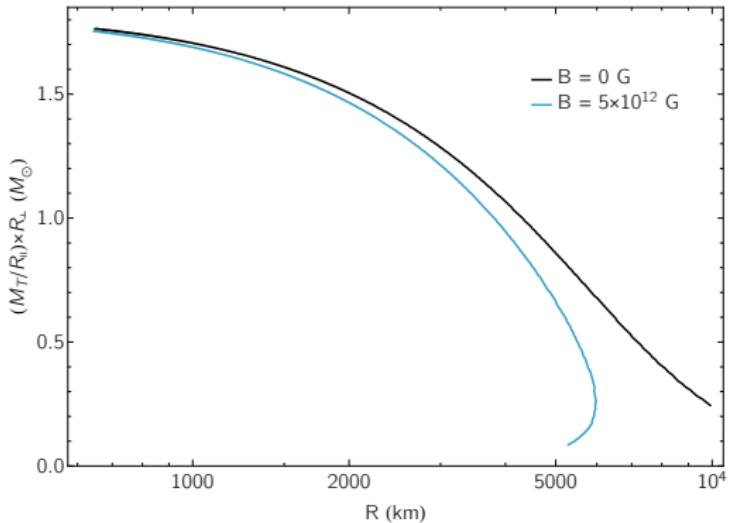
Precedents

- Metric in cylindrical coordinates

$$ds^2 = -e^{2\Phi(r)} dt^2 + e^{2\Lambda(r)} dr^2 + r^2 d\phi^2 + e^{2\Psi(r)} dz^2$$

- Mass by Tolman definition

$$\frac{M_T}{R_{||}} = 4\pi \int_0^{R_\perp} r e^{\Phi(r) + \Lambda(r) + \Psi(r)} (E - 2P_\perp - P_{||}) dr$$

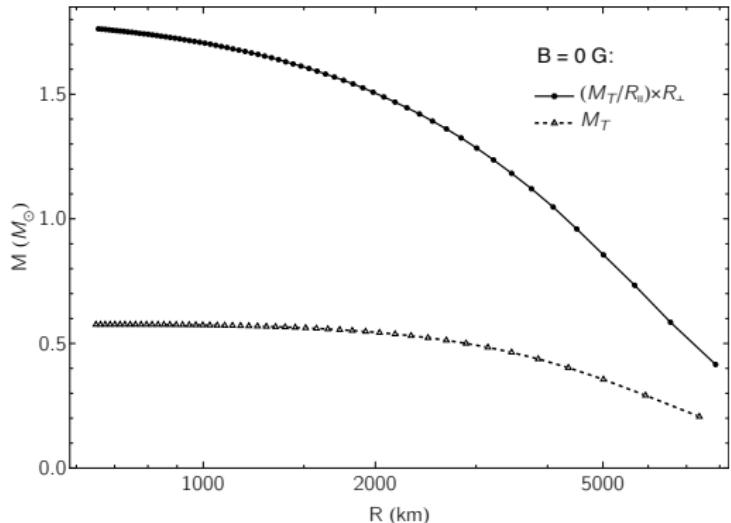


D. Manreza Paret et.al., RAA 15,
2015.

Precedents

- Metric in cylindrical coordinates
- Ansatz: $\frac{R_{\parallel}}{R_{\perp}} = \frac{P_{\parallel}}{P_{\perp}}$
- Mass by Tolman definition

$$M_T = 4\pi \int_0^{R_{\perp}} \int_0^{r^{\frac{P_{\parallel}}{P_{\perp}}}} r e^{\Phi(r) + \Lambda(r) + \Psi(r)} (E - 2P_{\perp} - P_{\parallel}) dz dr$$



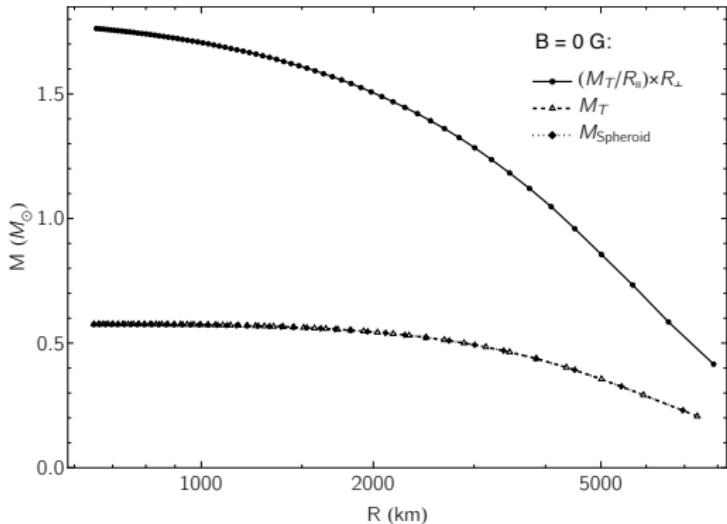
Precedents

- Metric in cylindrical coordinates

- Ansatz: $\gamma = \frac{P_{||0}}{P_{\perp0}} = \frac{z}{r}$

- Spheroid mass

$$\frac{dM}{dr} = 4\pi r^2 E \gamma$$



Anisotropic model for Compact Objects

- γ Metric in spherical coordinates (arXiv:gr-qc/9810079v1; Omair Zubairi et al 2015 J. Phys.: Conf. Ser. 615 012003)

$$ds^2 = - \left(1 - \frac{2M(r)}{r}\right)^\gamma dt^2 + \left(1 - \frac{2M(r)}{r}\right)^{-\gamma} dr^2 + r^2 \sin \theta d\phi^2 + r^2 d\theta^2$$

