

Constraints on the Dense Matter Equation of State from Neutron Stars

David Blaschke^{a,b,c}

Thanks to: M. Cierniak^a, O. Ivanytskyi^a, T. Fischer^a, M. Shahrbafe^a,
A. Ayriyan^b, A. Bauswein^d & S. Typel^d

a – University Wroclaw, b - JINR Dubna, c – NRNU (MEPhI) Moscow, d – GSI Darmstadt

1. Introduction: Recent relevant multi-messenger observations
2. New paradigm: Only hybrid stars fulfil new M-R constraints
3. Outlook: Supernovae & Mergers in the QCD phase diagram
→ Constraints for the Onset of Deconfinement?

STRONG-2020 Workshop of NA7-Hf-QGP, Hersonissos, 8.10.2021

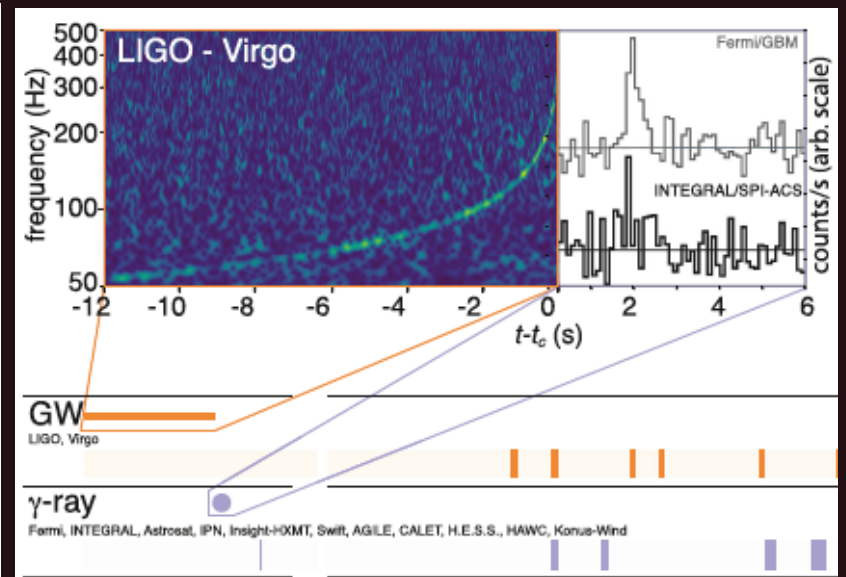
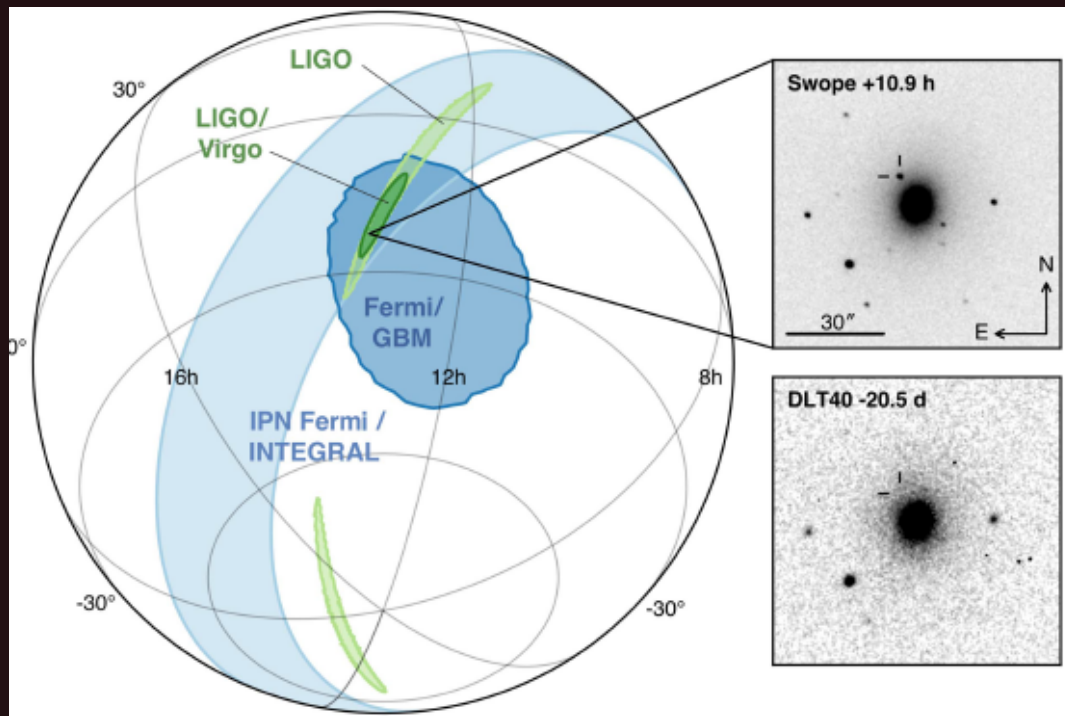
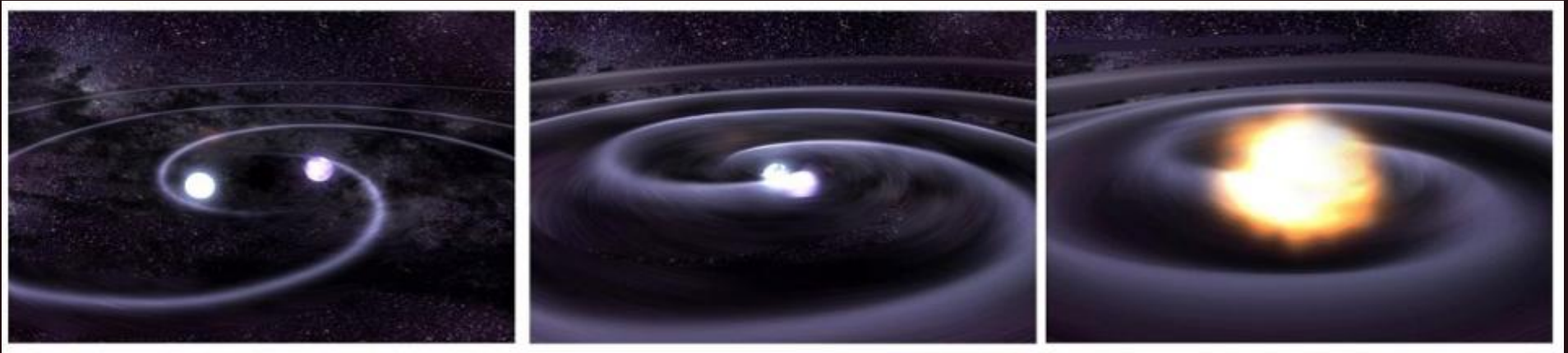


Grant No. UMO 2019 / 33 / B / ST9 / 03059

Grant No. 18-02-40137



Discovery: neutron star merger !



GW170817A , announced 16.10.2017 *)

*) B.P. Abbott et al. [LIGO/Virgo Collab.], PRL 119, 161101 (2017); ApJLett 848, L12 (2017)

NS-NS merger !

