"THERMAL" PHOTONS FROM IDEAL AND VISCOUS MEDIA

OUTLINE • Sources & EM emissivity

•Modelling the evolving system:

3D hydro

- 3D viscous hydro
- Fluctuating initial states
- •Are photons sensitive to all of the above?
- If so, can we quantify this?



INFO CARRIED BY THE RADIATION

$$dR = -\frac{g^{\mu\nu}}{2\omega} \frac{d^3k}{(2\pi)^3} \frac{1}{Z} \sum_{i} e^{-\beta K_i} \sum_{f} (2\pi)^4 \delta(p_i - p_f - k)$$

$$\times \langle j | J_{\mu} | i \rangle \langle i | J_{\nu} | j \rangle$$

Thermal ensemble average of the current-current correlator

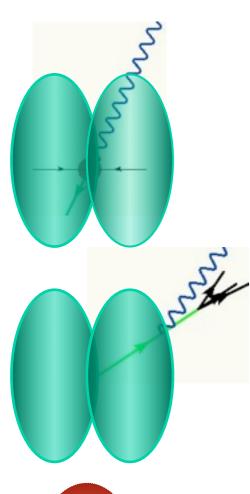
Emission rates:

$$\omega \frac{d^{3}R}{d^{3}k} = -\frac{g^{\mu\nu}}{(2\pi)^{3}} \operatorname{Im}\Pi_{\mu\nu}(\omega,k) \frac{1}{e^{\beta\omega}-1} \qquad \text{(photons)}$$

$$E_{+}E_{-}\frac{d^{6}R}{d^{3}p_{+}d^{3}p_{-}} = \frac{2e^{2}}{(2\pi)^{6}} \frac{1}{k^{4}} L^{\mu\nu} \operatorname{Im}\Pi_{\mu\nu}^{R}(\omega,k) \frac{1}{e^{\beta\omega}-1} \qquad \text{(dileptons)}$$

$$\operatorname{McLerran, Toimela (85), Weldon (90), Gale, Kapusta (91) \qquad 2}$$

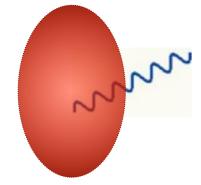




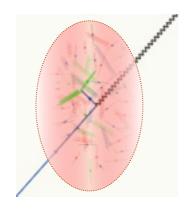
Sources of photons:

Hard direct photons. pQCD with shadowing Non-thermal

Fragmentation photons. pQCD with shadowing Non-thermal







Jet-plasma photons Thermal

Jet in-medium bremsstrahlung Thermal



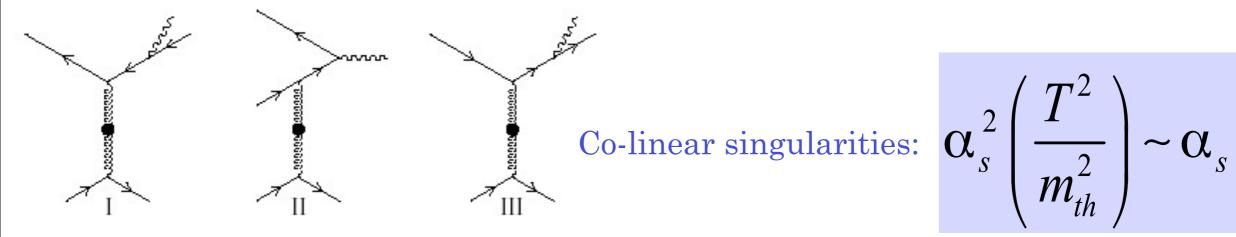


Thermal Photons from hot QCD: HTL program (Klimov (1981), Weldon (1982), Braaten & Pisarski (1990); Frenkel & Taylor (1990))

$$\mathrm{Im}\,\Pi^{\mu}_{R\mu}\sim\ln\left(\frac{\varpi T}{\left(m_{th}\left(\sim gT\right)\right)^{2}}\right)$$

Kapusta, Lichard, Seibert (1991) Baier, Nakkagawa, Niegawa, Redlich (1992)

Going to two loops: Aurenche, Kobes, Gelis, Petitgirard (1996) Aurenche, Gelis, Kobes, Zaraket (1998)



2001: Results complete at $O(\alpha_{c})$

Arnold, Moore, and Yaffe JHEP 12, 009 (2001); JHEP 11, 057 (2001) Incorporate LPM; Inclusive treatment of collinear enhancement, photon and gluon emission





ELECTROMAGNETIC RADIATION FROM HADRONS

Chiral, Massive Yang-Mills: O. Kaymakcalan, S. Rajeev, J. Schechter, PRD 30, 594 (1984)

$$\mathcal{L} = \frac{1}{8} F_{\pi}^{2} \operatorname{Tr} D_{\mu} U D^{\mu} U^{\dagger} + \frac{1}{8} F_{\pi}^{2} \operatorname{Tr} M \left(U + U^{\dagger} \right)$$
$$- \frac{1}{2} \operatorname{Tr} \left(F_{\mu\nu}^{L} F^{L\mu\nu} + F_{\mu\nu}^{R} F^{R\mu\nu} \right) + m_{0}^{2} \operatorname{Tr} \left(A_{\mu}^{L} A^{L\mu} + A_{\mu}^{R} A^{R\mu} \right)$$

+ non-minimal terms

Parameters and form factors are constrained by hadronic phenomenology:

•Masses & strong decay widths

•Electromagnetic decay widths

•Other hadronic observables:

•e.g. $a_1 \rightarrow \pi \rho \quad D/S$

(See also, Lichard and Vojik, arXiv:1006.2919)

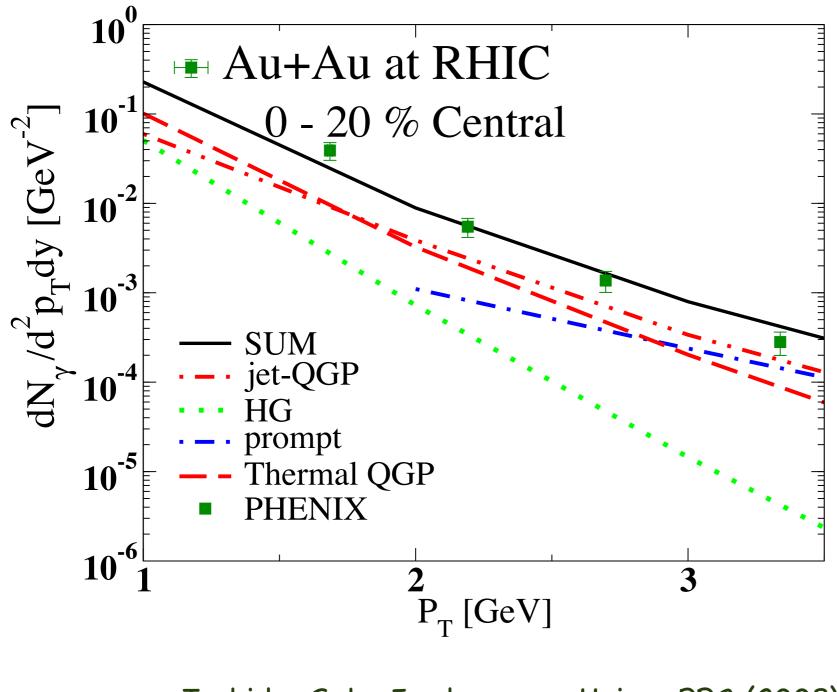


EM emissivities computed: Turbide, Rapp, Gale, PRC (2004); Turbide, McGill PhD (2006)



