

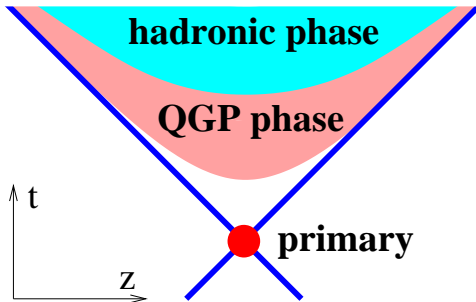
**10th International Symposium on Non-equilibrium Dynamics, 25 - 29 November, 2024,
Krabi, Thailand**

New concepts in EPOS4

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Space-time picture of pp, pA, AA at high energy

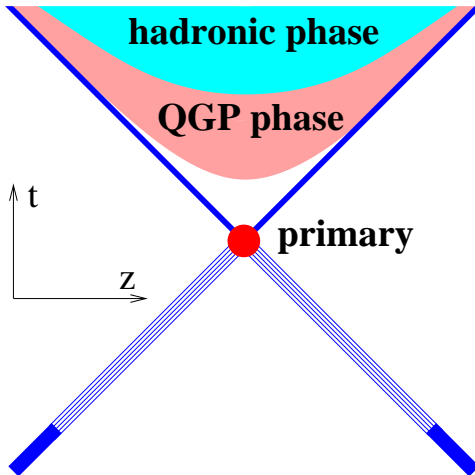


Primary interactions, pointlike overlap
(even in pA, AA scattering: large γ factors)

followed (later) by QGP formation

BUT the picture is not really correct...

More realistic space-time picture



splitting into
multiple partons
(parton evolution)
long in advance,
takes a long time
(large γ factors)

**but the interaction
region (red point)
is pointlike**

multiple scatterings must happen in parallel

EPOS4 philosophy

concerning primary interactions

- **Avoid sequential scatterings,**
 - **concerning both parton-parton**
 - **and nucleon-nucleon interactions**

- **Do multiple scatterings**
rigorously in parallel

- **Respect the rule “MC = theory”**

Some technical remarks:

Graphs are nice ... but it should be clear what is behind. We use

symbol	meaning
$T(s, t)$	elastic scattering T-matrix; s, t Mandelstam variables
$T(s, b)$	Fourier transformation of $T(s, t)$ with respect to the momentum transfer, divided by $2s$ (impact parameter representation)
G	$2 \operatorname{Im} T$ – representing inelastic scattering (cut diagram)
$\tilde{\sigma}$	pp cross sections: $\sigma^{pp} = \int d^2b \tilde{\sigma}^{pp}(s, b)$ $A+B$ cross sections: $\sigma^{AB} = \int db_{AB} \tilde{\sigma}^{AB}(s, b, \{b_i^A\}, \{b_i^B\})$

$$\int db_{AB} = \int d^2b \int \prod_{i=1}^A d^2b_i^A T_A(b_i^A) \int \prod_{j=1}^B d^2b_j^B T_B(b_j^B),$$

with transv. nucleon coordinates b_i^A and b_j^B , with the nuclear thickness function

$$T_A(b) = \int dz \rho_A(\sqrt{b^2 + z^2})$$

where ρ_A is the (normalized) nuclear density for nucleus A .

Warmup: Gribov-Regge (GR) approach

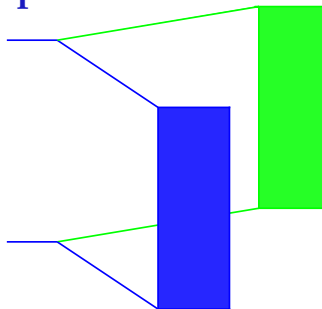
V. A. Abramovsky, V. N. Gribov, O. V. Kancheli, L. N. Lipatov (1967-1973)

Multiple scattering in pp strictly parallel

**Box: inelastic subscattering G ,
producing chain of particles**

(precise structure: unknown)