GW170817A , announced 16.10.2017 *)

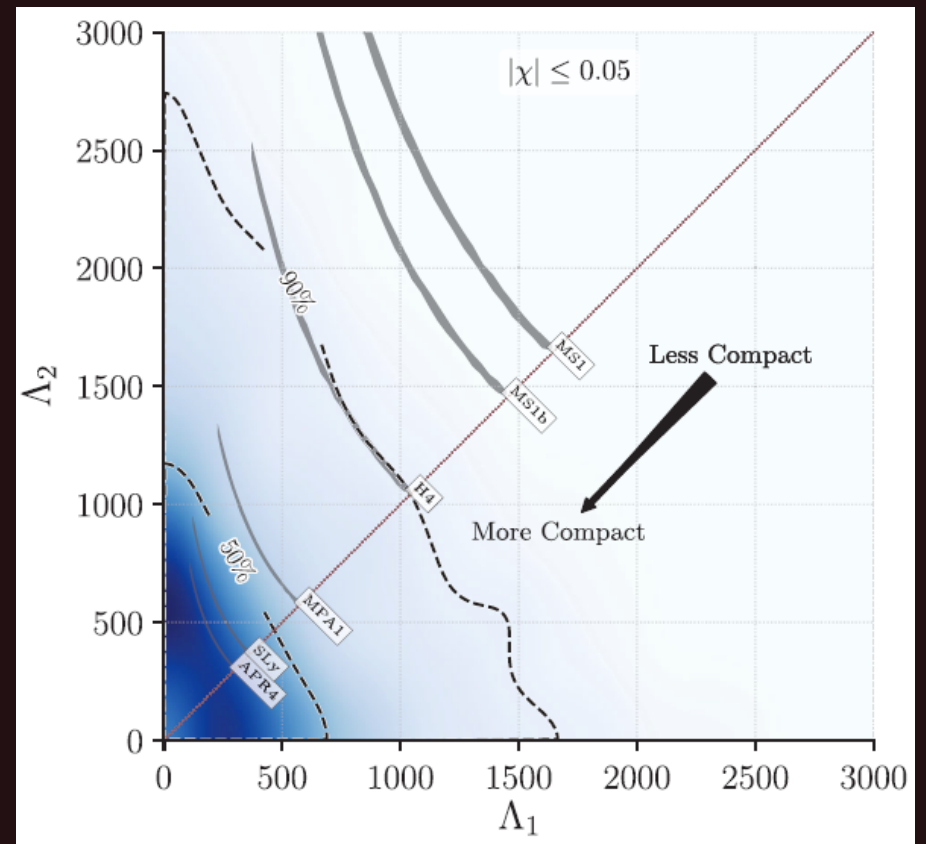
Multi-Messenger Astrophysics !!

Low-spin priors ($ \chi \leq 0.05$)	
Primary mass m_1	$1.36\text{--}1.60 M_\odot$
Secondary mass m_2	$1.17\text{--}1.36 M_\odot$
Chirp mass \mathcal{M}	$1.188^{+0.004}_{-0.002} M_\odot$
Mass ratio m_2/m_1	$0.7\text{--}1.0$
Total mass m_{tot}	$2.74^{+0.04}_{-0.01} M_\odot$
Radiated energy E_{rad}	$> 0.025 M_\odot c^2$
Luminosity distance D_L	40^{+8}_{-14} Mpc

Constraint on neutron star maximum mass

$$M_{\text{TOV}} < 2.17 M_{\text{sun}}$$

(Margalit & Metzger, arxiv:1710.05938)



Constraint on parameter ($\Lambda < 800$)

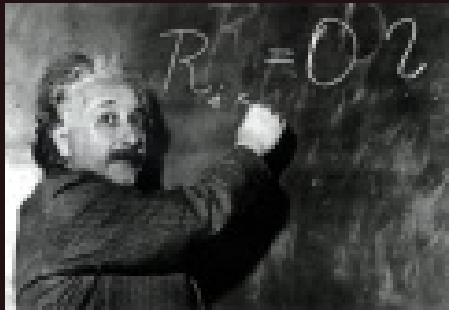
$$\tilde{\Lambda} = \frac{16(m_1 + 12m_2)m_1^4\Lambda_1 + (m_2 + 12m_1)m_2^4\Lambda_2}{(m_1 + m_2)^5}$$

Dimensionless tidal deformability

$$\Lambda = (2/3)k_2[(c^2/G)(R/m)]^5$$

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Compact stars and black holes in Einstein's General Relativity theory

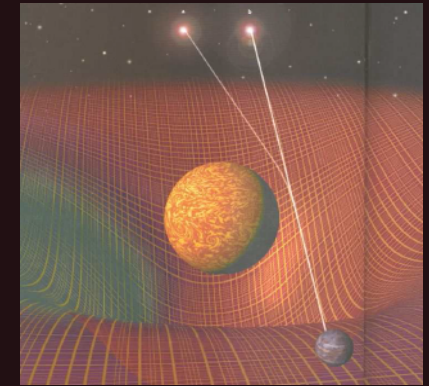


Space-Time

$$G_{\mu\nu} = 8\pi G T_{\mu\nu}$$

Matter

Massive objects curve the Space-Time



Non-rotating, spherical masses \rightarrow Schwarzschild Metrics

$$ds^2 = -\left(1 - \frac{2M}{r}\right)dt^2 + \left(1 - \frac{2M}{r}\right)^{-1}dr^2 + r^2d\Omega^2$$

Einstein eqs. \rightarrow Tolman-Oppenheimer-Volkoff eqs.*)

For structure and stability of compact stars

$$\frac{dP(r)}{dr} = -G \frac{m(r)\epsilon(r)}{r^2} \left(1 + \frac{P(r)}{\epsilon(r)}\right) \left(1 + \frac{4\pi r^3 P(r)}{m(r)}\right) \left(1 - \frac{2Gm(r)}{r}\right)^{-1}$$