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APPLYING THIS TO THE SOFT SECTOR @ RHIC



• At low p_T, spectrum dominated by thermal components (HG, QGP)

- At high p_T, spectrum dominated by pQCD
- Window for jet-QPG contributions at mid-

 p_{T}

Turbide, Gale, Frodermann, Heinz, PRC (2008); Higher p_T : G. Qin et al., PRC (2009)

(e)

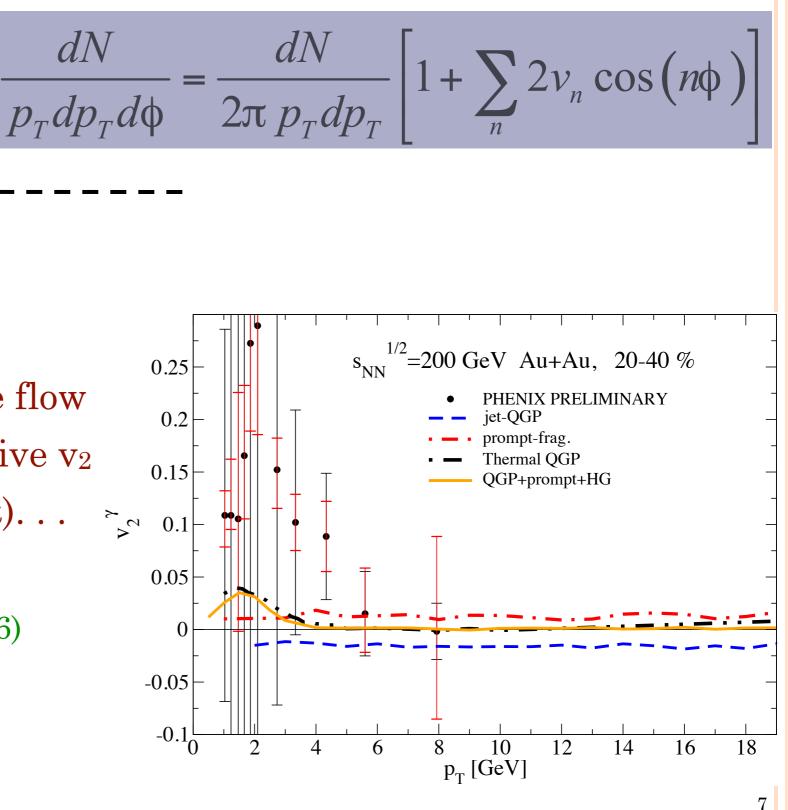


BEYOND ONE-BODY DATA: FLOW AND CORRELATIONS



<u>Details will matter</u>: flow, T(t). . . $\succ_{S^{n}}$

Turbide, Gale, Fries PRL (2006) Low p_T : Chatterjee *et al.*, PRL (2006) All p_T : Turbide *et al.*, PRC (2008)

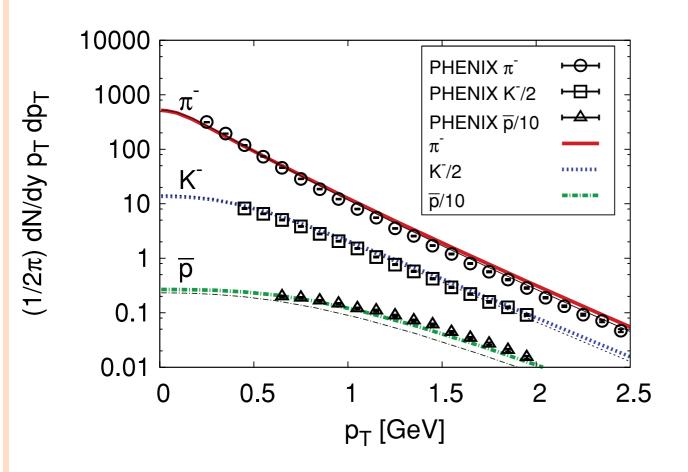




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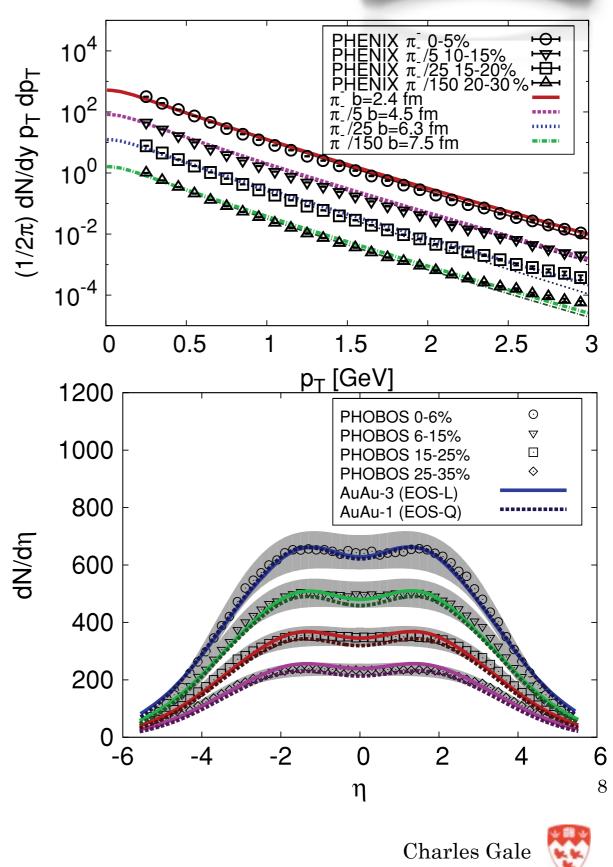
3D RELATIVISTIC HYDRODYNAMICS:





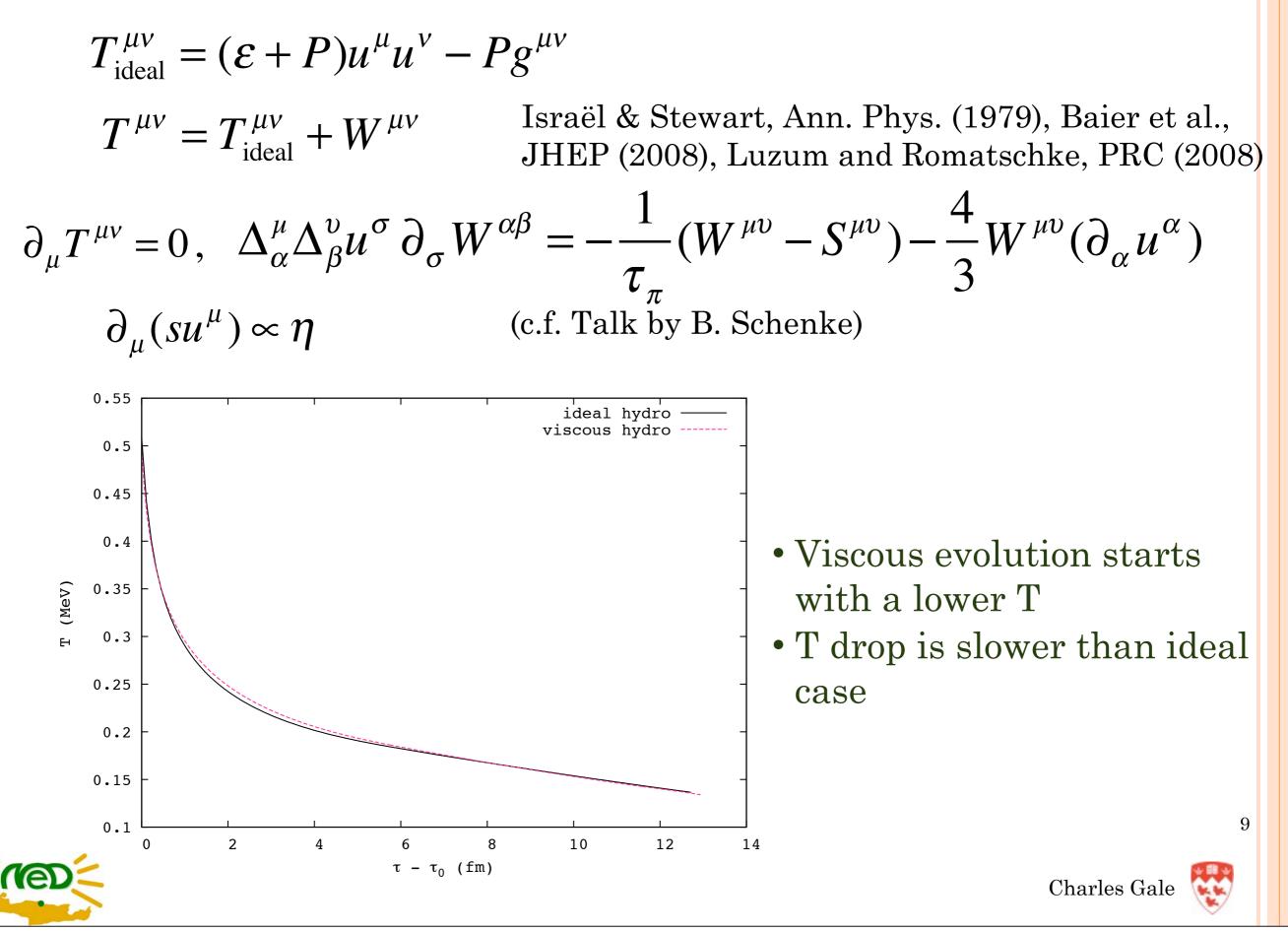
•MUSIC: 3D relativistic hydro

- Ideal: Schenke, Jeon, and Gale, PRC (2010)
- FIC and Viscous: Schenke, Jeon, Gale, PRL (2011)





THE EFFECTS OF SHEAR VISCOSITY ON BULK DYNAMICS



THE EFFECTS OF SHEAR VISCOSITY ON BULK DYNAMICS

$$T_{\text{ideal}}^{\mu\nu} = (\varepsilon + P)u^{\mu}u^{\nu} - Pg^{\mu\nu}$$

$$T^{\mu\nu} = T_{\text{ideal}}^{\mu\nu} + W^{\mu\nu} \qquad \text{Israël & Stewart, Ann. Phys. (1979), Baier et al., JHEP (2008), Luzum and Romatschke, PRC (2008)}$$

$$\partial_{\mu}T^{\mu\nu} = 0, \quad \Delta^{\mu}_{\alpha}\Delta^{\nu}_{\beta}u^{\sigma}\partial_{\sigma}W^{\alpha\beta} = -\frac{1}{\tau_{\pi}}(W^{\mu\nu} - S^{\mu\nu}) - \frac{4}{3}W^{\mu\nu}(\partial_{\alpha}u^{\alpha})$$

$$\partial_{\mu}(su^{\mu}) \propto \eta \qquad (c.f. Talk by B. Schenke)$$

THE EFFECTS OF SHEAR VISCOSITY ON THE PHOTON DISTRIBUTION

In-medium **hadrons**:

$$f_{0}(u^{\mu}p_{\mu}) = \frac{1}{(2\pi)^{3}} \frac{1}{\exp[(u^{\mu}p_{\mu} - \mu)/T] \pm 1}$$

$$f \to f_{0} + \delta f, \quad \delta f = f_{0}(1 \pm f_{0})p^{\alpha}p^{\beta}W_{\alpha\beta}\frac{1}{2(\varepsilon + P)T^{2}}$$

$$\frac{d^{3}R}{d^{3}q} = \int \frac{d^{3}p_{1}}{2(2\pi)^{3}E_{1}} \frac{d^{3}p_{2}}{2(2\pi)^{3}E_{2}} \frac{d^{3}p_{3}}{2(2\pi)^{3}E_{3}}(2\pi)^{4}|M|^{2}\delta^{4}(...)\frac{f(E_{1})f(E_{2})[1 \pm f(E_{3})]}{2(2\pi)^{3}}$$

One considers all the reaction and radiative decay channels of external state combinations of:

 $\{\pi, K, \rho, K^*, a_1\}$ With hadronic form factors

+ QGP Photons



 q_0



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THE EFFECTS OF SHEAR VISCOSITY ON THE PHOTON DISTRIBUTION

In-medium **hadrons**:

$$\begin{split} f_0(u^{\mu}p_{\mu}) &= \frac{1}{(2\pi)^3} \frac{1}{\exp[(u^{\mu}p_{\mu} - \mu)/T] + 1} & \text{Moore, Teaney} \\ FRC (2010) \\ f \to f_0 + \delta f, \quad \delta f &= f_0(1 \pm f_0)p^{\alpha}p^{\beta}W_{\alpha\beta}\frac{1}{2(\varepsilon + P)T^2} \\ q_0 \frac{d^3R}{d^3q} &= \int \frac{d^3p_1}{2(2\pi)^3E_1} \frac{d^3p_2}{2(2\pi)^3E_2} \frac{d^3p_3}{2(2\pi)^3E_3} (2\pi)^4 |M|^2 \delta^4(...) \frac{f(E_1)f(E_2)[1 \pm f(E_3)]}{2(2\pi)^3} \end{split}$$

