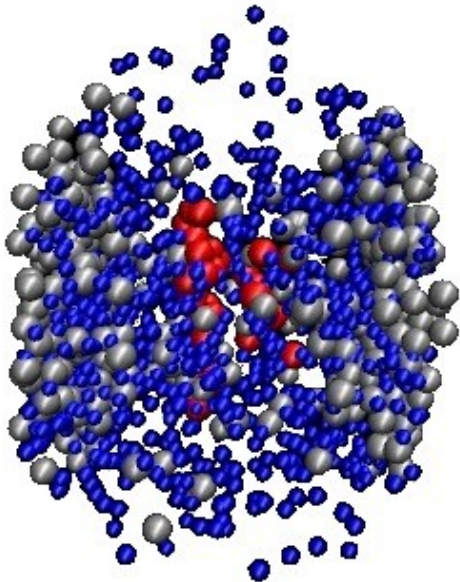


Chiral magnetic effect and evolution of electromagnetic field in relativistic heavy-ion collisions



Volodya Konchakovski



Network Workshop "TORIC"
Heraklion, Crete, Greece
6 September 2011

Outline

- **(P)HSD** – relativistic microscopic transport approach
- Comparison with the observables
- Chiral Magnetic Effect (**CME**)
- Model results
- Summary / Conclusions
- Coffee break

Basic Concept of HSD Transport Approaches

HSD – Hadron-String-Dynamics transport approach

Ehehalt, Cassing, Nucl.Phys. A602 (1996) 449; Cassing, Bratkovskaya, Phys. Rep.308 (1999) 65.

- the phase-space density f_i follows the **transport equations**

$$\left(\frac{\partial}{\partial t} + (\nabla_{\vec{p}} H) \nabla_{\vec{r}} - (\nabla_{\vec{r}} H) \nabla_{\vec{p}} \right) f_i(\vec{r}, \vec{p}, t) = I_{coll}(f_1, f_2, \dots, f_M)$$

with **collision terms** I_{coll} describing:

- **elastic and inelastic hadronic reactions:**

baryon-baryon, meson-baryon, meson-meson

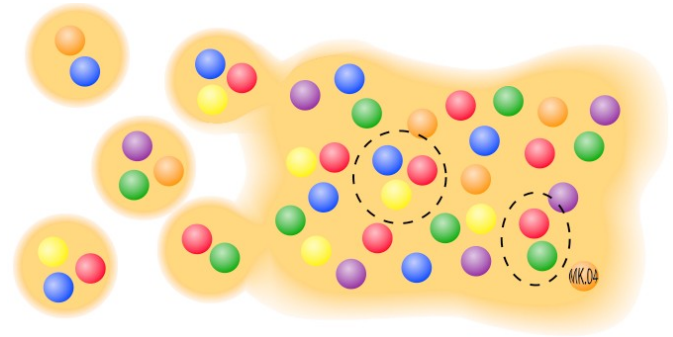
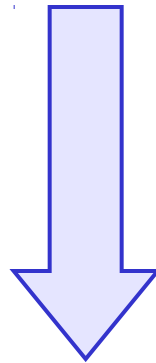
- **formation and decay of baryonic and mesonic resonances**

- **string formation and decay**

(for inclusive particle production: $BB \rightarrow X$, $mB \rightarrow X$, $X = \text{many particles}$)

- implementation of **detailed balance** on the level of $1 \leftrightarrow 2$ and $2 \leftrightarrow 2$ reactions (+ **$2 \leftrightarrow n$ multi-particle reactions in HSD !**)
- no explicit phase transition from hadronic to partonic degrees of freedom

Transport description of the **partonic** and **hadronic** phase



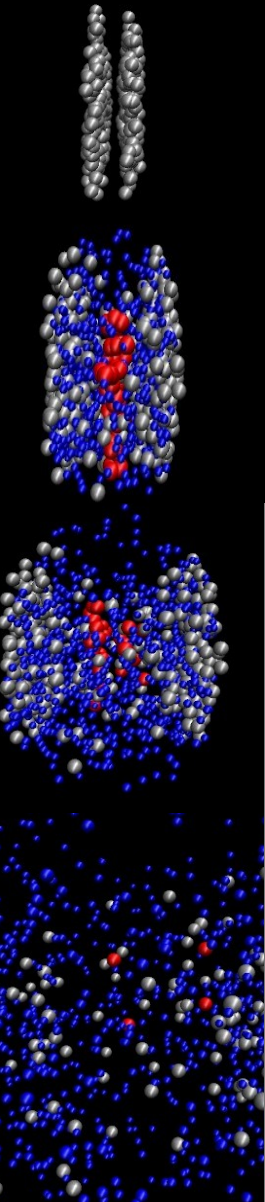
Parton-Hadron-
String-Dynamics
PHSD



PHSD - basic concepts

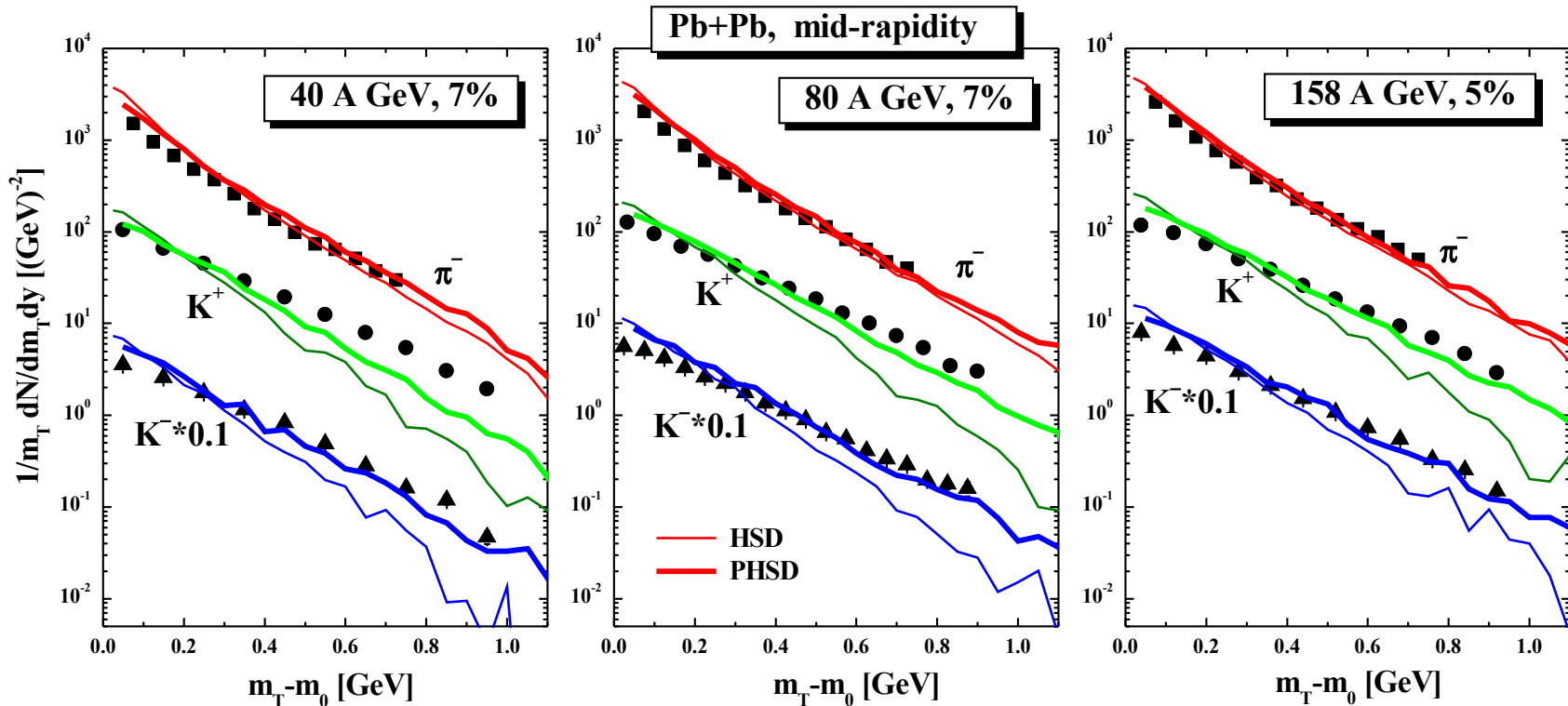
W. Cassing, E. Bratkovskaya,
PRC 78 (2008) 034919;
NPA831 (2009) 215;
EPJ ST 168 (2009) 3.