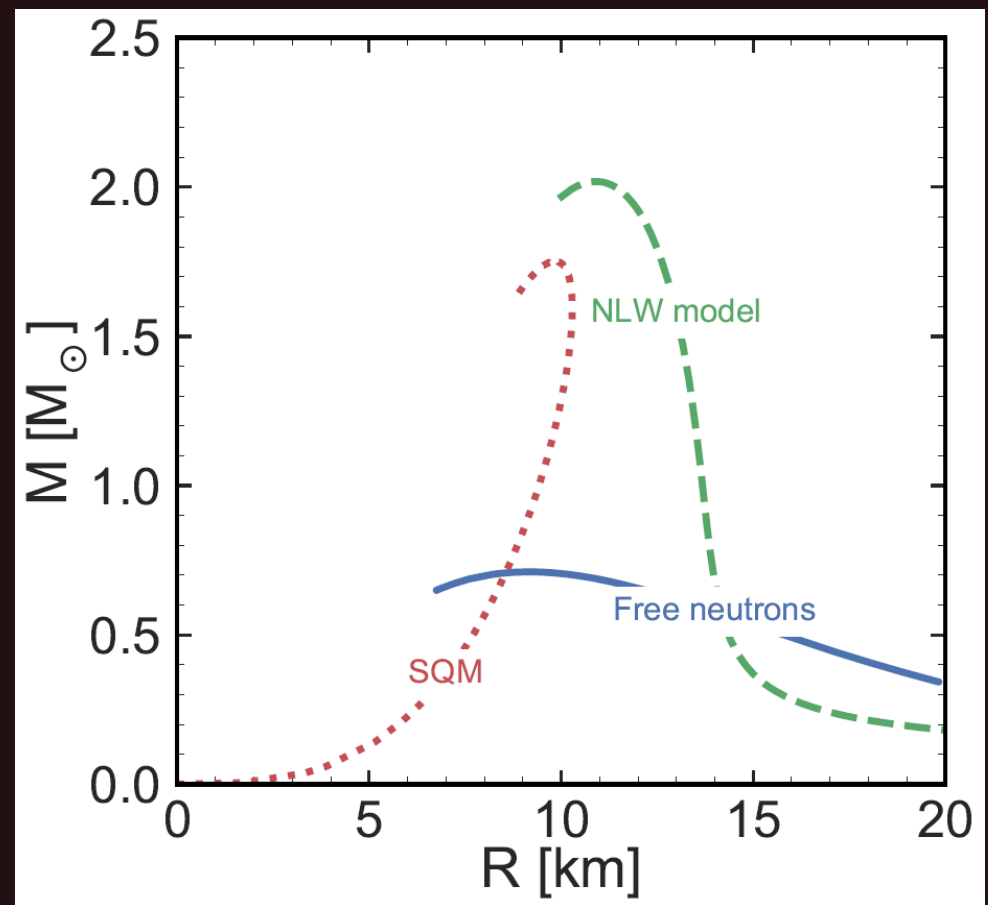
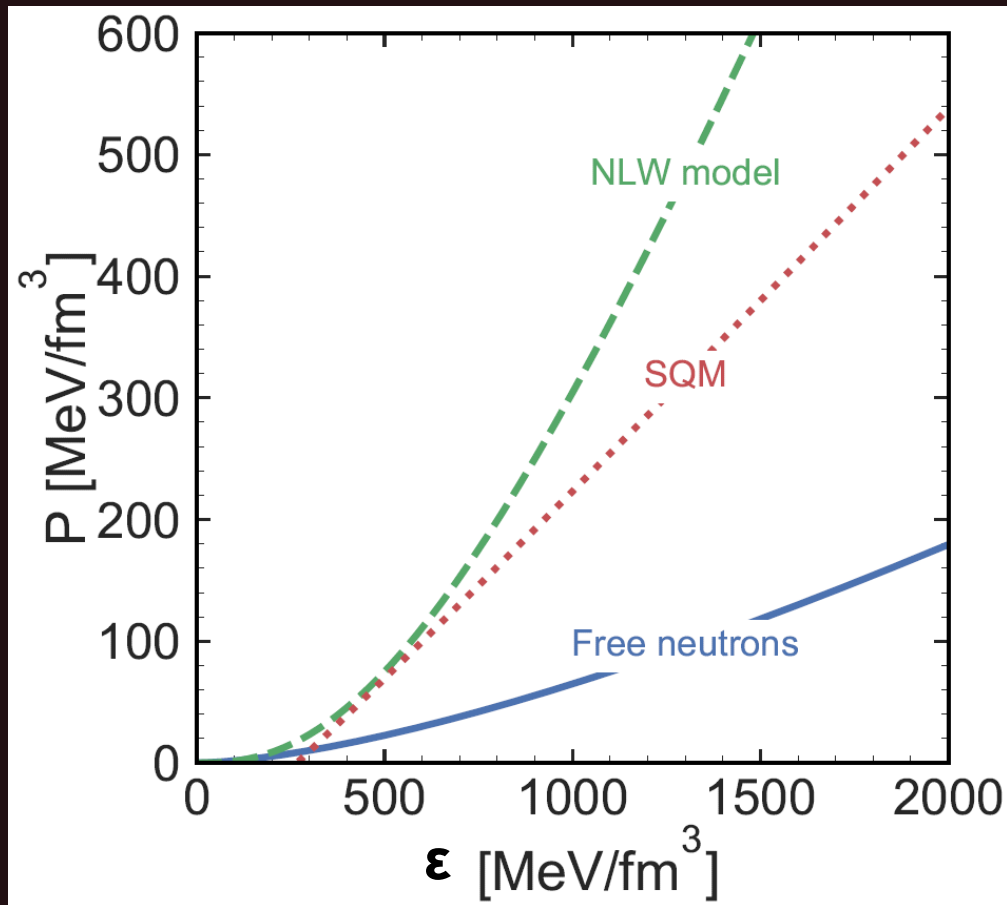
Newtonian case x GR corrections from EoS and metrics



*) R. C. Tolman, Phys. Rev. 55 (1939) 364 ; J. R. Oppenheimer, G. M. Volkoff, ibid., 374

The 1:1 relation $P(\varepsilon) \leftrightarrow M(R)$ via TOV

Simple examples*)



Free neutrons: Oppenheimer & Volkoff, Phys. Rev. 55 (1939) 374

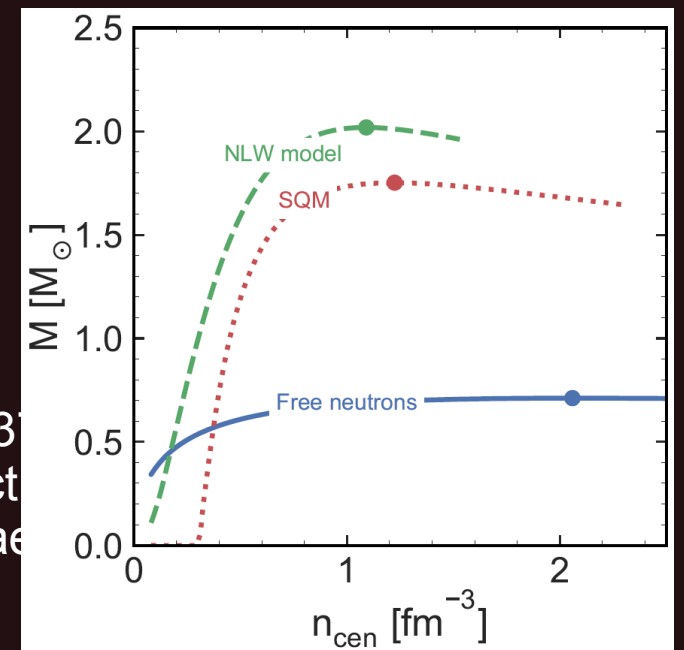
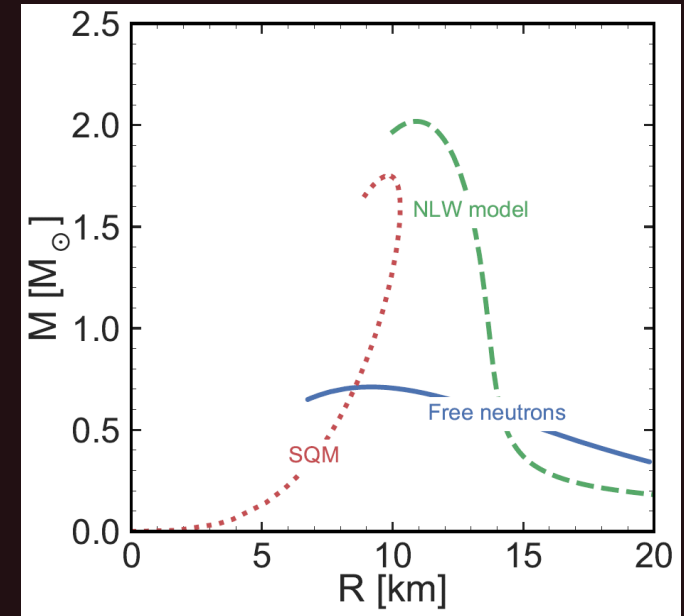
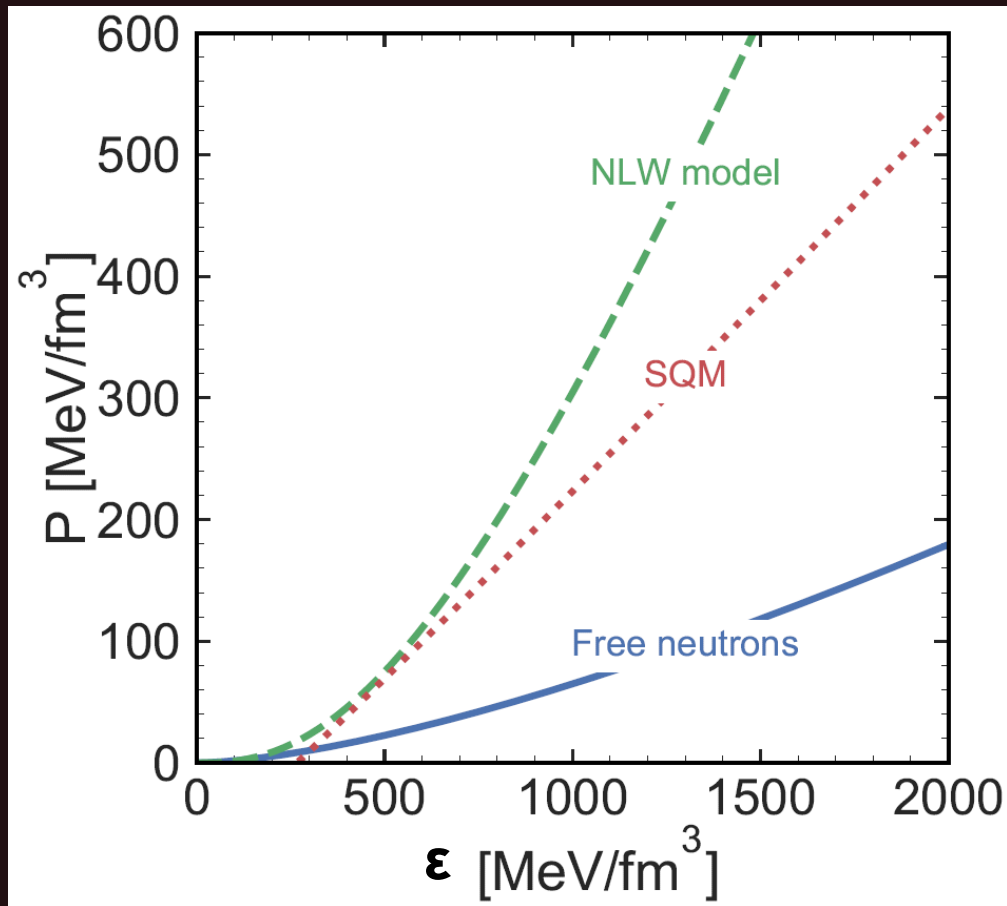
NLW (nonlinear Walecka) model: N. K. Glendenning, Compact Stars (Springer, 2000)

