

The 10th International Symposium on Non-equilibrium Dynamics (NeD-2024), Nov. 25-29, 2024

Elliptic flow in heavy-ion collisions at intermediate energies

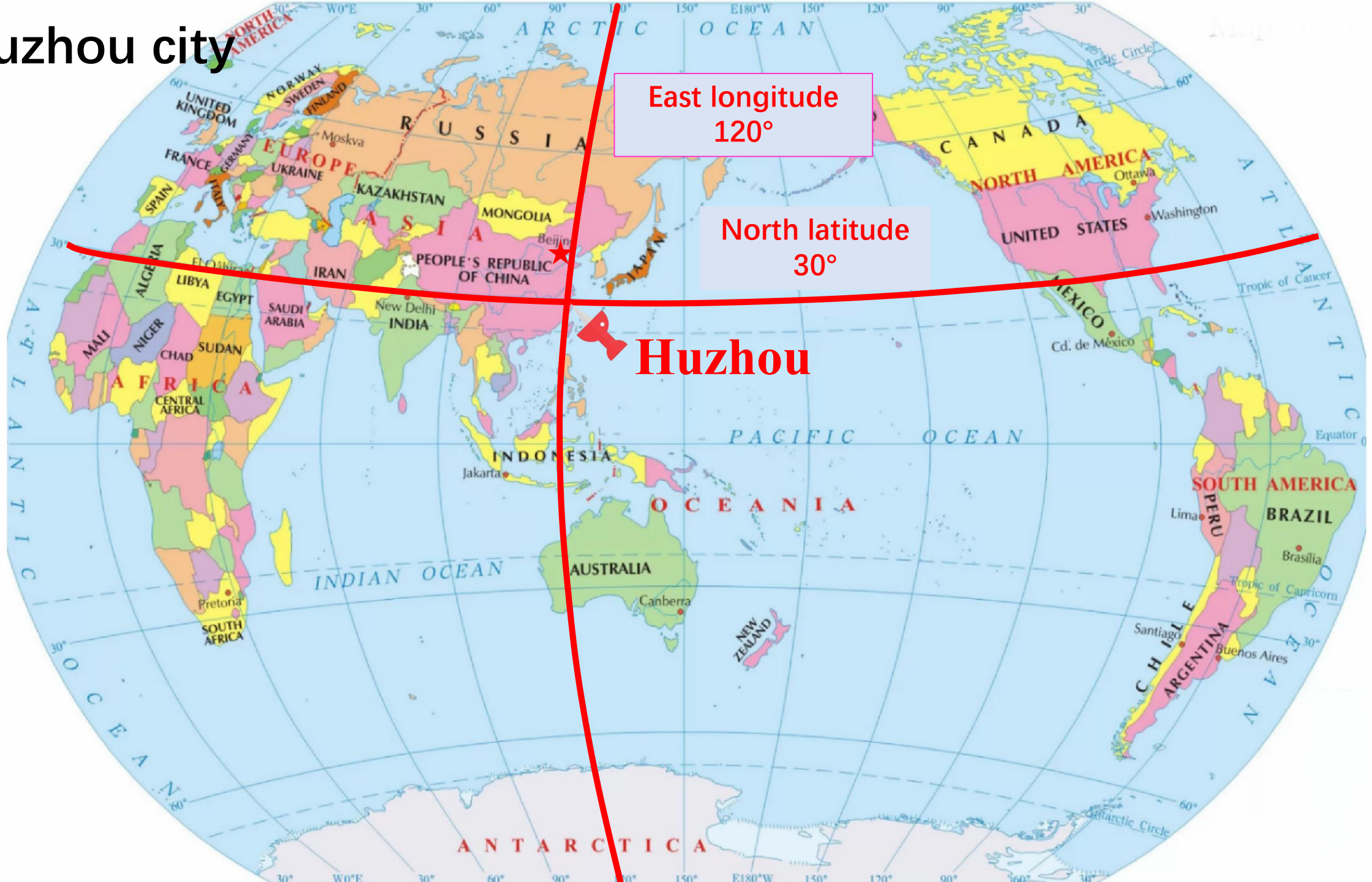
Yongjia Wang (Huzhou University, China)

Huzhou city

East longitude
120°

North latitude
30°

Huzhou





Hunter (Huzhou nuclear theory group)



于少英



沈彩万



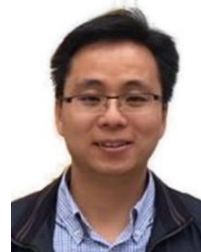
李庆峰



王建松



戴连荣



朱祥荣



孙宇梁



李可



宋红芳



李自立

22 faculties, including 8 Professors

Particle physics

Heavy-ion physics

Nuclear structure

Experimental nuclear physics



刘艳鑫



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董国香



穆成富



李柏青



赵卿



徐浩浩



王永佳



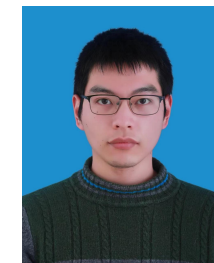
孙晓慧



张肇文



李鹏程



顾安

Outline

➤ 1. Background

➤ 2. Elliptic flow

- ✓ 2.1 The role of nuclear potential and collision term
- ✓ 2.2 The density interval at which the nuclear equation of state is probed by the elliptic flow
- ✓ 2.3 Effects of rapidity and transverse momentum cuts

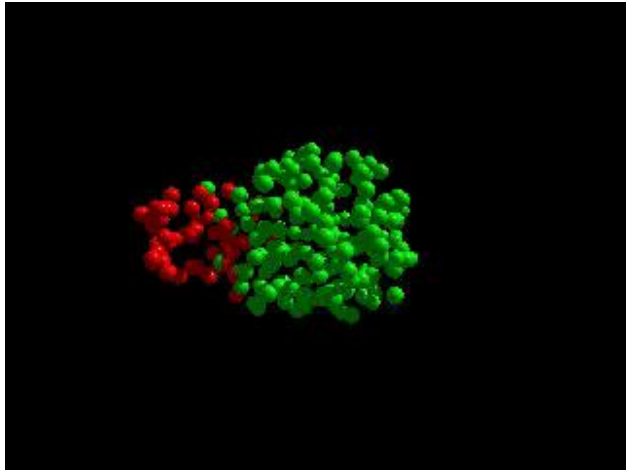
➤ 3. Bayesian analysis with FOPI data

➤ 4. Summary

01

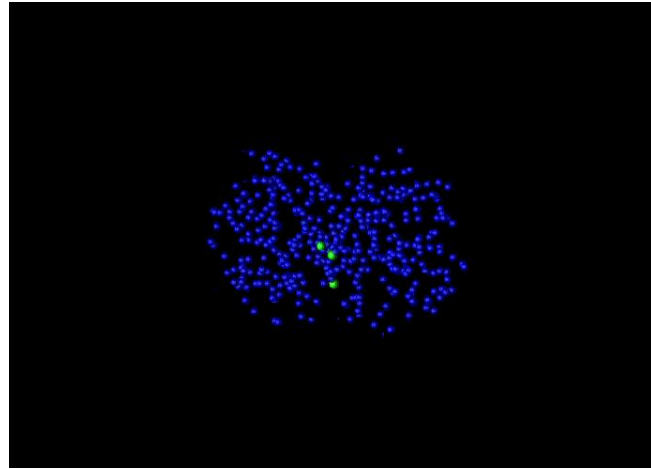
Overview of HIC

◆ intermediate energies HIC

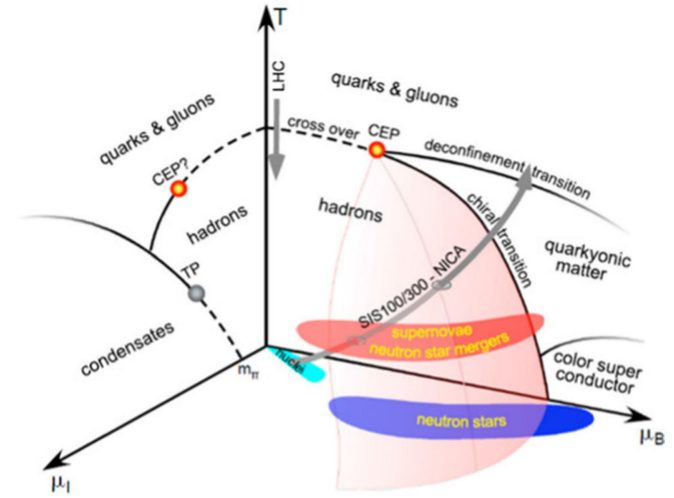


◆ Low energy fusion reaction

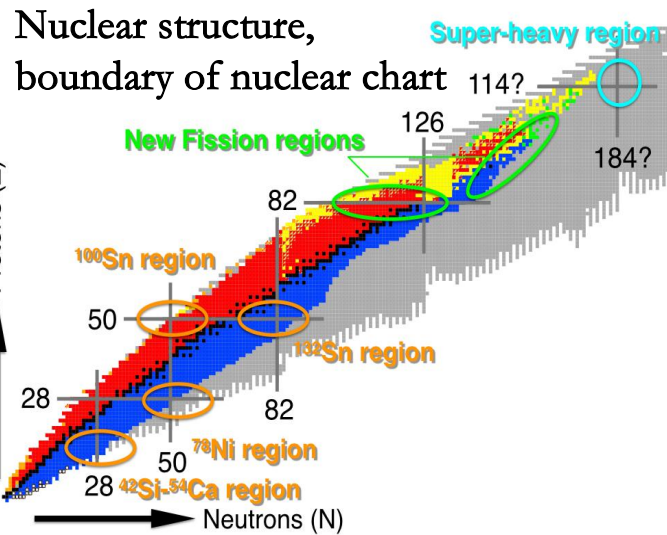
Nuclear equation of state



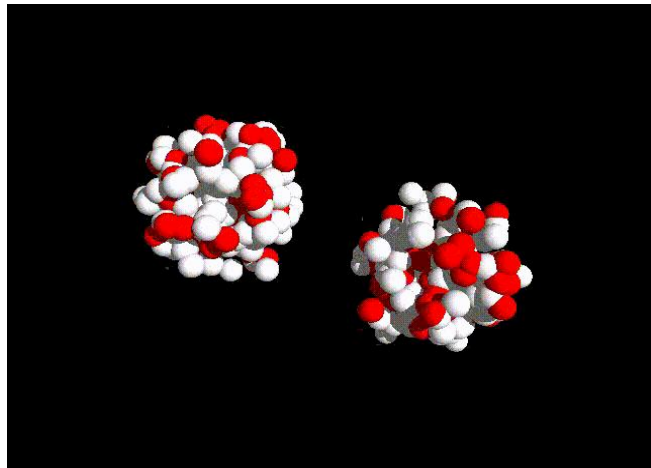
~100 MeV- a few GeV/nucleon



QCD phase diagram, QGP properties

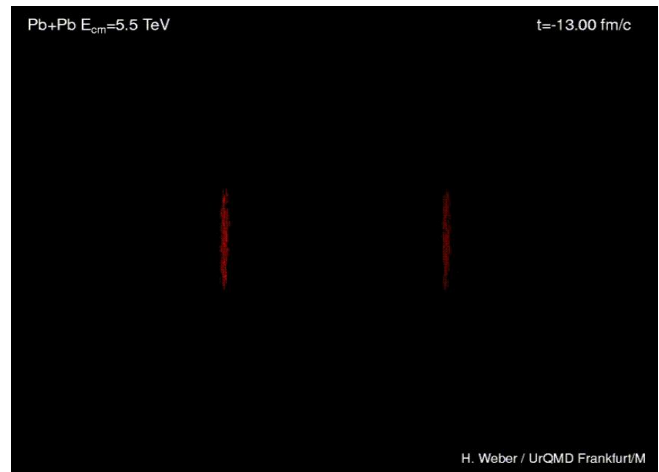


~100 MeV/nucleon



> hundreds of GeV/nucleon

beam energy



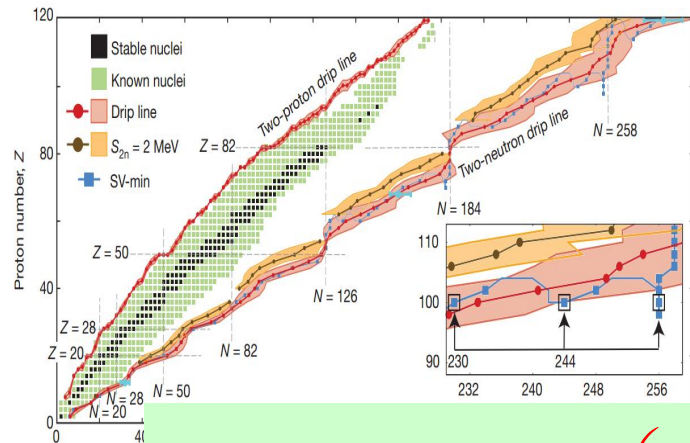
◆ Relativistic HIC

Nuclear equation of state (EOS)

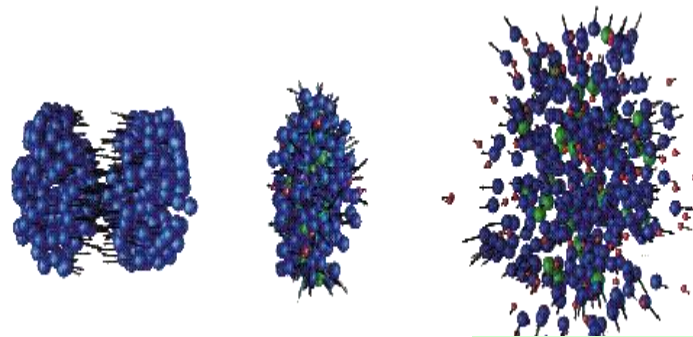
The thermodynamic relationship between the binding energy E (or pressure P) and density ρ , as well as the isospin asymmetry δ .

$$E(\rho, \delta) = E(\rho, 0) + E_{sym}(\rho) \delta^2 + \dots, \quad \delta = \frac{\rho_n - \rho_p}{\rho_n + \rho_p}$$

Nuclear landscape



Heavy ion collision



Neutron star



$$E(\rho, 0) = E_0 + \frac{K_0}{2} \left(\frac{\rho - \rho_0}{3\rho_0} \right)^2 + \dots,$$