**No contradictions even
when long “preparation”**



Xsections expressed in terms of weights P depending on G

$$\tilde{\sigma}_{\text{in}}^{pp} = \sum_{m=1}^{\infty} \underbrace{\frac{1}{m!} G^m e^{-G}}_{P(m)},$$

$$\tilde{\sigma}_{\text{in}}^{AB} = \sum_{\{m_k\}} \underbrace{\prod_{k=1}^{AB} \frac{1}{m_k!} (G_k)^{m_k} e^{-G_k}}_{P(\{m_k\}) \text{ basis of MC}}$$

EPOS4 improvement, step 1:

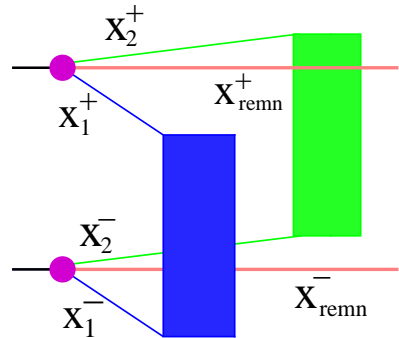
- **Implement energy-momentum conservation, referred to as GR⁺**
- **For certain observables not so important**
(total cross sections)
- **For others absolutely crucial**
(particle production)
- **Necessary as solid basis for MC**
(otherwise contradictions, but GR still widely used)

Energy-momentum sharing (GR⁺) in EPOS

In pp or for each NN scattering in $A+B$

New variables: lightcone momentum fractions x_m^+ and x_m^- of subscatterings, being conserved:

$$x_{\text{remn}}^\pm = 1 - \sum x_m^\pm \quad (pp)$$



Xsections still expressed in terms of weights P of configurations $K = \{ \{m_k\}, \{x_{k\mu}^\pm\} \}$ (m_k subscatterings per pair k , with $x_{k\mu}^\pm$)

Solid basis of Monte Carlo:

- one determines K according to $P(K)$,
- instantaneously, no sequences, in parallel!! Here: MC = theory

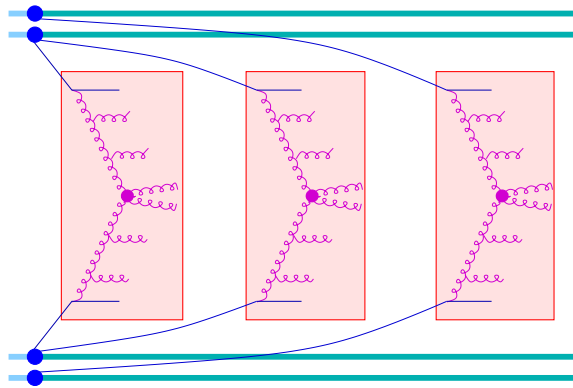
EPOS4 improvement, step 2:

- So far: general framework, based on “some G ”
- G represents a subscattering
- **Now: make link with QCD: $G = G_{\text{QCD}}$**
- **G_{QCD} represents parton-parton scattering, based on pQCD, including DGLAP evolution**

See: K. Werner and B. Guiot, PRC 108, 034904 (2306.02396)

Early work (no HF): H.J. Drescher, M. Hladik, S. Ostapchenko, T. Pierog,
K. Werner, Phys.Rept. 350 (2001) 93-289 (hep-ph/0007198)

