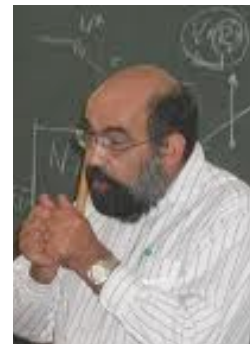




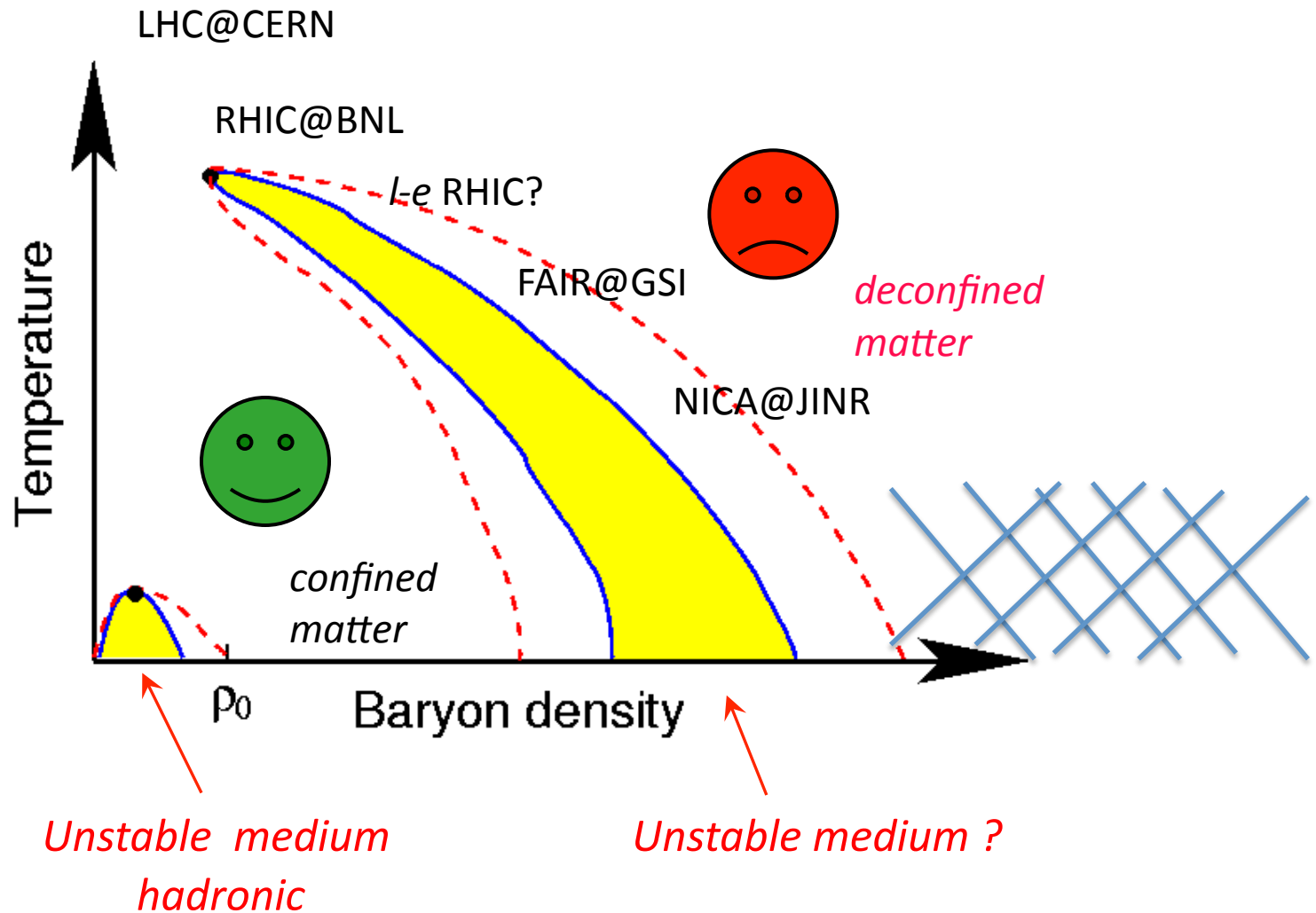
*Dynamics of density instabilities  
associated with the first-order  
confinement phase transition*



*Jørgen Randrup, LBNL, Berkeley*



*Schematic and simplified  
phase diagram of strongly interacting matter*

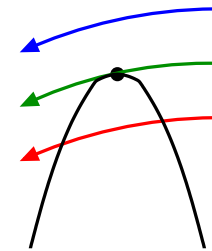
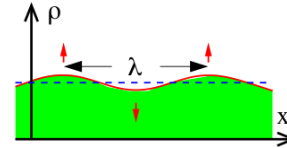


# Dynamics of collective modes in evolving many-body systems

Amplitude evolution:

$$\frac{d}{dt}A_\nu(t) = -i\omega_\nu A_\nu(t) + B_\nu(t)$$

$\omega_\nu = \epsilon_\nu + i\gamma_\nu$



Markovian noise:

$$\langle B_\nu(t) B_\mu(t')^* \rangle = 2\mathcal{D}_{\nu\mu} \delta(t - t')$$

Correlation function:

$$\sigma_{\nu\mu}(t_1, t_2) \equiv \langle A_\nu(t_1) A_\mu(t_2)^* \rangle$$

Evolution:

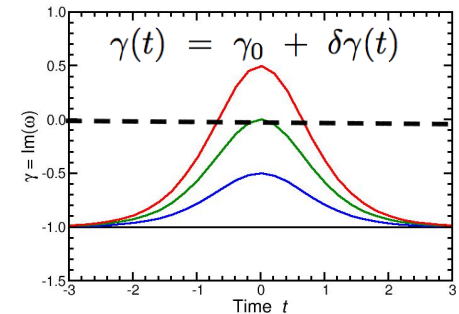
$$\frac{d}{dt}\sigma_{\nu\mu}(t) = 2\mathcal{D}_{\nu\mu} \overset{\text{seed}}{-} \overset{\text{feedback}}{-} i(\omega_\nu - \omega_\mu^*)\sigma_{\nu\mu}$$

Variance of a single mode:

$$\frac{d}{dt}\sigma_\nu^2 = 2\mathcal{D}_\nu + 2\gamma_\nu\sigma_\nu^2$$

*Lalime Equation*

$$\sigma_\nu^2(t) = \left[ 2\mathcal{D}_\nu \int_0^t e^{-2\Gamma_\nu(t')} dt' + \sigma_0^2 \right] e^{2\Gamma_\nu(t)}$$



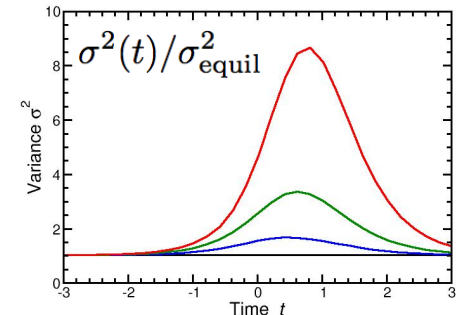
$$\Gamma_\nu(t) \equiv \int_0^t \gamma_\nu(t') dt'$$

*Amplification coefficient*

$$\gamma_\nu < 0: \sigma_\nu^2(t) \rightarrow -\mathcal{D}_\nu/\gamma_\nu$$

$$\gamma_\nu = 0: \sigma_\nu^2(t) = 2\mathcal{D}_\nu t$$

$$\gamma_\nu > 0: \sigma_\nu^2(t) \sim e^{2\Gamma_\nu(t)}$$

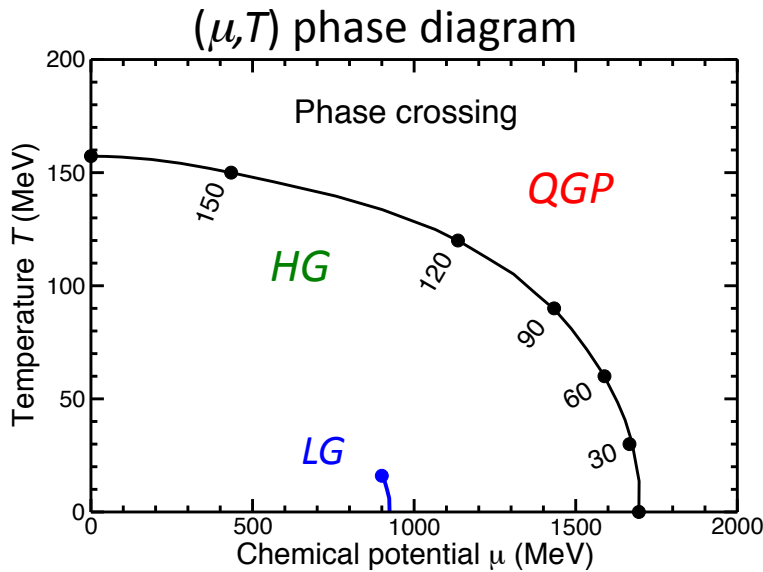
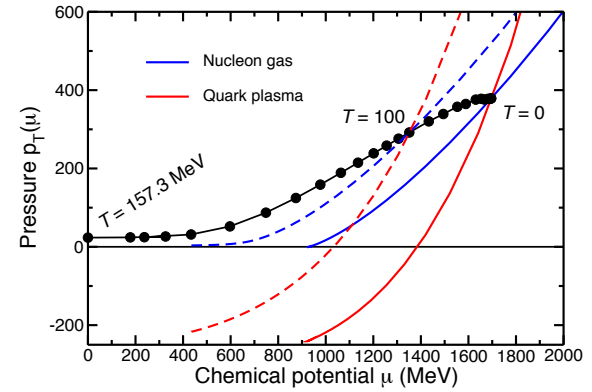


Colonna, Chomaz, Randrup,  
Nucl Phys A 567 (1994) 637

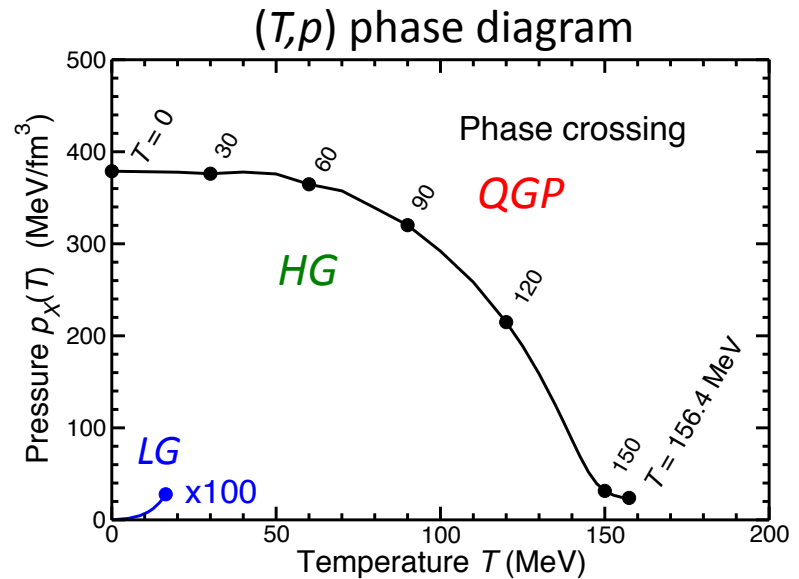
**Phase crossing:** Two different media have the same  $T, \mu, p$