SQM (strange quark matter): P. Haensel, J. L. Zdunik, R. Schaeffer, A&A 160 (1986) 121

*) courtesy: Konstantin Maslov

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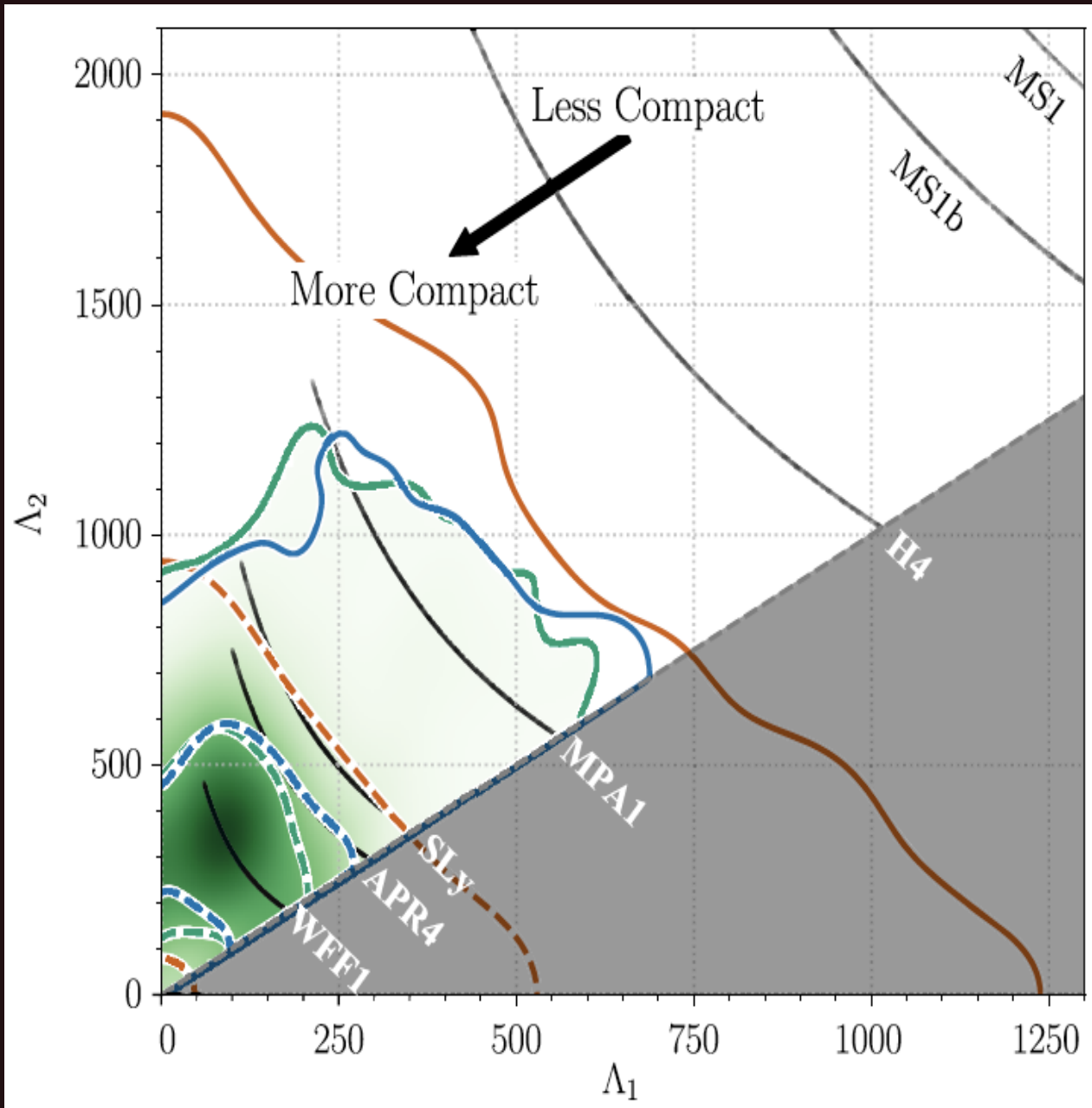
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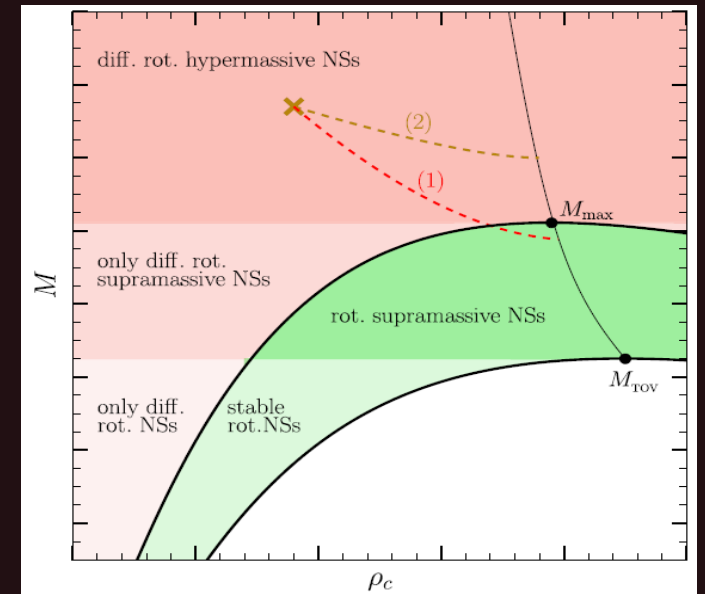
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Constraints on NS mass and radii !