Replace boxes
by QCD

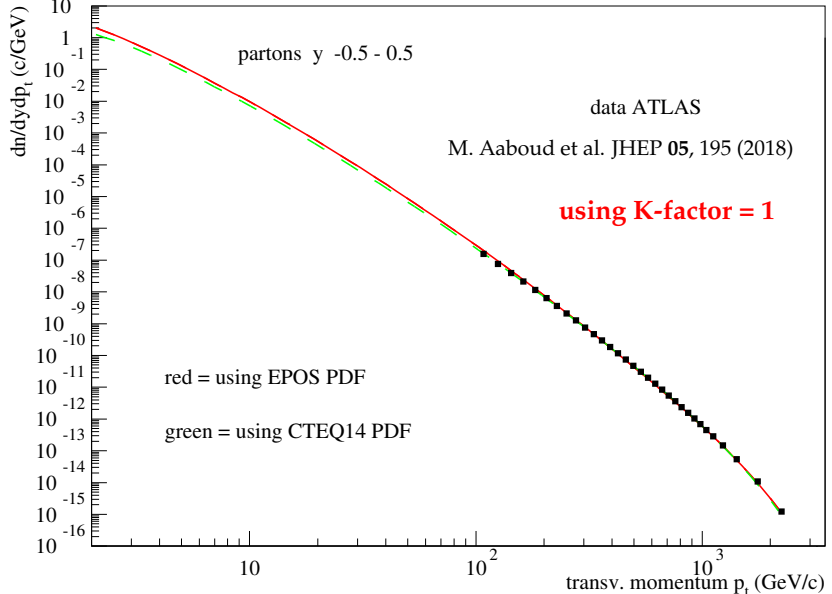


collision of two nuclei with three subscatterings

**We compute and tabulate “moduls”
(QCD evolution, Born cross sections, vertices)
which then allow to evaluate the diagram
Different ways to rearrange the modules...**

one may define (and tabulate) a PDF, allowing to compute the jet cross section vs p_t for pp at 13 TeV

$$E_3 E_4 \frac{d^6 \sigma}{d^3 p_3 d^3 p_4} = \sum_{klmn} \iint d\xi_1 d\xi_2 f_{\text{PDF}}^k(\xi_1, \mu_F^2) f_{\text{PDF}}^l(\xi_2, \mu_F^2) \frac{1}{32s\pi^2} \sum_i |\mathcal{M}^{kl \rightarrow mn}|^2 \delta^4(p_1 + p_2 - p_3 - p_4) \frac{1}{1 + \delta_{mn}}$$



Looks good, but

- Here we considered just one single subscattering
- In GR, the full multiple scattering scenario is equal to the single one for inclusive cross sections (AGK theorem)

$$\frac{d\sigma_{\text{incl}}^{AB}}{dp_t} = AB \times \frac{d\sigma_{\text{incl}}^{\text{single scattering}}}{dp_t}$$

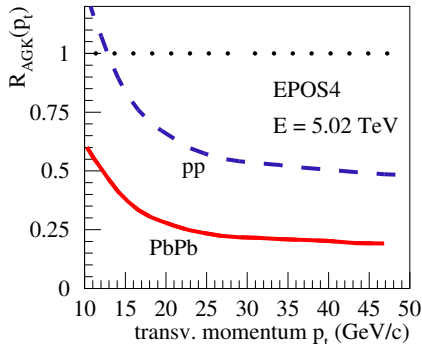
- Does AGK hold in our case (GR⁺) ?
- And does AGK hold for nuclear scattering (which would amount to binary scaling)?

Validity of AGK Check p_t of partons
for minimum bias PbPb and pp scatterings at 5.02 TeV.

Ratio

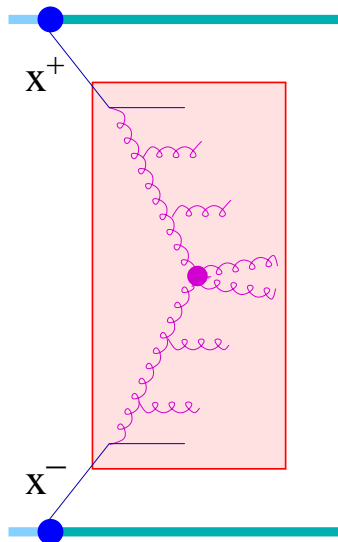
$$R_{AGK}(p_t) = \frac{d\sigma_{incl}^{AB}}{dp_t} / \left\{ AB \times \frac{d\sigma_{incl}^{single\ scattering}}{dp_t} \right\}$$

should be unity



AGK badly violated!!!

The problem is
the energy sharing
among subscatterings



Inclusive particle spectra (like p_t) are determined by the distribution of the LC momenta x^+ and x^- of the subscatterings.