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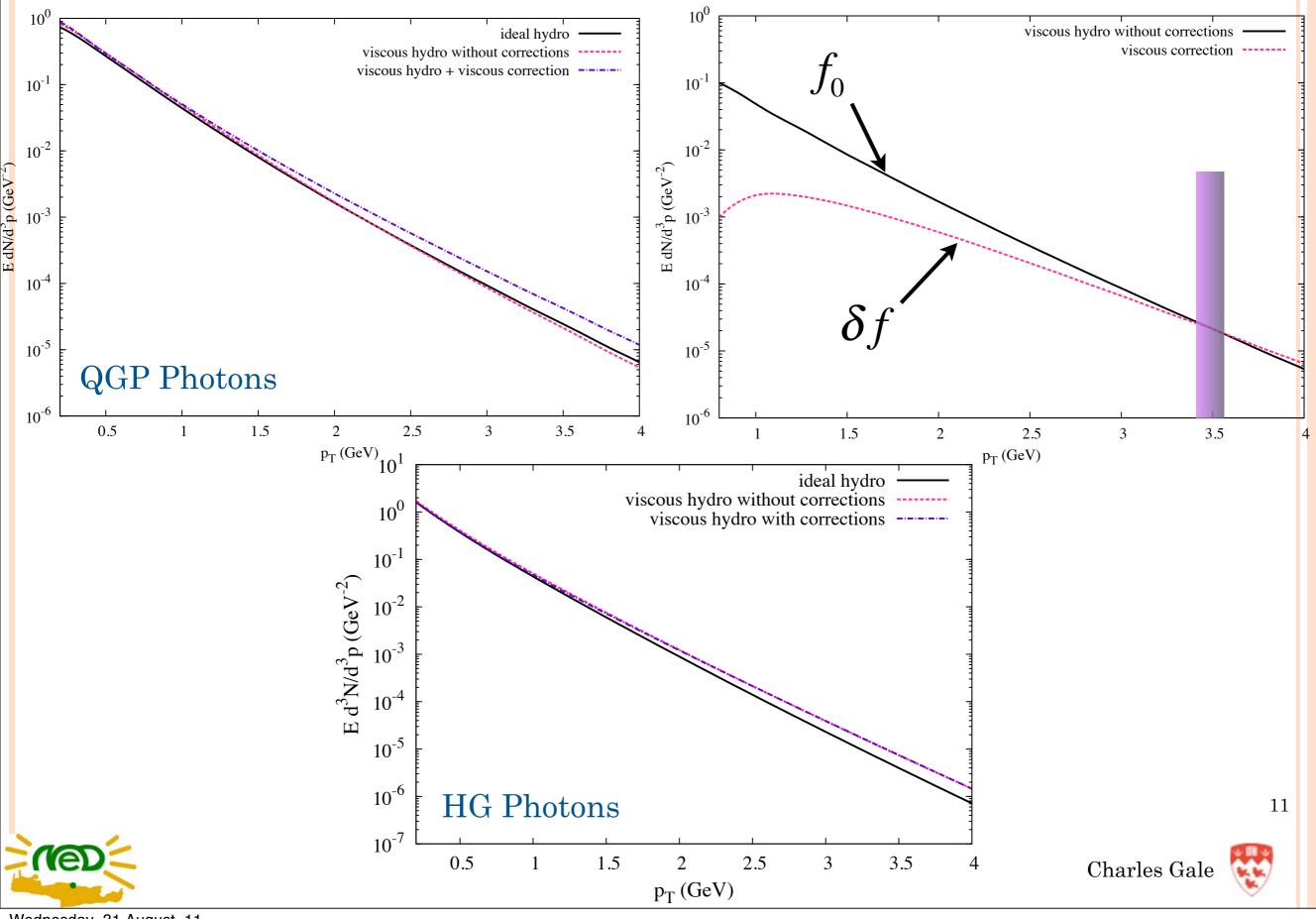
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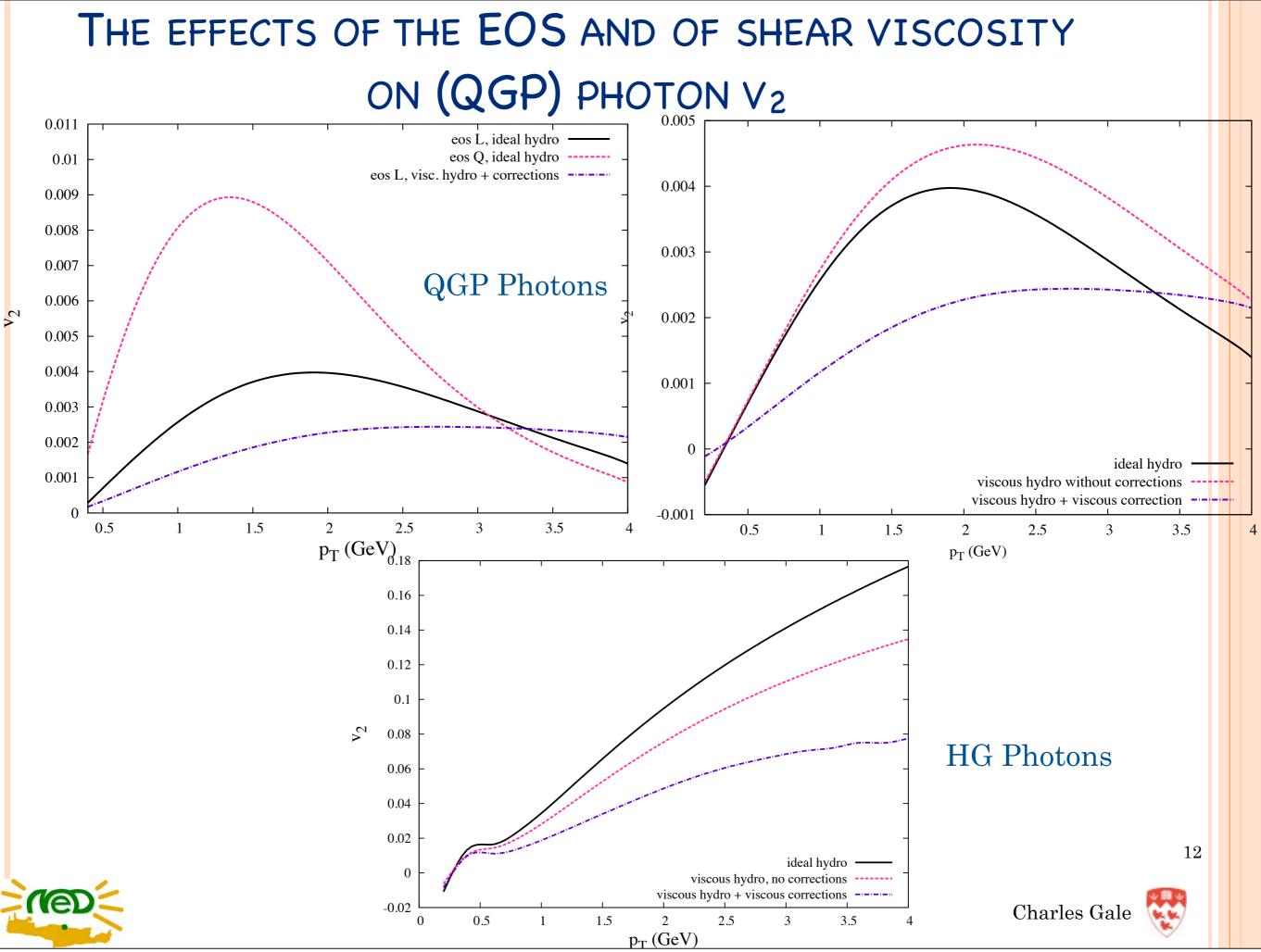