**HG:**  $p^H = p_\pi + p_N + p_{\bar{N}} + p_w$

**QGP:**  $p^Q = p_g + p_q + p_{\bar{q}} - B$



- qualitatively similar to LG

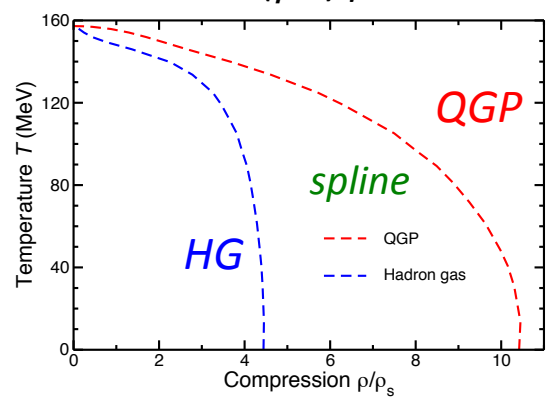


- qualitatively **different** from LG!

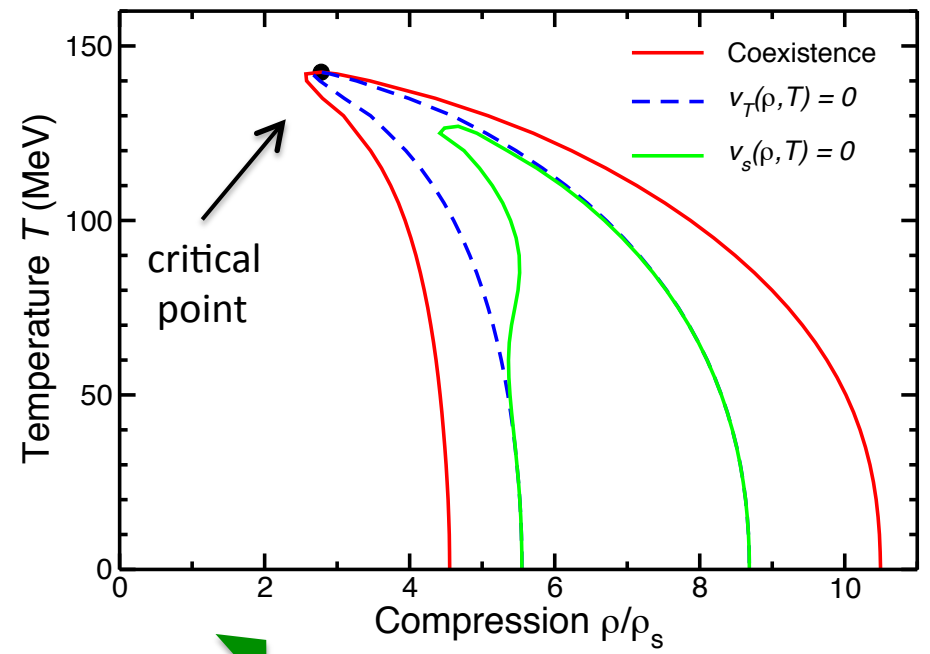
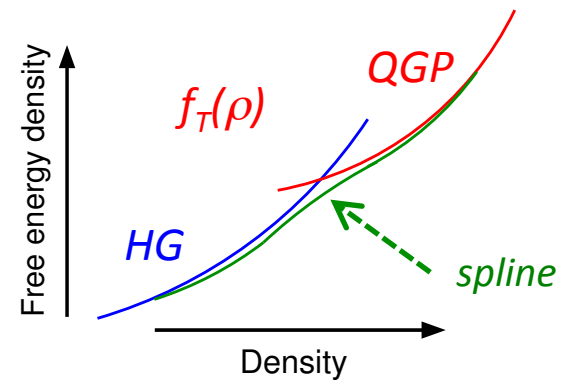
[Igor Iosilevsky, 2010]

# Construct a two-phase equation of state by splining between the HG and the QGP:

HG & QGP phase crossing  
in the  $(\rho, T)$  plane:



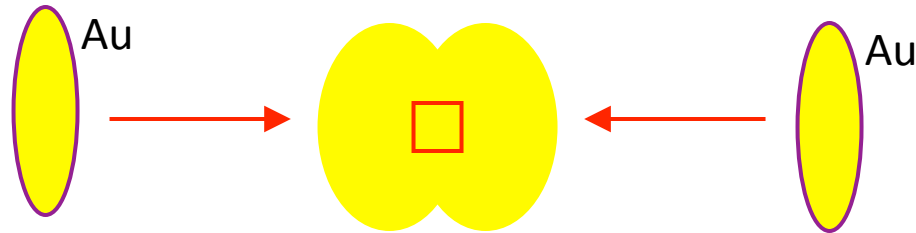
Spline between HG & QGP:



↔ ↔

Horizontal scale is rather uncertain!