- **Initial A+A collisions – HSD: string formation and decay to pre-hadrons**
- **Fragmentation of pre-hadrons into quarks:** using the quark spectral functions from the **Dynamical QuasiParticle Model (DQPM)** approximation to QCD
DQPM: Peshier, Cassing, PRL 94 (2005) 172301; Cassing, NPA 791 (2007) 365; NPA 793 (2007).
- **Partonic phase:** quarks and gluons (= 'dynamical quasiparticles') with **off-shell spectral functions** (width, mass) defined by the DQPM
- **Elastic and inelastic parton-parton interactions:** using the effective cross sections from the DQPM
 - ✓ $q + \bar{q}$ (flavor neutral) \Leftrightarrow gluon (colored)
 - ✓ gluon + gluon \Leftrightarrow gluon (possible due to large spectral width)
 - ✓ $q + \bar{q}$ (color neutral) \Leftrightarrow hadron resonances
- **Hadronization:** based on DQPM - massive, off-shell quarks and gluons with broad spectral functions hadronize to **off-shell mesons and baryons:**
 - ✓ gluons $\Rightarrow q + \bar{q}$
 - ✓ $q + \bar{q} \Rightarrow$ meson (or string)
 - ✓ $q + q + q \Rightarrow$ baryon (or string)
- **Hadronic phase:** hadron-string interactions – off-shell HSD



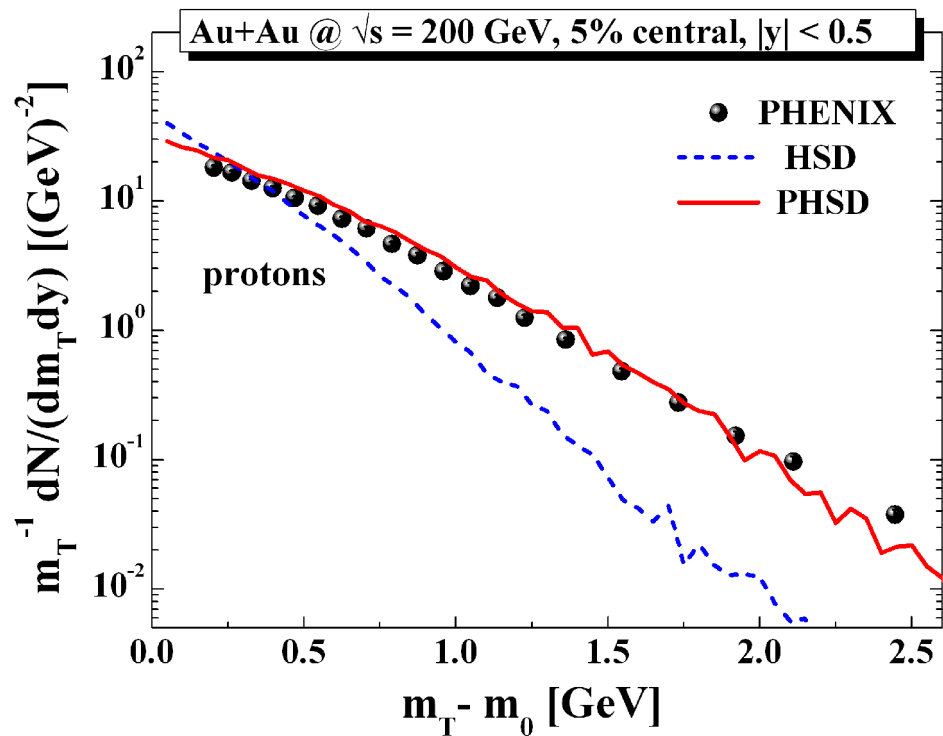
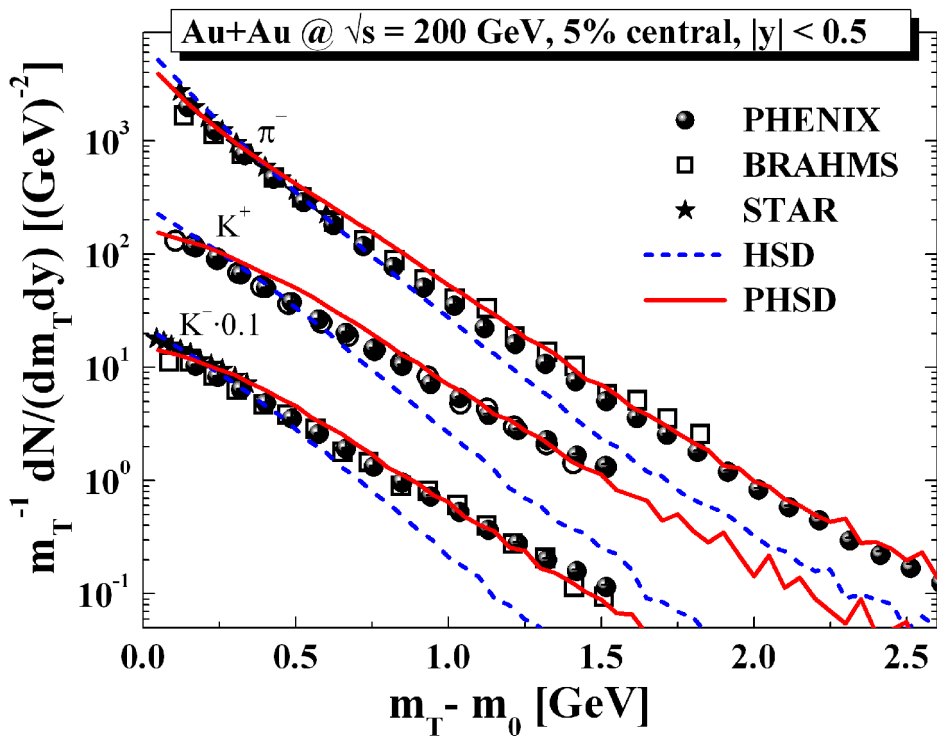
Transverse mass spectra at SPS energies

Central Pb + Pb at SPS energies



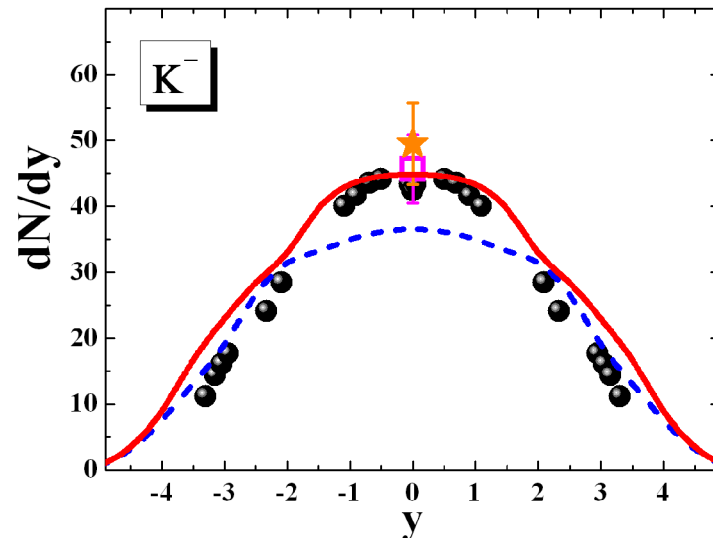
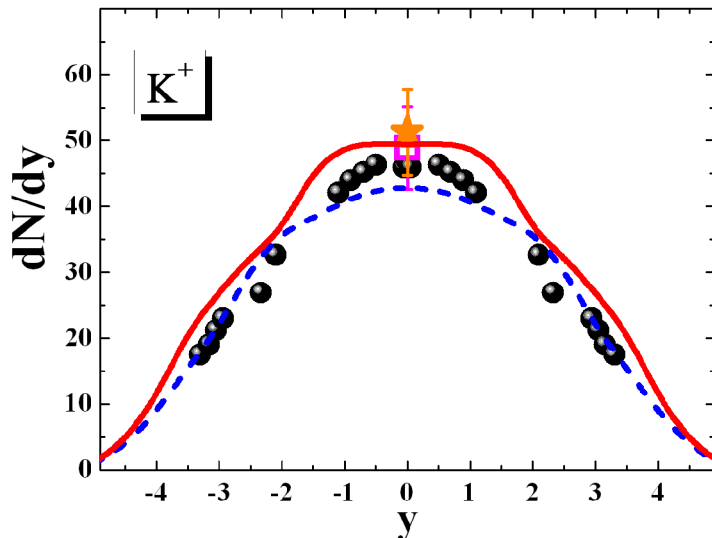
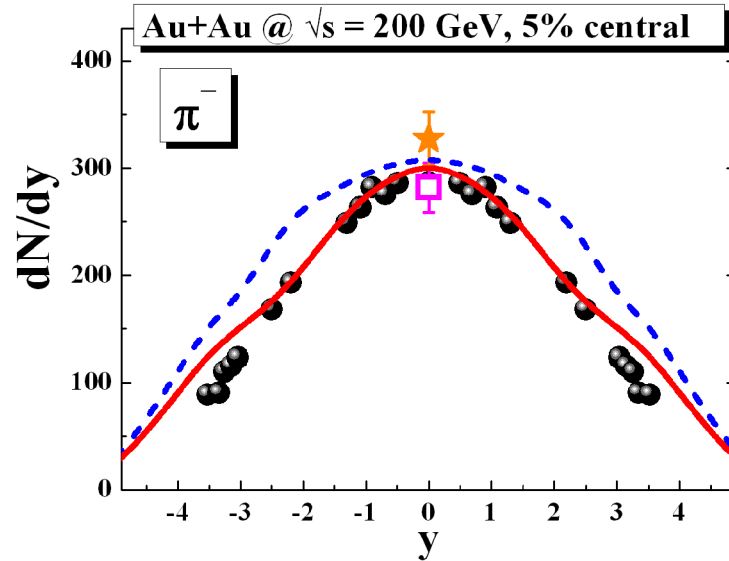
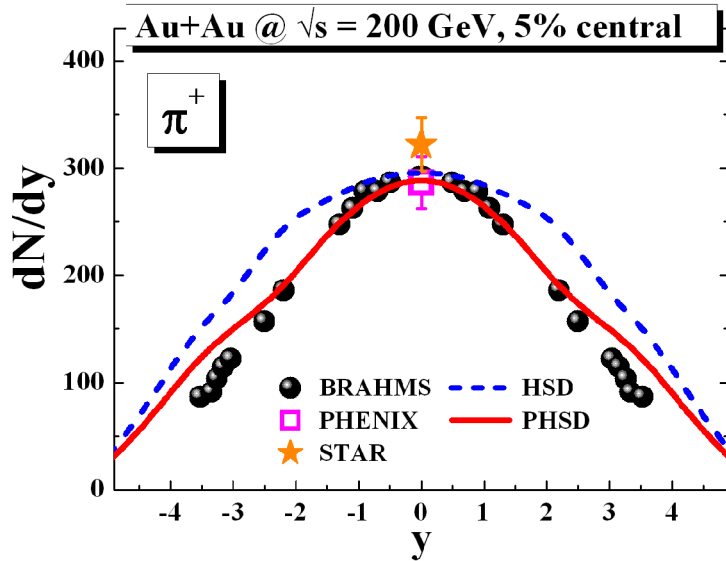
- PHSD gives harder spectra and works better than HSD at SPS (and top FAIR) energies
- At low SPS (and low FAIR) energies the effect of the partonic phase is less pronounced in rapidity distributions and m_T spectra

Transverse mass spectra at RHIC energies



PHSD improves significantly with respect to HSD (and the data) !