Crucial variable: the squared CMS energy fraction

$$x_{PE} = x^+ x^- \approx s / s_{\text{tot}}$$

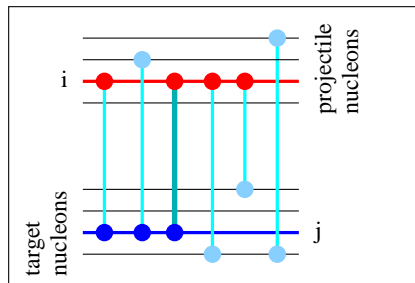
For a given scattering, involving
projectile nucleon i and
target nucleon j

define:

$$N_{\text{conn}} = \frac{N_P + N_T}{2}$$

N_P = number of scatterings involving i

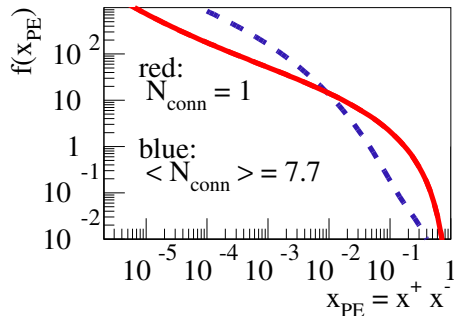
N_T = number of scatterings involving j



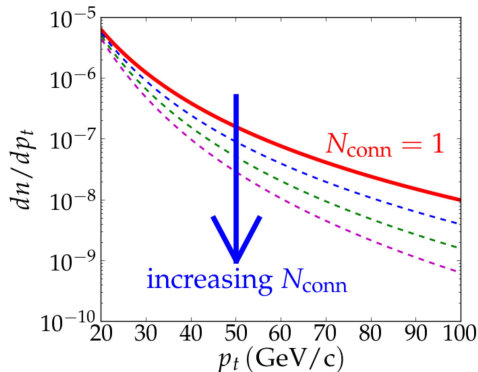
The x_{PE} distributions $f(x_{\text{PE}})$
depend on N_{conn}

Large $N_{\text{conn}} \Rightarrow$ large x_{PE} suppressed
small x_{PE} enhanced

We will use the notation $f^{(N_{\text{conn}})}(x_{\text{PE}})$



Large N_{conn} \Rightarrow large x_{PE} suppressed \Rightarrow large p_t suppressed



Min, bias pp or AA = superposition of different N_{conn} contributions

Cannot be equal to the single-scattering case ($N_{\text{conn}} = 1$)
 \Rightarrow violation of AGK

We define the “deformation” of $f^{(N_{\text{conn}})}(x_{\text{PE}})$ relative to the reference $f^{(1)}(x_{\text{PE}})$

$$R_{\text{deform}} = \frac{f^{(N_{\text{conn}})}(x_{\text{PE}})}{f^{(1)}(x_{\text{PE}})}$$

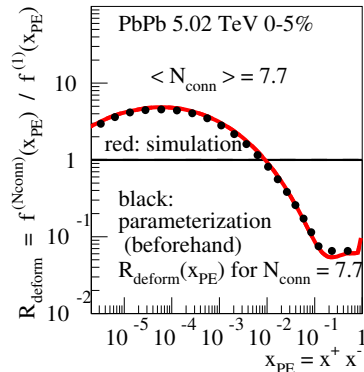
$R_{\text{deform}} \neq 1$ creates the problem

But we are able to parameterize R_{deform} and tabulate it, for all systems, all centrality classes

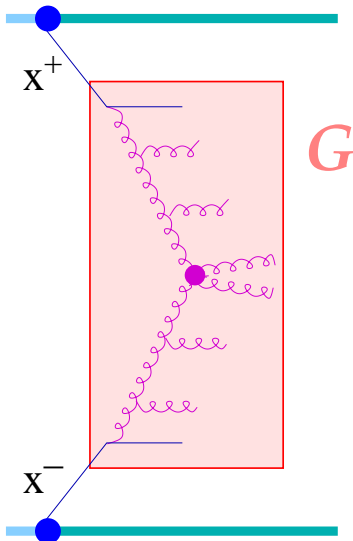
So:

$$R_{\text{deform}} = R_{\text{deform}}(N_{\text{conn}}, x_{\text{PE}})$$

can be considered to be known, it is tabulated and available via interpolation (to be used later).



