

Impact Parameter Dependent nPDFs Based on EKS98 and EPS09 Parametrizations TORIC

Ilkka Helenius

In collaboration with
Kari J. Eskola, Heli Honkanen

University of Jyväskylä
Department of Physics

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Outline

- 1 Introduction and Motivation
 - Nuclear Geometry in A+A collisions
- 2 Model
 - Assumptions
 - Fitting Procedure
 - Outcome
- 3 Application
 - Central-to-Peripheral Ratio
- 4 Summary & Outlook

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Nuclear Thickness Function

Amount of nuclear matter in beam direction

Thickness function

Woods-Saxon density profile:

$$T_A(\mathbf{s}) = \int_{-\infty}^{\infty} dz \frac{n_0}{1 + \exp\left[\frac{\sqrt{\mathbf{s}^2 + z^2} - R_A}{d}\right]}$$

$$d = 0.54 \text{ fm}$$

$$R_A = 1.12A^{1/3} - 0.86A^{-1/3} \text{ fm}$$

$$n_0 = \frac{3}{4\pi} \frac{A}{R_A^3} \frac{1}{\left(1 + \left(\frac{\pi d}{R_A}\right)^2\right)}$$

$$A = \int d^2\mathbf{s} T_A(\mathbf{s})$$

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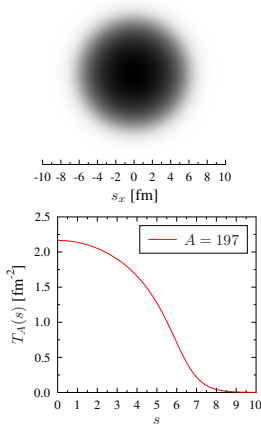
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A+A collision

From $T_A(\mathbf{s})$ we can construct the nuclear overlap function

$$T_{AB}(\mathbf{b}) = \int d^2\mathbf{s} T_A(\mathbf{s}) T_B(\mathbf{s} + \mathbf{b}),$$

Amount of the interacting matter at impact parameter \mathbf{b} .

In the absence of nuclear effects the hard cross section for given centrality class in $A + A$ collisions

$$d\sigma^{AB \rightarrow k+X} = \int_{b_1}^{b_2} d^2\mathbf{b} T_{AB}(\mathbf{b}) d\sigma^{NN}, \quad (1)$$

where

$$d\sigma^{NN} = \sum_{i,j} f_{i/N} \otimes f_{j/N} \otimes d\hat{\sigma}^{ij \rightarrow k+X}$$

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A+A collision

Nucleons in nucleus \neq free nucleons

\Rightarrow Modification to PDFs

$$f_{i/N}(x, Q^2) \rightarrow f_{i/A}(x, Q^2) = R_{i/A}(x, Q^2) \cdot f_{i/N}(x, Q^2)$$

Integrating the cross section formula (1) over the whole \mathbf{b} -space, we have

$$d\sigma^{AB \rightarrow k+X} = \sum_{i,j} \int d^2\mathbf{b} \int d^2\mathbf{s}_1 T_A(\mathbf{s}_1) R_{i/A}(x_1, Q^2) f_{i/N}(x_1, Q^2) \otimes \int d^2\mathbf{s}_2 T_B(\mathbf{s}_2) R_{j/B}(x_2, Q^2) f_{j/N}(x_2, Q^2) \otimes \delta(\mathbf{s}_2 - \mathbf{s}_1 - \mathbf{b}) d\hat{\sigma}^{ij \rightarrow k+X}$$

Now we want to replace

$$\int d^2\mathbf{s}_1 T_A(\mathbf{s}_1) R_{i/A}(x_1, Q^2) \rightarrow \int d^2\mathbf{s}_1 T_A(\mathbf{s}_1) r_{i/A}(x_1, Q^2, \mathbf{s}_1)$$

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Assumptions

Where $r_A(x, Q^2, \mathbf{s})$ is the nuclear modification depending on the transverse position of the nucleon \mathbf{s} .