*Collision energy  $\Leftrightarrow$  Phase region explored*



- Dynamical transport simulations (Au + Au at  $b=0$ )

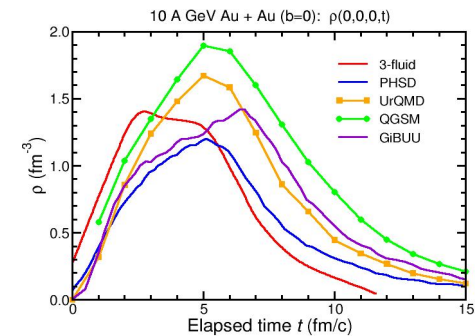
3-fluid: Yu.B. Ivanov, V.N. Ruusikikh, V.D. Toneev

PHSD: W. Cassing

UrQMD: I.C. Arsene, L.V. Bravina

QGSM: V.D. Toneev

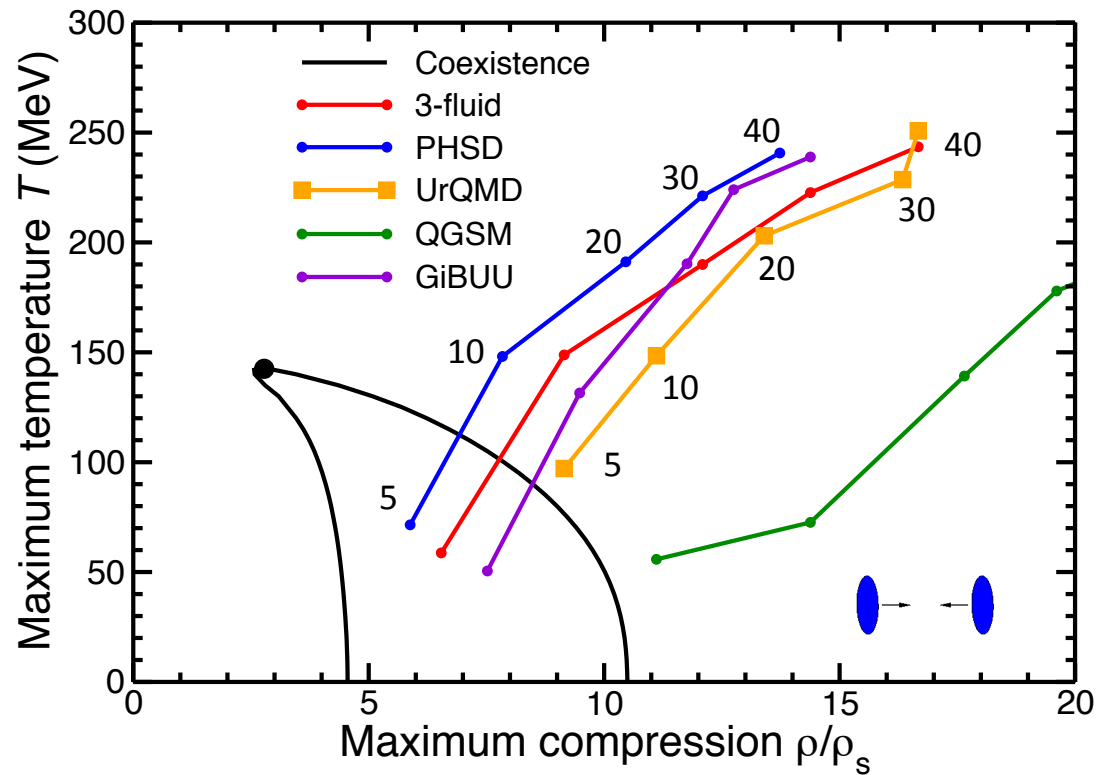
GiBUU: A. Larionov



- Extract  $\rho(t)$  and  $\epsilon(t)$  at the center

Arsene, Bravina, Cassing, Ivanov, Larionov, Randrup, Ruusikikh, Toneev, Zeeb, Zschesche: PRC75 (2007) 034902