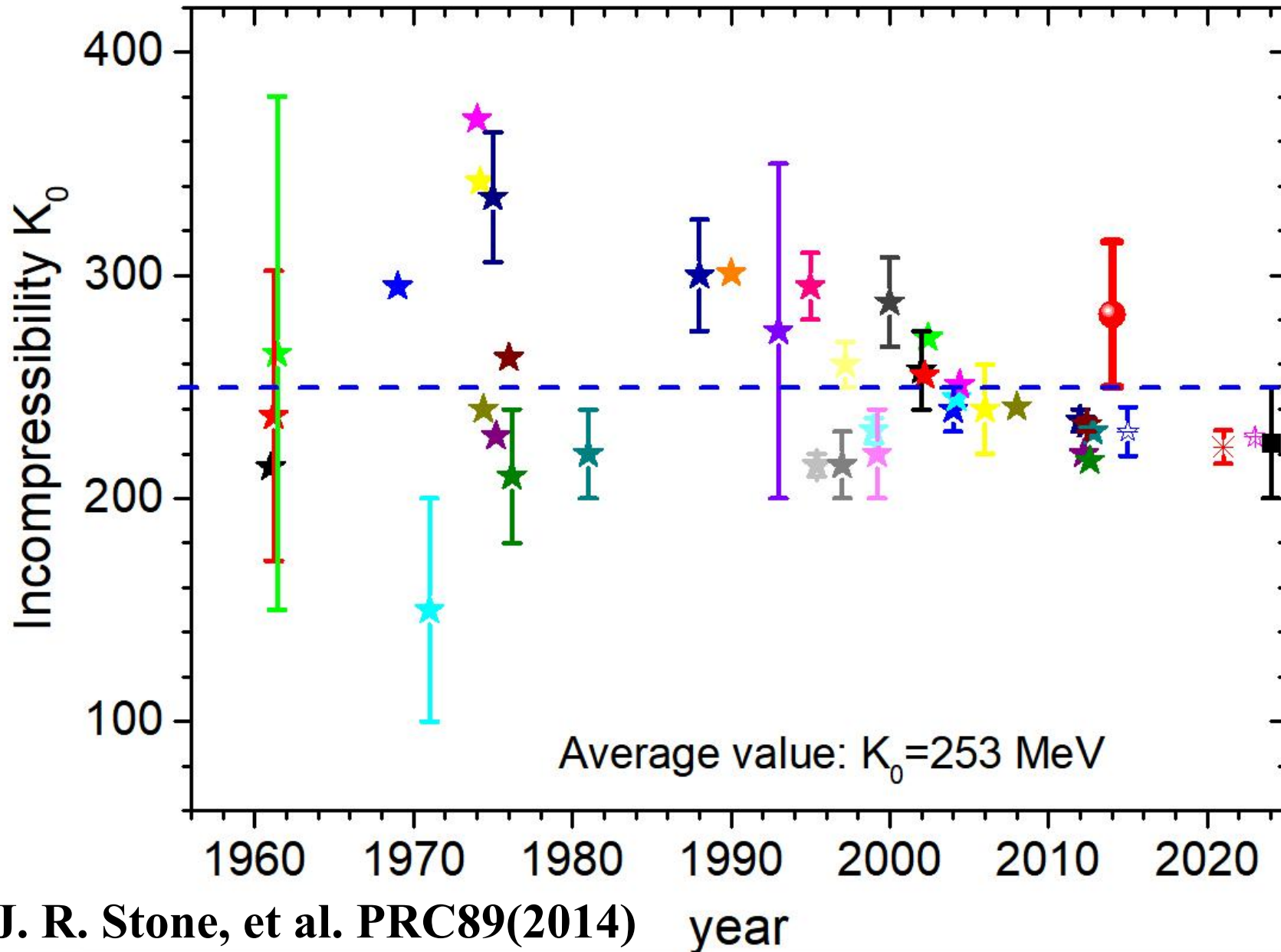
$$K_0 = 9\rho^2 \left(\frac{\partial^2 E}{\partial \rho^2} \right) \Big|_{\rho=\rho_0}$$

$$E_{sym}(\rho) = S_0 + L \left(\frac{\rho - \rho_0}{3\rho_0} \right) + \frac{K_{sym}}{2} \left(\frac{\rho - \rho_0}{3\rho_0} \right)^2 + \dots$$

$$L = 3\rho \frac{dE_{sym}(\rho)}{d\rho} \Big|_{\rho=\rho_0}$$

K_0 and L determine the EOS in the vicinity of the saturation density.

The incompressibility K_0 from properties of nuclei



Data: isoscalar giant monopole resonances (GMR), nuclear masses, nuclear decay...

Model: Skyrme-Hartree-Fock model, Skyrme-Hartree-Fock-Bogoliubov, ...

Ning Wang, et al. PLB(2015)

Jun Xu, et al. PRC(2021)

Zhenzhen Li, et al. PRL(2023)

Xiao Liu, et al. PRC(2024)

The incompressibility K_0 from heavy-ion physics

VOLUME 55, NUMBER 24

PHYSICAL REVIEW LETTERS

9 DECEMBER 1985

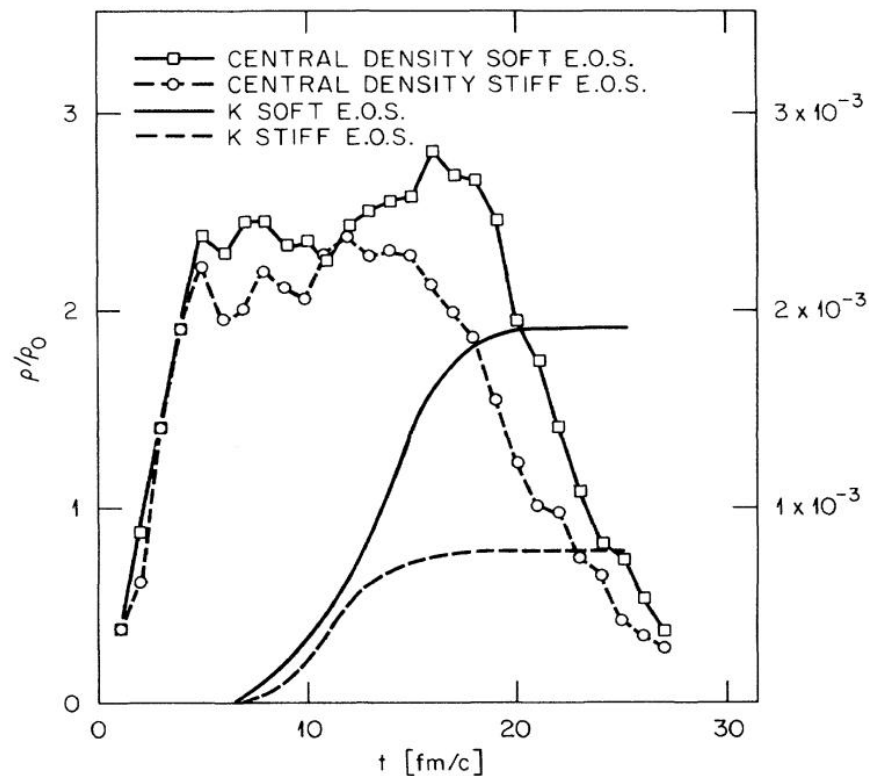
Data: pion and kaon production, collective flow

Subthreshold Kaon Production as a Probe of the Nuclear Equation of State

J. Aichelin and Che Ming Ko^(a)

Joint Institute for Heavy Ion Research, Holifield Heavy Ion Research Facility, Oak Ridge, Tennessee 37831
(Received 11 June 1985; revised manuscript received 23 September 1985)

The production of kaons at subthreshold energies from heavy-ion collisions is sensitive to the nuclear equation of state. In the Boltzmann-Uehling-Uhlenbeck model, the number of produced kaons from central collisions between heavy nuclei at incident energies around 700 MeV/nucleon can vary by a factor of ~ 3 , depending on the equation of state.



Model: QMD and BUU

Further evidence for a stiff nuclear equation of state from a transverse-momentum analysis of Ar(1800 MeV/nucleon) + KCl

[Joseph J. Molitoris](#) and [Horst Stöcker](#)

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Phys. Rev. C **32**, 346(R) – Published 1 July, 1985

Evidence for a Soft Nuclear Equation-of-State from Kaon Production in Heavy-Ion Collisions

[C. Sturm](#)¹, [I. Böttcher](#)⁴, [M. Dębowski](#)⁵, [A. Förster](#)¹, [E. Grosse](#)^{6,7}, [P. Koczoń](#)², [B. Kollmeier](#)⁴, [F. Laue](#)², and [M. Mang](#)^{2 et al.}

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Phys. Rev. Lett. **86**, 39 – Published 1 January, 2001

Hadronic Matter Is Soft

[Ch. Hartnack](#)¹, [H. Oeschler](#)², and [Jörg Aichelin](#)¹

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Phys. Rev. Lett. **96**, 012302 – Published 9 January, 2006

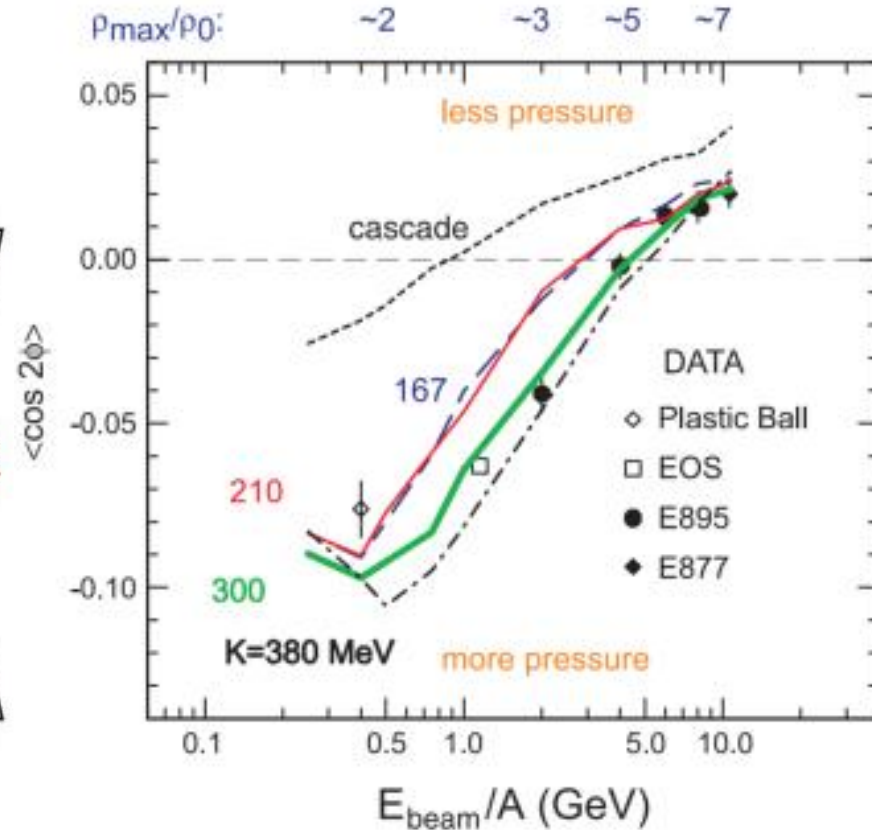
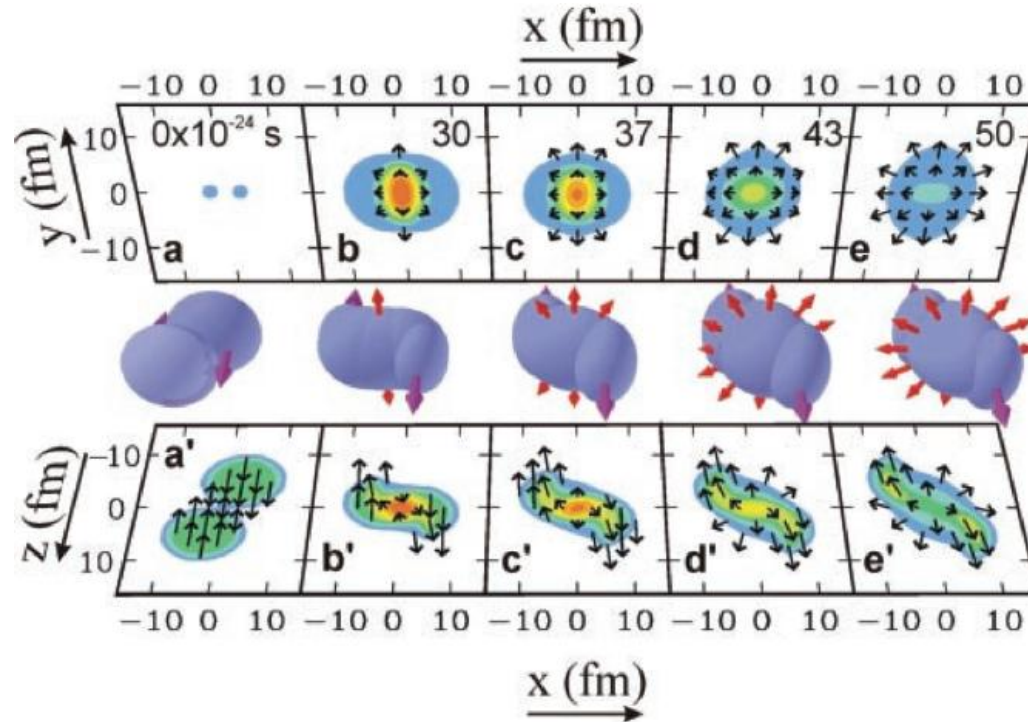
Stiff: $K_0 > 380$ MeV

Soft: $K_0 < 200$ MeV

HIC offers a unique way to create nuclear matter with high density and isospin asymmetry in laboratory.



Determination of the Equation of State of Dense Matter
 Pawel Danielewicz, *et al.*
Science **298**, 1592 (2002);
 DOI: 10.1126/science.1078070



EOS can be deduced from the comparison between experimental observables and transport model calculations.

Nuclear equation of state

nature

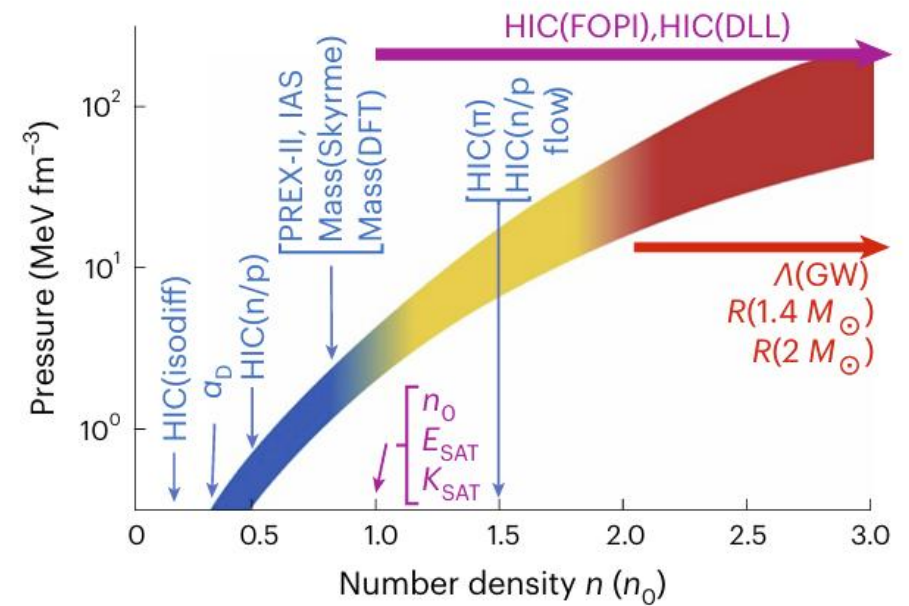
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Constraining neutron-star matter with microscopic and macroscopic collisions

Sabrina Huth, Peter T. H. Pang, Ingo Tews, Tim Dietrich, Arnaud Le Fèvre, Achim Schwenk, Wc Trautmann, Kshitij Agarwal, Mattia Bulla, Michael W. Coughlin & Chris Van Den Broeck



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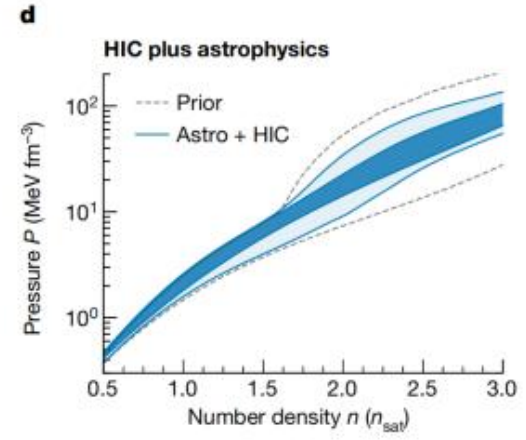
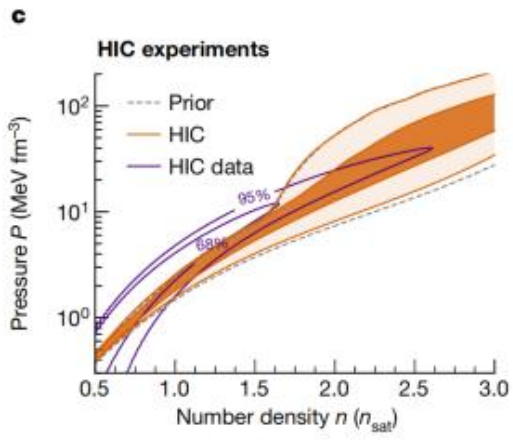
nature > nature astronomy > articles > article

Article | Published: 05 January 2024

Determination of the equation of state from nuclear experiments and neutron star observations

Chun Yuen Tsang, ManYee Betty Tsang, William G. Lynch, Rohit Kumar & Charles J. Horowitz