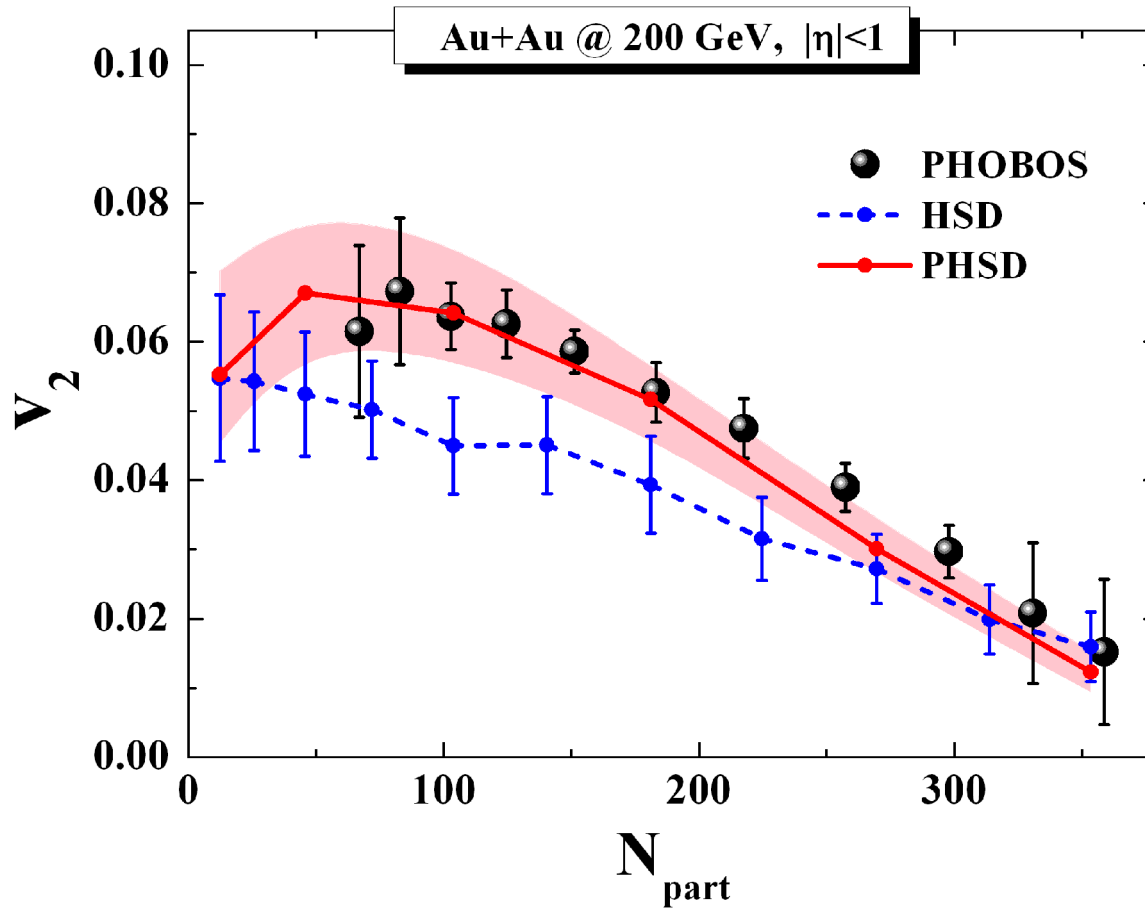
Rapidity distributions at RHIC energies



Look quite reasonable in comparison to data
from STAR, PHENIX and BRAHMS

Elliptic flow vs. centrality

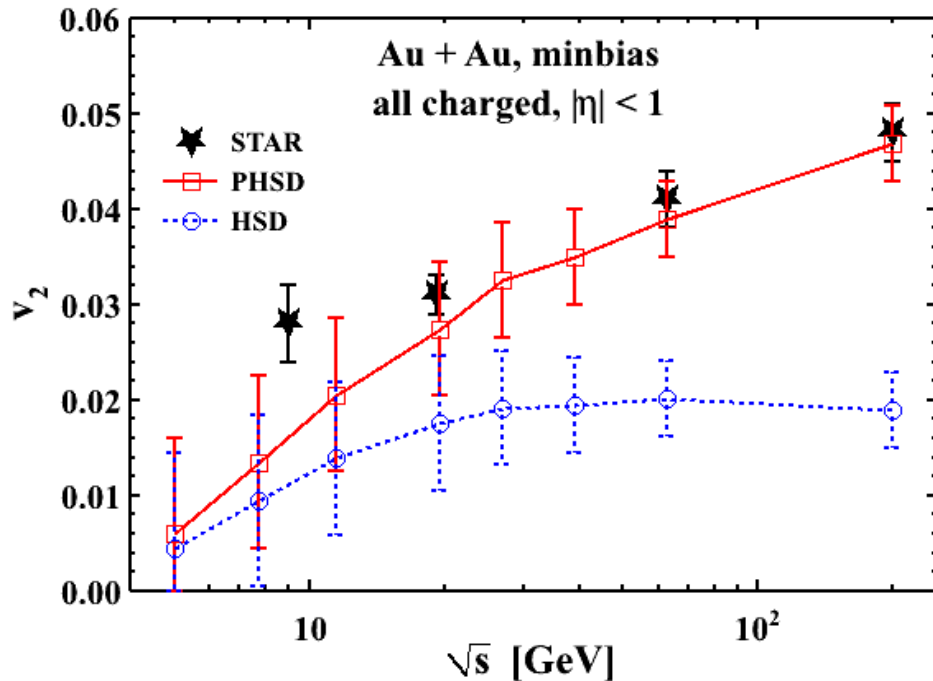
Au + Au collisions at $\sqrt{s} = 200$ GeV



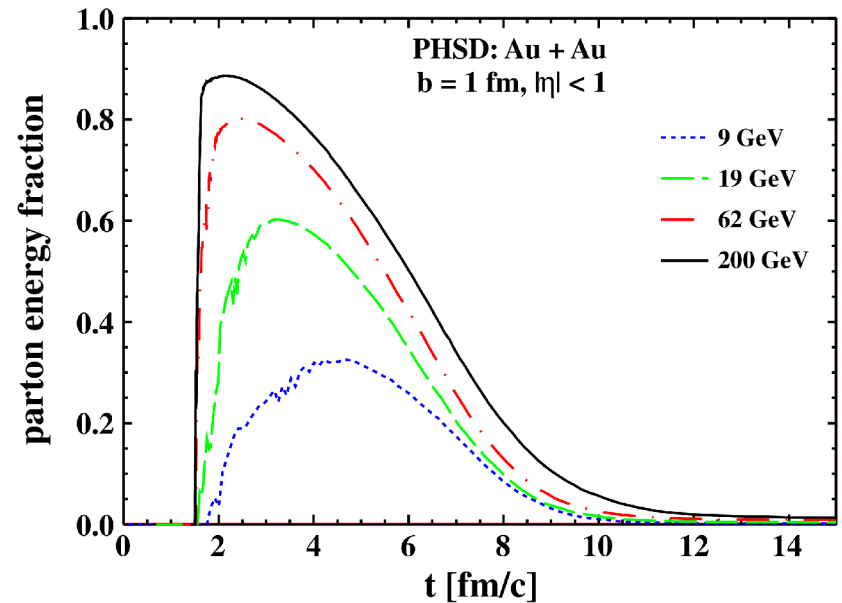
PHSD improves relative to HSD (in line with the data from PHOBOS)

Elliptic flow vs. collision energy

Elliptic flow v_2 at midrapidity



Parton energy fraction at midrapidity



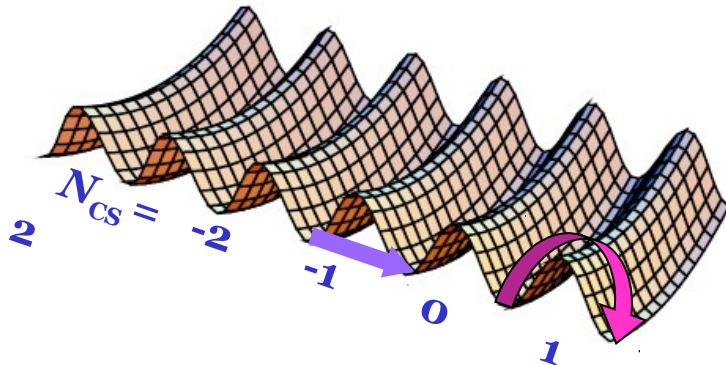
**Increase of parton fraction with energy
leads to increasing v_2**

Parity violation in strong interactions

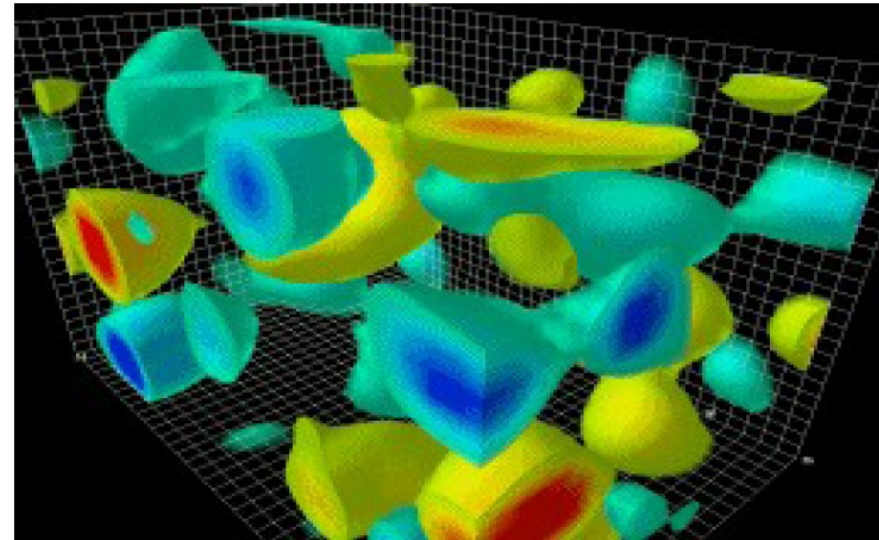
In QCD, chiral symmetry breaks due to a non-trivial topological effect; among the best evidence of this physics would be event-by-event strong parity violation.