## Maximum density and temperature achieved

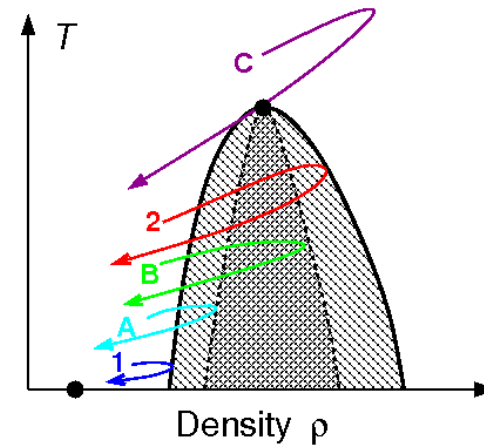
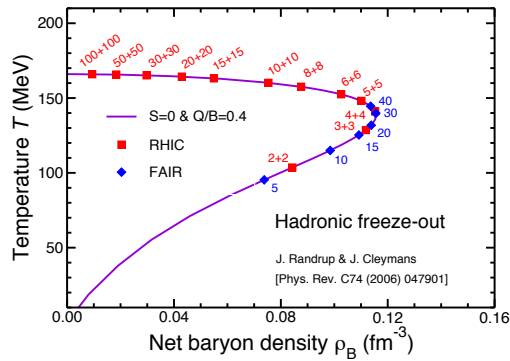
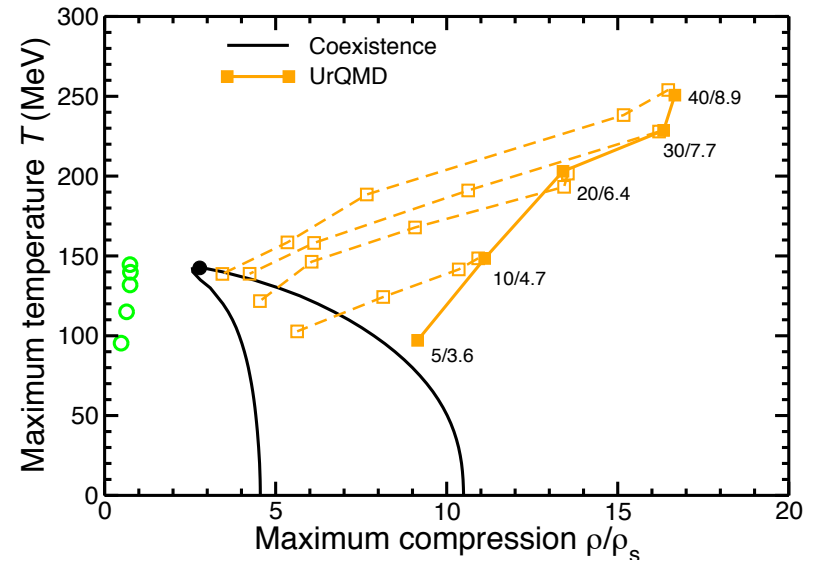
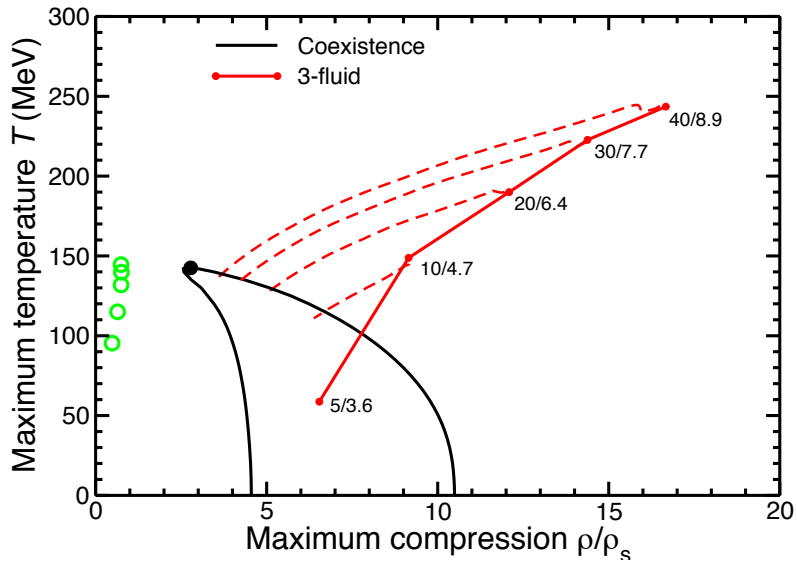


**QGP:**

$$\rho = 2\mu_q T^2 + \frac{2}{\pi^2} \mu_q^3$$

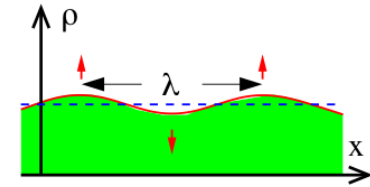
$$\varepsilon = \frac{37\pi^2}{30} T^4 + 3\mu_q^2 T^2 + \frac{3}{2\pi^2} \mu_q^4 + B \quad \Rightarrow (\rho, T)$$

# Expansion ( $\rho, T$ ) phase trajectories





# Evolution of density fluctuations with dissipative fluid dynamics



$$A_k \sim \exp(-i\omega_k t)$$

$$\omega_k = \epsilon_k + i\gamma_k$$

## Dispersion relations:

Ideal fluid dynamics:

$$\omega^2 \doteq v_s^2 k^2$$

+ gradient term:

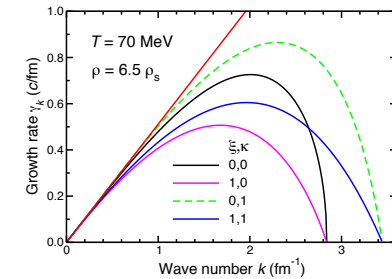
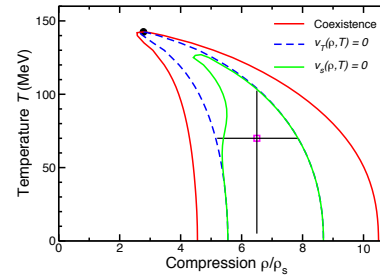
$$\omega^2 \doteq v_s^2 k^2 + C \frac{\rho_0^2}{h_0} k^4 \quad \Leftarrow \quad p_k \rightarrow p_k + C \rho_0 k^2 \rho_k$$

+ shear & bulk viscosity:

$$\omega^2 \doteq v_s^2 k^2 + C \frac{\rho_0^2}{h_0} k^4 - i\xi \frac{\omega}{h_0} k^2 \quad \Leftarrow \quad \xi \equiv \frac{4}{3}\eta + \zeta$$

+ heat conduction:

$$\omega^2 \doteq v_T^2 k^2 + C \frac{\rho_0^2}{h_0} k^4 - i\xi \frac{\omega}{h_0} k^2 + \frac{v_s^2 - v_T^2}{1 + i\kappa k^2 / \omega c_v} k^2 \quad \Leftarrow \quad \kappa$$



# Transport coefficients

$$\eta_0 \geq 1$$

$$\kappa_0 \geq 1$$

1) Bulk viscosity  $\zeta$ : Ignore  $\zeta \ll \eta \Rightarrow \xi \equiv \frac{4}{3}\eta + \zeta \approx \frac{4}{3}\eta$