Nature Astronomy 8, 328–336 (2024) | Cite this article





Progress in Particle and Nuclear Physics

Available online 19 September 2023, 104080

In Press, Journal Pre-proof [?](#) What's this? [?](#)



Review

Dense nuclear matter equation of state from heavy-ion collisions

[Agnieszka Sorensen](#)¹ , [Kshitij Agarwal](#)², [Kyle W. Brown](#)^{3,4}, [Zbigniew Chajecski](#)⁵, [Paweł Danielewicz](#)^{3,6}, [Christian Drischler](#)⁷, [Stefano Gandolfi](#)⁸, [Jeremy W. Holt](#)^{9,10}, [Matthias Kaminski](#)¹¹, [Che-Ming Ko](#)^{9,10}, [Rohit Kumar](#)³, [Bao-An Li](#)¹², [William G. Lynch](#)^{3,6}, [Alan B. McIntosh](#)¹⁰, [William G. Newton](#)¹², [Scott Pratt](#)^{3,6}, [Oleh Savchuk](#)^{3,13}, [Maria Stefaniak](#)^{14,15}, [Ingo Tews](#)⁸, [ManYee Betty Tsang](#)^{3,6}...[Yi Yin](#)^{9,4}

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Abstract

This [White Paper](#) highlights the essential role of hadronic transport simulations of heavy-ion collisions in studies involving the equation of state of nuclear matter. It also elucidates many connections between inferences of the equation of state from heavy-ion collision data and other efforts aiming to understand the properties of nuclear matter.

arXiv > nucl-th > arXiv:2211.02224

Nuclear Theory

[Submitted on 4 Nov 2022 (v1), last revised 8 Nov 2022 (this version, v2)]

Long Range Plan: Dense matter theory for heavy-ion collisions and neutron stars

[Alessandro Lovato](#), [Travis Dore](#), [Robert D. Pisarski](#), [Bjoern Schenke](#), [Katerina Chatziioannou](#), [Jocelyn S. Read](#), [Philippe Landry](#), [Pav Hannah Elfner](#), [Veronica Dexheimer](#), [Rajesh Kumar](#), [Michael Strickland](#), [Johannes Jahan](#), [Claudia Ratti](#), [Volodymyr Vovchenko](#), [Mikh Hippert](#), [Jacquelyn Noronha-Hostler](#), [Jorge Noronha](#), [Enrico Speranza](#), [Nicolas Yunes](#), [Chuck J. Horowitz](#), [Steven P. Harris](#), [Larry McStefano Gandolfi](#), [Ingo Tews](#), [M. Coleman Miller](#), [Cecilia Chirenti](#), [Zohreh Davoudi](#), [Jamie M. Karthein](#), [Krishna Rajagopal](#), [Salvatore Vladimír Skokov](#), [Ulrich Heinz](#), [Christian Drischler](#), [Daniel R. Phillips](#), [Madappa Prakash](#), [Zoltan Fodor](#), [David Radice](#), [Christopher Plu Fraga](#), [Aleksi Kurkela](#), [James M. Lattimer](#), [Andrew W. Steiner](#), [Jeremy W. Holt](#), [Bao-An Li](#), [Chun Shen](#), [Mark Alford](#), [Alexander Haber](#),

Since the release of the 2015 Long Range Plan in Nuclear Physics, major events have occurred that reshaped our understanding of quantum chromodynamics of equilibrium. The US nuclear community has an opportunity to capitalize on advances in astrophysical observations and nuclear experiments and engage matter that connects low- and high-energy nuclear physics, astrophysics, gravitational waves physics, and data science

Comments: 70 pages, 3 figures, [White Paper for the Long Range Plan for Nuclear Science](#)

Subjects: **Nuclear Theory (nucl-th)**; High Energy Astrophysical Phenomena (astro-ph.HE); High Energy Physics - Phenomenology (hep-ph)

Report number: LA-UR-22-31648

Cite as: arXiv:2211.02224 [nucl-th]

(or arXiv:2211.02224v2 [nucl-th] for this version)

<https://doi.org/10.48550/arXiv.2211.02224>

Submission history

From: [Jacquelyn Noronha-Hostler](#) [[view email](#)]

[v1] Fri, 4 Nov 2022 02:15:29 UTC (2,372 KB)

[v2] Tue, 8 Nov 2022 01:52:13 UTC (2,373 KB)

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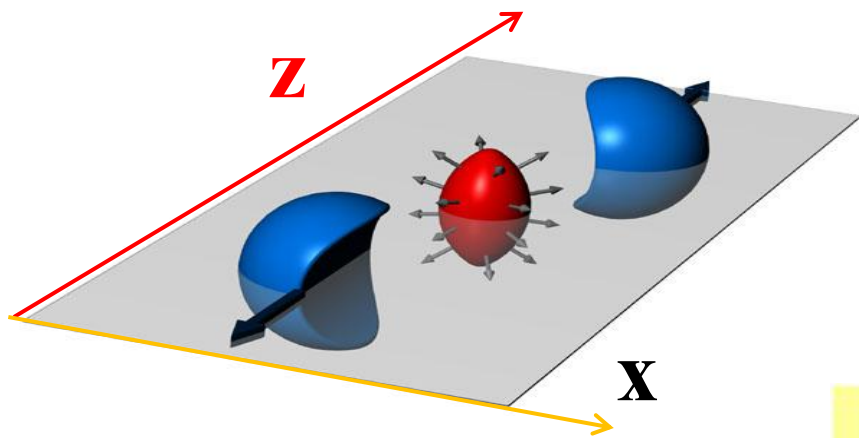
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01

Elliptic flow (v_2)

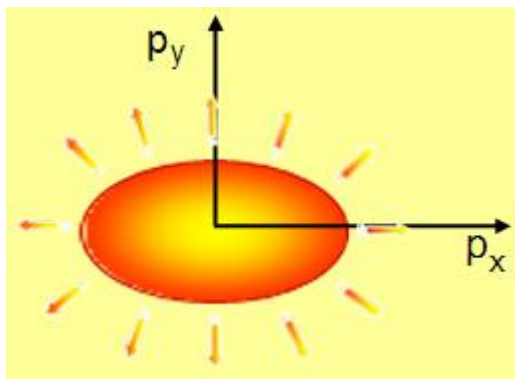
$$v_2 = \left\langle \frac{p_x^2 - p_y^2}{p_x^2 + p_y^2} \right\rangle \quad p_t = \sqrt{p_x^2 + p_y^2}; y = \frac{1}{2} \ln \left[\frac{E + p_z}{E - p_z} \right]$$

Coordinate-space anisotropy

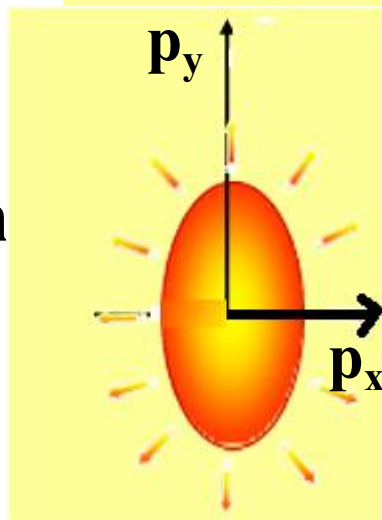


v_2 is a multivariate function of rapidity (y_z) and transverse momentum (p_t)

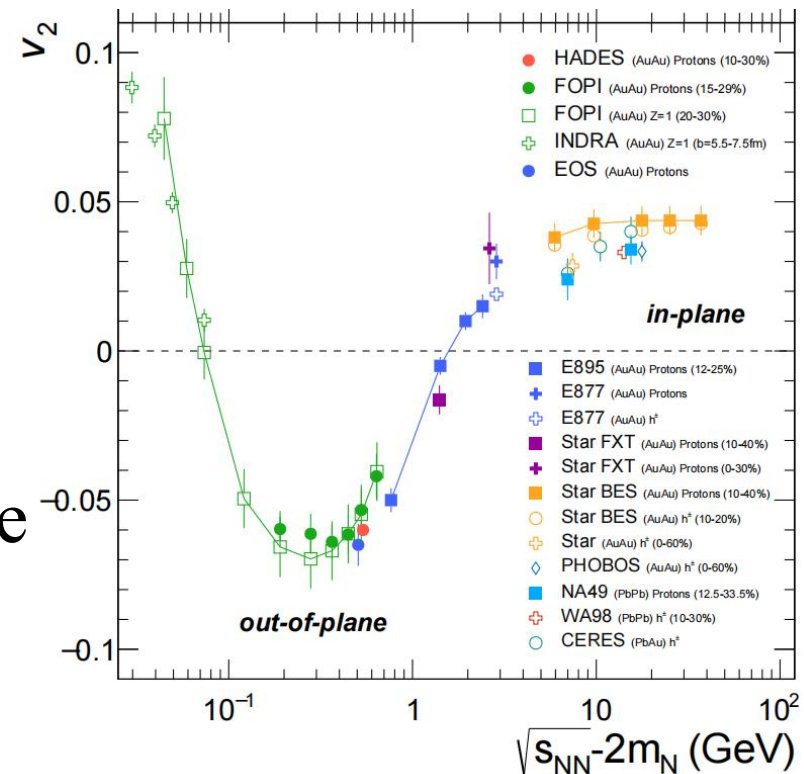
Momentum-space anisotropy



In-plane
 $v_2 > 0$



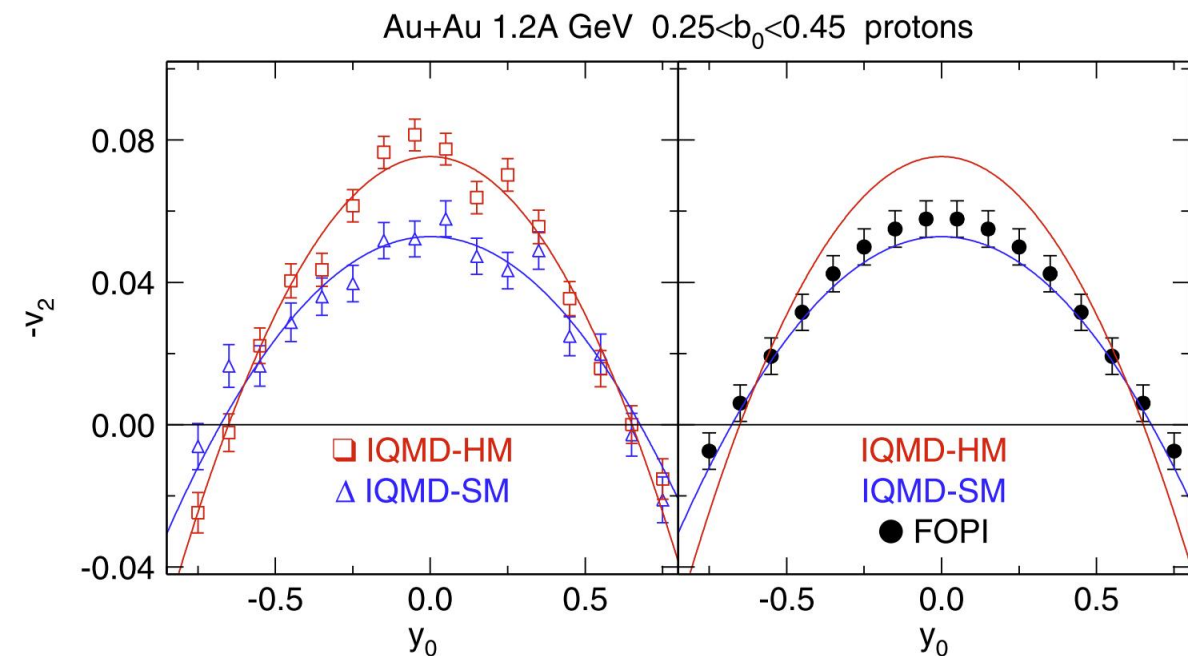
Out-of-plane
Squeeze out
 $v_2 < 0$



v_2 has been widely used to probe the nuclear EOS, the nucleon effective mass,...

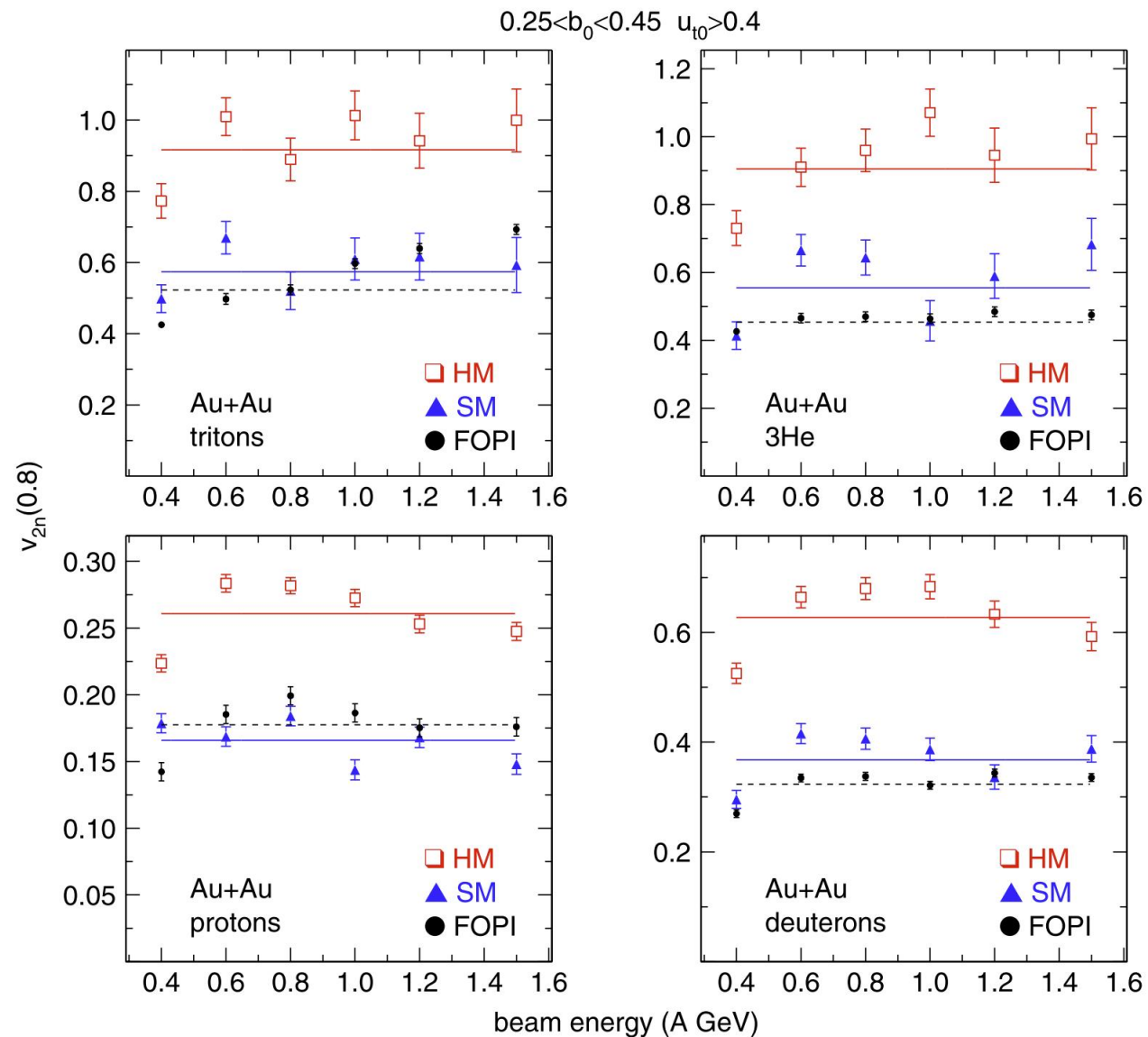
Elliptic flow (v_2)

Le Fevre, A., Leifels, Y., Reisdorf, W., Aichelin, J., & Hartnack, C. (2016). Constraining the nuclear matter equation of state around twice saturation density. *Nuclear Physics A*, 945, 112-133.