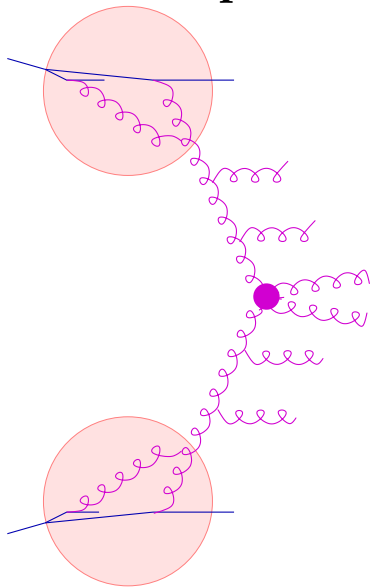
There are actually two problems



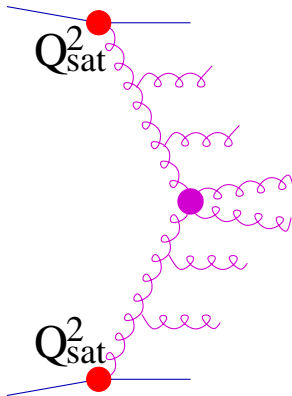
concerning the single scattering expression G , the fundamental building block of the multiple scattering formalism

- **The assumption $G = G_{\text{QCD}}$ seems to be wrong (AGK problem)**
- **Nonlinear effects are completely missing**

EPOS4 improvement, step 3: Add saturation



Saturation phenomena
(nonlinear effects, inside circles)
may be “summarized”
by saturation scales



Saturation phenomena
(nonlinear effects, inside circles)
may be “summarized”
by saturation scales

suggesting to treat nonlinear effects
by introducing saturation scales Q_{sat}^2
as the lower limits Q_0^2 of the virtualities
for DGLAP evolutions

**We compute and tabulate
 $G_{\text{QCD}}(Q_0^2, x^+, x^-, s, b)$ for a large
range of Q_0^2 values**

K. Werner and B. Guiot, PRC 108, 034904 (2306.02396)

For the connection between the basic multiple scattering building block G and the QCD expression G_{QCD} one postulates:

For each subscattering, for given x^\pm , s , b , and N_{conn} :

$$G(x^+, x^-, s, b) = n \frac{G_{\text{QCD}}(Q_{\text{sat}}^2, x^+, x^-, s, b)}{R_{\text{deform}}(N_{\text{conn}}, x_{\text{PE}})}$$

**such that G does not depend on N_{conn} ,
whereas Q_{sat}^2 does depend on x^+ , x^- , N_{conn}**

n is a normalization constant

Early attempts in this direction:

K. Werner, F.-M. Liu, and T. Pierog, Phys. Rev. C 74, 044902 (2006), hep-ph/0506232

K. Werner, B. Guiot, I. Karpenko, and T. Pierog, J. Phys. Conf. Ser. 458, 012020 (2013)

T. Pierog and K. Werner, Acta Phys. Polon. Supp. 8, 1031 (2015)

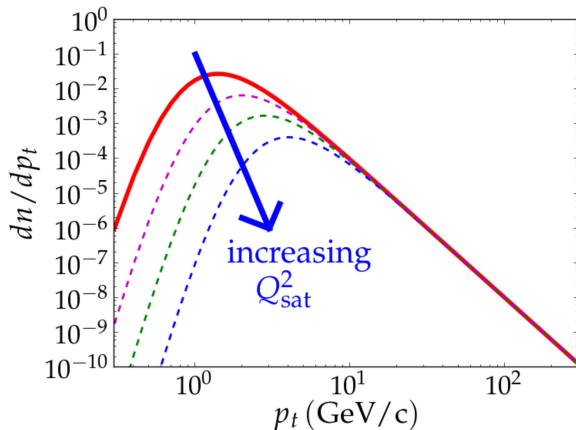
On can show see K. Werner, PRC 109, 034918 (2310.09380)

$$\frac{d^2\sigma_{\text{incl}}^{AB(N_{\text{conn}})}}{dx^+ dx^-} \propto \frac{d\sigma_{\text{incl}}^{\text{single scattering}}}{dx^+ dx^-} [Q_{\text{sat}}^2(N_{\text{conn}}, x^+, x^-)]$$

i.e., the A+B cross section (given N_{conn})

- is equal to the single scattering case,**
- but with Q_{sat}^2 corresponding to N_{conn}**