The volume of the box is 2.4 by 2.4 by 3.6 fm.
The topological charge density of 4D gluon field configurations.
(Lattice-based animation by *Derek Leinweber*)

Energy of gluonic field is periodic in N_{cs} direction (\sim a generalized coordinate)



Instantons and sphalerons are localized (in space and time) solutions describing transitions between different vacua via tunneling or crossing the barrier



Dynamics is a random walk between states with different topological charges.

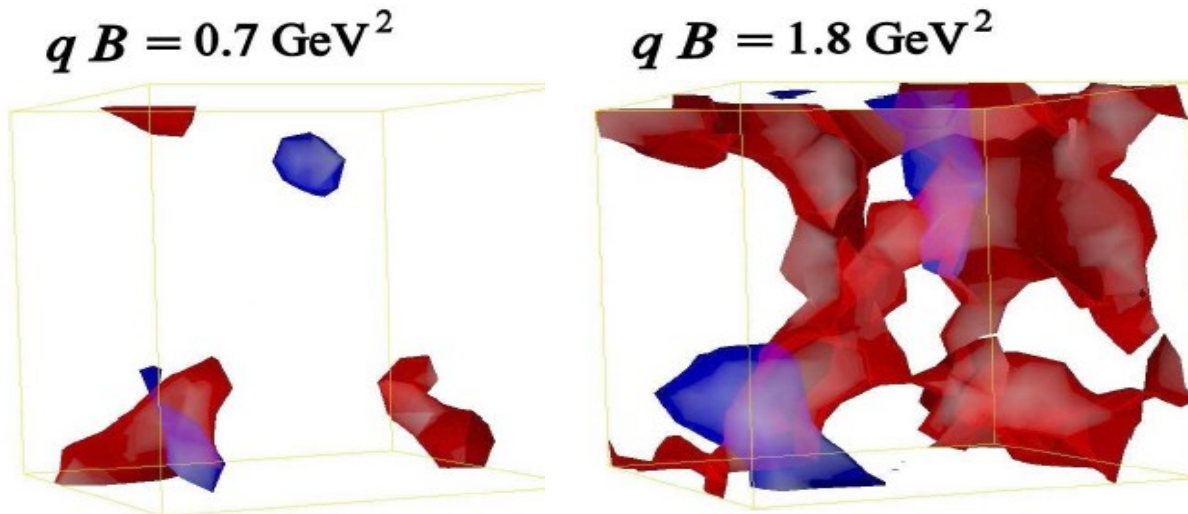
Charge separation: CP violation signal

Dynamics is a random walk between states with different topological charges. In this states **a balance** between left-handed and right-handed quarks **is destroyed**, $N_R - N_L = Q_T \rightarrow$ **violation** of P-, CP- symmetry.

Average total topological charge **vanishes** $\langle n_w \rangle = 0$ but its **variance** is equal to the total number of transitions $\langle n_w^2 \rangle = N_t$

Fluctuation of topological charges **in the presence** of magnetic field induces electric current which will separate different charges

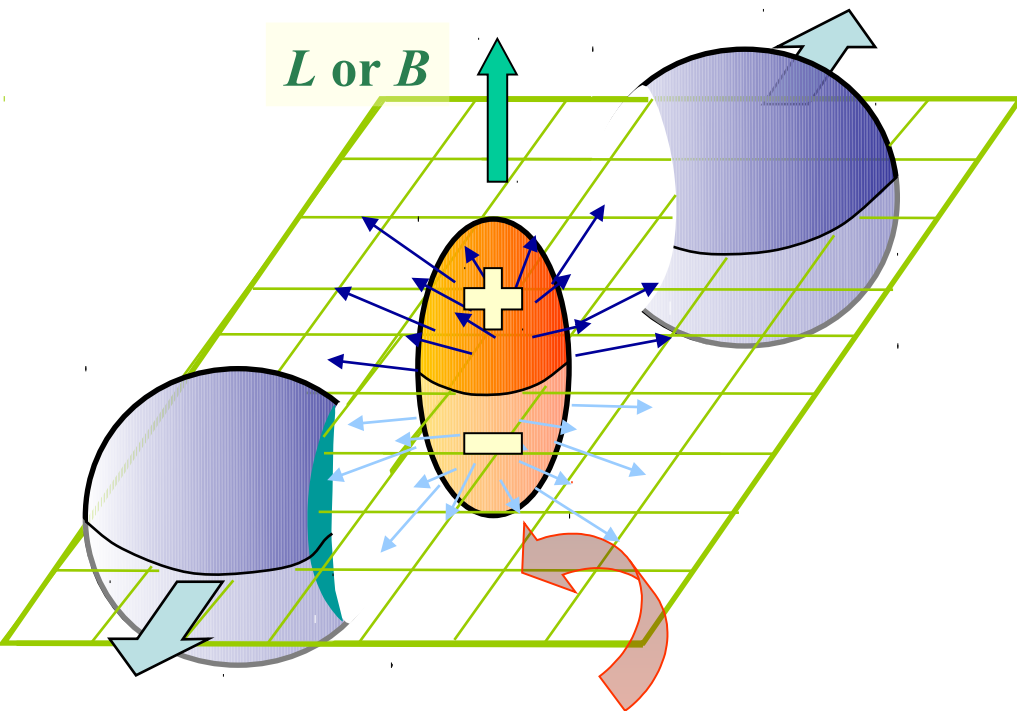
Lattice gauge theory



The excess of electric charge density due to the applied magnetic field.
Red — positive charges,
blue — negative charges.

P.V. Buividovich et al.,
PRD80 (2009) 054503

Charge separation in HIC



Non-zero angular momentum
(or equivalently magnetic field)
in heavy-ion collisions make it
possible for P- and CP-odd
domains to induce charge
separation

D.Kharzeev, PLB 633 (2006) 260.

Electric dipole moment of QCD matter !

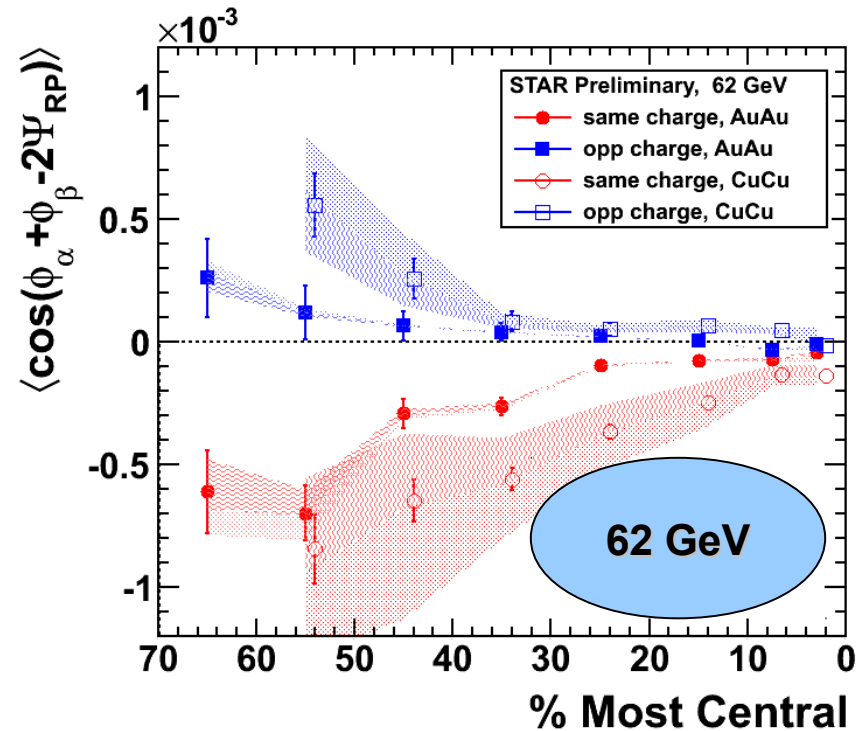
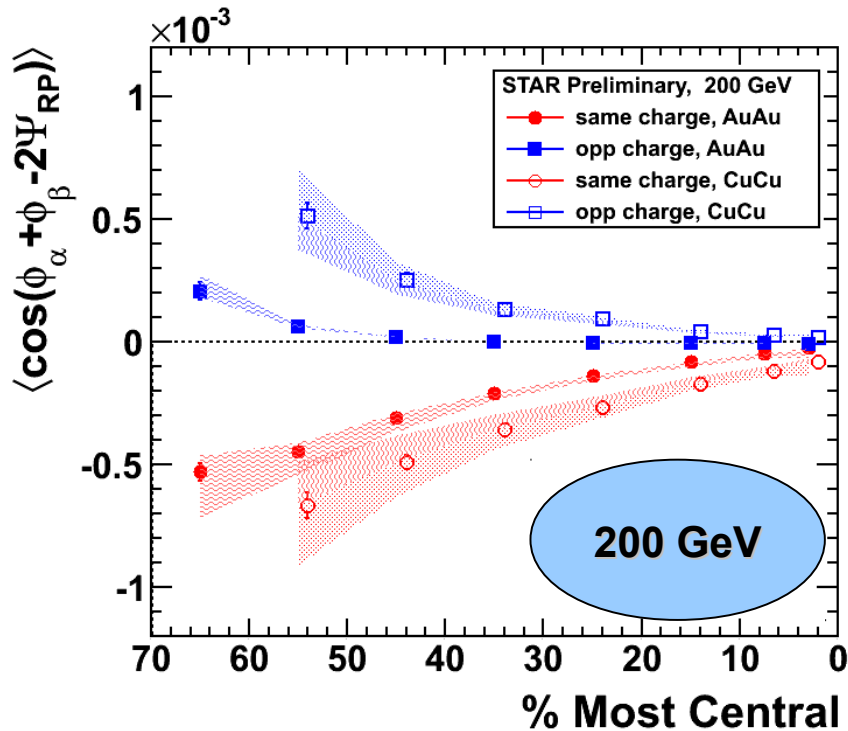
$$\langle \cos(\varphi_a + \varphi_b - 2\psi_{RP}) \rangle$$