$$v_2 = v_{20} + v_{22} \cdot y_0^2.$$

$$v_{2n} = |v_{20}| + |v_{22}|$$



01

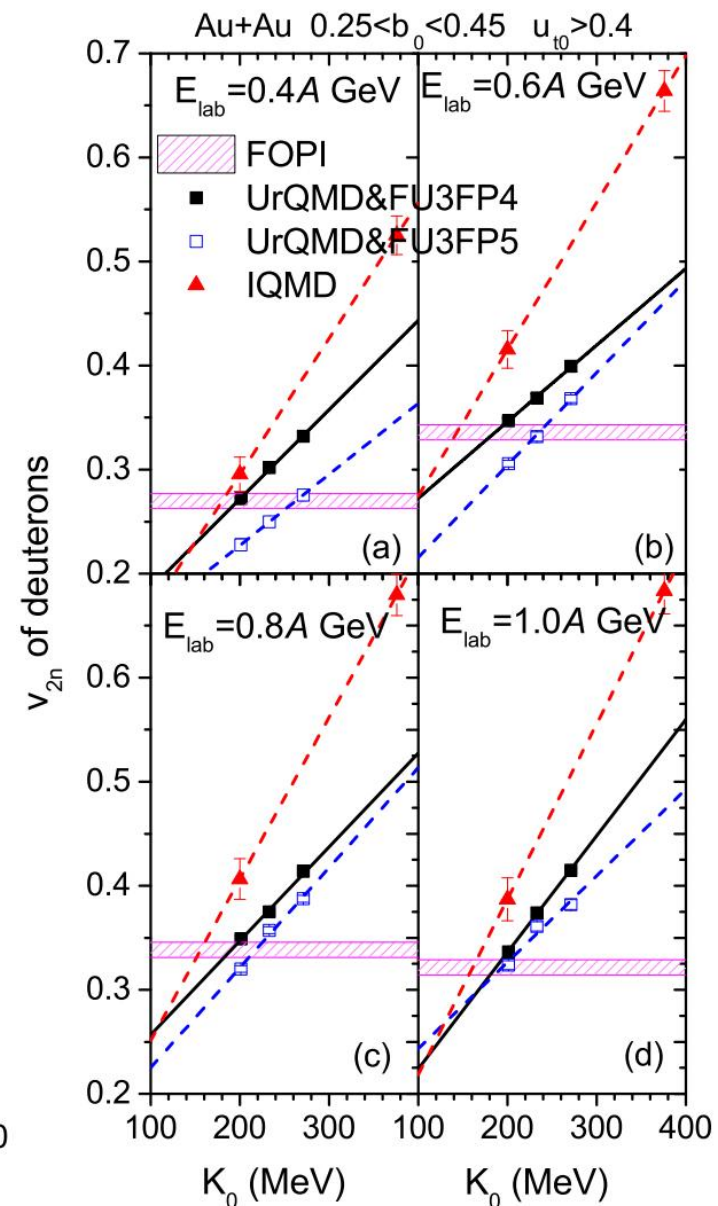
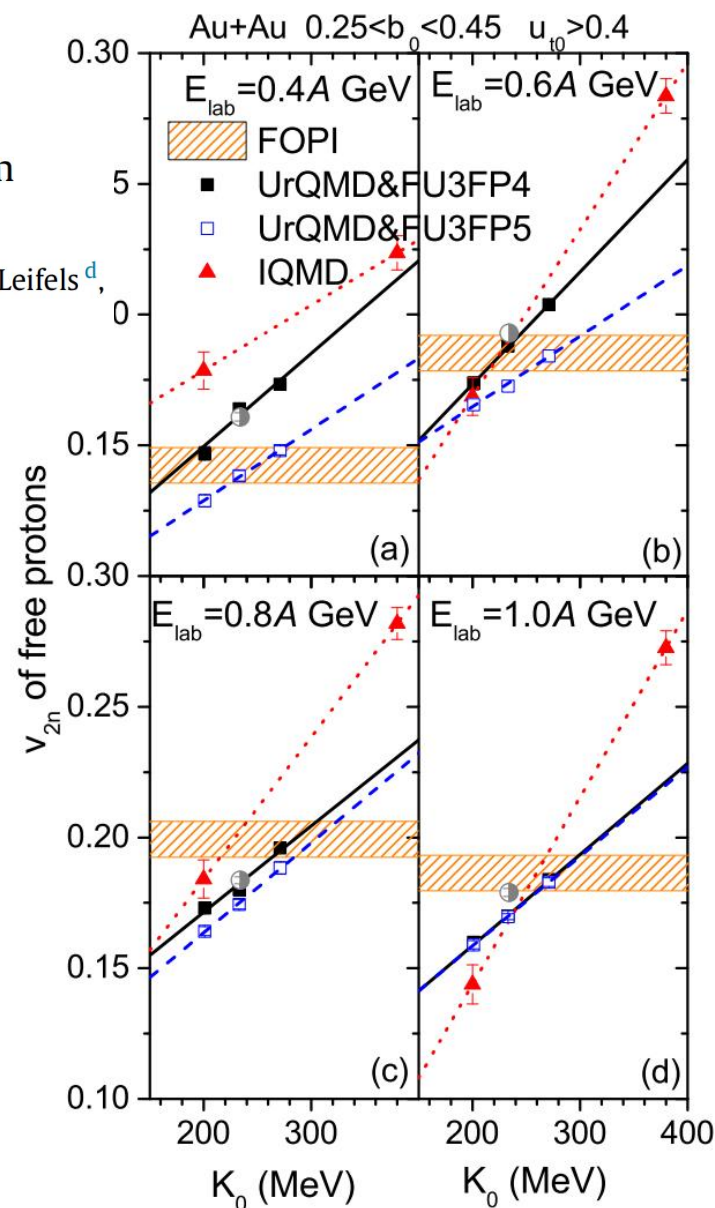
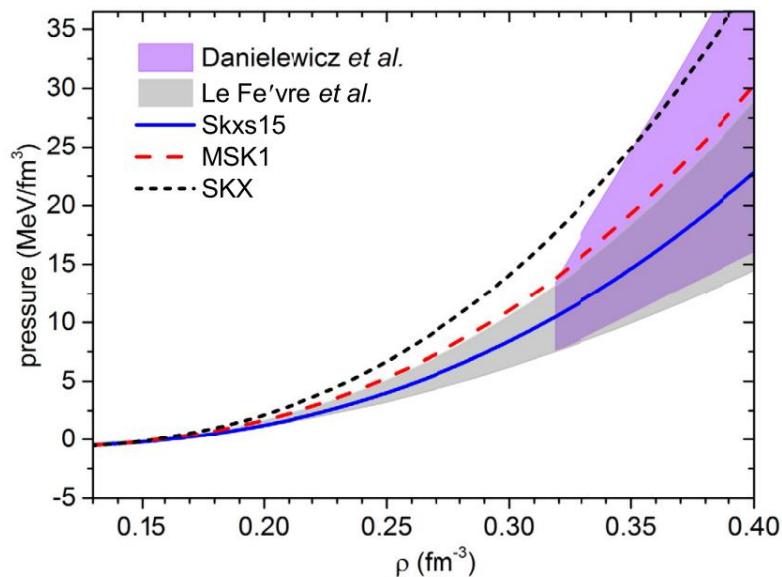
Elliptic flow (v_2)

Determination of the nuclear incompressibility from the rapidity-dependent elliptic flow in heavy-ion collisions at beam energies 0.4A–1.0A GeV

Yongjia Wang^{a,*}, Chenchen Guo^b, Qingfeng Li^{a,c,**}, Arnaud Le Fèvre^d, Yvonne Leifels^d, Wolfgang Trautmann^d

Physics Letters B 778 (2018) 207–212.

	K_0 (MeV)	S_0 (MeV)	L (MeV)
Skxs15	201	31.88	34.79
MSK1	234	30.00	33.92
SKX	271	31.10	33.18
SV-sym34	234	34.00	80.95



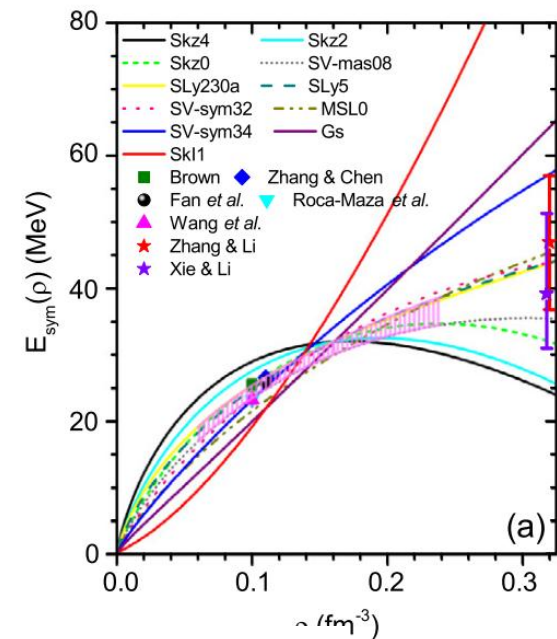
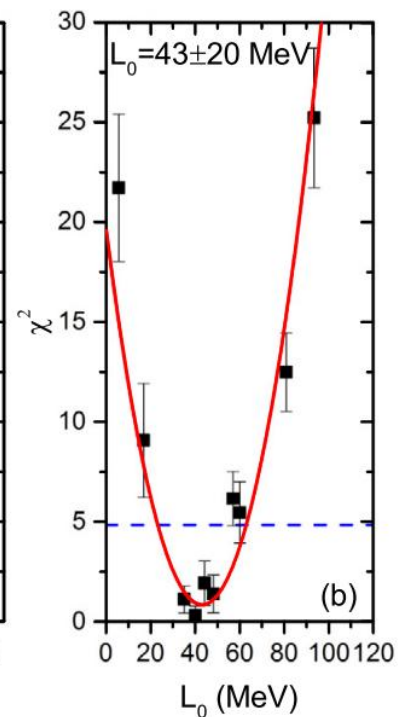
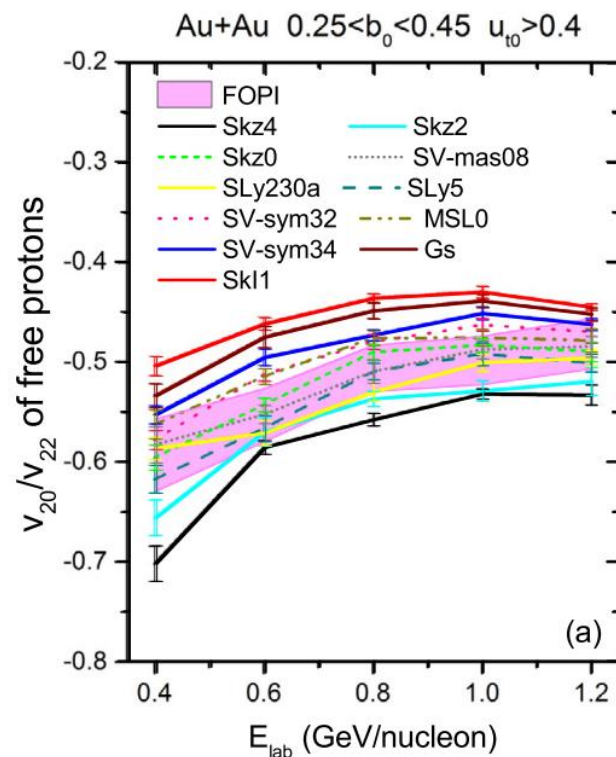
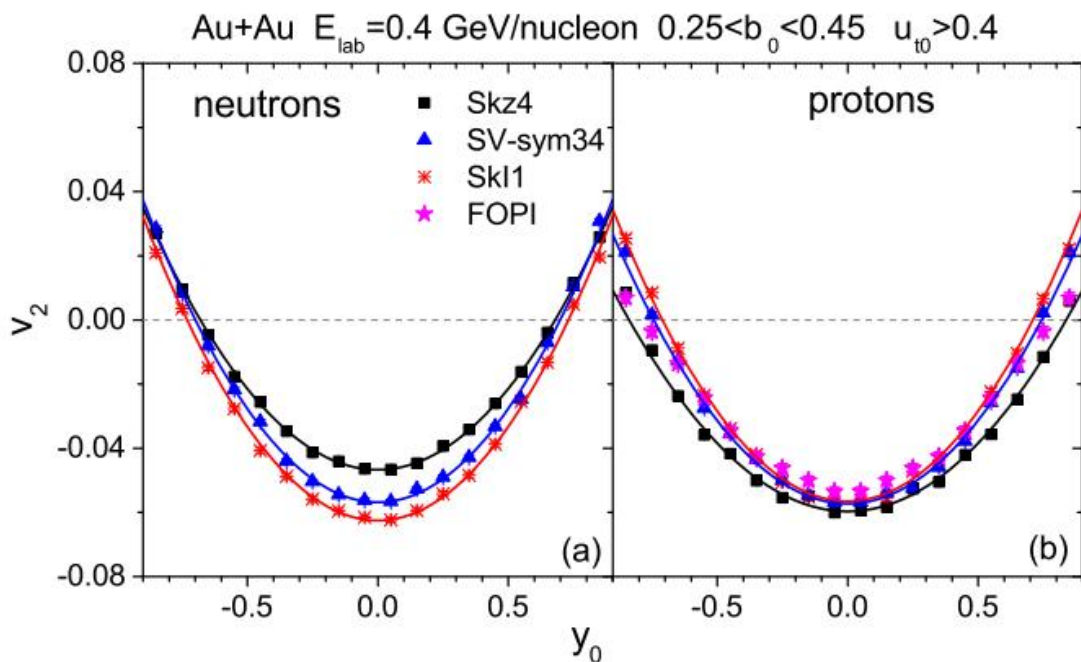
01

Elliptic flow (v_2)

Study of the nuclear symmetry energy from the rapidity-dependent elliptic flow in heavy-ion collisions around 1 GeV/nucleon regime

Yongjia Wang^a, Qingfeng Li^{a,b,*}, Yvonne Leifels^c, Arnaud Le Fèvre^c

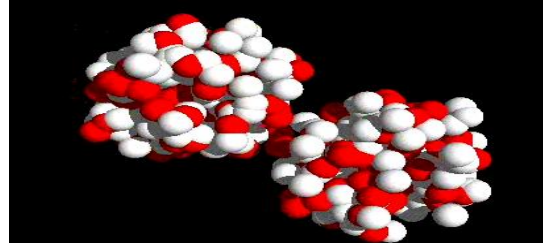
Physics Letters B 802 (2020) 135249



Ultrarelativistic Quantum Molecular Dynamics (UrQMD) Model

1). Initialization

Get the coordinate \mathbf{r}
and the momentum \mathbf{p}



$$\phi_i(\vec{r}_i; t) = \frac{1}{(2\pi)^{3/4}(\Delta x)^{3/2}} \exp \left\{ -\frac{[\vec{r}_i - \vec{R}_i(t)]^2}{(2\Delta x)^2} + i\vec{r}_i \cdot \vec{P}_i(t) \right\}.$$

2). Propagation

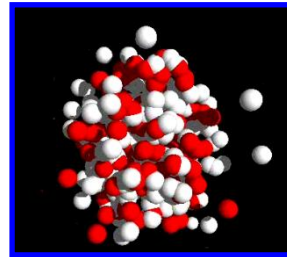
Nucleon moves in nuclear potential.
Density, momentum, isospin-dependent.

$$\dot{\mathbf{p}}_i = -\frac{\partial \langle H \rangle}{\partial \mathbf{r}_i} \quad \text{and} \quad \dot{\mathbf{r}}_i = \frac{\partial \langle H \rangle}{\partial \mathbf{p}_i},$$

Input nuclear potential
EOS, Symmetry energy

3). Collision term

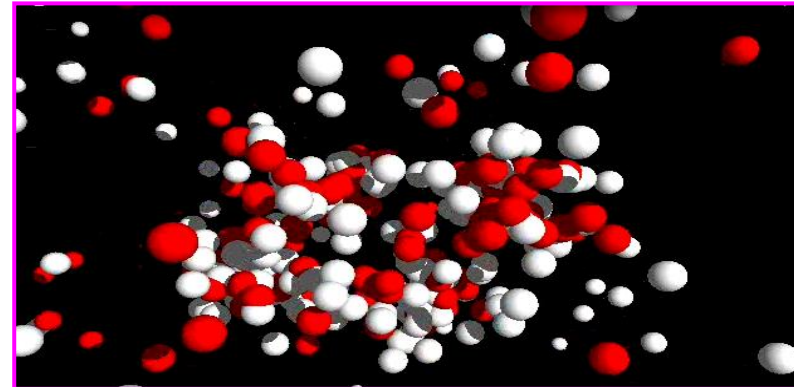
Medium modified cross section. Density,
momentum, isospin-dependent. Pauli
blocking.



