

Properties of strange hadrons in vacuum and cold nuclear matter

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Analysis of the data and tuning of the transport codes:





Energy Regime

| Proton Energy [GeV] | 1-4 | 30 | |
|--|----------|-----------|--------------------------|
| Heavy Ion Energy [AGeV] | 1-2 | 2-10 | $\sqrt{s} = 4 - 11 GeV$ |
| | | HADES | NICA |
| | HADES | + | |
| | At SIS18 | + | |
| Purpose: Creation of Compressed baryonic matter ρ up to 5-6 ρ_0 ?? | | СВМ | |
| | | At SIS100 | |



P+P, P+NB AT 3.5GEV



- Role of resonances in strangeness production in p+p collisions at 3.5 GeV
- KOs properties in cold nuclear matter p+Nb at 3.5 GeV



How to understand HIC in the GeV regime

Examine strangeness in nuclear matter in the lab (HADES experiment, GSI, Darmstadt)

Heavy-ion collisions $\rho_B \leq 2 - 3 \rho_0$



Vienna University of Technology

N-N primary reactions π N, Δ N secondary reactions Resonance Excitation Propagation of particles through nuclear matter Scattering and Absorption In-Medium Spectral Function Mean field Approaches



Role Played by Resonances at intermediate Energies: I

$$p + p(@3.5 \, GeV) \rightarrow$$

$$\Sigma(1385)^+ + K^+ + n$$

$$\Lambda + \pi^+$$

Chinowsky, W. et al. Phys.Rev. 165 (1968) 1466-1478

$$\frac{\Delta(1900 - 2000)^{++}}{\Gamma} \rightarrow \Sigma(1385)^{+} + K^{+}$$

$$\Gamma = 150 - 200 \text{MeV/c}^{2}$$

~ **30%** Δ⁺⁺

ТШ

MU

Role of the Δ *(1940) in the π +p- \rightarrow K+ Σ +(1385) and pp- \rightarrow nK+ Σ +(1385) reactions Ju-Jun Xie, En Wang, Bing-Song Zou

arxiv.1405.5586



Role Played by Resonances at intermediate Energies: II



HADES, arXiv:1403.6662



 $\sigma(p+p \to \Lambda + p + \pi^+ + K_S^0) = 2.57 \pm 0.02^{+0.21}_{-1.98} \pm 0.18 \ \mu b$

Role Played by Resonances at intermediate Energies: II



 $\sigma(p+p \to \Lambda + \Delta^{++} + K_S^0) = 26.27 \pm 0.64^{+2.57}_{-2.13} \pm 1.84 \text{ }\mu\text{b}$ $\overline{z}\sigma(p+p \to \Lambda + p + \pi^+ + K_S^0) = 2.57 \pm 0.02^{+0.21}_{-1.98} \pm 0.18 \text{ }\mu\text{b}$

ue with the Resonances in hadron-hadron collisions: Interferences



with the Resonances in hadron-hadron collisions: Interferences

 $p + p(@3.5 \, GeV) \rightarrow$ $p + K^+ + \Lambda$ $p + N^* \rightarrow p + K^+ + \Lambda$

Bonn-Gatchina Partial Wave Analysis

http://pwa.hiskp.uni-bonn.de/

A.V. Anisovich, V.V. Anisovich, E. Klempt, V.A. Nikonov and A.V. Sarantsev Eur. Phys. J. A 34, 129152 (2007)

What we included to model the $PK^+\Lambda$ process:

N* Resonances in the PDG with measured decay into $K^{\scriptscriptstyle +}\Lambda$

N(1650), N(1710), N(1720), N(1875), N(1880), N(1895), N(1900) Non-resonant PK⁺Λ production waves Interferences Issue with the Resonances in hadron-hadron collisions: Interferences



ue with the Resonances in hadron-hadron collisions: Interferences



$$p + N^* \to p + K^+ + \Lambda$$



included resonances: Non-resonant waves:

N(1650), N(1710), N(1720), N(1900), N(1895) $(pL)({}^{1}S_{0}) - K \quad (pL)({}^{3}S_{1}) - K \quad (pL)({}^{1}P_{1}) - K$ $(pL)({}^{3}P_{0}) - K \quad (pL)({}^{3}P_{2}) - K \quad (pL)({}^{3}P_{1}) - K$ $(pL)({}^{3}D_{1}) - K \quad (pL)({}^{1}D_{2}) - K \quad (pL)({}^{3}D_{2}) - K$



Intermediate Conclusions I

- Decay chains for Resonances can be determined by exclusive/semi-exclusive measurements of final state in p+p and p+n collision for energies between 2-10 GeV
- Interferences among resonances play an important role for some final states (for dileptons too btw) but those can be estimated using PWA on elementary collisions like p+N and π +N. A 4π distribution can be extracted which one could implement into transport models.
- Still... A systematic analysis of different data sets at different energies is needed to pin down quantitatively the contribution of N* to pK Λ e.g.
- **Upcoming** measurements of π +N and π +A for p=0.8-1.65 GeV/c and planned PWA analysis to determine N* and Δ yields and their contribution to the strange and dilepton final states.