Measuring the charge separation with
respect to the reaction plane was proposed
by S.Voloshin, Phys. Rev. C 70 (2004) 057901.

Charge separation in RHIC experiments

STAR Collaboration, PRL 103 (2009) 251601

$$\langle \cos(\varphi_a + \varphi_b - 2\psi_{RP}) \rangle$$



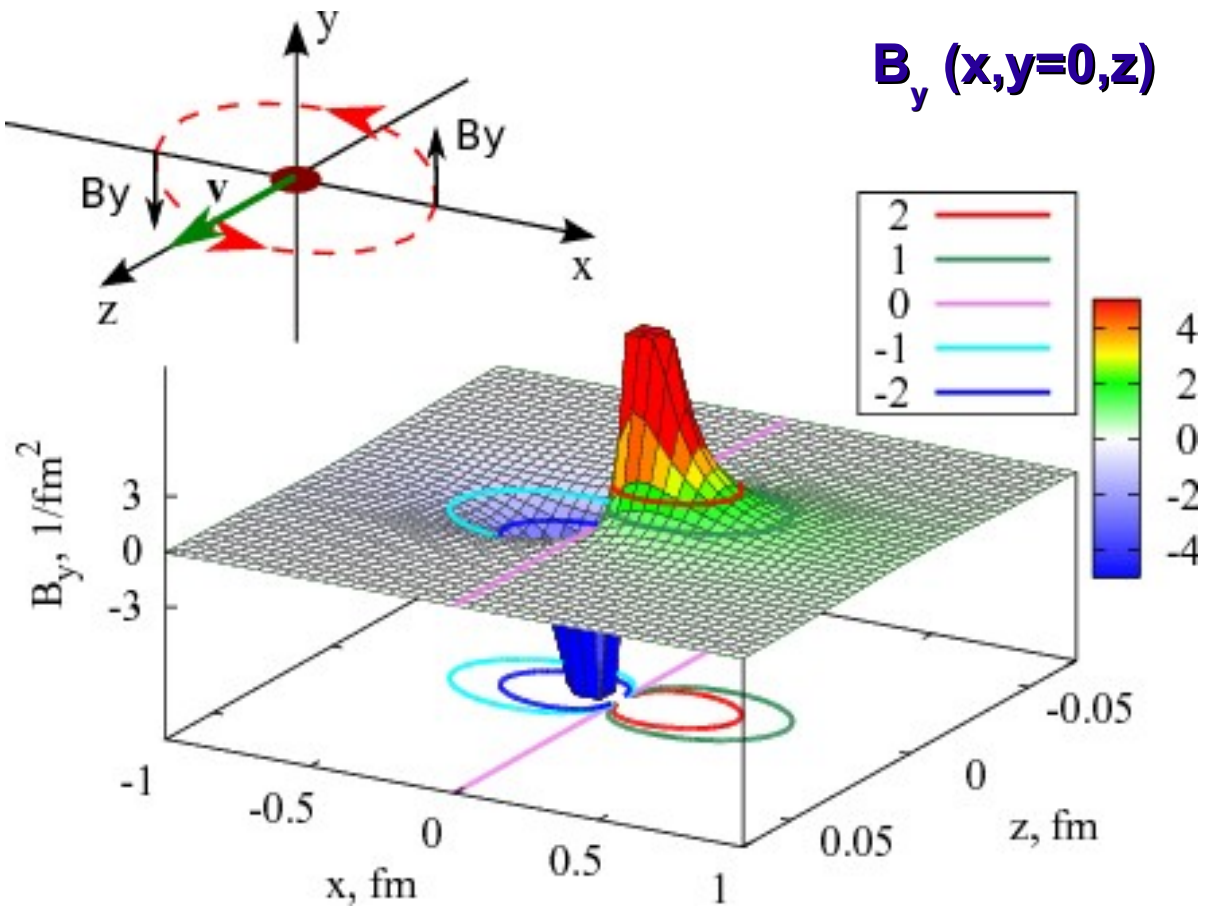
Combination of intense B and deconfinement is needed for a spontaneous parity violation signal