Input: $F * \sigma_{NN}^{free}$

4). Cluster recognition

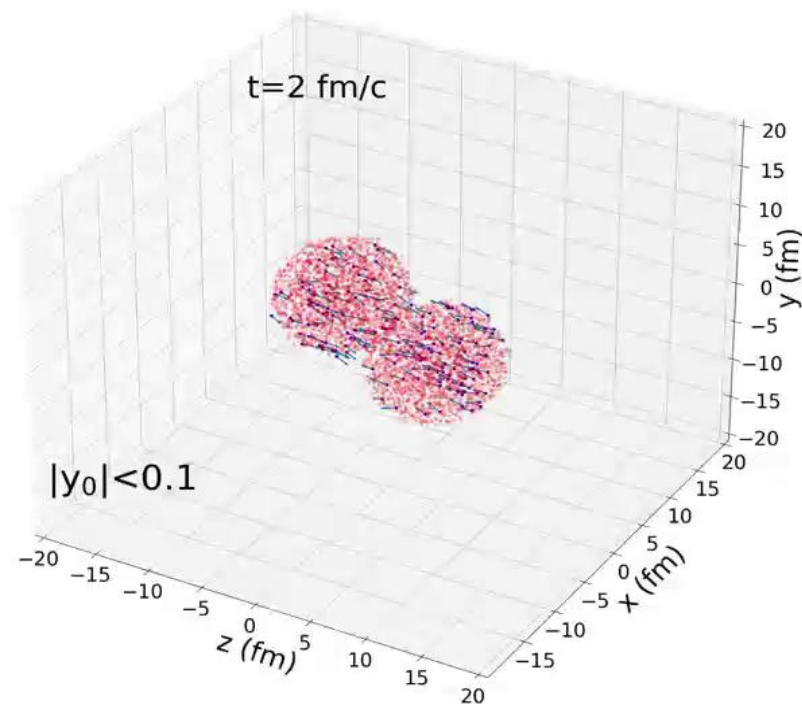
an isospin-dependent Minimum Spanning Tree



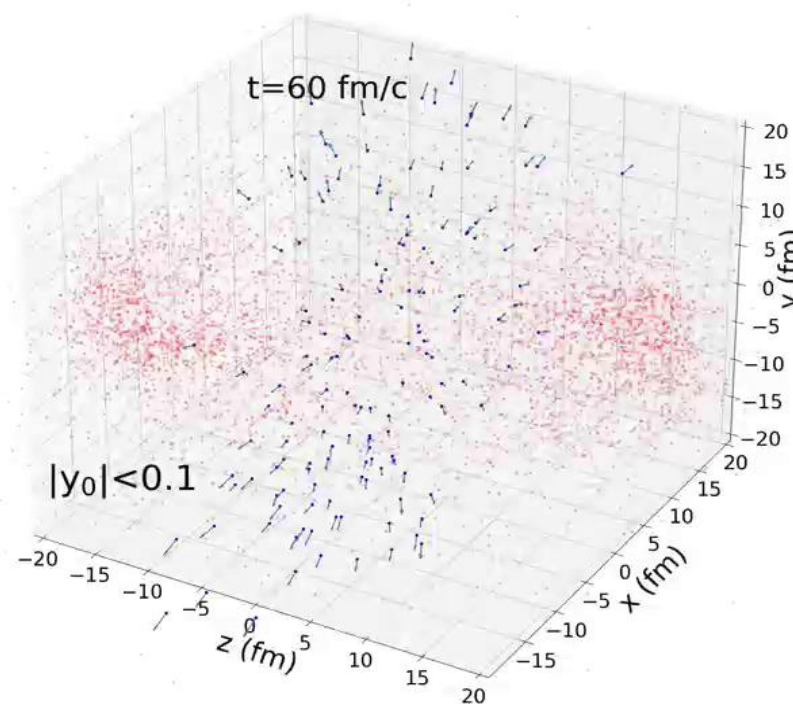
Elliptic flow

The momentum of each nucleon can be changed either by the force caused by the nuclear mean field (MF) potential or by nucleon-nucleon (NN) scattering.

Au+Au $E_{lab}=0.4A$ GeV $b = 6$ fm



Au+Au $E_{lab}=0.4A$ GeV $b = 6$ fm



$$v_2 = \langle V_2 \rangle = \left\langle \frac{p_x^2 - p_y^2}{p_t^2} \right\rangle.$$

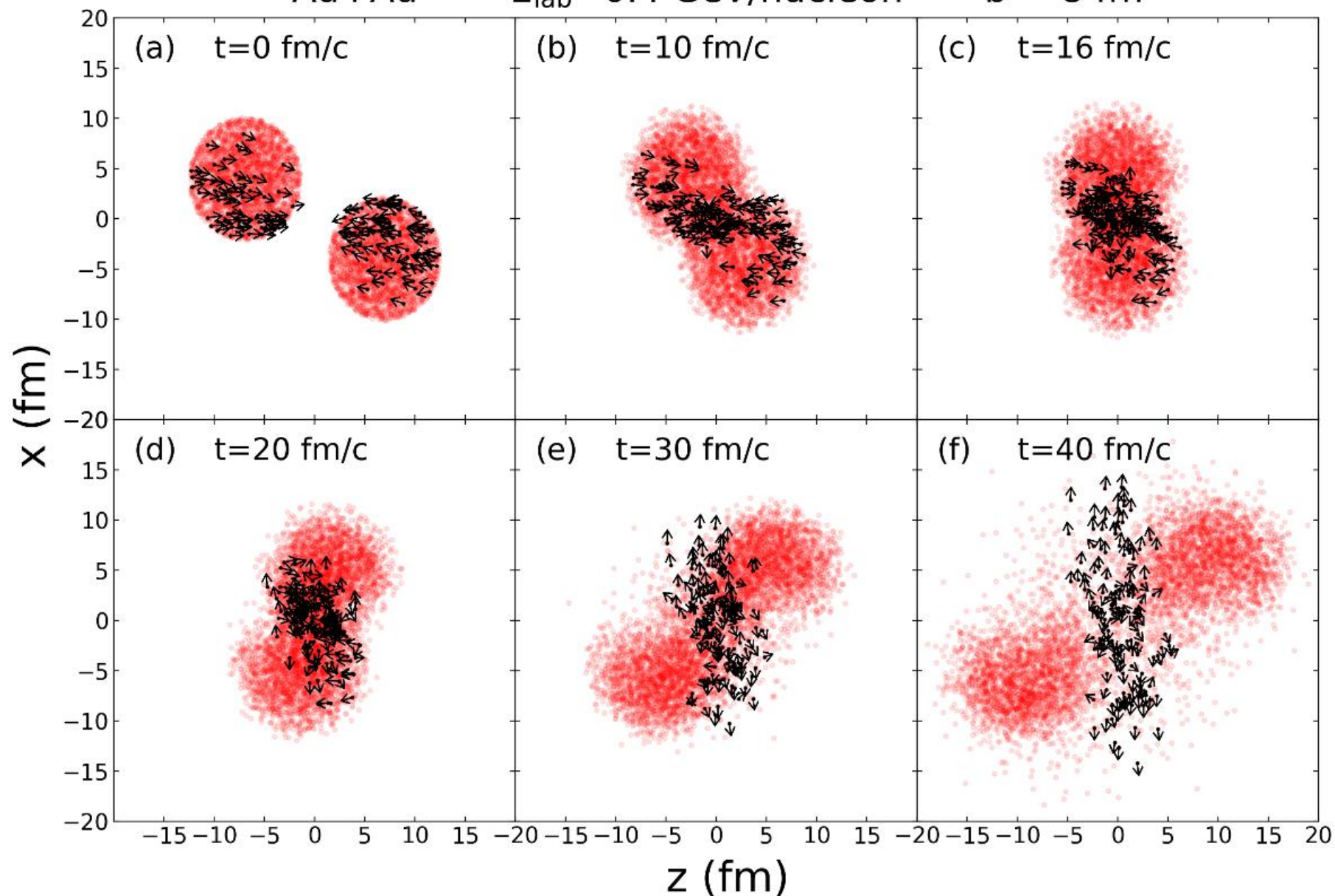
$$\Delta v_2^{\text{coll}}(t) = \left\langle V_2^{\text{aft.coll}}(t) - V_2^{\text{bef.coll}}(t) \right\rangle.$$

$$\Delta v_2^{\text{mf}}(t) = \left\langle V_2^{\text{aft.mf}}(t + \Delta t) - V_2^{\text{bef.mf}}(t) \right\rangle.$$

Reversely tracing nucleons that are finally emitted at mid-rapidity ($-0.1 < y_0 < 0.1$).
Recording the momentum changes caused by MF and NN collisions.

Elliptic flow

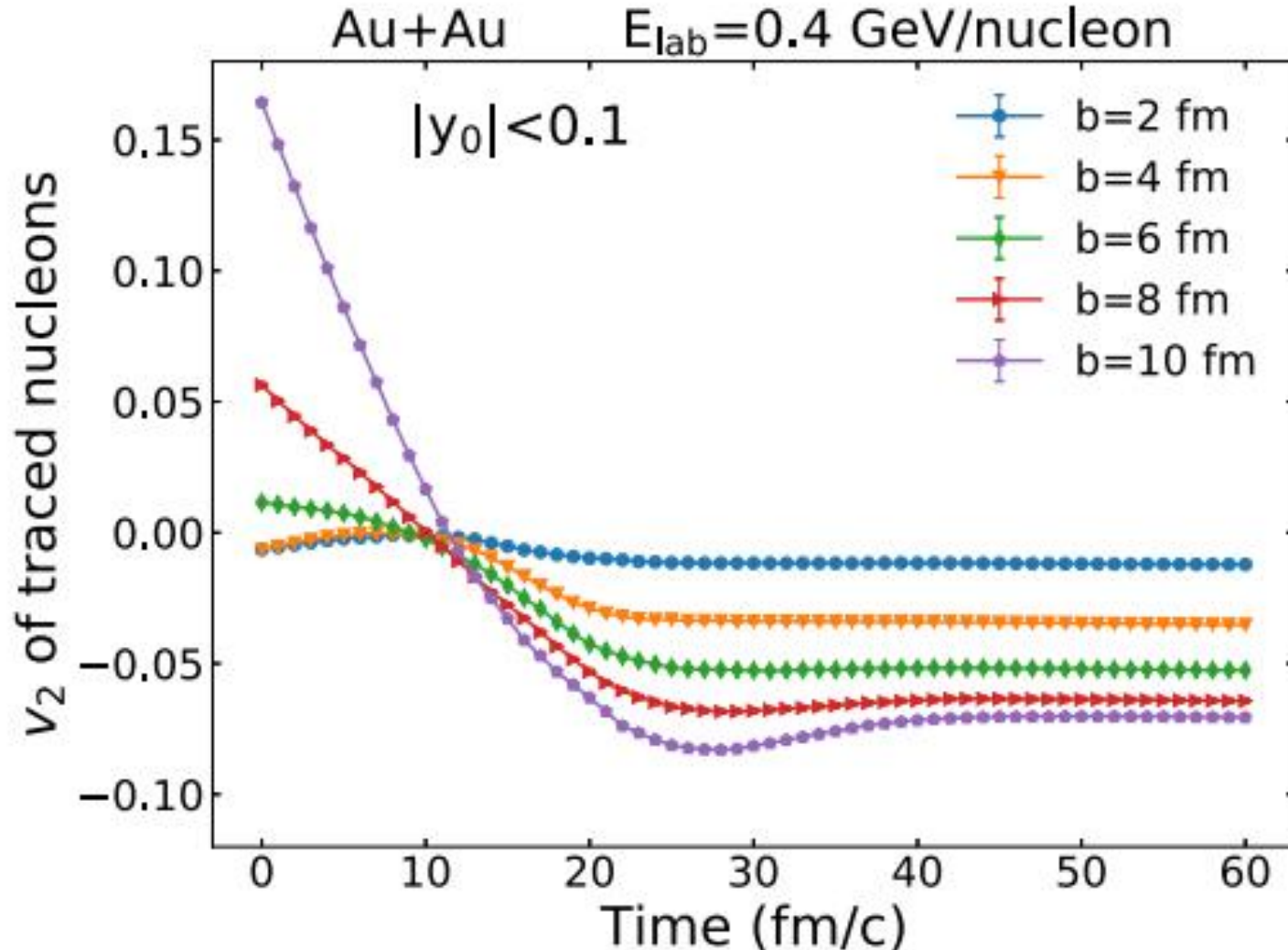
Au+Au

 $E_{\text{lab}} = 0.4 \text{ GeV/nucleon}$
 $b = 8 \text{ fm}$


Solid dots with arrows represent the traced nucleons (those will be appeared at mid-rapidity at the final stage), while circles denote other nucleons.

Elliptic flow

Time evolution of v_2 of the **traced nucleons finally emitted at mid-rapidity** ($|y_0| < 0.1$) with different impact parameters.



At the initial time, v_2 is non-zero, indicating that those nucleons **are not randomly “selected”** out of the uniform distribution in the initial state. This is especially so for large b where $v_2(t=0 \text{ fm/c})$ is positive because nucleons with larger $|p_x|$ and smaller $|p_y|$ tend to have higher probability to experience NN scattering and finally emerge at midrapidity.