Same relation for p_t distributions (deduced from $x^+ x^-$)



one expects with increasing N_{conn}

- an increasing Q_{sat}^2
- and a reduction at $p_t^2 < Q_{\text{sat}}^2$ compared to $N_{\text{conn}} = 1$ (red)

But no change for large p_t . If interested in large p_t :

One replaces Q_{sat}^2 by some constant $Q_0^2 = \max\{Q_{\text{sat}}^2\}$

One gets finally

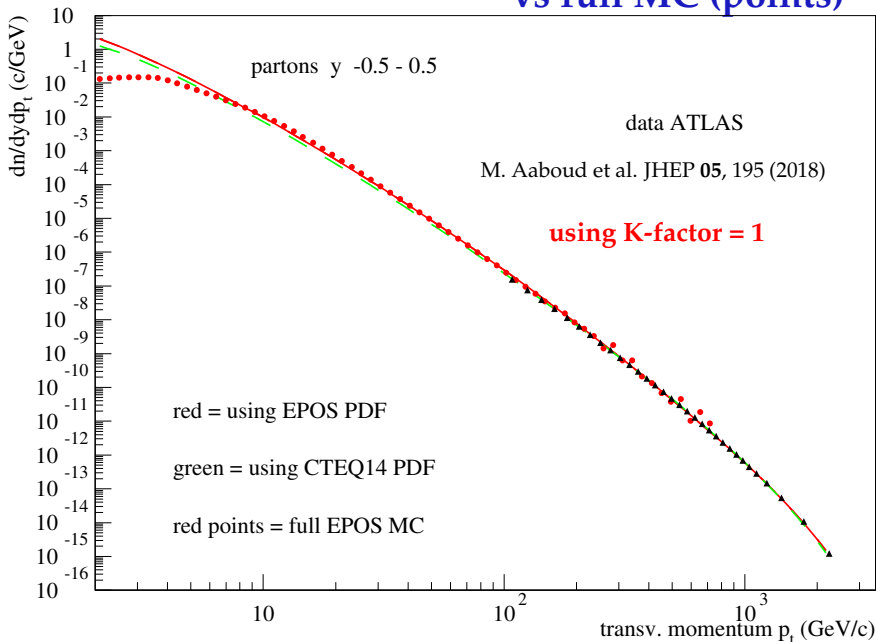
$$\frac{d\sigma_{\text{incl}}^{AB} (mb)}{dp_t} = AB \frac{d\sigma_{\text{incl}}^{\text{single scattering}}}{dp_t} [Q_0^2]$$

but only for p_t^2 bigger than the relevant Q_{sat}^2 values
(gAGK theorem)

Extremely important: One gets factorization (in pp and $A+B$) for inclusive cross sections at high p_t in a fully self-consistent (*) multiple (parallel) scattering scheme.

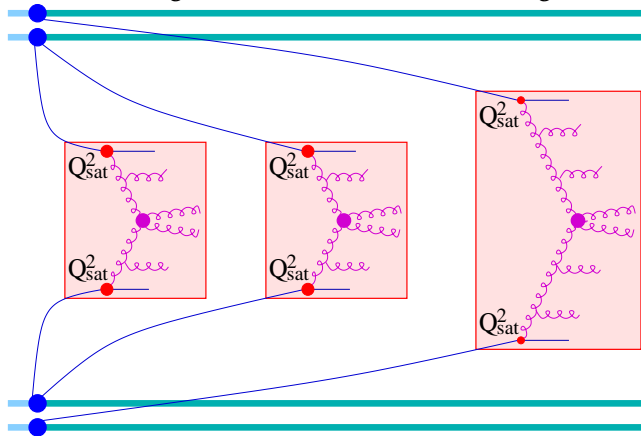
(*) Mandatory: (A) energy-mom. conservation, (B) parallel scattering,
(C) MC = theory, (D) factorization

Jet cross section vs p_t for pp at 13 TeV, factorization result vs full MC (points)