Hadron-String-Dynamics HSD



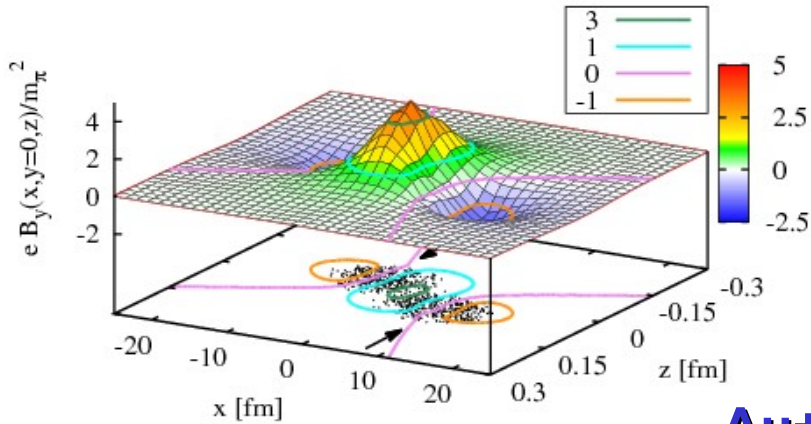
Retarded electromagnetic field

Magnetic field for a single moving charge

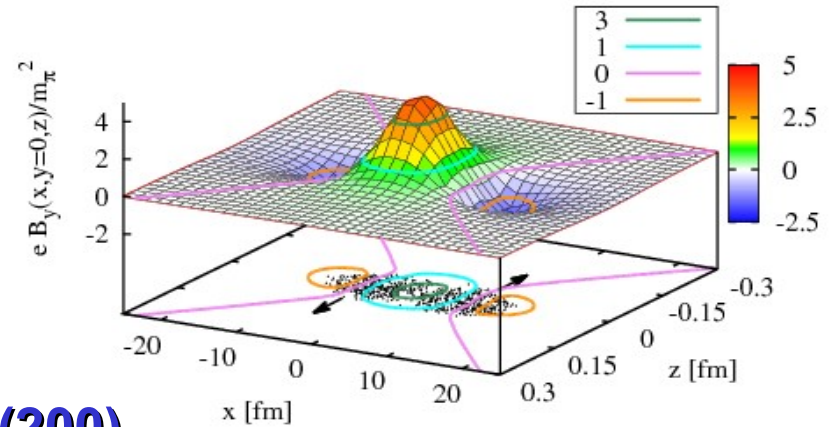


Magnetic field evolution

AuAu, $\sqrt{s_{NN}} = 200$ GeV, $b=10$ fm, $t=0.01$ fm/c

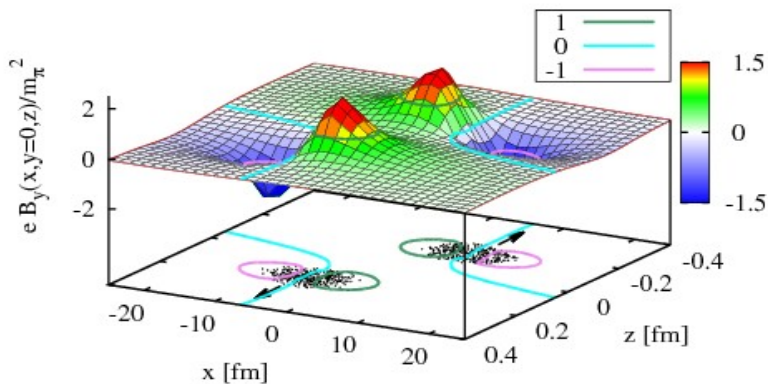


AuAu, $\sqrt{s_{NN}} = 200$ GeV, $b=10$ fm, $t=0.05$ fm/c

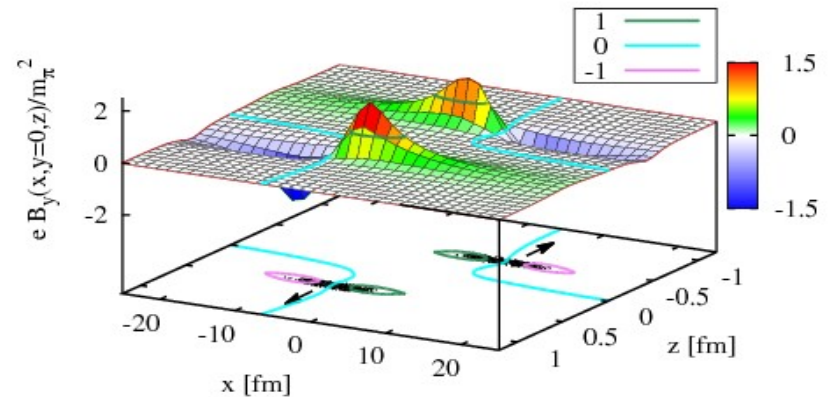


Au+Au (200)
b=10 fm

AuAu, $\sqrt{s_{NN}} = 200$ GeV, $b=10$ fm, $t=0.2$ fm/c

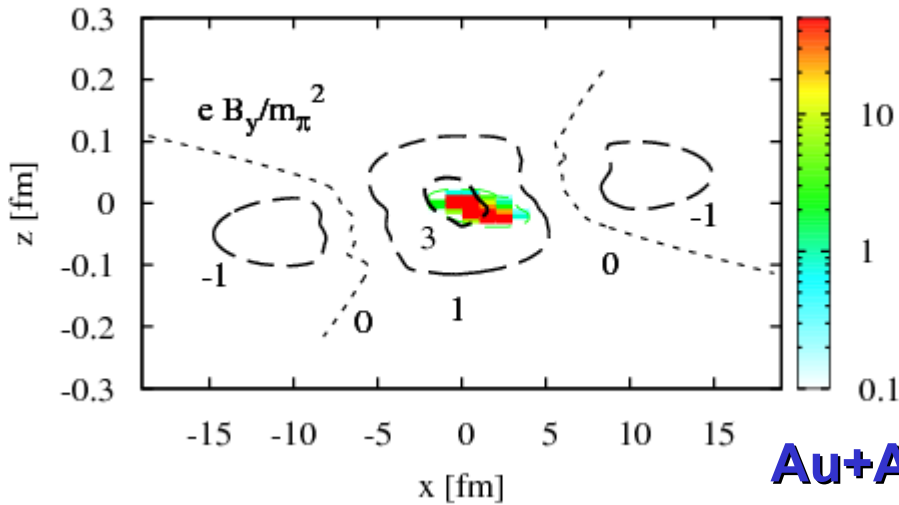


AuAu, $\sqrt{s_{NN}} = 200$ GeV, $b=10$ fm, $t=0.5$ fm/c

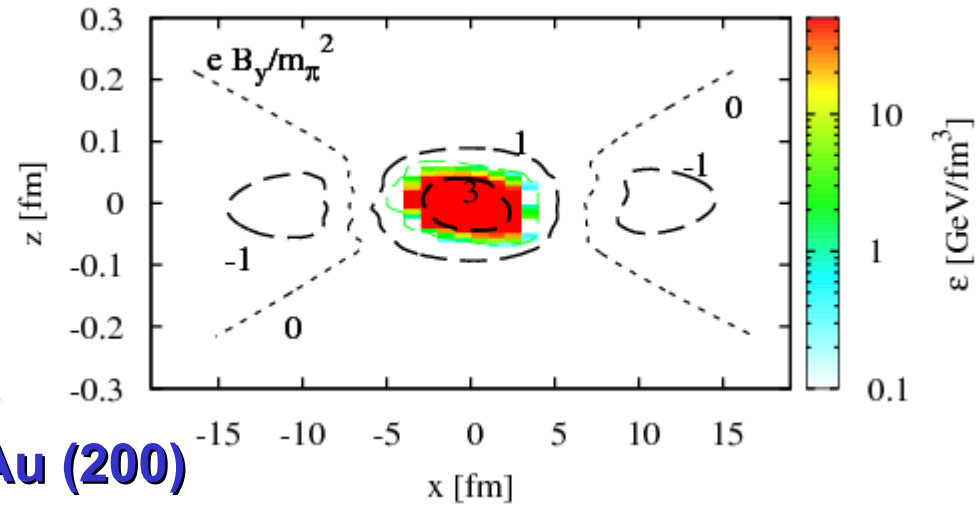


Magnetic field and energy density correlation

AuAu, $\sqrt{s_{NN}} = 200$ GeV, $b=10$ fm, $t=0.01$ fm/c

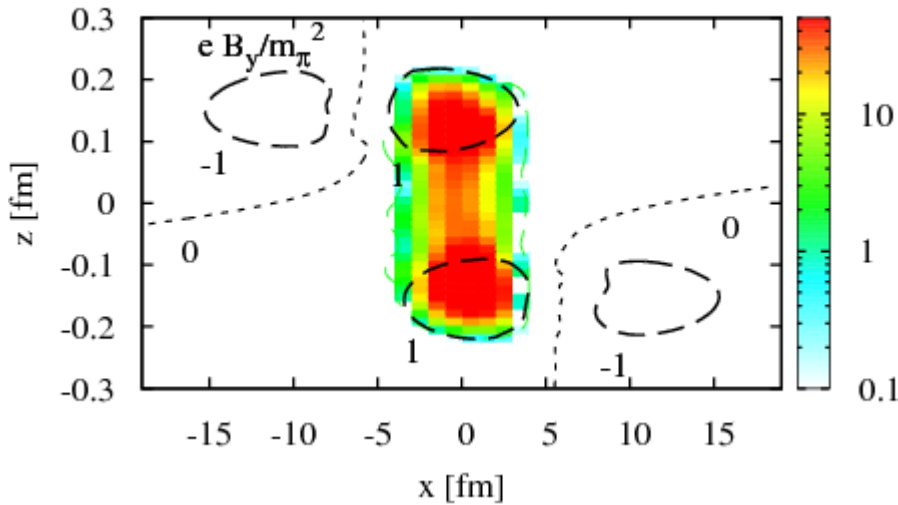


AuAu, $\sqrt{s_{NN}} = 200$ GeV, $b=10$ fm, $t=0.05$ fm/c

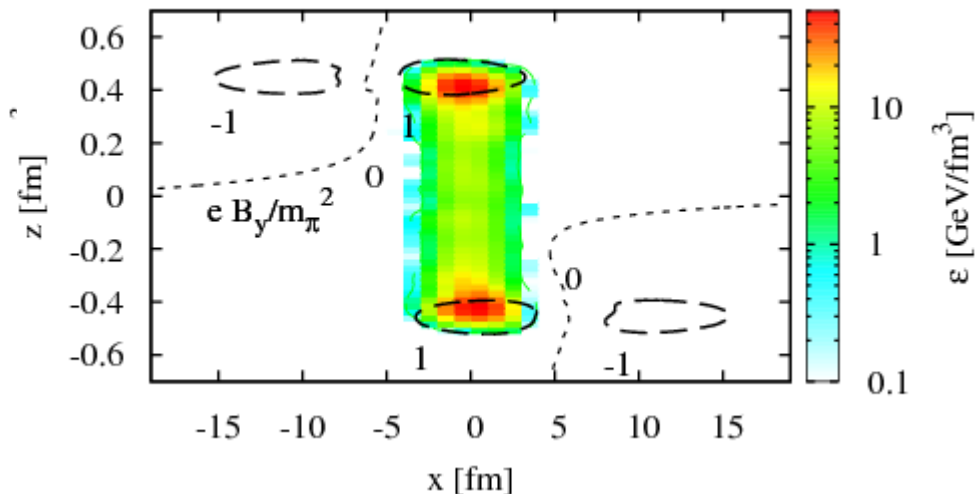


Au+Au (200)
b=10 fm

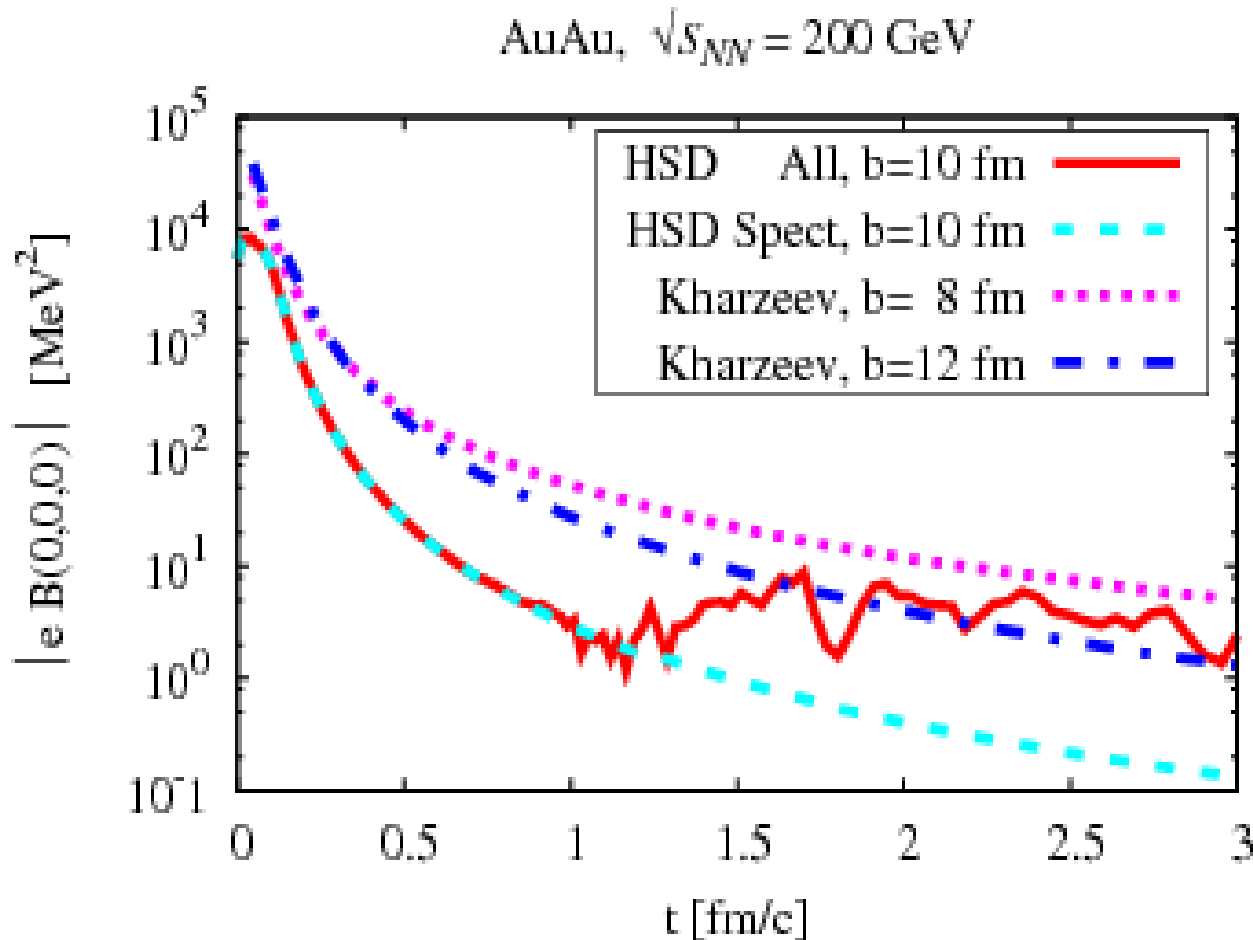
AuAu, $\sqrt{s_{NN}} = 200$ GeV, $b=10$ fm, $t=0.2$ fm/c