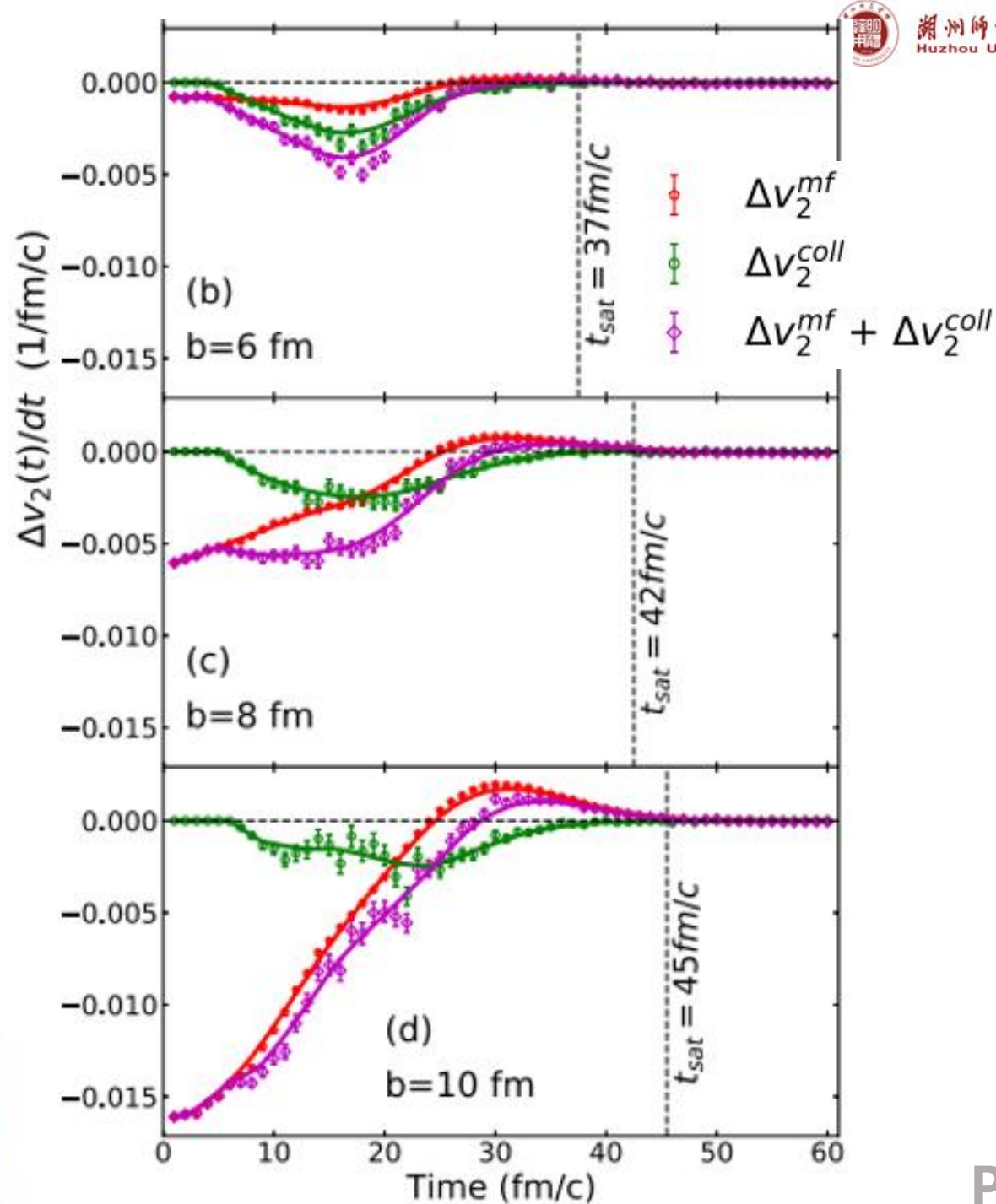
Machine learning

$$v_2 = \langle V_2 \rangle = \left\langle \frac{p_x^2 - p_y^2}{p_t^2} \right\rangle.$$

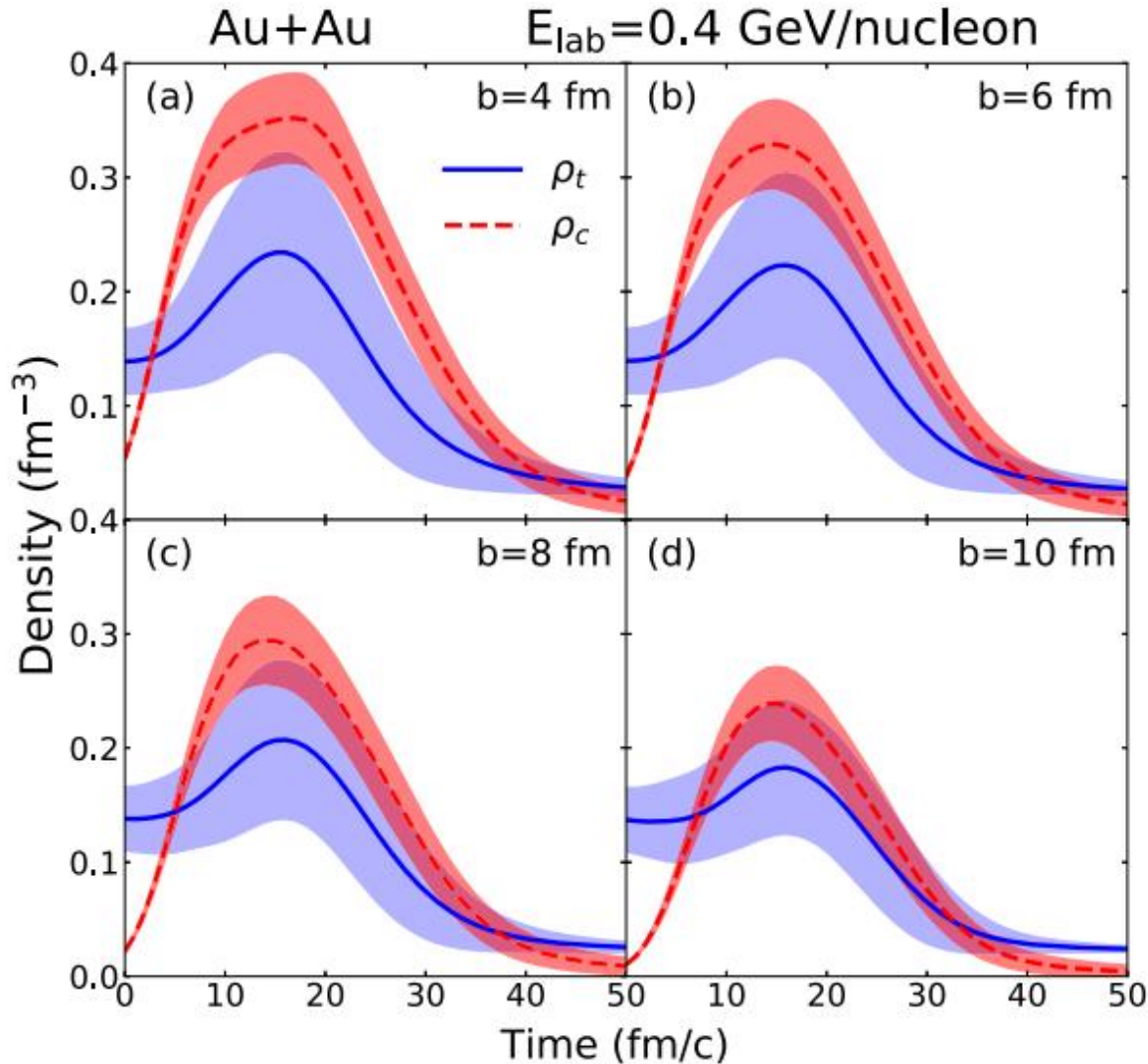
$$\Delta v_2^{\text{coll}}(t) = \left\langle V_2^{\text{aft.coll}}(t) - V_2^{\text{bef.coll}}(t) \right\rangle.$$

$$\Delta v_2^{\text{mf}}(t) = \left\langle V_2^{\text{aft.mf}}(t + \Delta t) - V_2^{\text{bef.mf}}(t) \right\rangle.$$

1. At the initial time, $\Delta v_2^{\text{coll}}(t)$ is zero because it takes approximately 5 fm/c for target and projectile nucleons to encounter each other.
2. $\Delta v_2^{\text{coll}}(t)$ is always negative which indicates NN collision enhances the signal of out-of-plane elliptic flow (decreases the value of v_2). This is because the outgoing nucleons along the reaction plane will be blocked.
3. The value of $\Delta v_2^{\text{mf}}(t)$ can be either negative or positive, depending on the impact parameter and time.



The density interval at which the nuclear equation of state is probed by the elliptic flow



ρ_c denotes the density at the central region.
 ρ_t denotes the density experienced by the traced nucleons.

$$\langle \rho \rangle_{v_2} = \frac{\int_0^{t_{sat}} |\Delta v_2(t)| \rho_t(t) dt}{\int_0^{t_{sat}} |\Delta v_2(t)| dt}$$

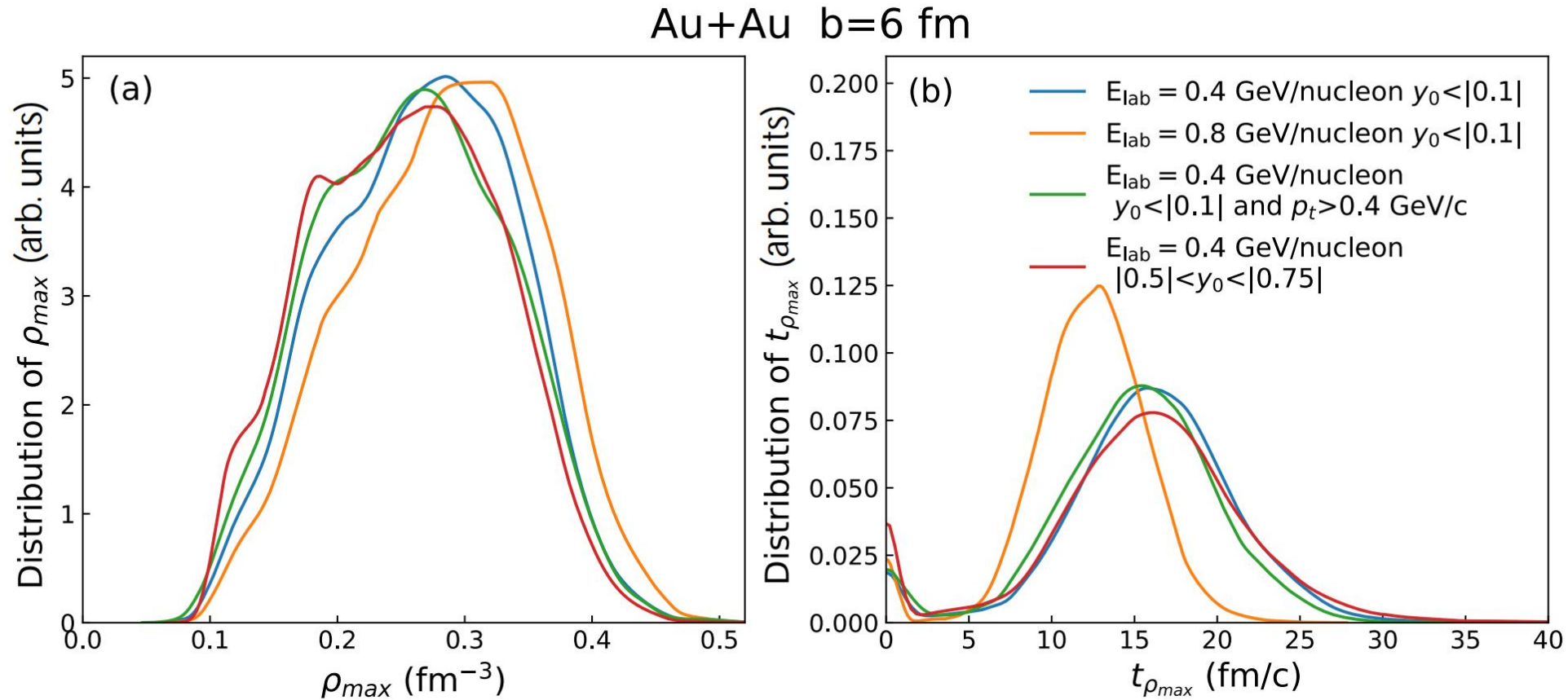
The most relevant density probed by v_2 of nucleons at mid-rapidity lies in $(0.9-1.3)\rho_0$ and dependent of the impact parameter.

Their values are found to be **60%** of the maximum density reached during the collisions.

The density interval at which the nuclear equation of state is probed by the elliptic flow

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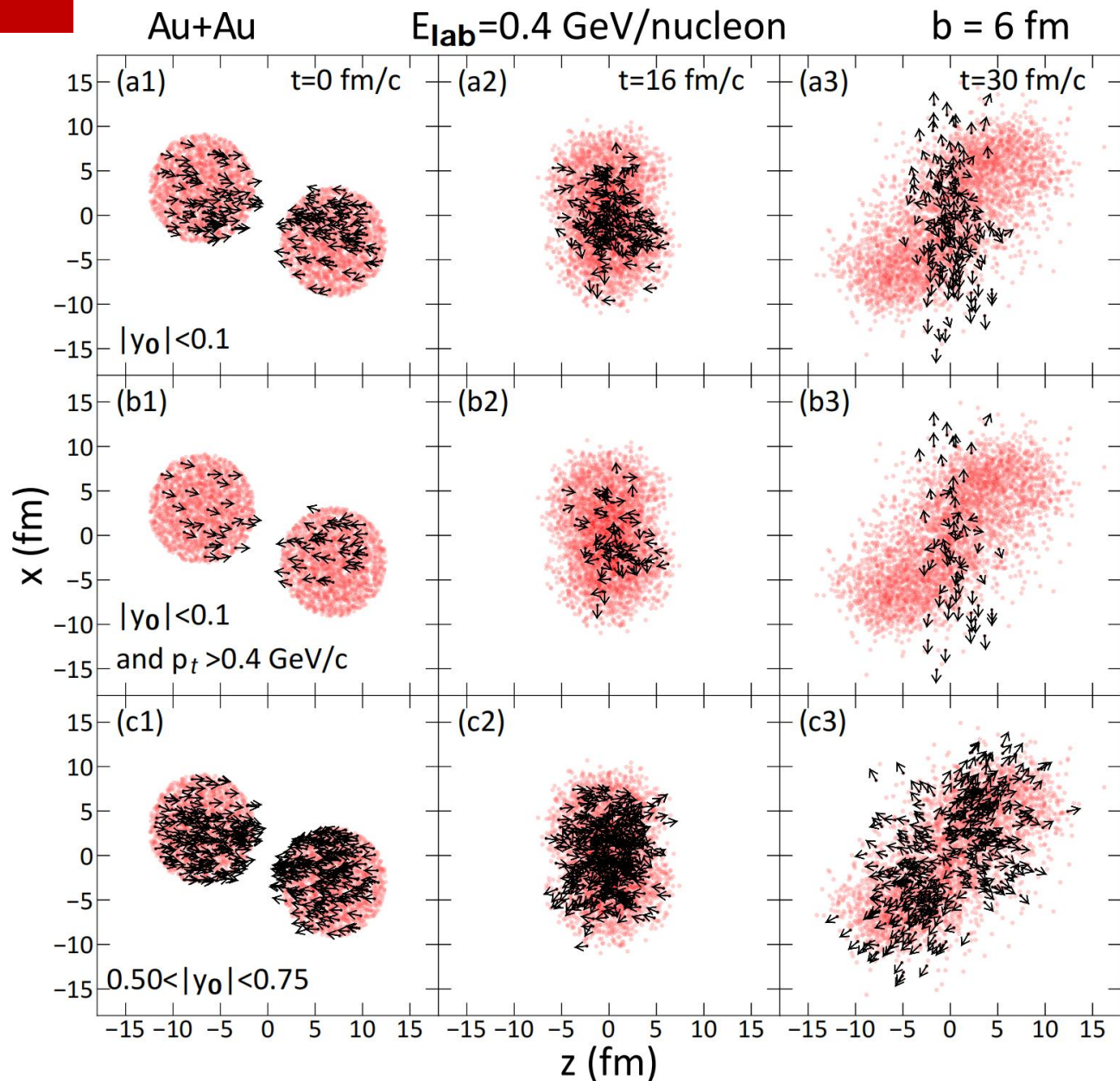
Anonymous referee: different nucleons reach their maximum densities at different times.



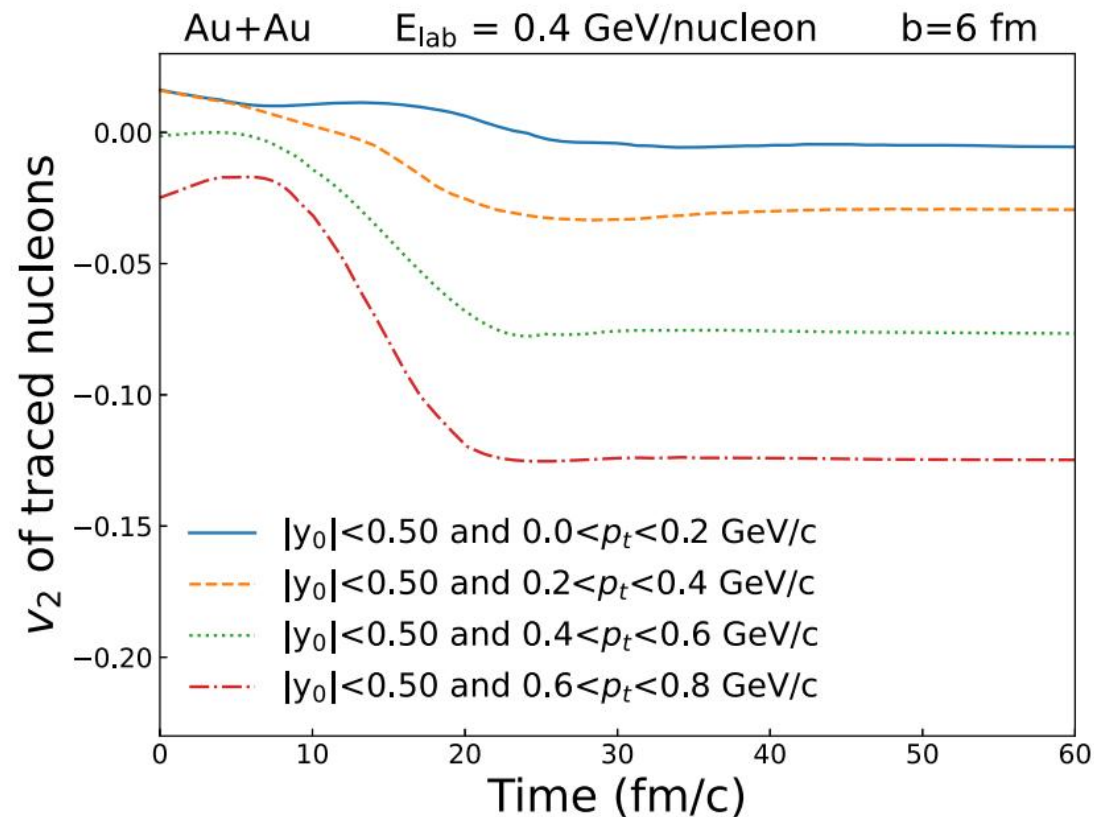
The distributions of the maximum density (ρ_{max}) defined for each traced nucleon.