Kaons in really cold nuclear matter

p+Nb, 3.5 GeV







Inside the Nucleus:

The K⁰ experiences a potential due to the surrounding nucleons. Modification of these processes.

K⁰_s in cold nuclear matter

Neutral kaons measured by HADES in p+p and p+⁹³Nb collisions at 3.5 GeV:





Resonance model of kaon production

All production channels:

| No | Reaction |
|--------|--|
| 1 | $nn \rightarrow n\Lambda K^+$ |
| 9 | $pp \rightarrow pMK^+$ |
| 2 | $pn \rightarrow nMK^{+}$ |
| 3 | $pp \rightarrow p \angle K$ $\Sigma - V^+$ |
| 4 | $nn \rightarrow n\Sigma \Lambda$ |
| о С | $pn \rightarrow n\Sigma^{-}K^{+}$ |
| 0 | $np \rightarrow p\Sigma^-K^+$ |
| 7 | $pp \rightarrow n \Sigma^+ K^+$ |
| 8 | $nn \rightarrow \Delta^- \Lambda K^+$ |
| 9 | $pp \rightarrow \Delta^{++} \Sigma^- K^+$ |
| 10 | $\Delta^{++}n \rightarrow p\Lambda K^+$ |
| 11 | $\Delta^- p \rightarrow n \Sigma^- K^+$ |
| 12 | $\Delta^{++}p \rightarrow \Delta^{++}\Lambda K^+$ |
| 13 | $\Delta^+ n \rightarrow \Delta^0 \Lambda K^+$ |
| 14 | $\Delta^+ p \rightarrow \Delta^+ \Lambda K^+$ |
| 15 | $\Delta^{++}n\to \Delta^{++}\Sigma^-K^+$ |
| 16 | $\Delta^0 p \rightarrow \Delta^+ \Sigma^- K^+$ |
| 17 | $\Delta^+ n \to \Delta^+ \Sigma^- K^+$ |
| 18 | $\Delta^{++}p \rightarrow \Delta^{++}\Sigma^0 K^+$ |
| 19 | $\Delta^+ n \rightarrow \Delta^0 \Sigma^0 K^+$ |
| 20 | $\Delta^+ p \rightarrow \Delta^+ \Sigma^0 K^+$ |
| 21 | $\Delta^+ p \rightarrow \Delta^0 \Sigma^+ K^+$ |
| 22 | $\Delta^+ \Delta^{++} \rightarrow \Delta^{++} \Lambda K^+$ |
| 23 | $\Delta^0 \Delta^{++} \rightarrow \Delta^+ \Lambda K^+$ |
| 24 | $\Delta^0 \Delta^+ \rightarrow \Delta^0 \Lambda K^+$ |
| 25 | $\Delta^{++}\Delta^0 \rightarrow \Delta^{++}\Sigma^-K^+$ |
| 26 | $\Delta^-\Delta^0 \to \Delta^-\Sigma^-K^+$ |
| 27 | $\Delta^0 \Delta^{++} \rightarrow \Delta^+ \Sigma^0 K^+$ |
| 28 | $\Delta^- \Delta^+ \to \Delta^0 \Sigma^- K^+$ |

Cross section parameterization:

$$\sigma(B_1B_2 \to B_3YK) = a\,\left(\frac{s}{s_0}-1\right)^b\,\left(\frac{s_0}{s}\right)^c,$$

Note: this is what's inside transport code

np-reactions isospin interrelations (one example):

$$\begin{split} \sigma(nn \to \Delta^{-}\Lambda K^{+}) &= \sigma(pp \to \Delta^{++}\Lambda K^{0}) \\ &= 3\sigma(pn \to \Delta^{0}\Lambda K^{+}) &= 3\sigma(np \to \Delta^{+}\Lambda K^{0}) \\ &= 3\sigma(pp \to \Delta^{+}\Lambda K^{+}) &= 3\sigma(nn \to \Delta^{0}\Lambda K^{0}), \end{split}$$

almost no experimental data for np!