2) Shear viscosity  $\eta$ :

$$\rho = 0 : h \equiv p + \varepsilon = T\sigma$$

$$\rho > 0, T \ll mc^2 : h \asymp mc^2 n \gg T\sigma$$

\*)  $\rho = 0 : \eta \geq \frac{\hbar}{4\pi} \sigma = \frac{\hbar}{4\pi} \frac{h}{T}$

$$\rho = 0 : n \sim T^3 \Rightarrow \frac{\hbar c}{T} = 4\pi c_0 d \quad d \equiv n^{1/3}$$

$$\eta(\rho, T) = \eta_0 \frac{c_0}{c} d(\rho, T) h(\rho, T)$$

$$\lambda_{\text{visc}} \equiv \frac{1}{c} \frac{\xi(\rho, T)}{h(\rho, T)/c^2} \approx \frac{4}{3} \eta_0 c_0 d(\rho, T)$$

3) Heat conductivity  $\kappa$ :

$$\eta \approx \frac{1}{3} n \bar{p} \ell$$

$$\frac{\kappa}{\eta} \approx \frac{c_v}{\hbar/c^2}$$

$$\bar{p} = m\bar{v}$$

$$h \asymp mc^2 n$$

$$\kappa \approx \frac{1}{3} \bar{v} \ell c_v$$

$$c_v \equiv \partial_T \varepsilon_T(\rho)$$

$$c_v \asymp \frac{3}{2} n$$

$$\kappa(\rho, T) = \kappa_0 c_0 c d(\rho, T) c_v(\rho, T)$$

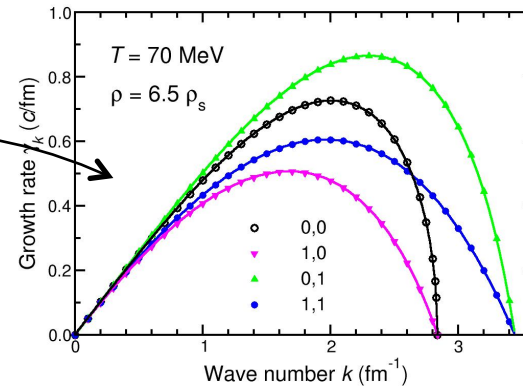
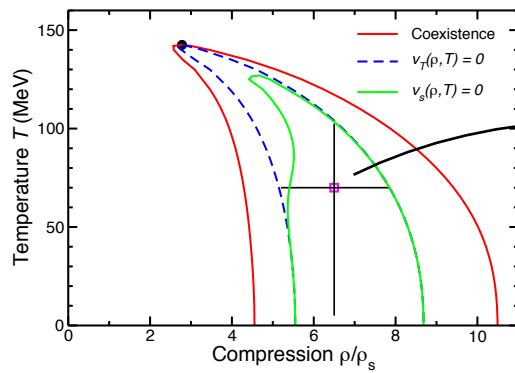
$$\lambda_{\text{heat}} \equiv \frac{1}{c} \frac{\kappa(\rho, T)}{c_v(\rho, T)} = \kappa_0 c_0 d(\rho, T)$$

\*) V Koch & J Liao, PRC81, 014902 (2010)

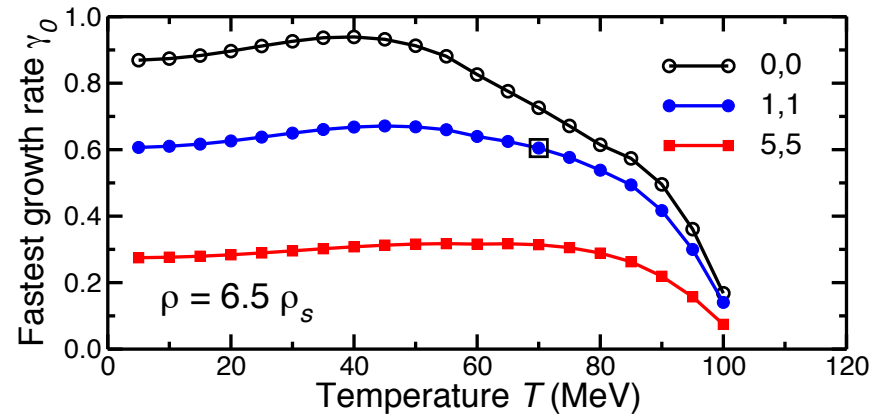
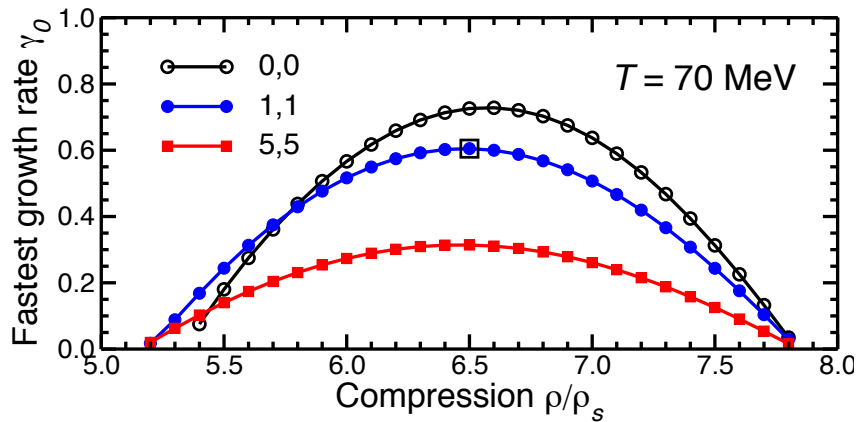
$$c_0 = \frac{1}{4\pi} \left[ (g_g + \frac{3}{3} g_q) \frac{\zeta(3)}{\pi^2} \right]^{\frac{1}{3}} \approx 0.12779$$