How to understand N_{conn} -dependent saturation scales

$A + B$ scattering ($A = B = 2$) with 3 scatterings



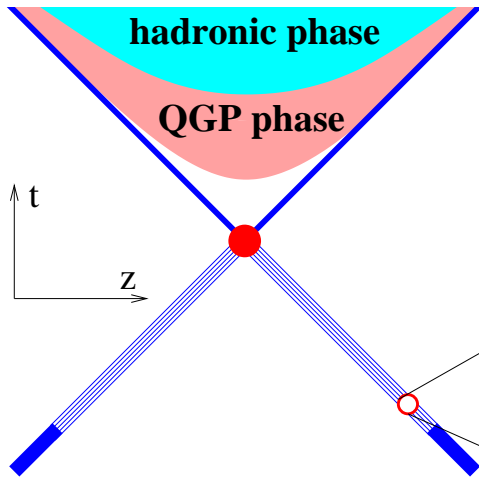
Two left scatterings compared to right one:

- N_{conn} bigger
- energy (\sqrt{s}) smaller
- Q_{sat}^2 bigger (bigger dots)
- evolution shorter
- central part identical
- particle production harder

formulas: K. Werner, B. Guiot, PRC 108, 034904 (2306.02396)

Implementing both energy-mom conservation and saturation is needed, the one compensates the other, such that the central part remains unchanged

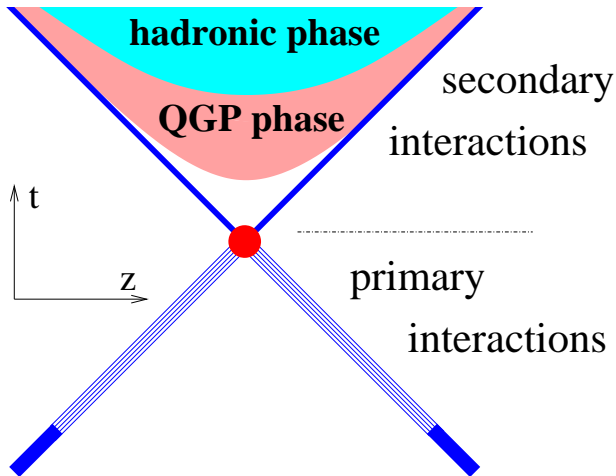
**Nonlinear effects
happen very early**



dynamical process
happening long before
the actual projectile-target
overlap (red point)

**initial state phenomenon,
contrary to final state
"string manipulations"**
(like color reconnections)

Secondary interactions



EPOS4:

Solid basis for the primary interactions, setting the stage for what follows

Secondary interactions:

core-corona separation¹
 core hydro evolution²
 microcanonical decay¹
 hadronic cascade³

¹) K. Werner, PRC 109, 014910 (2024), arXiv:2306.10277

²) I. Karpenko et al, Computer Physics Communications 185, 3016 (2014),
 K. Werner, B. Guiot, I. Karpenko, and T. Pierog, PRC 89, 064903 (2014), 1312.1233

³) S. A. Bass et al., Prog. Part. Nucl. Phys. 41, 225 (1998),
 M. Bleicher et al., J. Phys. G25, 1859 (1999)

**Core, corona, full
pp at 7 TeV**

pions, kaons, protons,
lambdas (top to bottom)

Green: $\frac{\text{core}}{\text{core} + \text{corona}}$

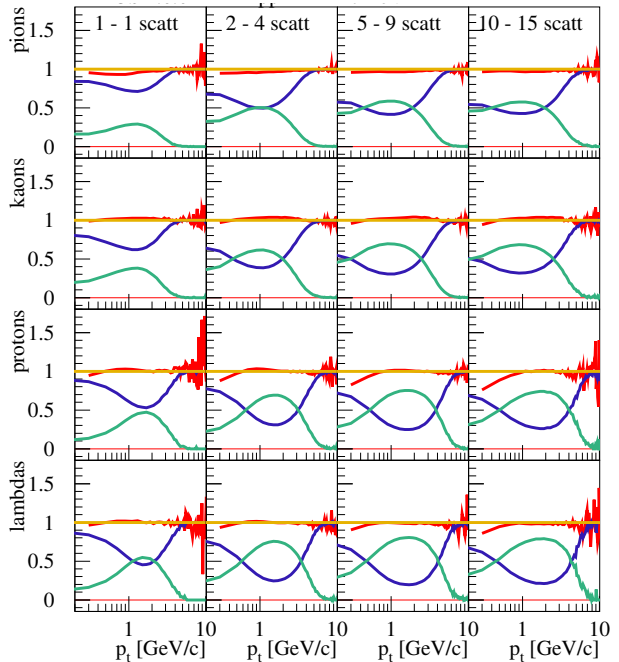
Blue: $\frac{\text{corona}}{\text{core} + \text{corona}}$