K⁰ production channels:

| Number of particles | Final state | What is in the model | |
|---------------------|--|------------------------------|--|
| 3-body | p Σ+ K ⁰ | p Σ+ K ⁰ | |
| 4-body | p π+ Λ K ⁰ | $\Delta^{++} \wedge K^0$ | |
| | p π ⁺ Σ ⁰ K ⁰ | $\Delta^{++} \Sigma^0 \ K^0$ | |
| | $n \pi^+ \Sigma^+ K^0$ | $\Delta^+ \Sigma^+ K^0$ | |
| | $p \pi^0 \Sigma^+ K^0$ | | |

Note:

1. Pion production goes exclusively through Δ .

2. No angular anisotropies in production.

K⁰_S from p+p@ 3.5GeV

HADES arXiv:1403.6662 [nucl-ex]., accepted by PRC



pNb data vs. tuned resonance model

p+Nb@3.5 GeV

GiBUU w/o pot.



▶ KN potential is OFF.

▶ 3-body reactions in np (np \rightarrow NYK) poorly constrained, scale factor 0.5 is applied to the Tsushima parameterizations.

• GiBUU simulations based on tuned resonance model describe data.



Rapidity distribution and role of secondary reactions

p+Nb@3.5 GeV



▶ Significant contribution of secondary reactions at backward rapidities (~70%).

- Three main sources:
- ΔN -reactions. Rely on the resonance model (Tsushima et al.).
- ▶ πN-reactions.
- KN scattering.

How well the two last processes are known?

Secondary processes: pion-nucleon reactions



- Elementary cross sections are known well and parametrized in the model.
- No angular distributions implemented in the model.

Pictures from K. Tsushima, A. Sibirtsev, A.W. Thomas, PRC62 (2000) 064904



Kaon-nucleon scattering



Parametrization: M. Effenberger, PhD. Giessen, 1999.

- Vacuum cross sections are well known.
- K⁰N scattering from isospin considerations.
- No angular distributions implemented in the model (some data are available).

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Kaons properties in Matter

How to describe the properties of Kaon in nuclear matter? Effective Chiral Lagrangian with Kaons and Nucleons as degree of Freedom

Mean Field Dynamics

C. Fuchs Progr. In Part and Nucl. Phys. 56 (2006) 1-103

$$\begin{bmatrix} \left(\partial_{\mu} \pm iV_{\mu}\right)^{2} + m_{K}^{*2} \end{bmatrix} \phi_{K^{\pm}}(x) = 0$$

$$m_{K}^{*} = \sqrt{m_{K}^{2} - \frac{\Sigma_{KN}}{f_{\pi}^{*2}}} \rho_{S} + V_{\mu}V^{\mu}$$

$$\boxed{V_{\mu}} = \frac{3}{8f_{\pi}^{*2}} j_{\mu} \text{ Vector potential attractive for } K^{*} \text{ repulsive for } K^{+}$$

$$\boxed{\Sigma_{KN}} = \text{Scalar potential, } \sim 300\text{-}450 \text{ MeV, same for } K^{+} \text{ and } K^{-}$$

$$E({\bf k}) = \sqrt{{\bf k}^2 + m_{\rm K}^2 - \frac{\Sigma_{\rm KN}}{f_\pi^2} \rho_{\rm s} + V_0^2} \pm V_0$$



 $IM \leq Re \geq No K^+$ absorption

T. Gaitanos, K. Lapidus -> GiBUU

HADES, arXiv:1404.7011

Effect of the potential in pNb: $p-\theta$ spectra







Systematic check

HADES, arXiv:1404.7011



• w/o potential

with potential ("U = +35 MeV")

$$\chi^2 = \sum_i \frac{\left((d\sigma/dp_t)_i^{\text{exp.}} - (d\sigma/dp_t)_i^{\text{sim.}} \right)^2}{\sigma_{\text{exp.syst.}}^2}$$







Intermediate Conclusions II

p+A at 3.5 GeV : ρ 0

detailed comparison of the experimental data for the inclusive KOs to ONE Transport model (GiBUU)

Pirimide Approach:

- -> Elementary cross-sections
- -> Check scattering effects via Rapidity density distribution
- -> Secondary processes included
- -> Test of the Chiral potential

Result for KOs:

Repulsive potential around 40 MeV for k=0 and $\rho = \rho 0$

Factor 2 larger than extracted from Flow in HIC and previous π +A measurement

-> new Measurements for π +A at high rate with HADES in 3 weeks.



Correlation function depends on scattering length: Lednicky model $C(k) = 1 + \sum_{s} \rho_s \left[\frac{1}{2} \left| \frac{f^s(k)}{r_0} \right|^2 \left(1 - \frac{d_0^s}{2\sqrt{\pi}r_0} \right) + \frac{2\Re(f^s(k))}{\sqrt{\pi}r_0} F_1(Qr_0) - \frac{\Im(f^s(k))}{r_0} F_2(Qr_0) \right]$



Thank you