# Spinodal growth rates

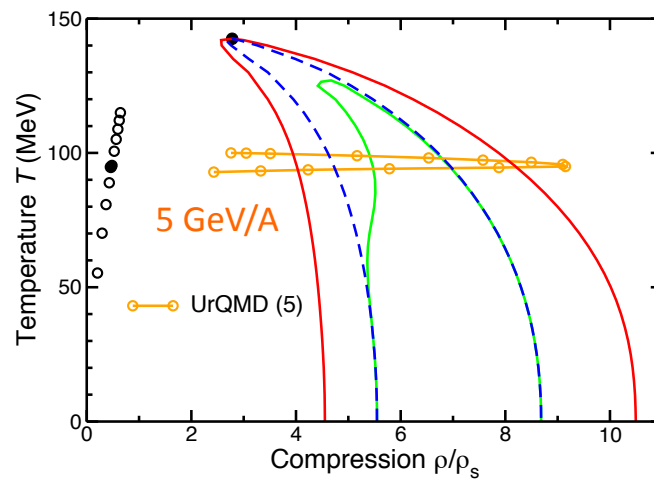
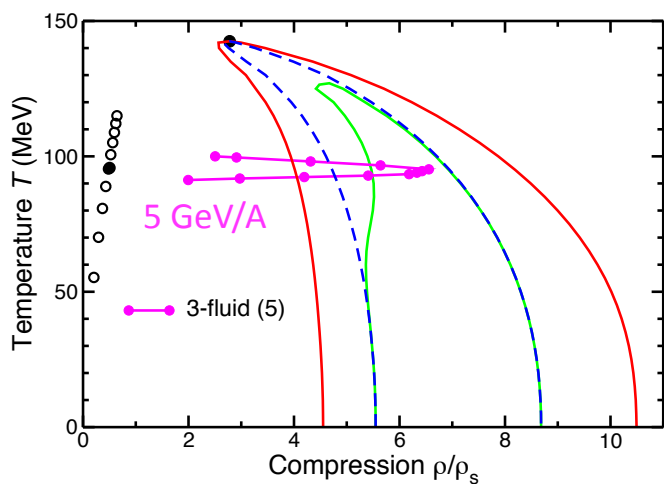
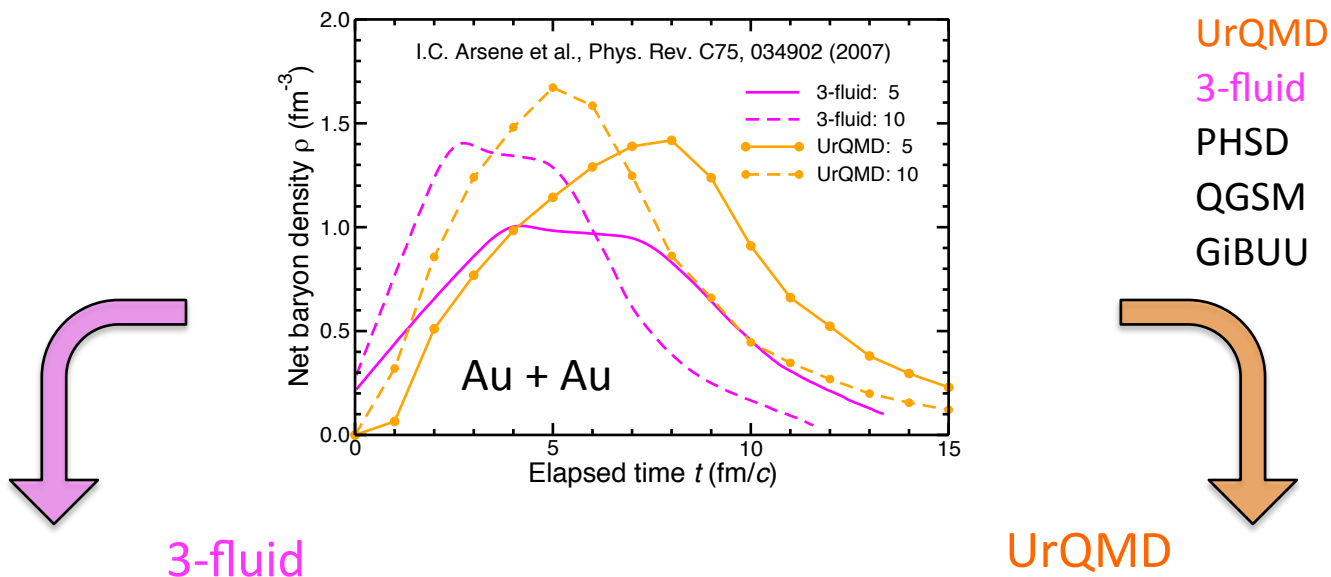
$$(\rho, T): \quad \omega^2 \doteq v_T^2 k^2 + C \frac{\rho_0^2}{h_0} k^4 - i\xi \frac{\omega}{h_0} k^2 + \frac{v_s^2 - v_T^2}{1 + i\kappa k^2 / \omega c_v} k^2 \quad \Rightarrow \quad \gamma_k(\rho, T)$$