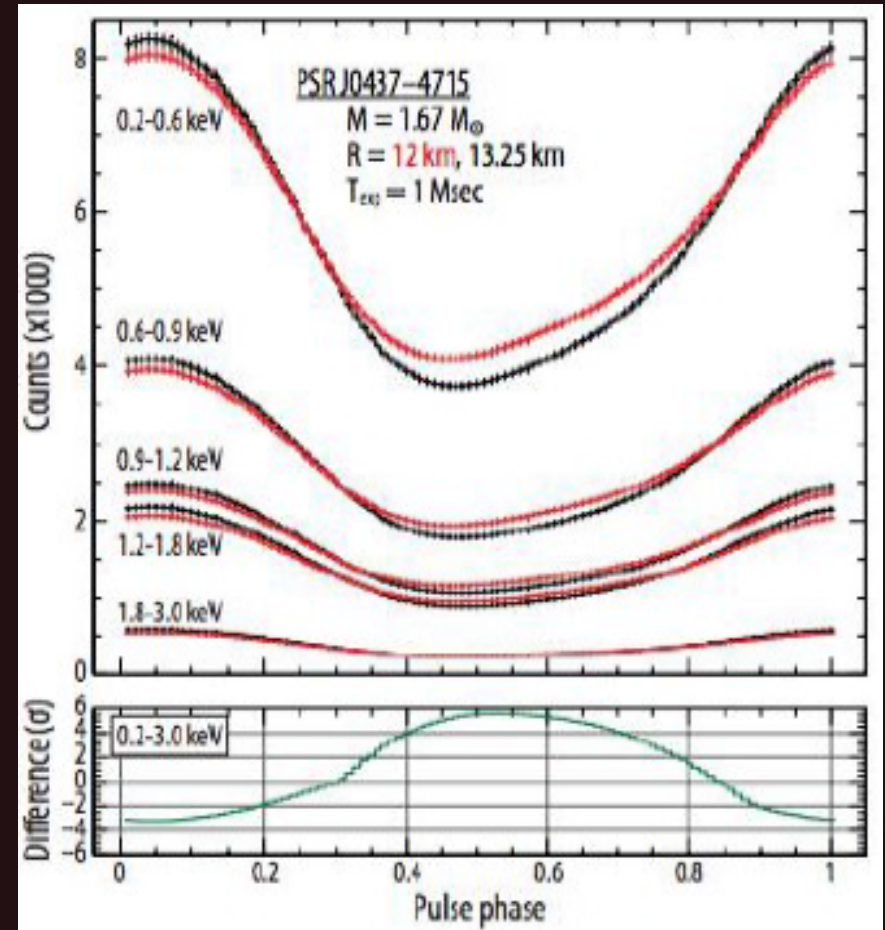
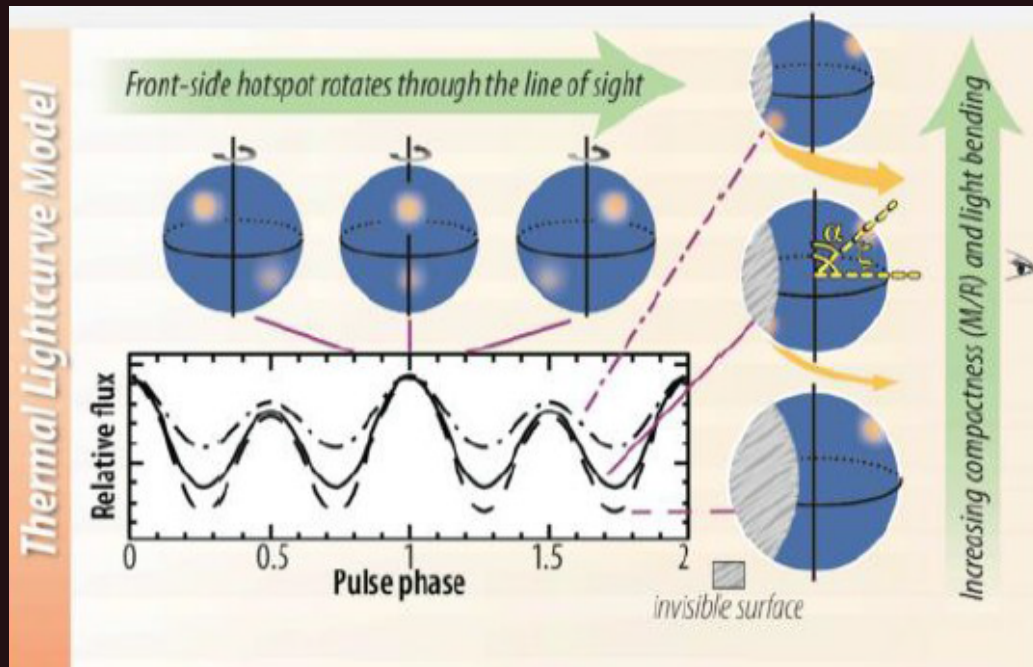


Constraint on maximum mass
 $2.01 < M_{\text{TOV}}/M_{\odot} < 2.16$
 (Rezzolla et al., arxiv:1710.05938)



LVC radius constraint
 GW170817
 (Abbott et al., PRL (2018))
 GW190425
 (Abbott et al., arxiv:2001.01761)
 NICER mass -radius constraint
 PSR J0030+0451
 (Miller et al., ApJLett. (2019))

Measure NS Radii ...