For $r_A(x, Q^2, \mathbf{s})$ we now assume

- Normalization:

$$R_A(x, Q^2) = \frac{1}{A} \int d^2\mathbf{s} T_A(\mathbf{s}) r_A(x, Q^2, \mathbf{s}),$$

where $R_A(x, Q^2)$ from

- EKS98 [*Eur. Phys. J.*, C9:61-68, 1999]
- EPS09 [*JHEP*, 04:065, 2009]
- Spatial dependence through $T_A(\mathbf{s})$

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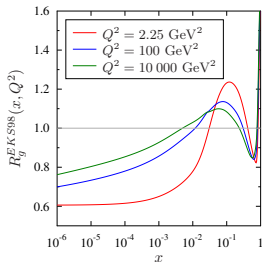
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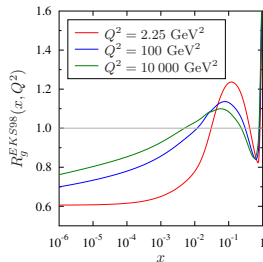
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Fitting Procedure

Functional form used for EKS98

$$r_A(x, Q^2, s) = 1 + c_1(x, Q^2)[T_A(s)] + c_2(x, Q^2)[T_A(s)]^2 \\ + c_3(x, Q^2)[T_A(s)]^3 + c_4(x, Q^2)[T_A(s)]^4$$

Important: no A dependence in fit parameters $c_i(x, Q^2)$!

Parameters $c_i(x, Q^2)$ obtained by minimizing the χ^2

$$\chi^2(x, Q^2) = \sum_A \left[\frac{R_A(x, Q^2) - \frac{1}{A} \int d^2s T_A(s) r_A(x, Q^2, s)}{R_A(x, Q^2) - 1} \right]^2,$$

where $A = 16, 20, 24, \dots, 300$.

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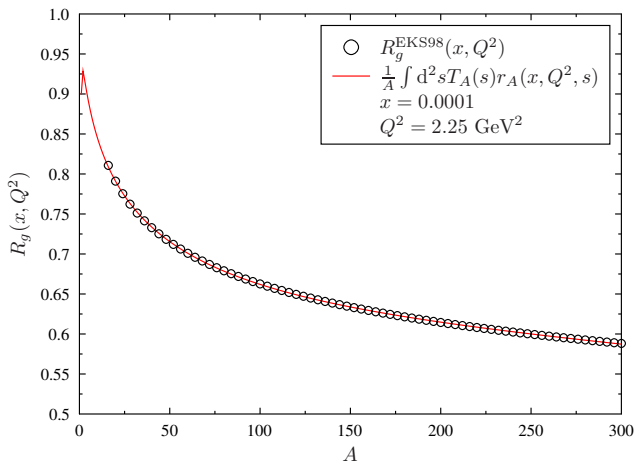
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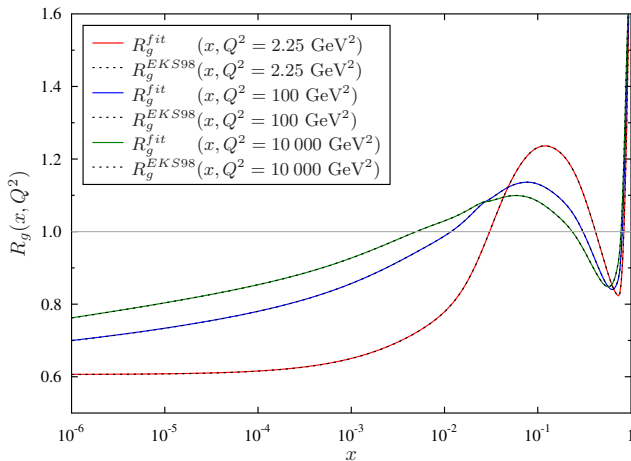
Fit Outcome