AuAu, $\sqrt{s_{NN}} = 200$ GeV, $b=10$ fm, $t=0.5$ fm/c



Time dependence of eB_y



D.E. Kharzeev et al.,
Nucl. Phys. A803, 227 (2008)

Collision of two infinitely thin layers (pancake-like)

V.Voronyuk, et al.,
PRC83 (2011) 054911

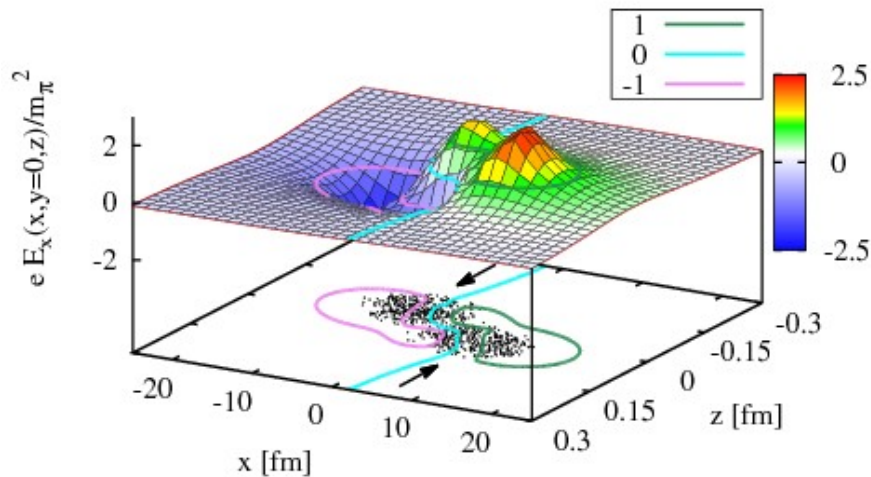
- Until $t \sim 1$ fm/c the induced magnetic field is defined by spectators only.
- Maximal magnetic field is reached during nuclear overlapping time $\Delta t \sim 0.2$ fm/c, then the field goes down exponentially.

Electric field evolution

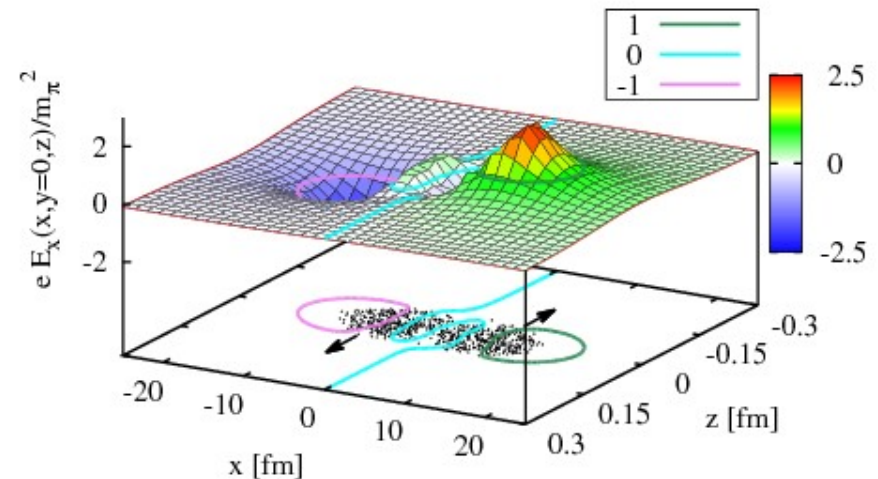


Electric field of a single moving charge has a “hedgehog” (or see urchin) shape

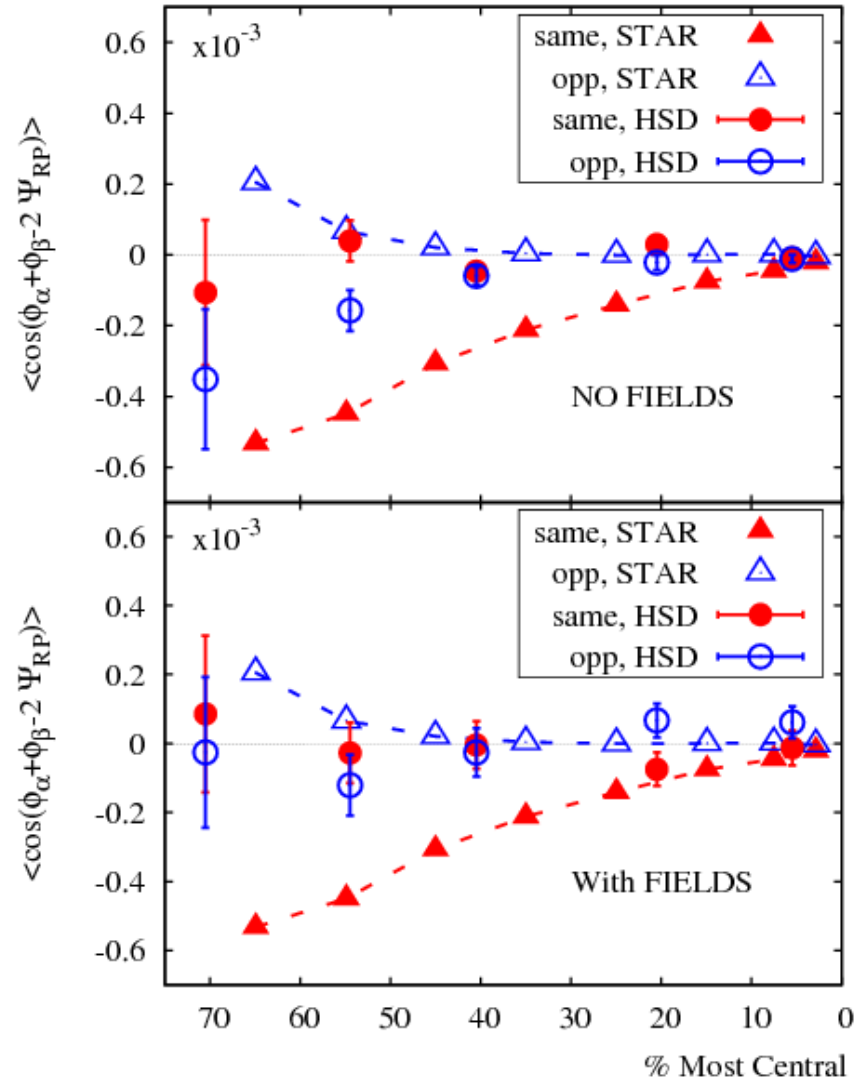
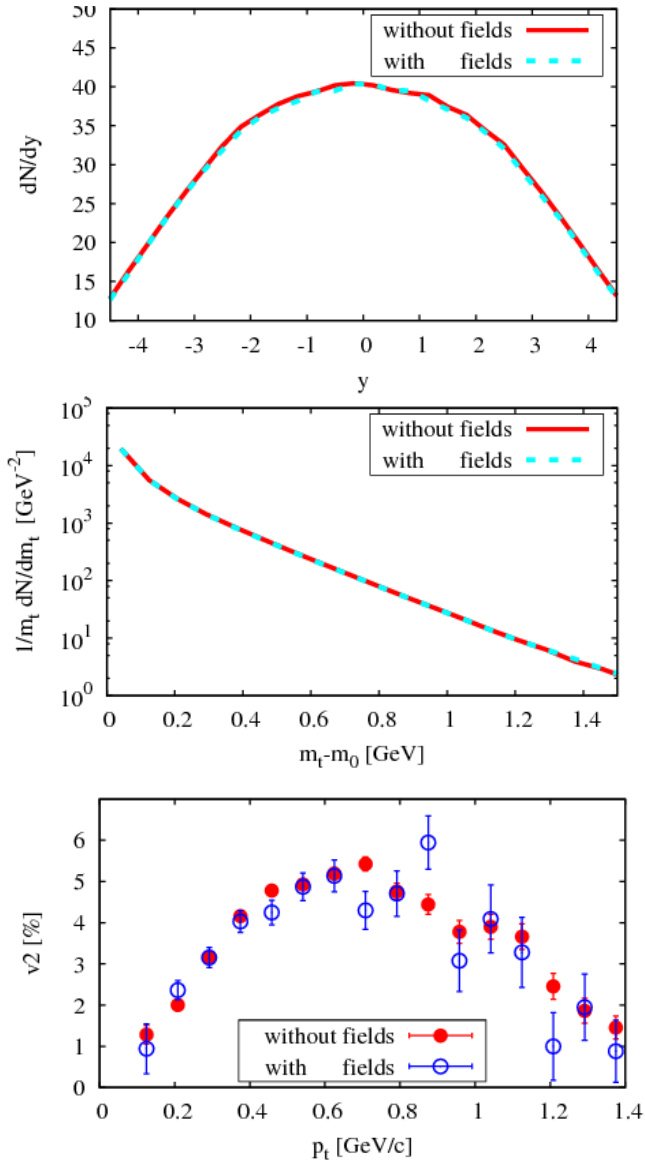
AuAu, $\sqrt{s_{NN}} = 200$ GeV, $b=10$ fm, $t=0.01$ fm/c