The average of ρ_{max} is about **80%** of the maximum average density created in the central region.

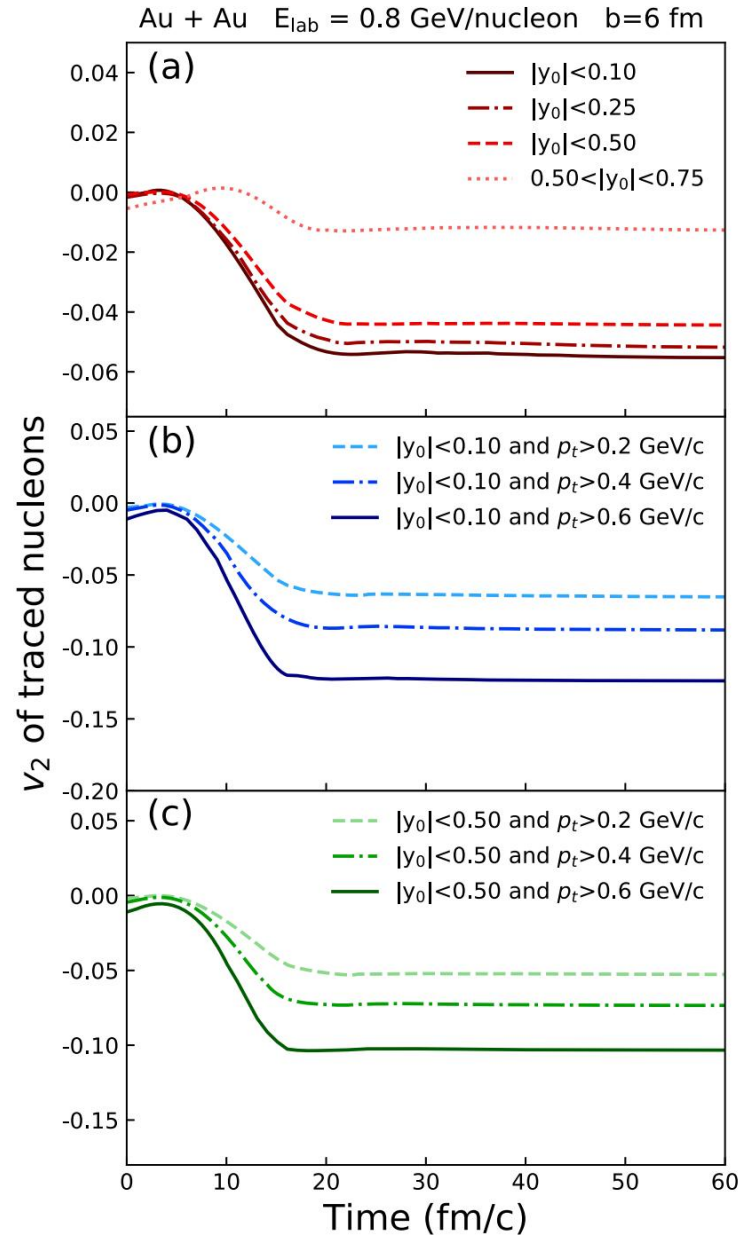
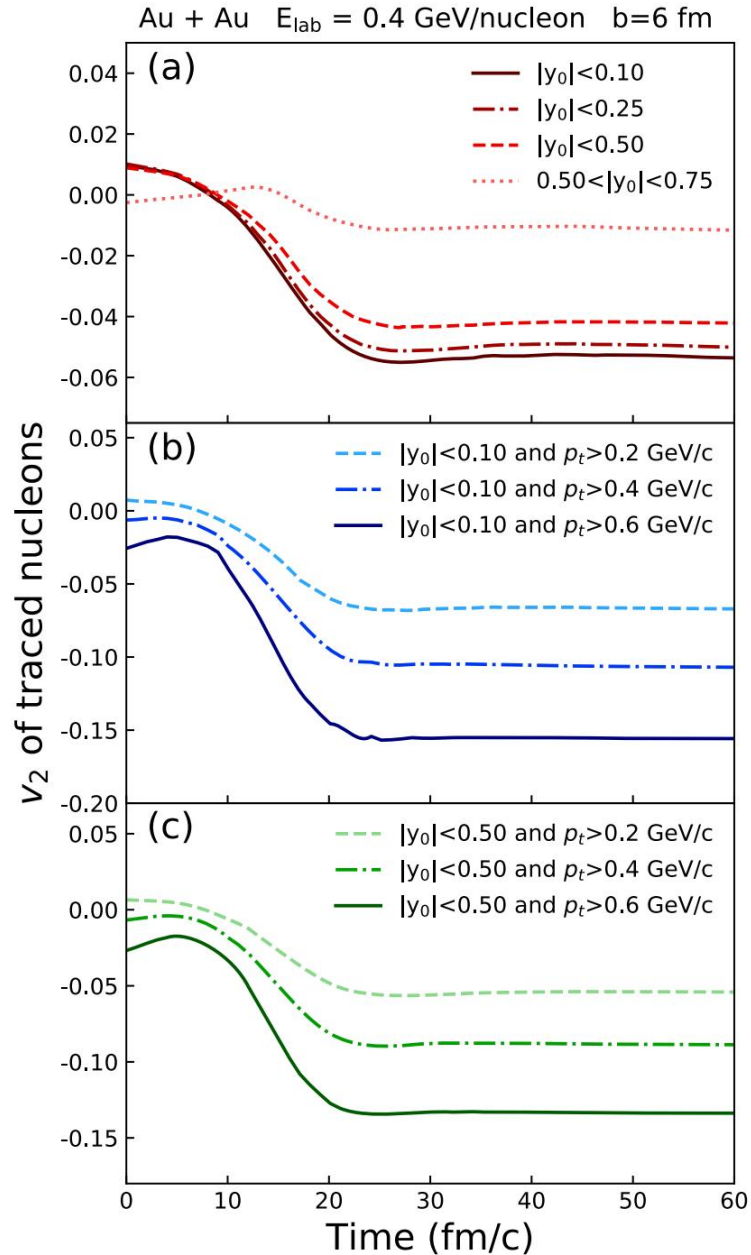
Effects of rapidity and transverse momentum cuts



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Effects of rapidity and transverse momentum cuts



Initial-final state correlation
 in the momentum space.
 This correlation becomes
 rather weak at
 $E_{\text{lab}} = 0.8$ GeV/nucleon.

Bayesian analysis of properties of nuclear matter with FOPI data

150 runs of different parameter sets in UrQMD to train emulator, and 20 runs for validate the trained emulator. 100, 000 events for each run.

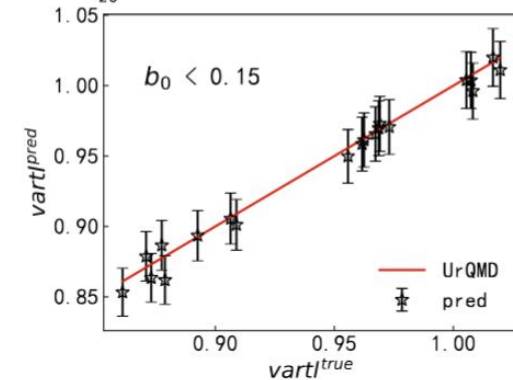
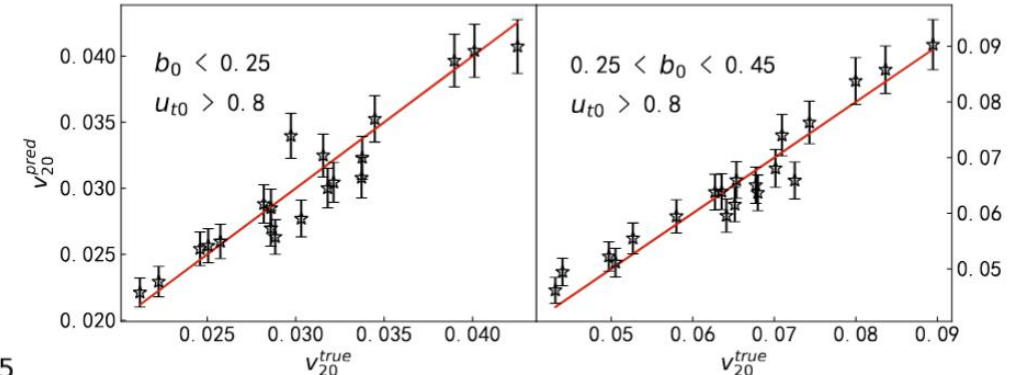
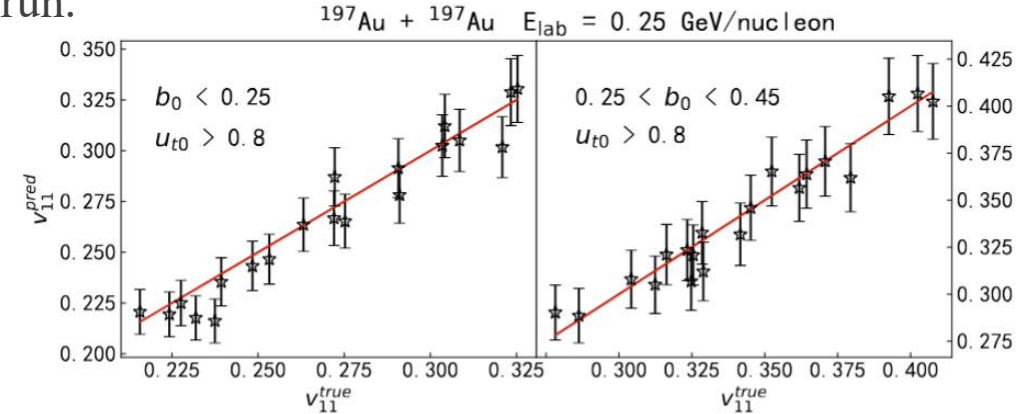
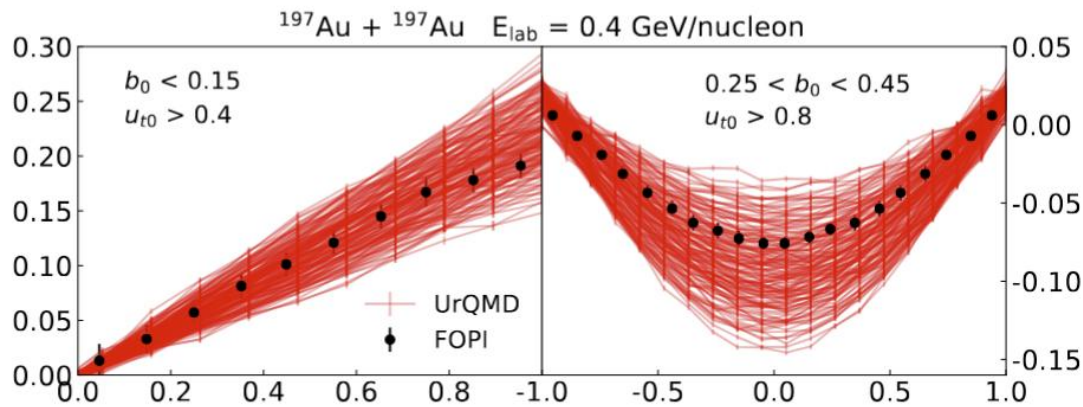
Table I. List of observables used in the analysis in $^{197}\text{Au} + ^{197}\text{Au}$ collisions with a beam energy of 0.25 GeV/nucleon

Observable	b_0	u_{t0}	value
v_{11}	$b_0 < 0.25$	$u_{t0} > 0.8$	0.23 ± 0.01
	$0.25 < b_0 < 0.45$	$u_{t0} > 0.8$	0.37 ± 0.01
$-v_{20}$	$b_0 < 0.25$	$u_{t0} > 0.8$	0.026 ± 0.001
	$0.25 < b_0 < 0.45$	$u_{t0} > 0.8$	0.046 ± 0.005
$vartl$	$b_0 < 0.15$	None	0.891 ± 0.041