Red: $\frac{\text{full}}{\text{core} + \text{corona}}$

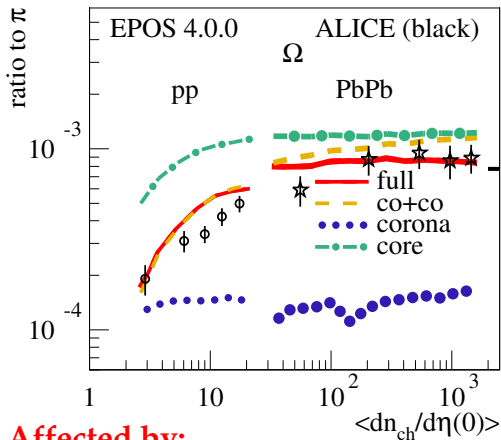
Core reaches to higher p_t for
baryons

Core has maximum at inter-
mediate p_t (flow)

Rescattering not very impor-
tant



continuous curve



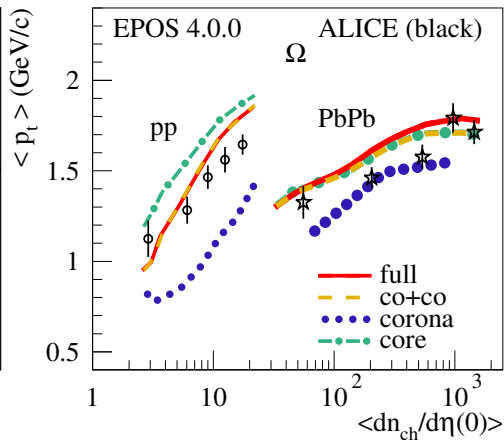
Affected by:

core-corona

microcanonical

hadronic cascade (UrQMD)

jump



saturation

flow

core-corona

EPOS4 web page

<https://klaus.pages.in2p3.fr/epos4/>

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<p style="text-align: center;">EPOS4: A Monte Carlo tool for</p> <p style="text-align: center;">from electron-positron annihilation</p> <p>Whereas elementary processes such as electron-proton and electron-nucleus scattering, EPOS4 are following standard procedures found in the literature. The parallel scattering process factorization approach is quite unique at high energies, multiple (nucleonic or partonic) primary scatterings are treated sequentially.</p> <p>The EPOS4 approach brings together ancient knowledge (multiple scattering theory and Regge theory) and modern concepts of perturbative QCD (k_T factorization approach). The parallel scattering process factorization approach is quite unique at high energies, multiple (nucleonic or partonic) primary scatterings are treated sequentially.</p>	

thanks Laurent Aphecetche
and Damien Vintache

Physics

Papers

Basis of EPOS4

<https://klaus.pages.in2p3.fr/epos4/physics/papers#basis>

EPOS4HQ

<https://klaus.pages.in2p3.fr/epos4/physics/papers#eposhq>

Code

very soon:

EPOS4.0.3

EPOS4HQ first public version

Summary (concerning primary scattering in EPOS4)

- **There are ad hoc assumptions, details may be questioned, but the overall multiple-scattering picture seems mandatory, since based on very fundamental principles:**
 - * **parallel scattering formalism**
 - * **energy-momentum sharing (EMS)**
 - * **implementing saturation**
 - Seems mandatory to implement “environment dependent” factorization scales, which compensate exactly the effect of the EMS
 - * **validity of factorization at high pt**
 - * **MC = theory**

- **Solid basis for further activities:**
 - EPOS4 systematic improvements**
 - EPOS4HQ, including quarkonia and HF collectivity,**
 - EPOS4JET ...**