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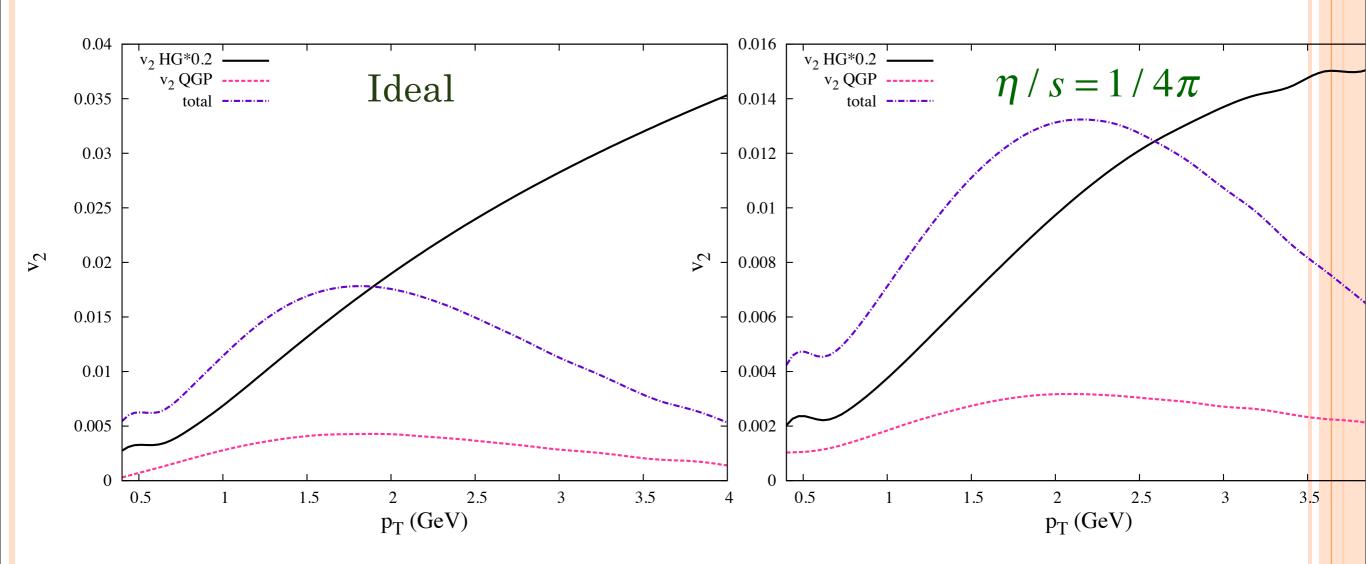
Ansatz: Dusling,

THE EFFECTS OF SHEAR VISCOSITY ON THE PHOTON DISTRIBUTION





Photons from the QGP + those from the HADRONIC GAS: Net V_2

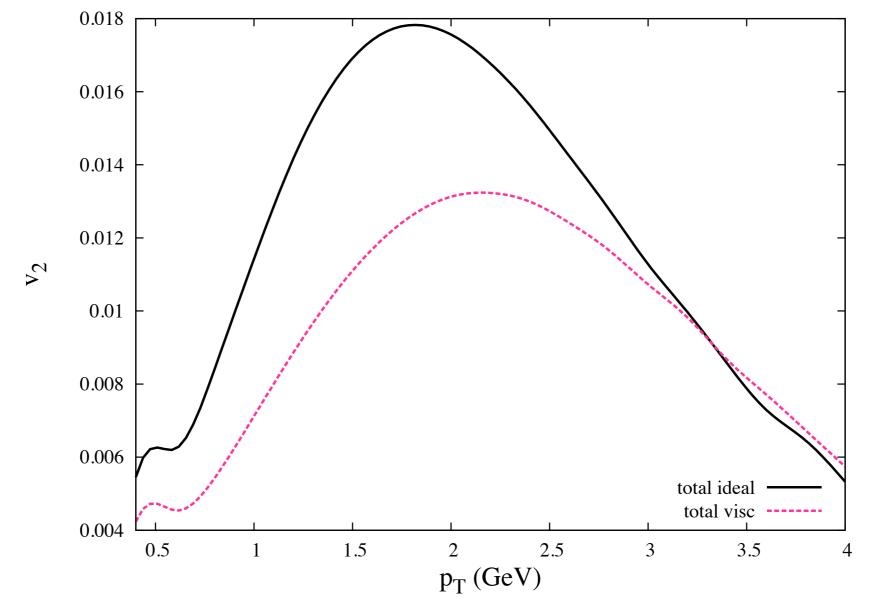




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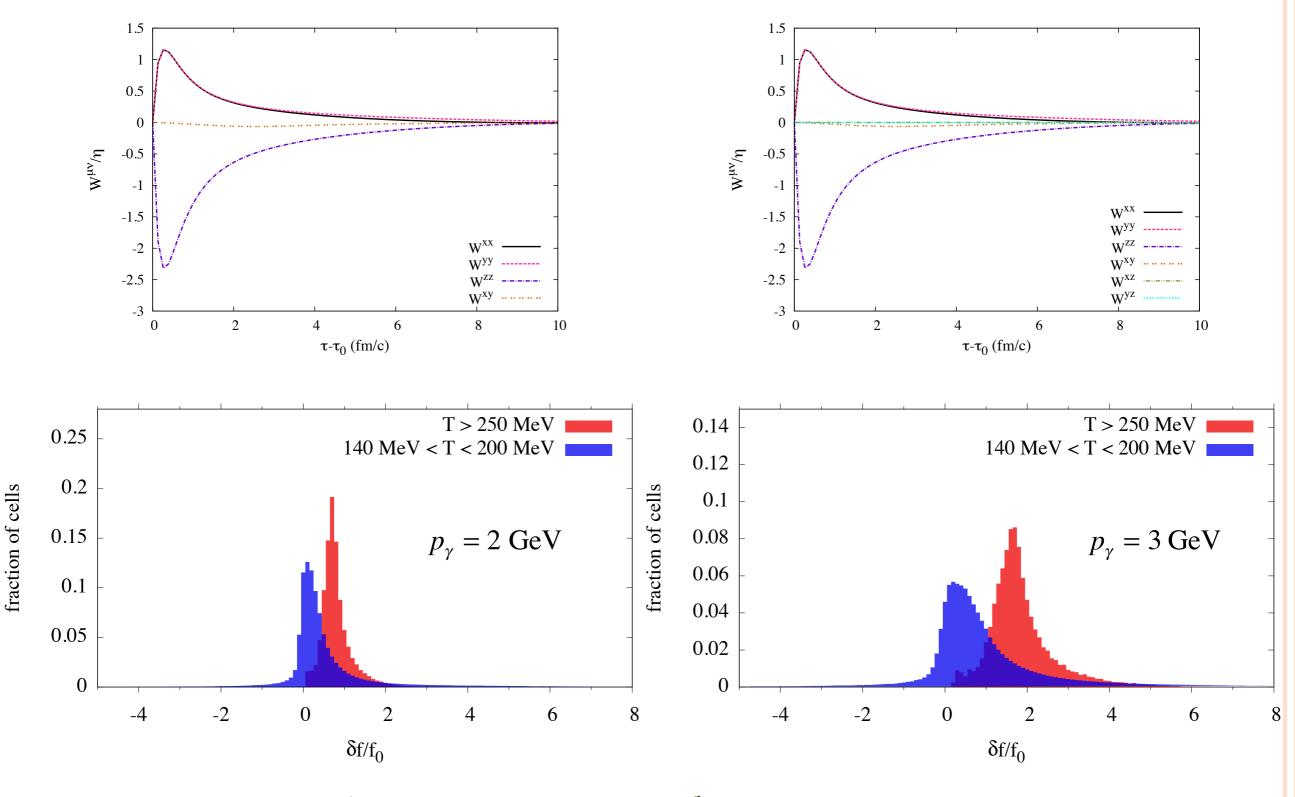
PHOTONS FROM THE QGP + THOSE FROM THE HADRONIC GAS: NET V₂



•Viscous corrections reduce v_2 (as for hadrons). At ~2 GeV ~40% reduction.