AuAu, $\sqrt{s_{NN}} = 200$ GeV, $b=10$ fm, $t=0.05$ fm/c



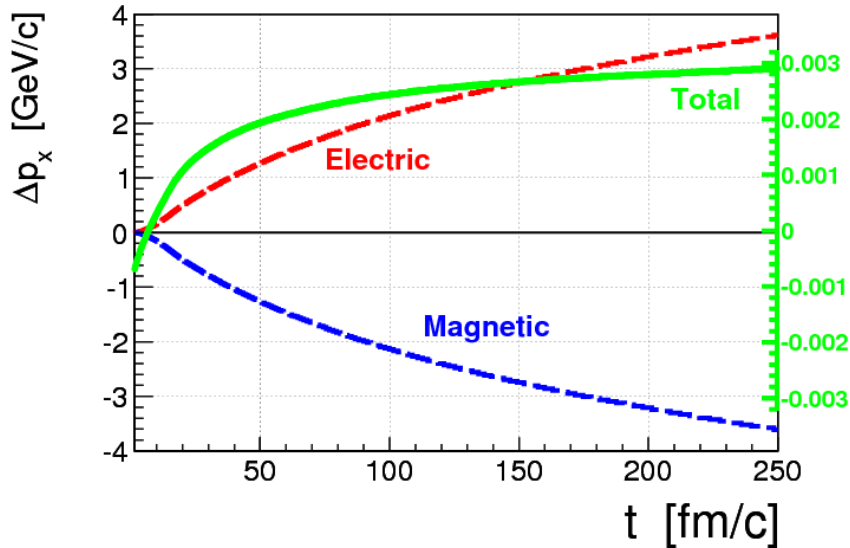
Observables



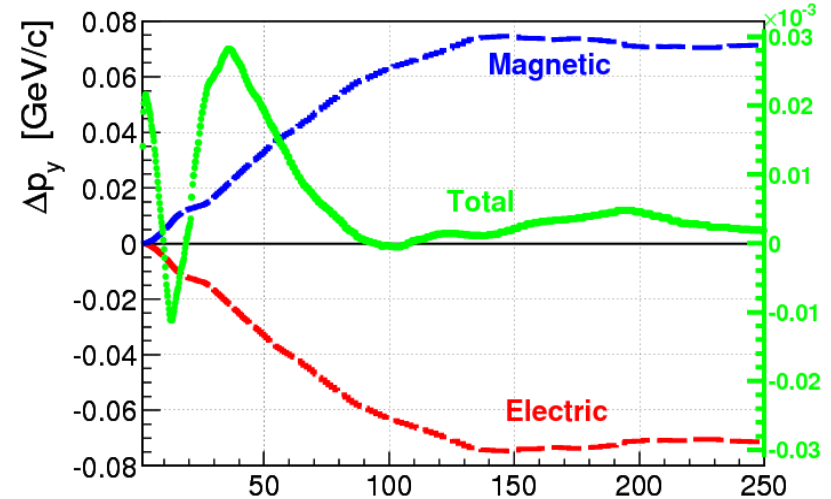
No electromagnetic field effects on observables !

Average momentum increment

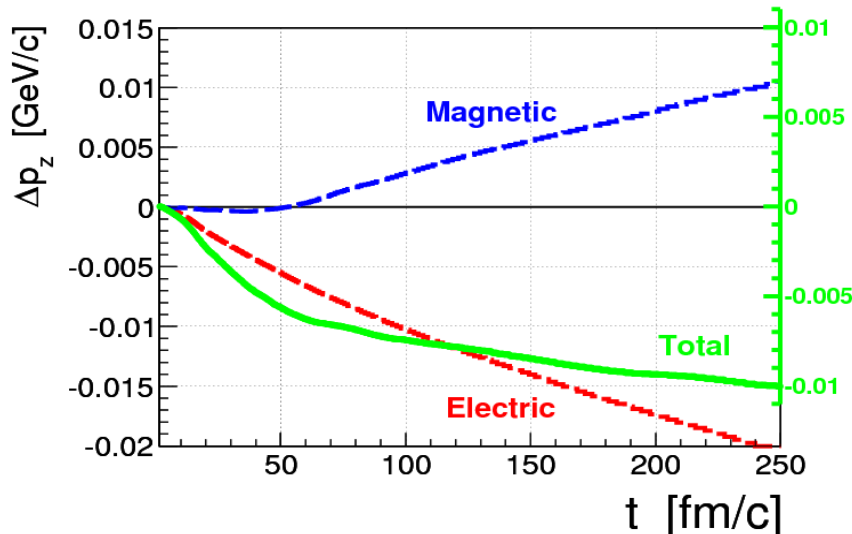
AuAu 200GeV, b=10fm



AuAu 200GeV, b=10fm



AuAu 200GeV, b=10fm



$$\dot{\vec{p}} \rightarrow e\vec{E} + e\vec{v} \times \vec{B}$$

$$\Delta\vec{p} = \sum_i \langle \delta\vec{p} \rangle_i \quad \text{for } p_z > 0$$

Transverse momentum increments Δp due to electric and magnetic fields compensate each other ! (worrying & hope)

Summary



- Some exp. data are not well reproduced in terms of the hadron-string picture => evidence for **nonhadronic degrees of freedom** => **PHSD** which provides a consistent description of off-shell parton dynamics in line with a lattice QCD equation of state
- PHSD provides a reasonable description of the **rapidity spectra** and **m_T -slopes** for A+A collisions at the SPS and RHIC energies
- The collective properties as expressed in terms of the **elliptic flow v_2** are reasonably described contrary to HSD calculations
- The **HSD transport model with retarded electromagnetic fields** has been developed. Actual calculations show no noticeable influence of the created electromagnetic fields on observables. It is due to a compensating effect between electric and magnetic fields
- Direct inclusion of quarks and gluons in evolution is needed (**PHSD model**)
- Experiments on the CME planned at RHIC by the low-energy scan program are of great interest since they hopefully will allow to infer the critical magnetic field eB_{crit} governing the spontaneous local CP violation

Thanks to

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Olena Linnyk

Vitalii Ozvenchuk



backup

Transport model with electromagnetic field

The Boltzmann equation is the basis of QMD like models:

$$\left\{ \frac{\partial}{\partial t} + \dot{\vec{r}} \cdot \vec{\nabla}_{\vec{r}} + \dot{\vec{p}} \cdot \vec{\nabla}_{\vec{p}} \right\} f(\vec{r}, \vec{p}, t) = I_{coll}(f, f_1, \dots, f_N)$$

Generalized on-shell transport equations in the presence of electromagnetic fields can be obtained formally by the substitution:

$$\begin{aligned} \dot{\vec{r}} &\rightarrow \frac{\vec{p}}{p_0} + \vec{\nabla}_{\vec{p}} U, \\ \dot{\vec{p}} &\rightarrow -\vec{\nabla}_{\vec{r}} U + e\vec{E} + e\vec{v} \times \vec{B} \end{aligned}$$

$$\begin{aligned} \left\{ \frac{\partial}{\partial t} + \left(\frac{\vec{p}}{p_0} + \vec{\nabla}_{\vec{p}} U \right) \vec{\nabla}_{\vec{r}} - \left(\vec{\nabla}_{\vec{r}} U - e\vec{E} - e\vec{v} \times \vec{B} \right) \vec{\nabla}_{\vec{p}} \right\} f(\vec{r}, \vec{p}, t) \\ = I_{coll}(f, f_1, \dots, f_N) \end{aligned} \quad U \sim \text{Re}(\Sigma^{ret})/2p_0$$

A general solution of the wave equations

$$\left\{ \begin{aligned} \vec{B} &= \vec{\nabla} \times \vec{A} \\ \vec{E} &= -\vec{\nabla} \Phi - \frac{\partial \vec{A}}{\partial t} \end{aligned} \right.$$

is as follows

$$\vec{A}(\vec{r}, t) = \frac{1}{4\pi} \int \frac{\vec{j}(\vec{r}', t') \delta(t - t' - |\vec{r} - \vec{r}'|/c)}{|\vec{r} - \vec{r}'|} d^3r' dt'$$

$$\Phi(\vec{r}, t) = \frac{1}{4\pi} \int \frac{\rho(\vec{r}', t') \delta(t - t' - |\vec{r} - \vec{r}'|/c)}{|\vec{r} - \vec{r}'|} d^3r' dt'$$

For point-like particles $\rho(\vec{r}, t) = e \delta(\vec{r} - \vec{r}(t)); \quad \vec{j}(\vec{r}, t) = e \vec{v}(t) \delta(\vec{r} - \vec{r}(t)) \quad \vec{\nabla} \times \vec{A} \rightarrow \text{LW eq.}$