Example fit: gluon modification



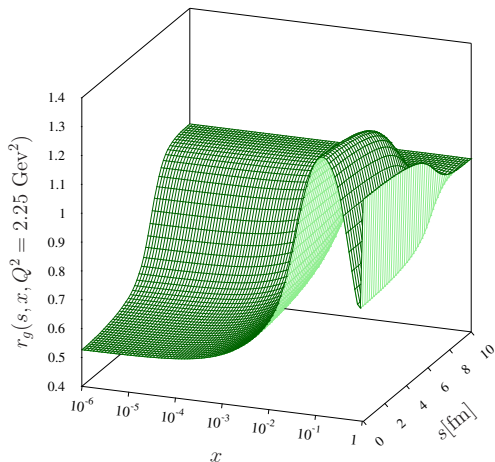
Fitted $R(x, Q^2)$ vs. old $R(x, Q^2)$

$$R^{fit}(x, Q^2) = \frac{1}{A} \int d^2\mathbf{s} T_A(s) \left[1 + \sum_{i=1}^4 c_i(x, Q^2) [T_A(s)]^i \right]$$



Spatial Dependence of Nuclear Modification for Au

$$r^{EKS98s}(x, Q^2, s) = 1 + \sum_{i=1}^4 c_i(x, Q^2) [T_A(s)]^i$$



Observations

- The shape in x is similar to EKS98
- Effects are slightly stronger in small s compared to EKS98
- Nuclear effects die out when $s > R_A$

Nuclear Modification Factor R_{AA}^{1-jet}

- The 1-jet distribution for a centrality class with $b \in [b_1, b_2]$ can be calculated from

$$\left\langle \frac{d^2 N_{AA}^{1-jet}}{dp_T dy} \right\rangle_{b_1, b_2} = \frac{\int_{b_1}^{b_2} d^2 b \frac{d^2 N_{AA}^{1-jet}(b)}{dp_T dy}}{\int_{b_1}^{b_2} d^2 b p_{AA}^{inel}(b)}$$

- $p_{AA}^{inel}(b) = 1 - e^{-T_{AA}(b)\sigma_{inel}^{NN}}$ (optical Glauber model)
- Comparison with proton-proton (pp) collision gives the R_{AA}^{1-jet} for each centrality class

$$R_{AA}^{1-jet} = \frac{\left\langle \frac{d^2 N_{AA}^{1-jet}}{dp_T dy} \right\rangle_{b_1, b_2}}{\frac{\langle N_{bin} \rangle_{b_1, b_2}}{\sigma_{inel}^{NN}} \frac{d^2 \sigma_{pp}^{1-jet}}{dp_T dy}}$$

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Central-to-Peripheral Ratio R_{CP}^{1-jet}

The 1-jet central-to-peripheral ratio

$$R_{CP}^{1-jet} = \frac{\left\langle \frac{d^2 N_{AA}^{1-jet}}{dp_T dy} \right\rangle \frac{1}{\langle N_{bin} \rangle} (central)}{\left\langle \frac{d^2 N_{AA}^{1-jet}}{dp_T dy} \right\rangle \frac{1}{\langle N_{bin} \rangle} (peripheral)}$$

Parameters from Optical Glauber Model

	<i>central</i> = 0 – 5%			<i>peripheral</i> = 60 – 80%		
	b_1 [fm]	b_2 [fm]	$\langle N_{bin} \rangle$	b_1 [fm]	b_2 [fm]	$\langle N_{bin} \rangle$
RHIC	0.0	3.41	1079	11.82	13.65	12.06
LHC	0.0	3.47	1717	12.03	13.89	18.81