NON-EQUILIBRIUM EFFECTS

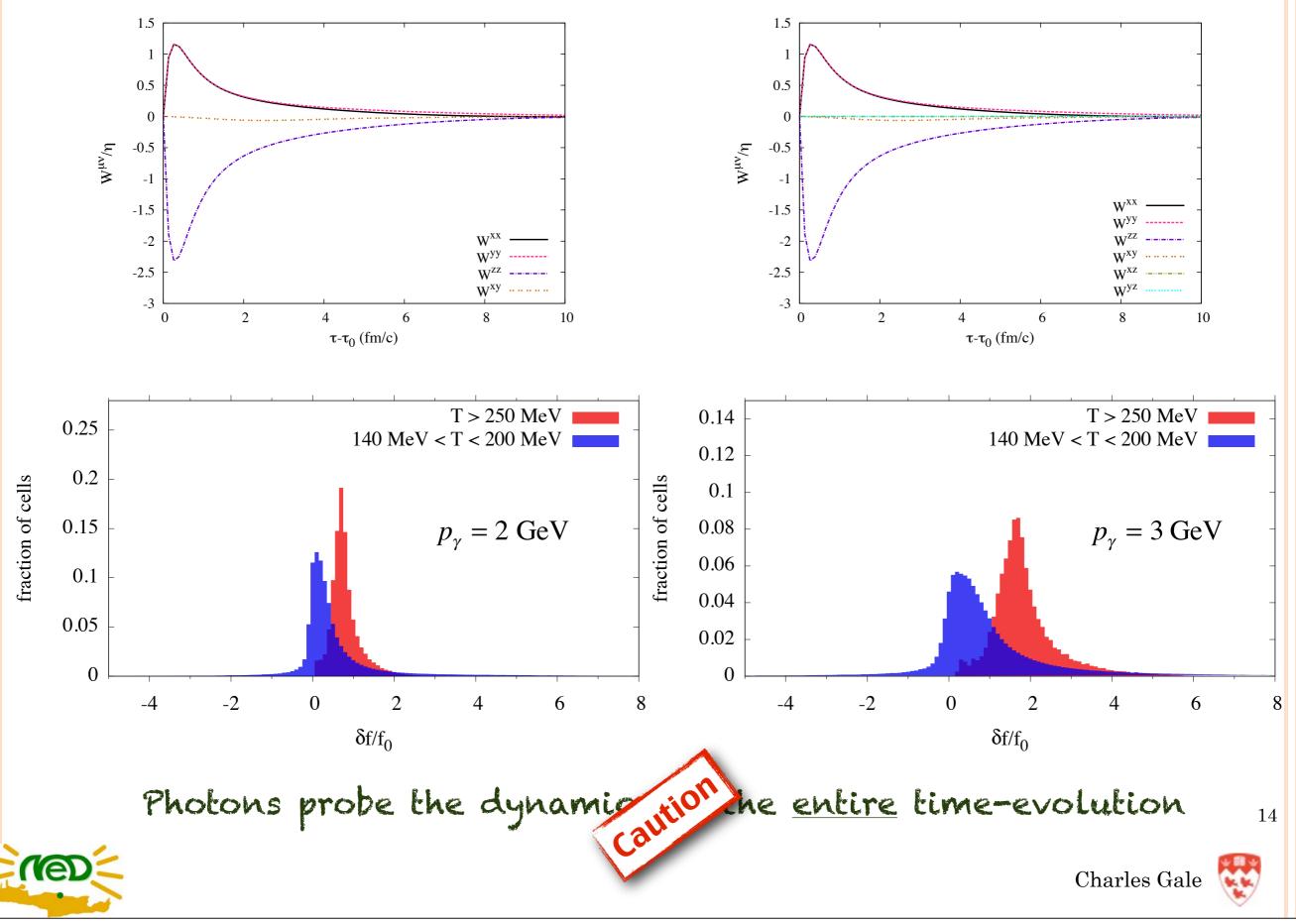


Photons probe the dynamics of the <u>entire</u> time-evolution

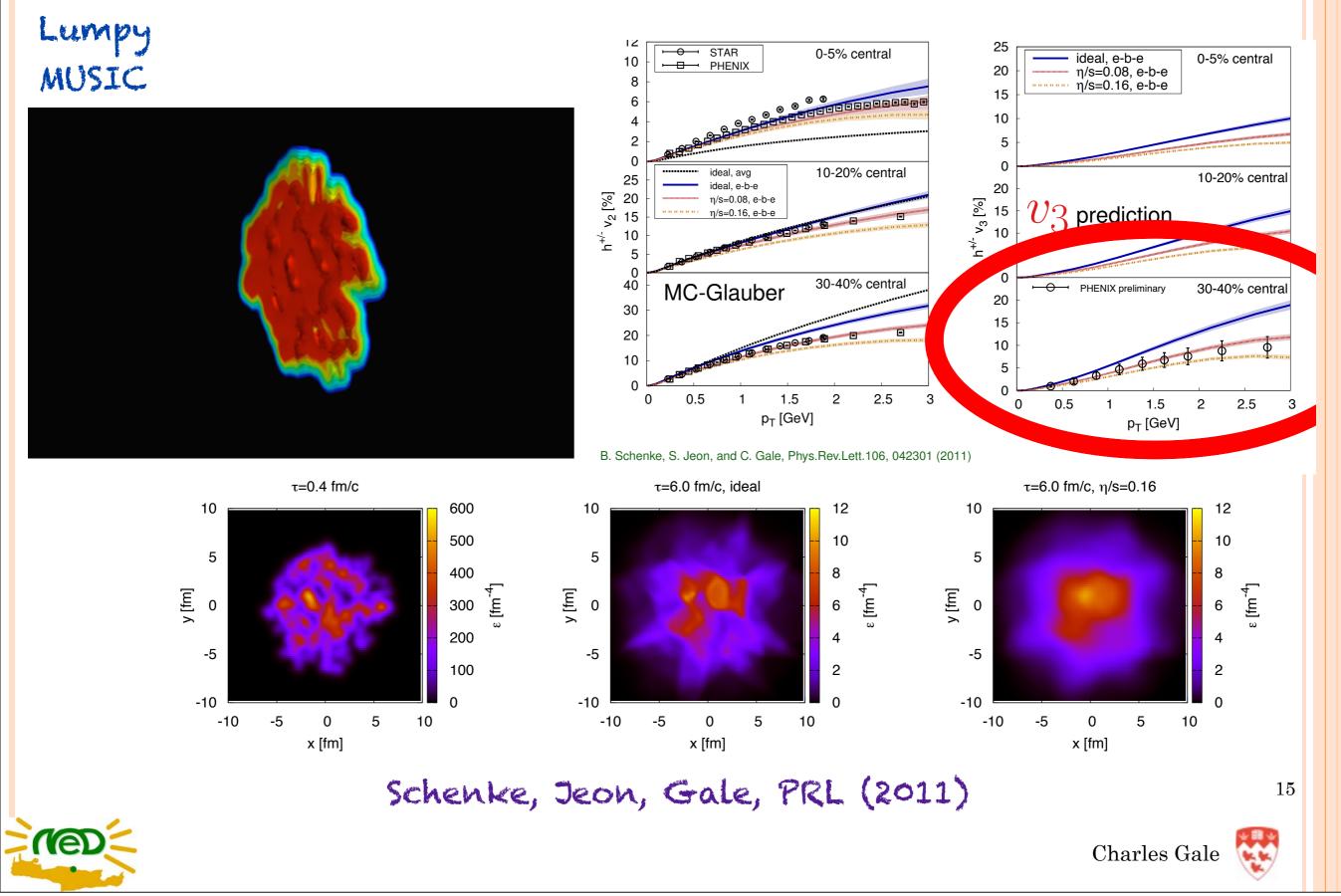




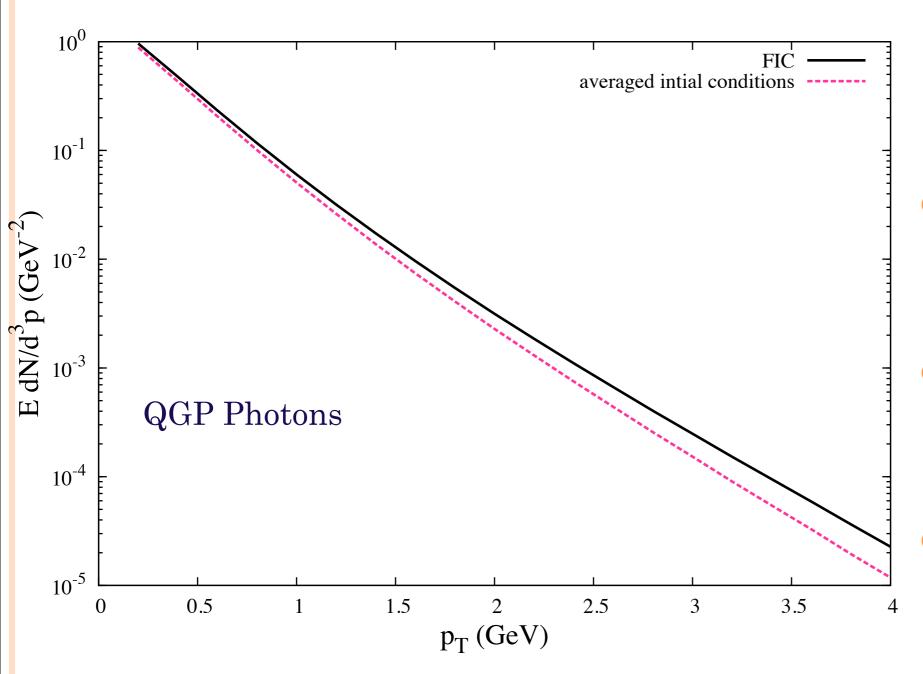
NON-EQUILIBRIUM EFFECTS



INITIAL STATE FLUCTUATIONS: A PARADIGM SHIFT IN HEAVY ION ANALYSES



THE EFFECT ON THE THERMAL PHOTON SPECTRUM



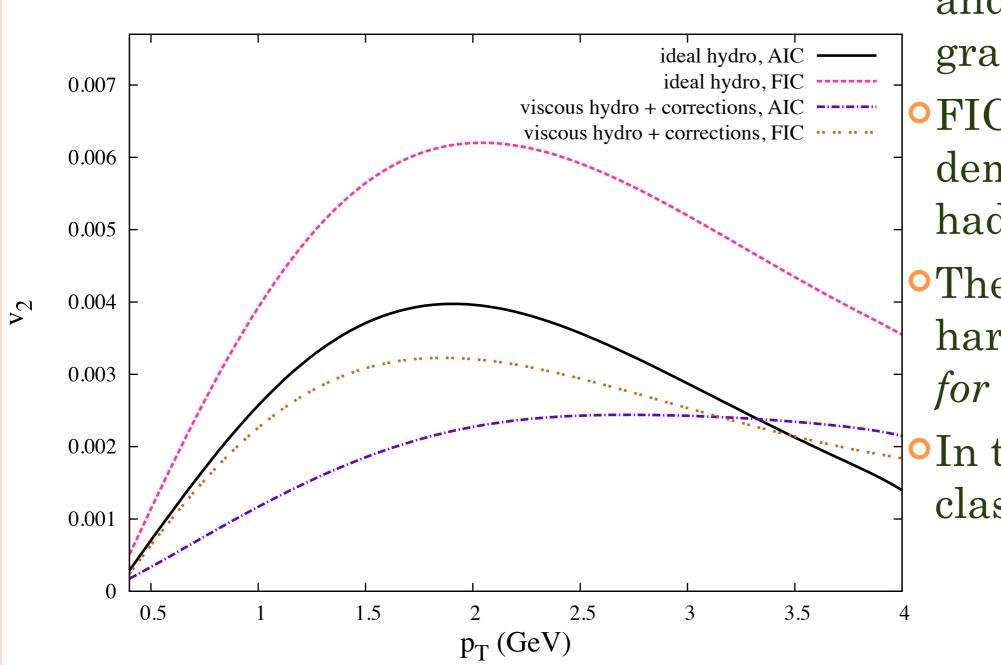
 FIC produces higher initial T (hot spots), and higher initial gradients
 FIC conditions are demanded by hadronic data (vodd)