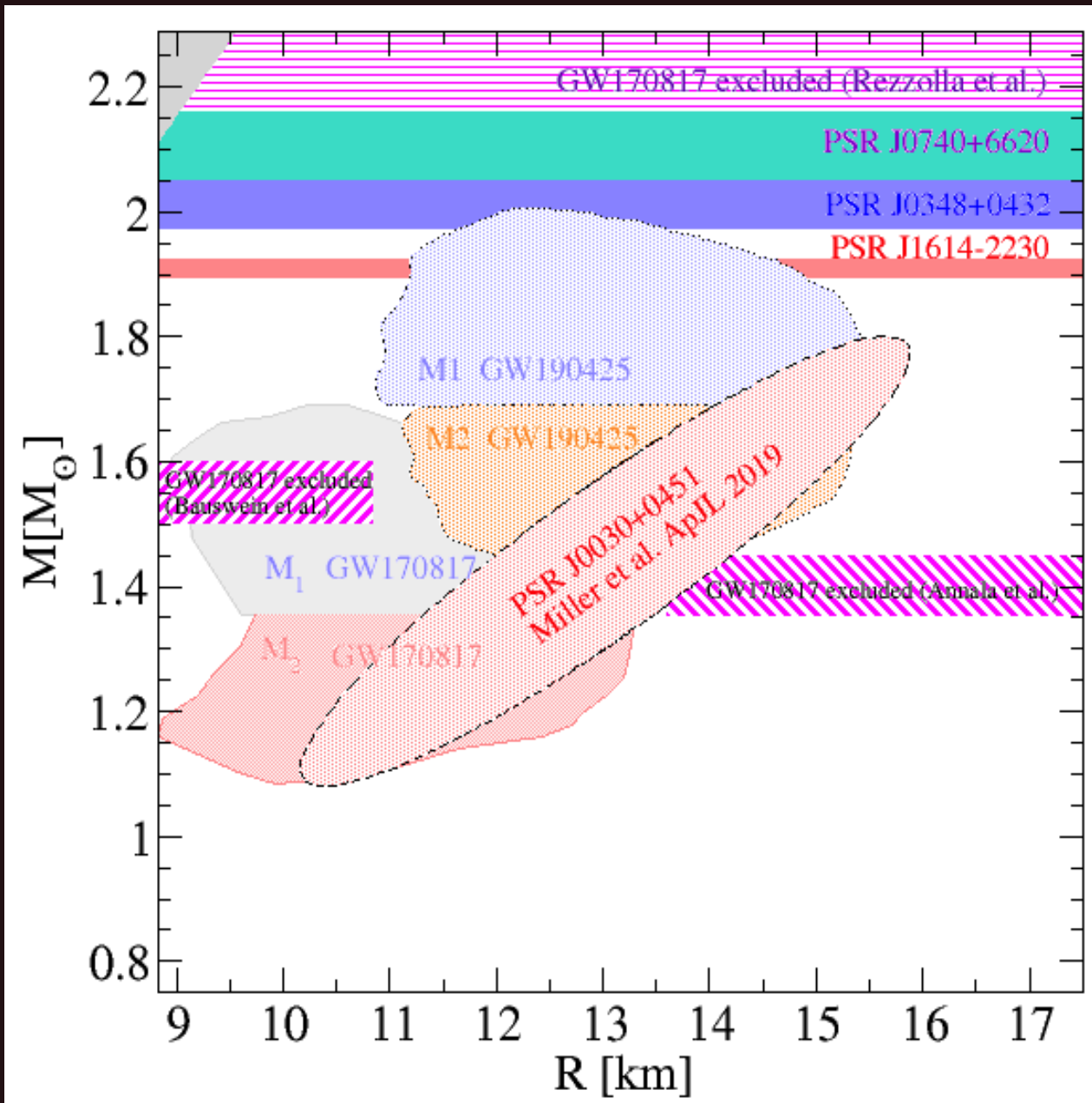


Thermal lightcurves: NS with “hot spots”



K.C. Gendreau et al., Proc. SPIE 8443 (2012) 844313 – first results end of 2019 !!

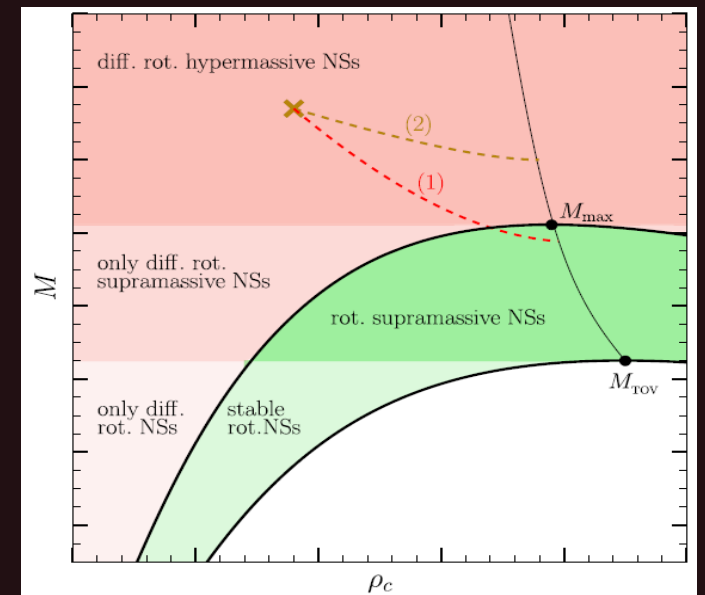
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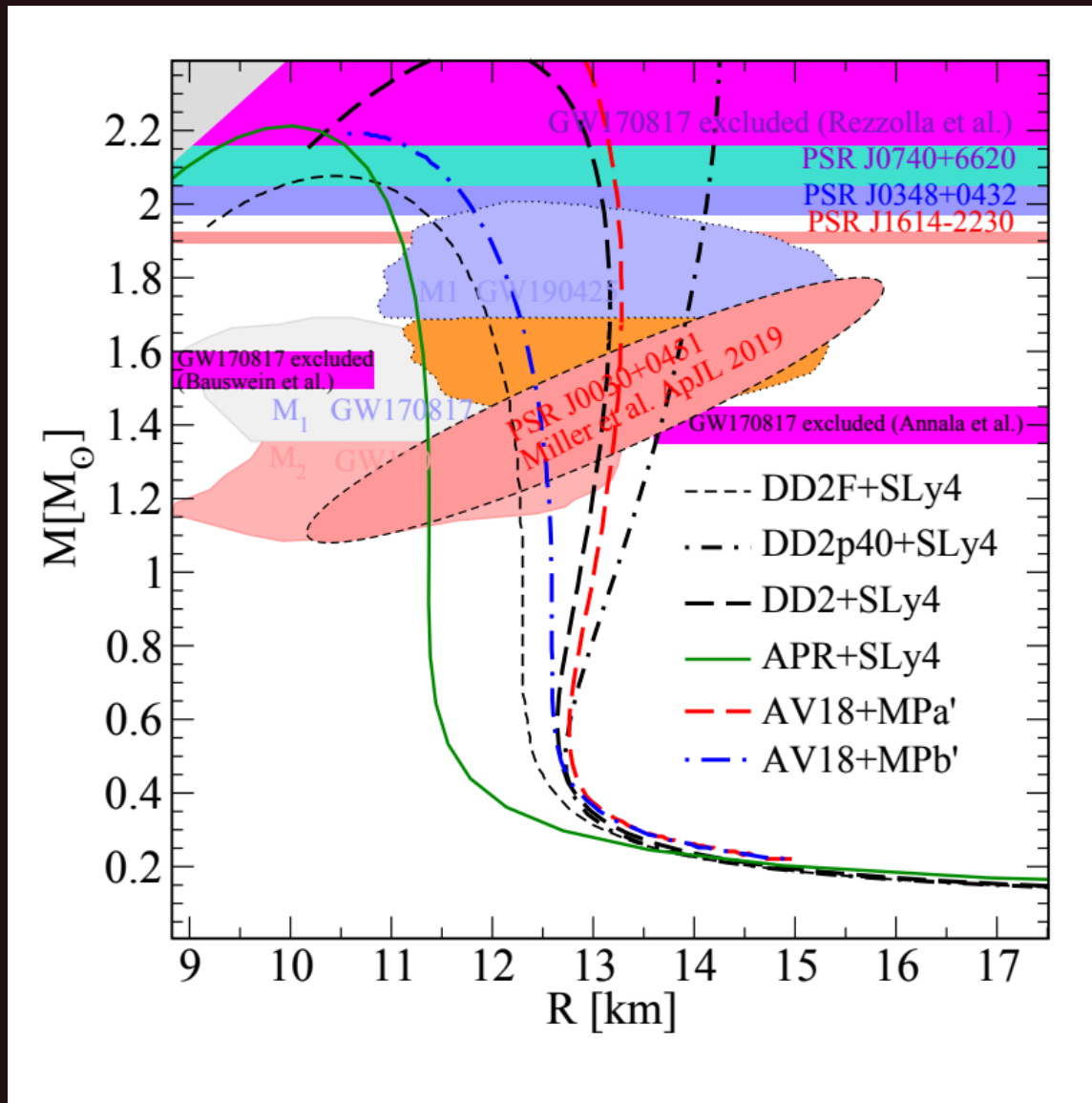
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Constraints on NS mass and radii !