- Ansatz: $\gamma = \frac{P_{\parallel 0}}{P_{\perp 0}} = \frac{z}{r}$
- Structure equations:

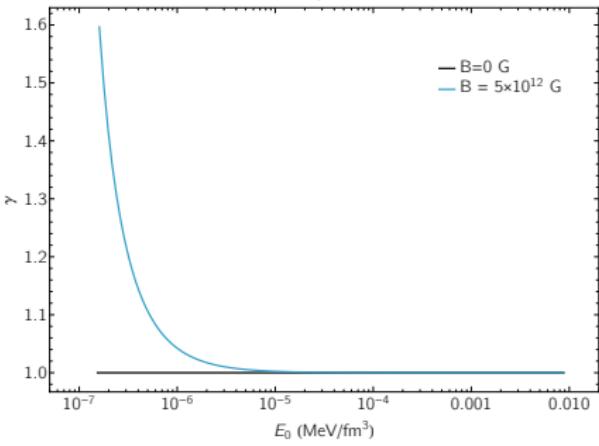
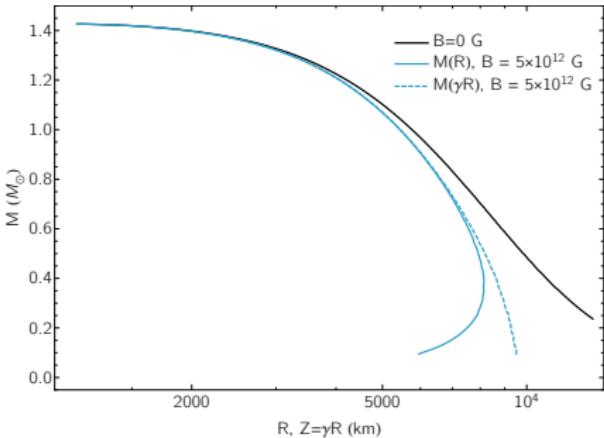
$$\frac{dM}{dr} = r^2 E \gamma$$

$$\frac{dP_{\parallel}}{dr} = - \frac{(E + P_{\parallel})[\frac{r}{2} + r^3 P_{\parallel} - \frac{r}{2}(1 - \frac{2M}{r})^{\gamma}]}{r^2(1 - \frac{2M}{r})^{\gamma}}$$

$$\frac{dP_{\perp}}{dz} = \frac{1}{\gamma} \frac{dP_{\perp}}{dr} = - \frac{(E + P_{\perp})[\frac{r}{2} + r^3 P_{\perp} - \frac{r}{2}(1 - \frac{2M}{r})^{\gamma}]}{\gamma r^2(1 - \frac{2M}{r})^{\gamma}}$$

Reduces to TOV equations at $P_{\parallel} \equiv P_{\perp}$ ($B = 0$ case)

Preliminary results for White Dwarfs



- Magnetic fields deform the star.
- The effects become important in the low density regime.

Summary and Future Work

- We have given the first steps towards a generalized structure equations that can account for the anisotropies present in magnetized compact stars and reduce to the standard TOV equations when $B = 0$.
- Magnetic fields seem to affect the shape of the star, and therefore, the radii, but not the maximum mass.

The near future:

- Study stability in more detail.
- Obtain solutions for different types of compact stars.
- Set $\gamma = P_{\parallel}/P_{\perp}$ varying throughout the star.



Structure equations in cylindrical coordinates

- Metric in cylindrical coordinates

$$ds^2 = -e^{2\Phi(r)} dt^2 + e^{2\Lambda(r)} dr^2 + r^2 d\phi^2 + e^{2\Psi(r)} dz^2$$

- Mass by Tolman definition

$$M_T = \int \sqrt{-g} (T_0^0 - T_1^1 - T_2^2 - T_3^3) dV$$

- Einstein equations

$$P'_\perp = -\Phi'(E + P_\perp) - \Psi'(P_\perp - P_\parallel),$$

$$4\pi e^{2\Lambda}(E + P_\parallel + 2P_\perp) = \Phi'' + \Phi'(\Psi' + \Phi' - \Lambda') + \frac{\Phi'}{r},$$

$$4\pi e^{2\Lambda}(E + P_\parallel - 2P_\perp) = -\Psi'' - \Psi'(\Psi' + \Phi' - \Lambda') - \frac{\Psi'}{r},$$

$$4\pi e^{2\Lambda}(P_\parallel - E) = \frac{1}{r}(\Psi' + \Phi' - \Lambda').$$