- RHIC: $\sigma_{inel}^{NN} = 42$ mb
- LHC: $\sigma_{inel}^{NN} = 62$ mb

Central-to-Peripheral Ratio R_{CP}^{1-jet}

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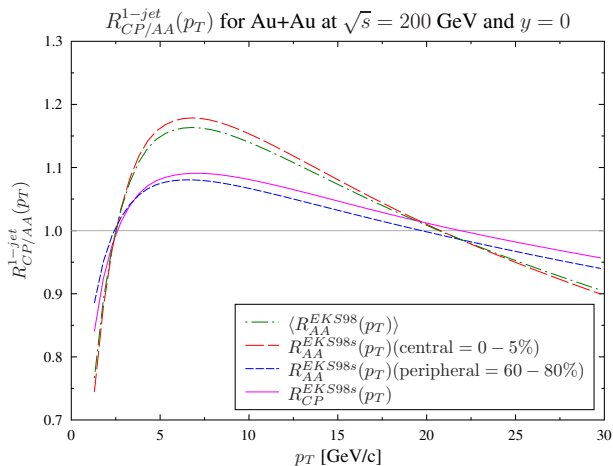
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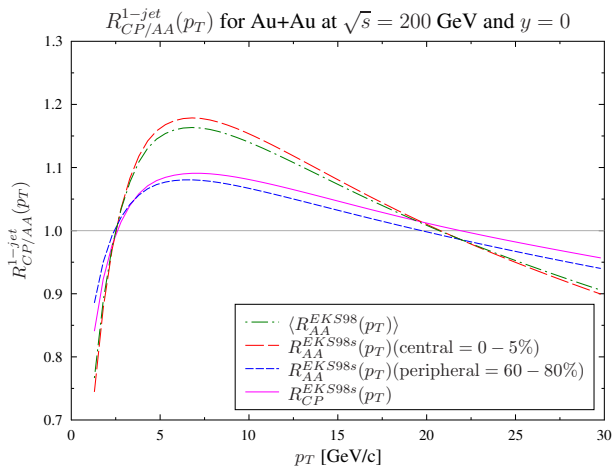
Observations

$$R_{AA}(\text{central}) \approx \langle R_{AA} \rangle$$

$$R_{AA}(\text{peripheral}) \neq 1$$

$$R_{CP} \neq \langle R_{AA} \rangle$$

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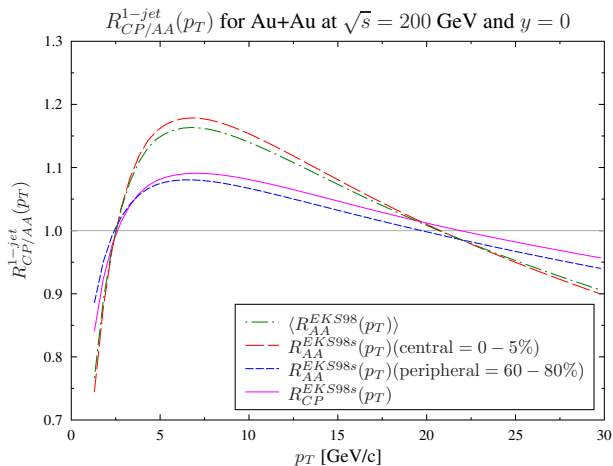
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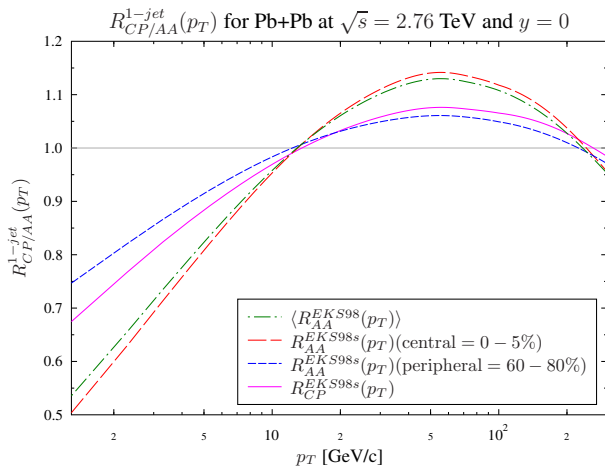
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We have

- Developed a model for spatial dependence of nuclear modification based on
 - the A dependence of the EKS98/EP09 parametrizations (= data!)
 - the nuclear thickness function $T_A(s)$
- Program to calculate $r_A(x, Q^2, s)$
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- ⇒ Nuclear modifications of any hard process in different centrality classes can be computed!

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Backup

Centrality Classes for RHIC