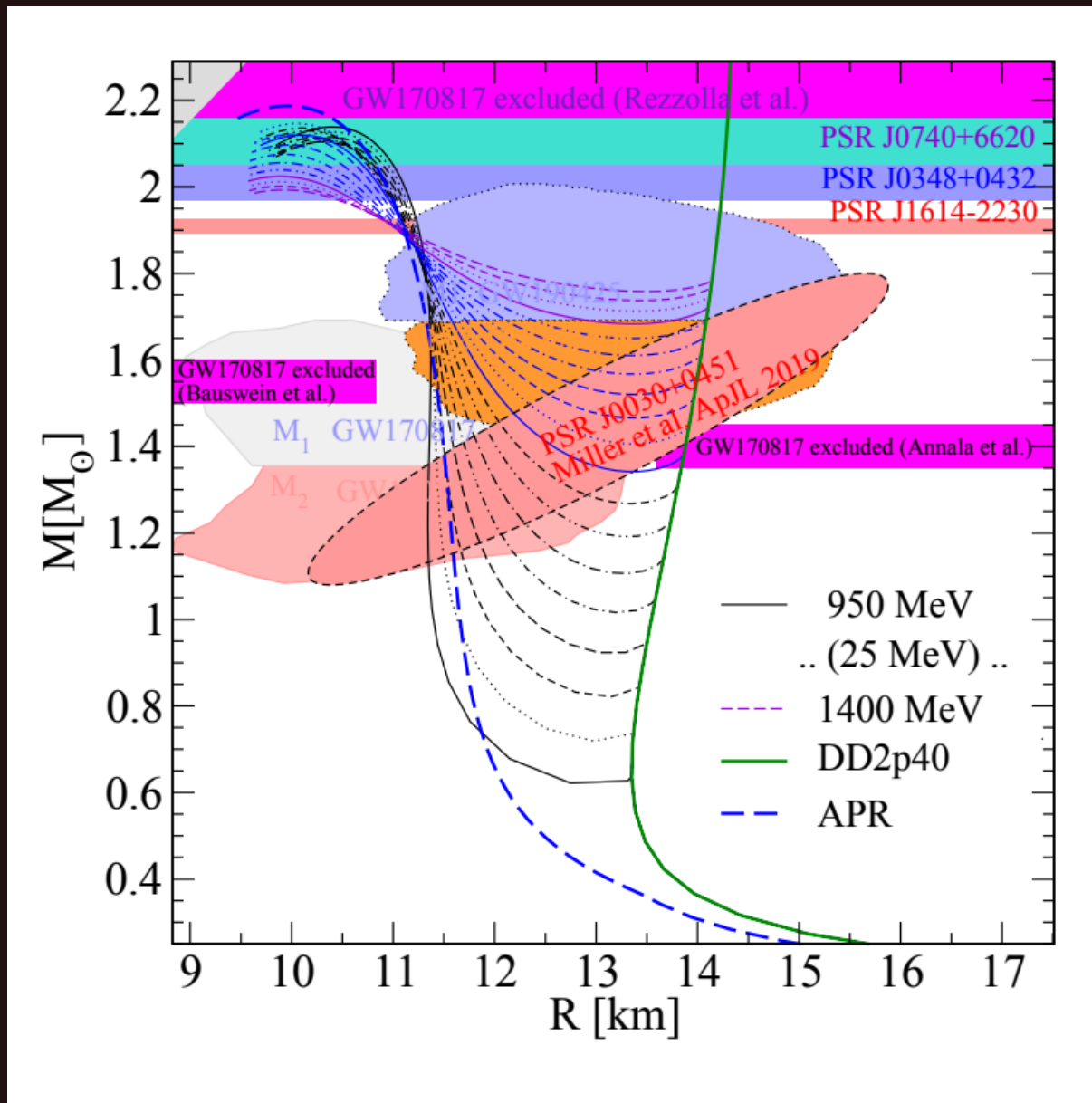


Examples of hadronic EoS
all do fulfill the constraints
 but none of them is applicable for
 Massive stars ($M > 1.5 M_{\text{sun}}$),
 Because of missing hyperons etc.

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AV18*: Yamamoto, Togashi et al., Phys. Rev C 96 (2017) 065804
 DD2*: Typel, Röpke, Klähn, et al., Phys. Rev. C 81 (2010) 015803

Constraints on NS mass and radii !



Examples of hadronic EoS
all do fulfill the constraints
 but **none of them is applicable** for
 Massive stars ($M > 1.5 M_{\text{sun}}$),
 Because of missing hyperons etc.

Which ways out?

- stiff hypernuclear matter
- early onset of deconfinement
 ($M_{\text{onset}} < 1.5 M_{\text{sun}}$)

Old quark matter paradigm:

- deconfinement softens EoS
- hybrid stars compacter, lower M_{max}

LVC radius constraint

GW170817

(Abbott et al., PRL (2018))

GW190425

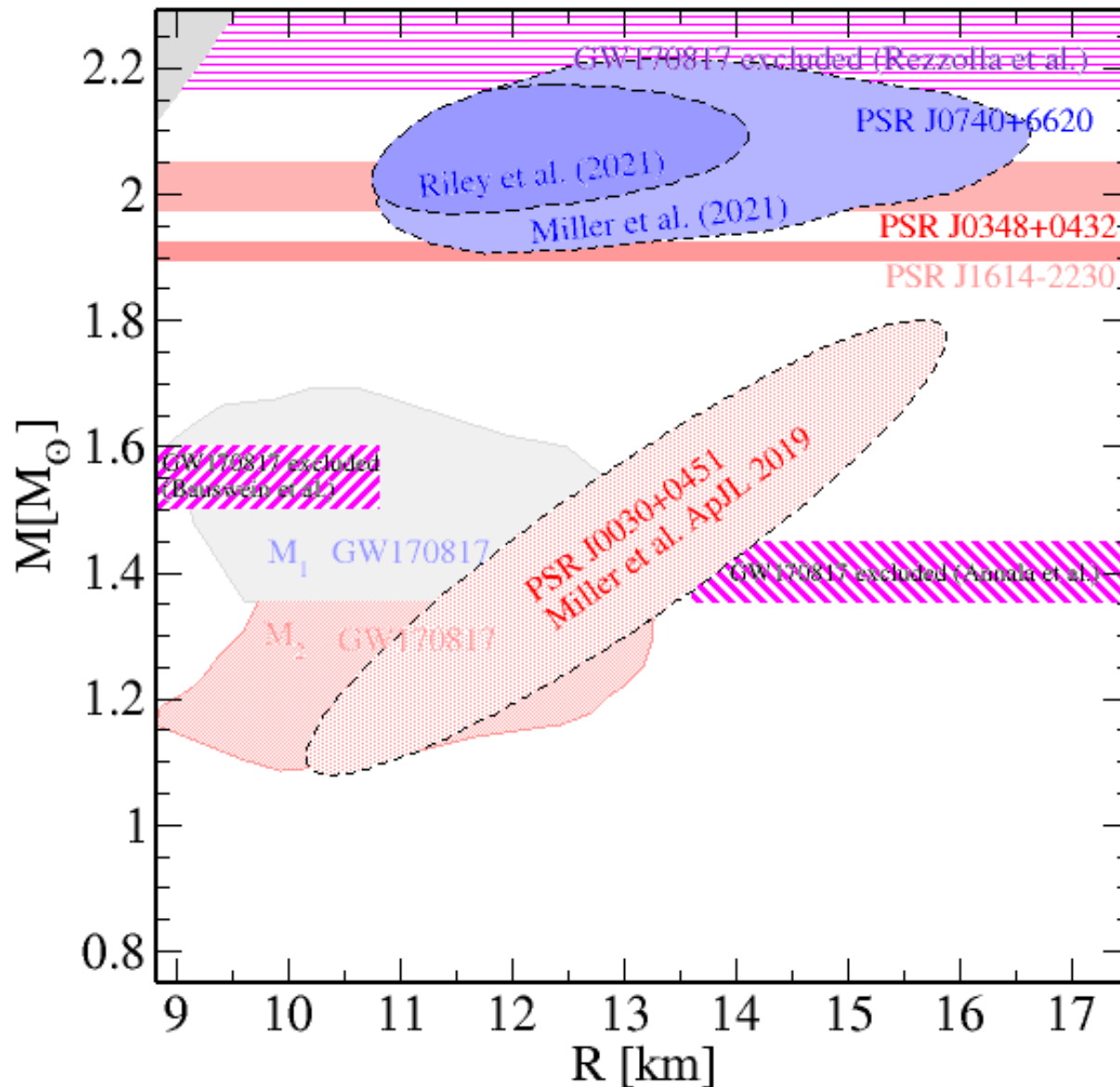
(Abbott et al., arxiv:2001.01761)

NICER mass -radius constraint

PSR J0030+0451

(Miller et al., ApJLett. (2019))

Constraints on NS mass and radii !