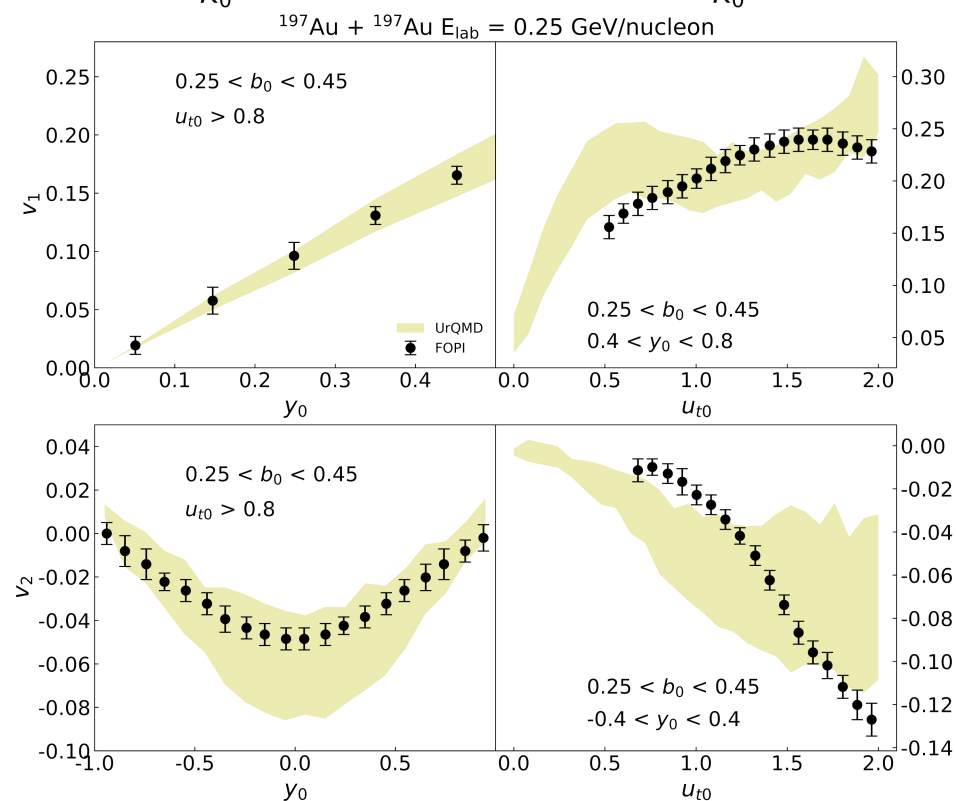
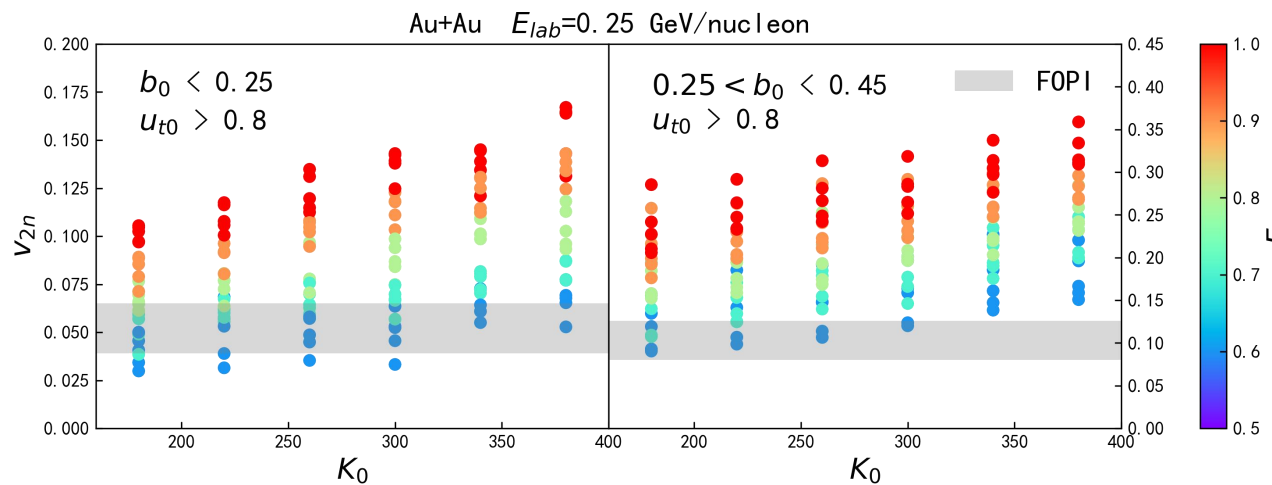
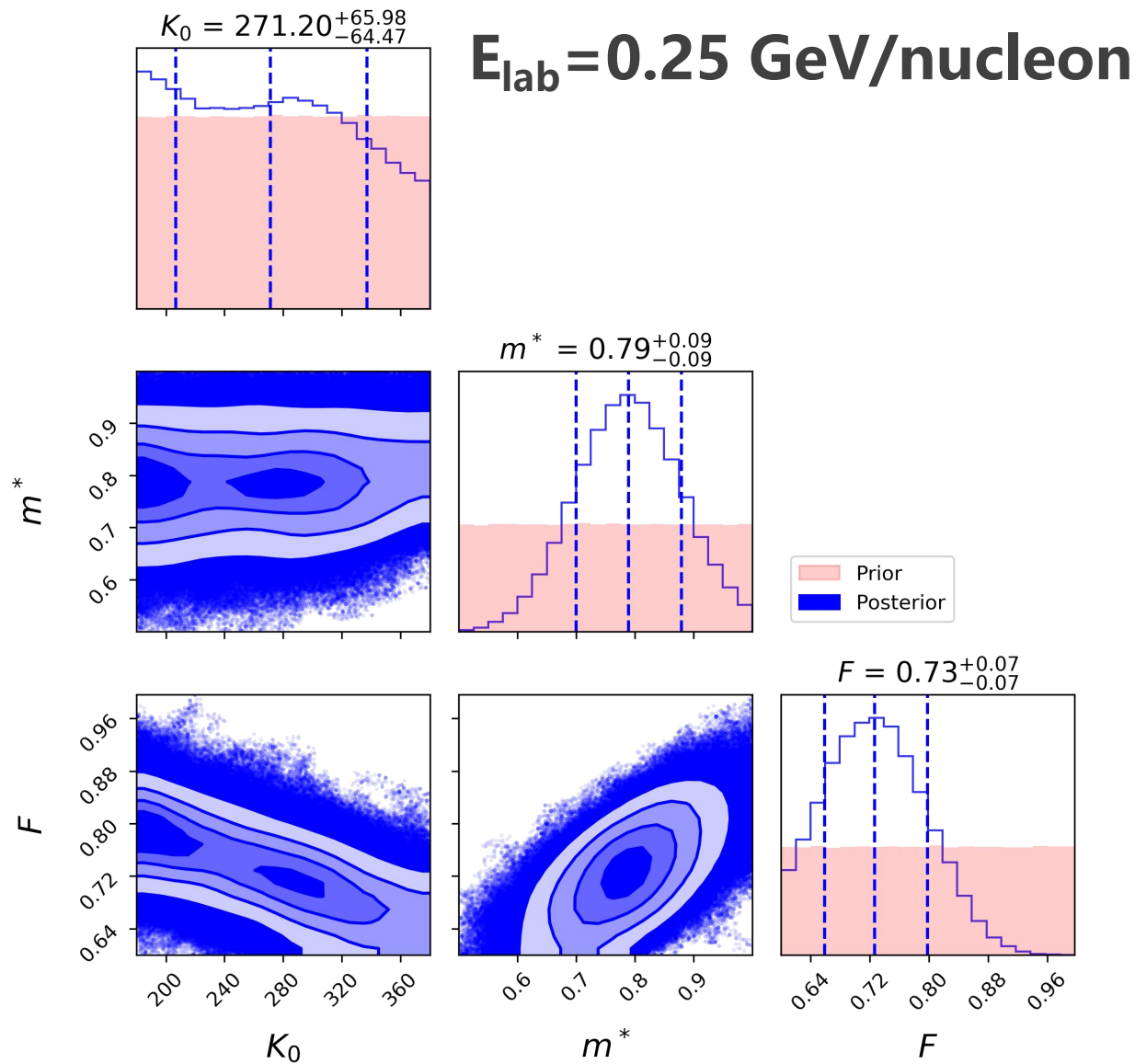
Table III. Parameters used in the present work

Para. Name	Description	Prior ranges
K_0	Incompressibility	[180, 380]
m^*	Isoscalar effective mass	[0.6, 0.95]
F	In-medium correction factor	[0.5, 1.0]

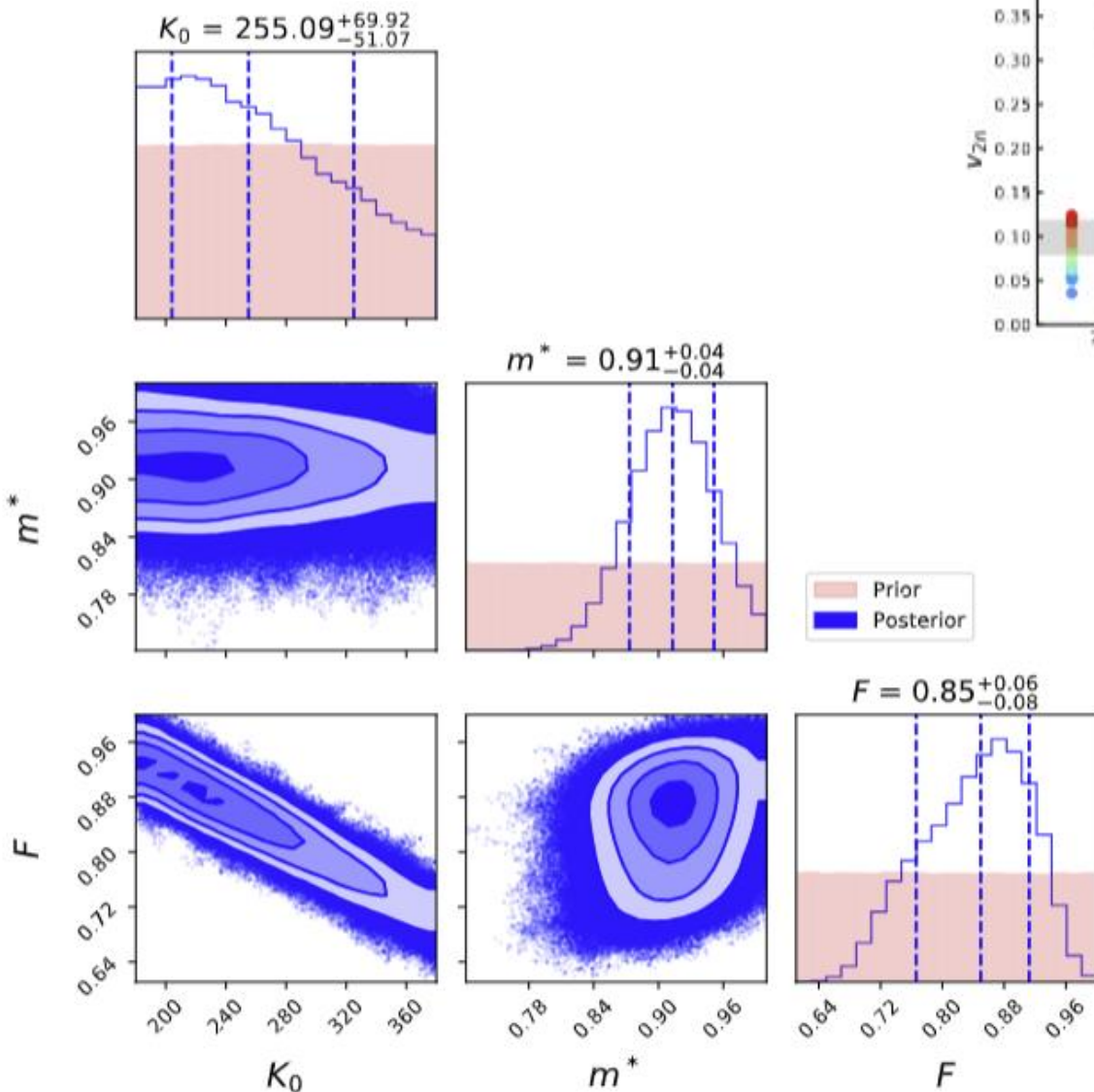
Gaussian process (GP) model is trained as an **emulator** of UrQMD model to interpolate the simulation results in the parameter space.



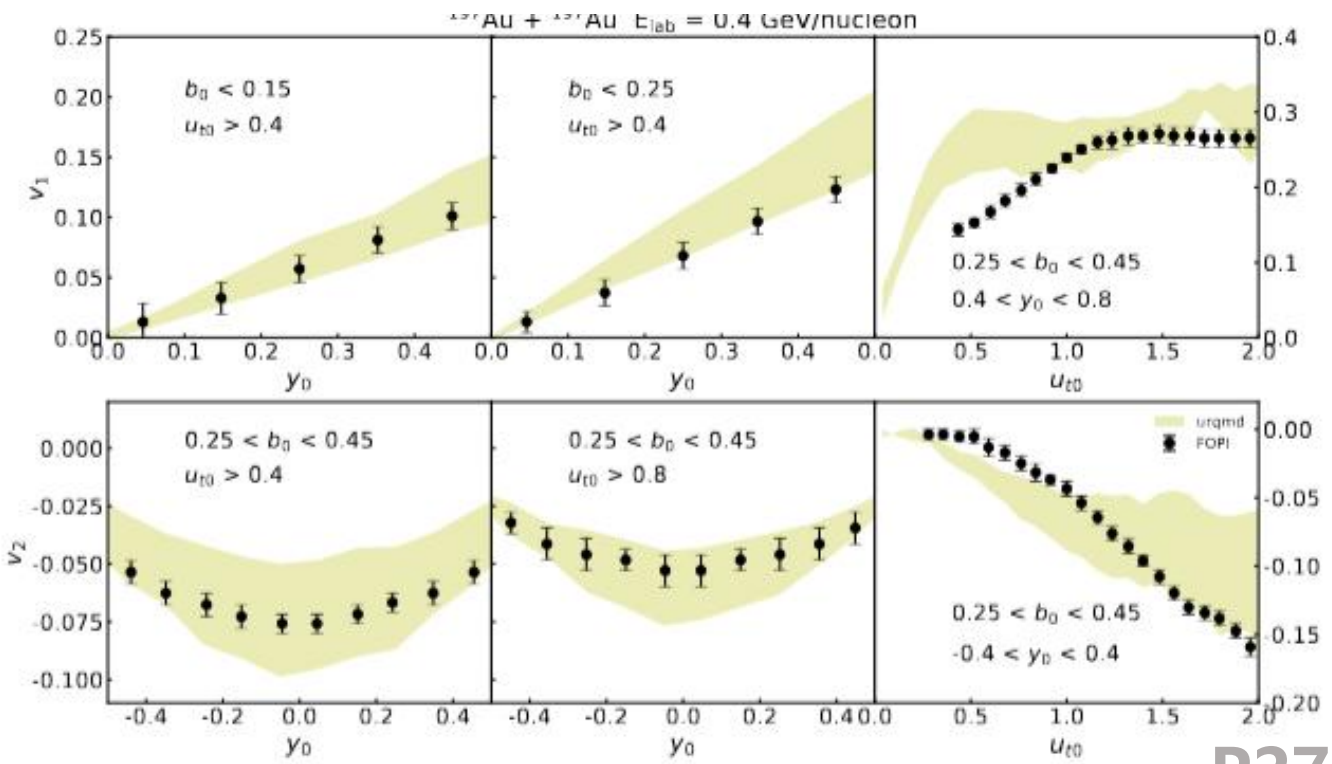
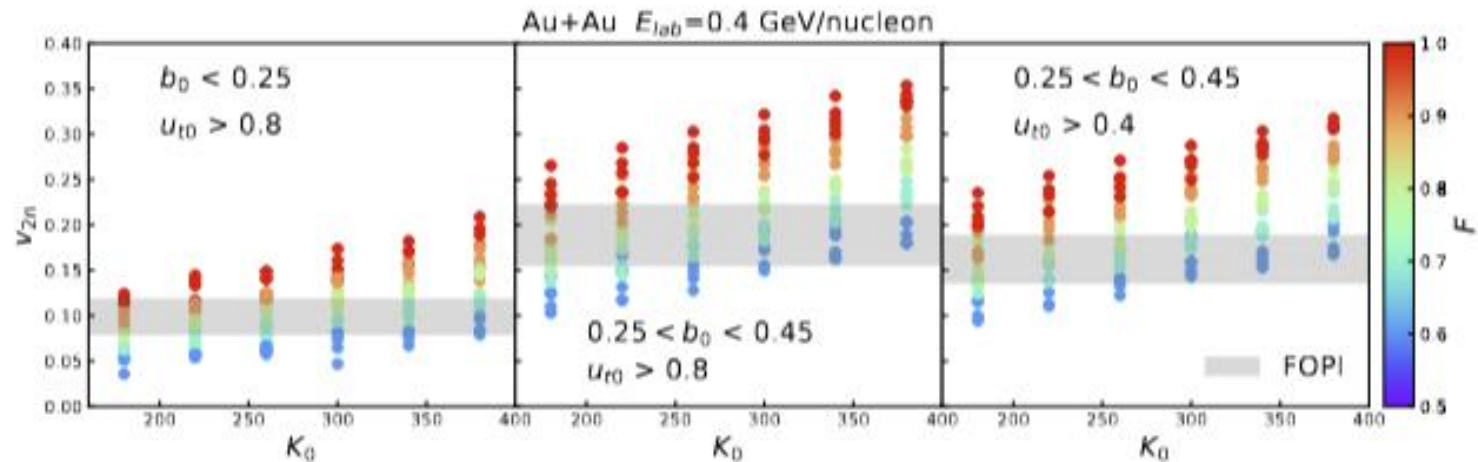
Bayesian analysis of properties of nuclear matter with FOPI data



Bayesian Inference of properties of nuclear matter with FOPI data



$E_{lab} = 0.4$ GeV/nucleon



Summary

1. Nucleon-nucleon collisions always suppress the value of v_2 (enhance the out-of-plane emission), while the nuclear mean field potential effect is more complex, which depends on both impact parameter and reaction time.
2. The density interval at which the nuclear equation of state is probed by the elliptic flow is about **60–80%** of the maximum average density created in the central region of heavy-ion collision.
3. $m^*/m=0.79\pm0.09$ and $F=\sigma^{\text{in-med.}}/\sigma^{\text{free}}=0.73\pm0.07$ with FOPI data at $E_{\text{lab}}=0.25$ GeV/nucleon, and the obtained values increased to $m^*/m=0.91\pm0.04$ and $F=0.85\pm0.07$ at $E_{\text{lab}}=0.4$ GeV/nucleon. **No tight constraint on K_0 !**

Thanks!