•These lead to a harder spectrum, *as for hadrons*

•In this centrality class, v₂ goes up



THE EFFECT ON THE THERMAL PHOTON SPECTRUM



•FIC produces higher initial T (hot spots), and higher initial gradients •FIC conditions are demanded by hadronic data (v_{odd}) •These lead to a harder spectrum, as for hadrons •In this centrality class, v₂ goes up



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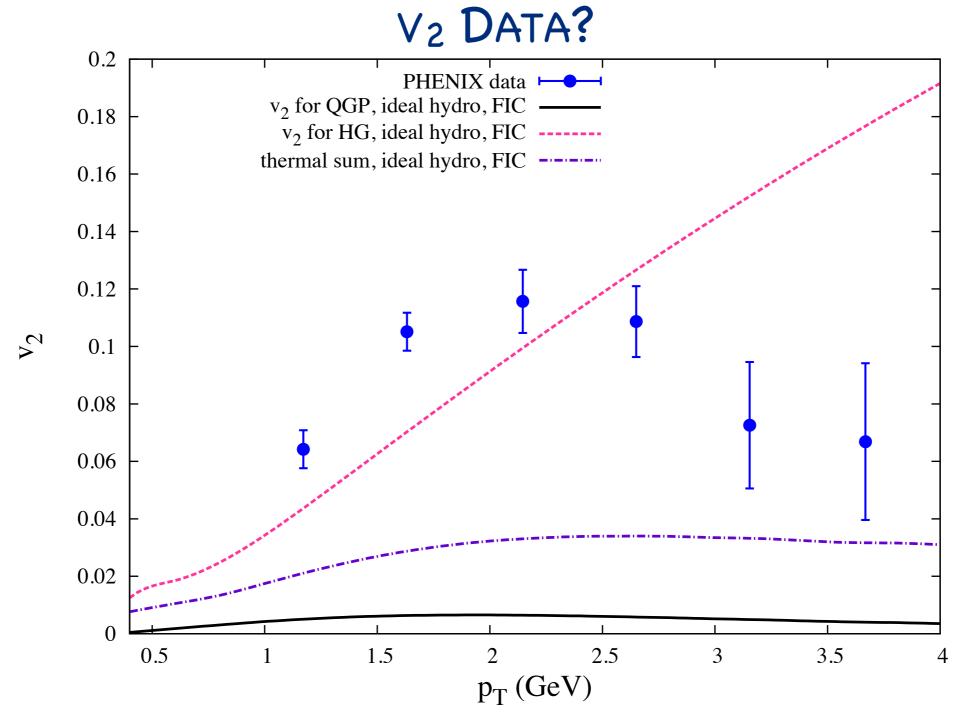
QUANTIFYING THE EFFECT ON THE THERMAL PHOTON SPECTRUM