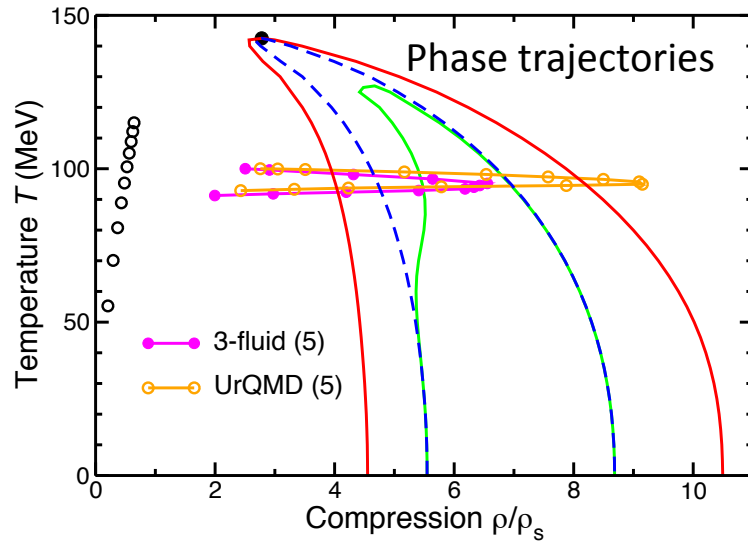
Fastest growth rates:



# Dynamical phase trajectories ( $\rho(t), T(t)$ ):



# Spinodal amplification

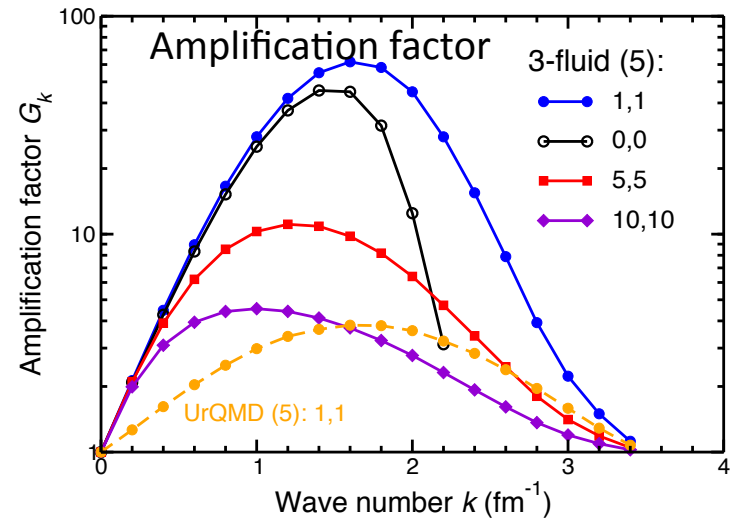
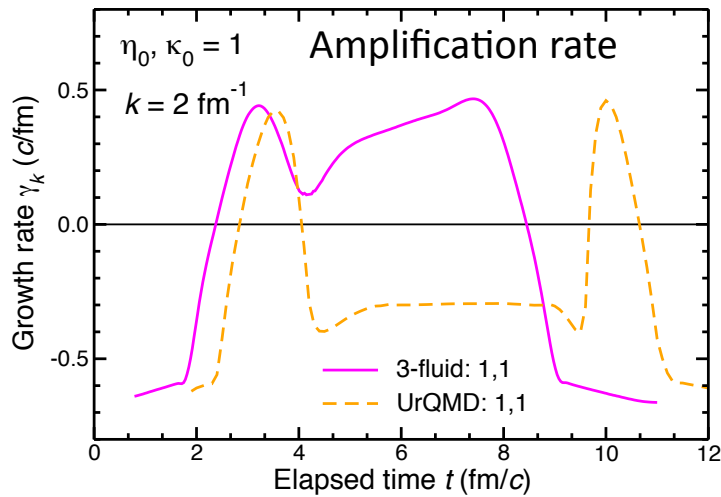


Amplification coefficient:

$$\Gamma_\nu(t) \equiv \int_0^t \gamma_\nu(t') dt'$$

Amplification factor:

$$G_k = e^{2\Gamma_k}$$



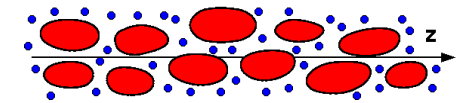
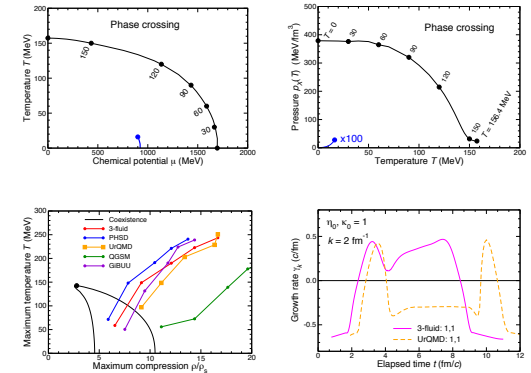
# Summary & outlook

Elementary considerations suggest the existence of a deconfined QGP phase at high  $T$  and/or  $\mu$ .

Dynamical model simulations suggest that this new phase is entered already at  $E_{beam} \approx 10$  GeV/A

The energy range *below* this threshold is optimal for exploring the character of the phase transition

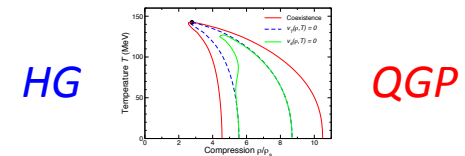
Spinodal decomposition (clumping) *may* occur in this range of collision energies



**BUT: Full dynamical simulations are needed!**



Two phases: hadron gas & quark-gluon plasma  
Equation of State with phase *transition*



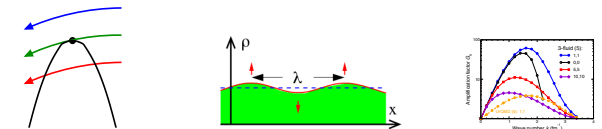
(SOME) MODEL REQUIREMENTS:

Phase coexistence: interface & mixed phase\*



Dynamics with instabilities\*

Spinodal modes\*



\* Requires finite range