P _T (GeV)	Viscosity	FIC	Viscosity + FIC
1	18%	18%	41%
2	30%	45%	82%
3	30%	77%	126%



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•New data is higher than calculation, even with e-b-e initial state fluctuations, and ideal hydro



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Work done with:



Maxime Dion
Jean-François Paquet
Björn Schenke
Clint Young
Sangyong Jeon

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19

 Bad news: Photon v₂ is very sensitive to the EOS, and to various hydro parameters such as viscosity, and initial state fluctuations

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19

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19

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- Good news: Photon v₂ is very sensitive to the EOS, and to various hydro parameters such as viscosity, and initial state fluctuations
- Photons are penetrating probes: at high p_T, they will reveal jet-medium interactions, at low p_T, they will reveal details of out-of-equilibrium dynamics

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- Photons are penetrating probes: at high $p_{\rm T},$ they will reveal jet-medium interactions, at low $p_{\rm T},$ they will reveal details of out-of-equilibrium dynamics
- New data is intriguing. . .

Work done with:

MED

- Maxime Dion
- Jean-François Paquet
- Björn Schenke
- Clint Young
- Sangyong Jeon



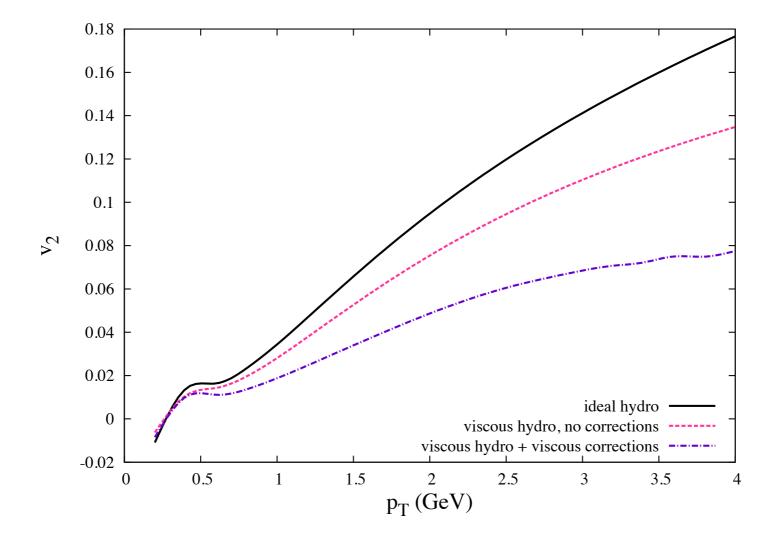
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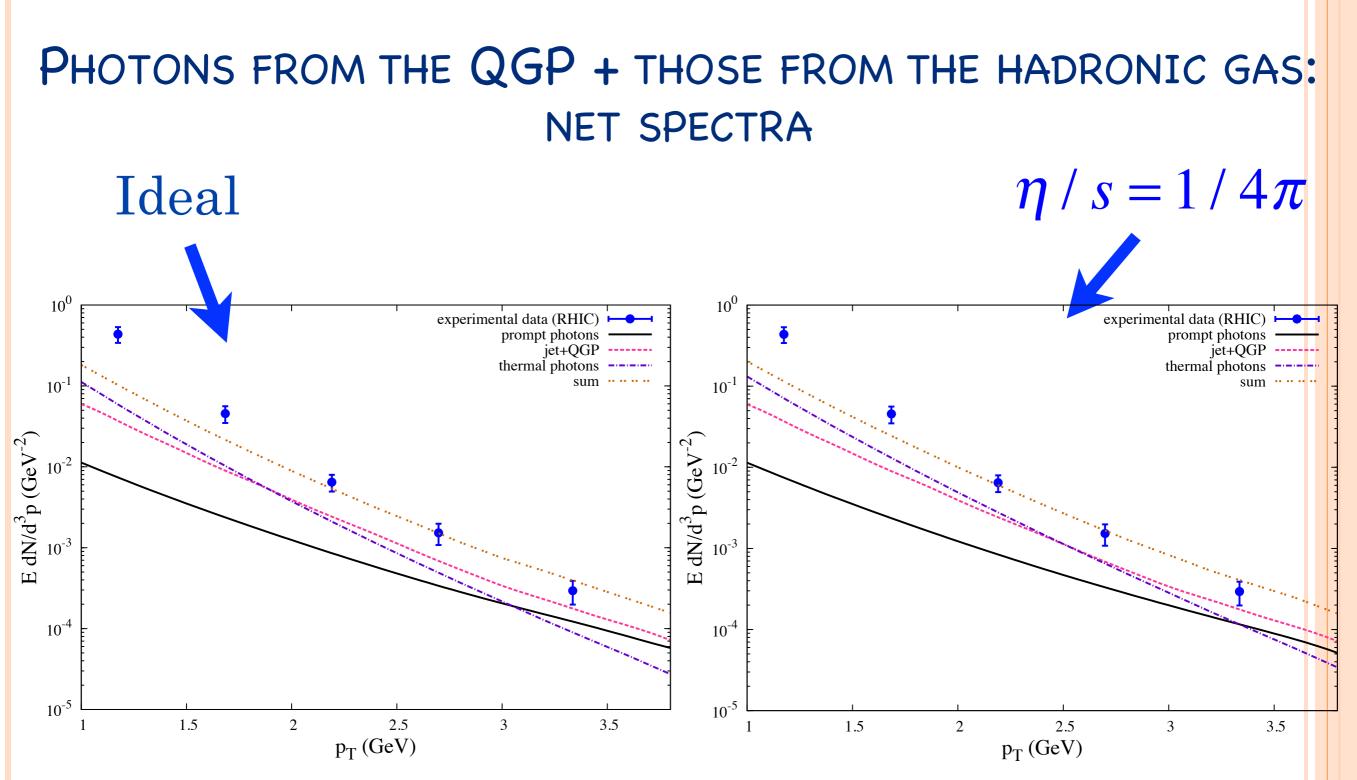


Viscous effects on HG anisotropy





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It is important to compare situations that yield *the same* (hadronic) final state: no contradictions between photon and hadron analyses
 Viscous effects on the net photon yield are small (a few %)
 ²² Charles Gale Charles Gale