New NICER mass-radius data

PSR J0740+6620

(Riley et al., arxiv:2105.06980

Miller et al., arxiv:2105.06979)

Hypernuclear EoS out ?!

→ stiff hypernuclear matter

→ early onset of deconfinement

($M_{\text{onset}} < 1.5 M_{\text{sun}}$)

New quark matter paradigm:

→ deconfinement to stiff QM EoS

→ hybrid stars larger, higher M_{max}

LVC radius constraint

GW170817

(Abbott et al., PRL (2018))

NICER mass -radius constraint

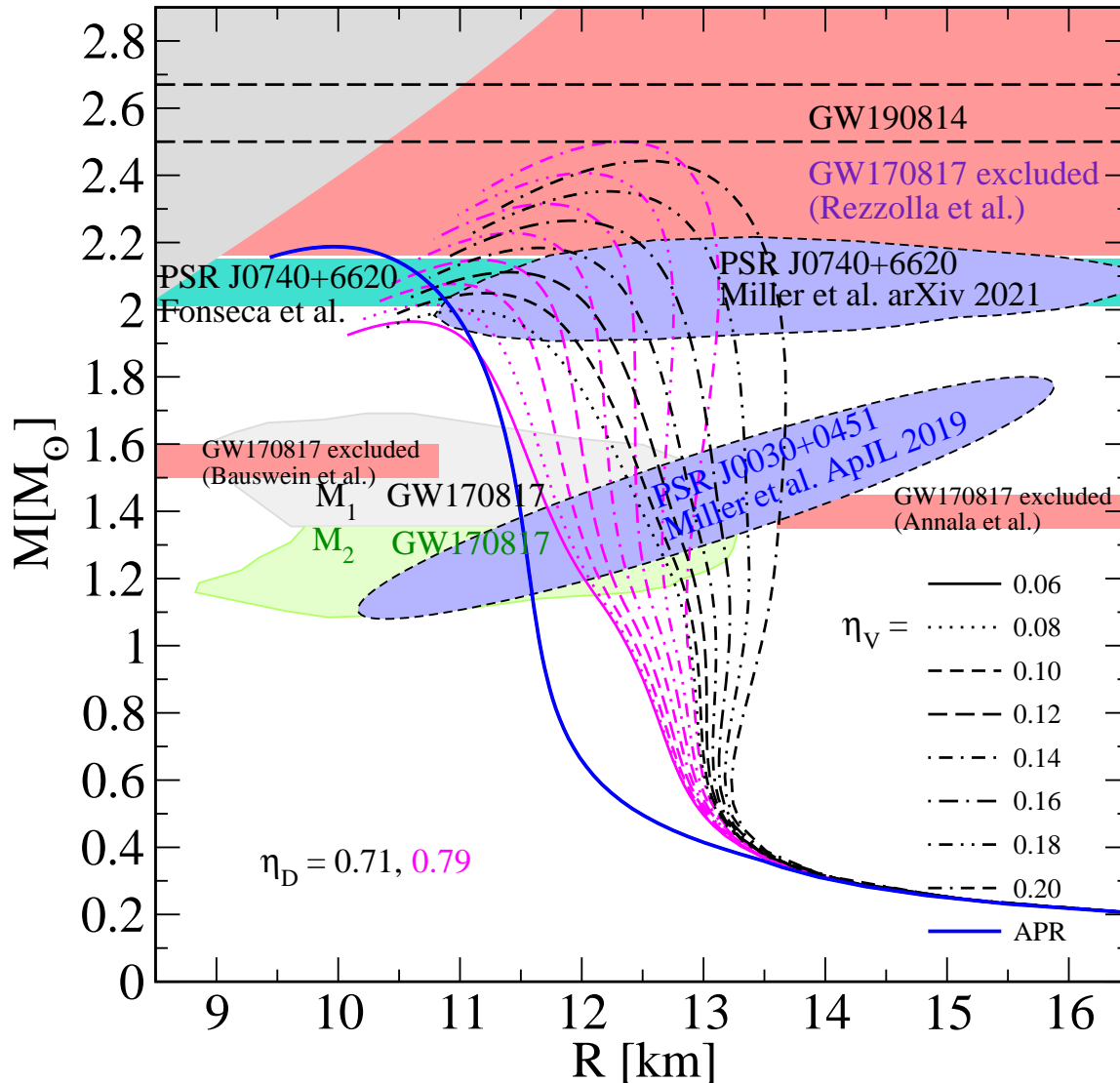
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Constraints on NS mass and radii !



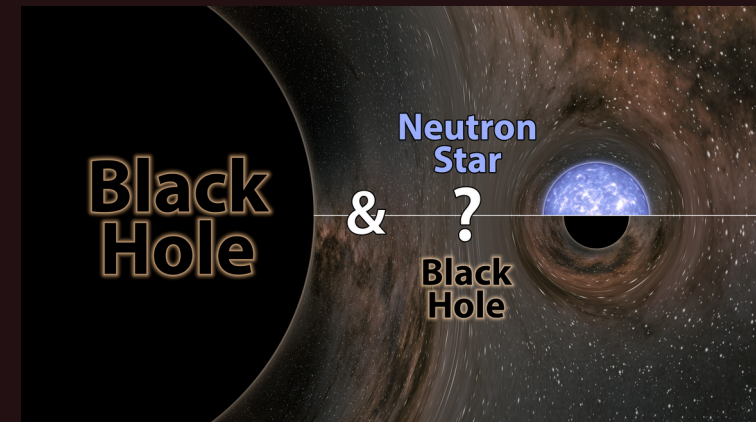
GW190814 - Enigma

Heaviest NS or Lightest BH ??

$$M_1 = 22.2 - 24.3 M_{\odot}$$

$$M_2 = 2.50 - 2.67 M_{\odot} \leftarrow$$

(Abbott et al., ApJL 896:L44 (2020))



LVC radius constraint

GW170817

(Abbott et al., PRL (2018))

NICER mass -radius constraint

PSR J0030+0451

(Miller et al., ApJLett. (2019))

PSR J0740+6620

(Miller et al., arxiv:2105.06979)

NICER radius measurement on PSR J0740+6620

New, large NICER radius for J0740: Riley et al., 2105.06980; Miller et al., 2105.06979

Attention:

Above $\sim 1.5 M_{\text{sun}}$ hyperons
Appear in the center of neutron stars.

Non-hyperonic nuclear EoS (APR)
Are no longer applicable for
High-mass neutron stars $\sim 2M_{\text{sun}}$!

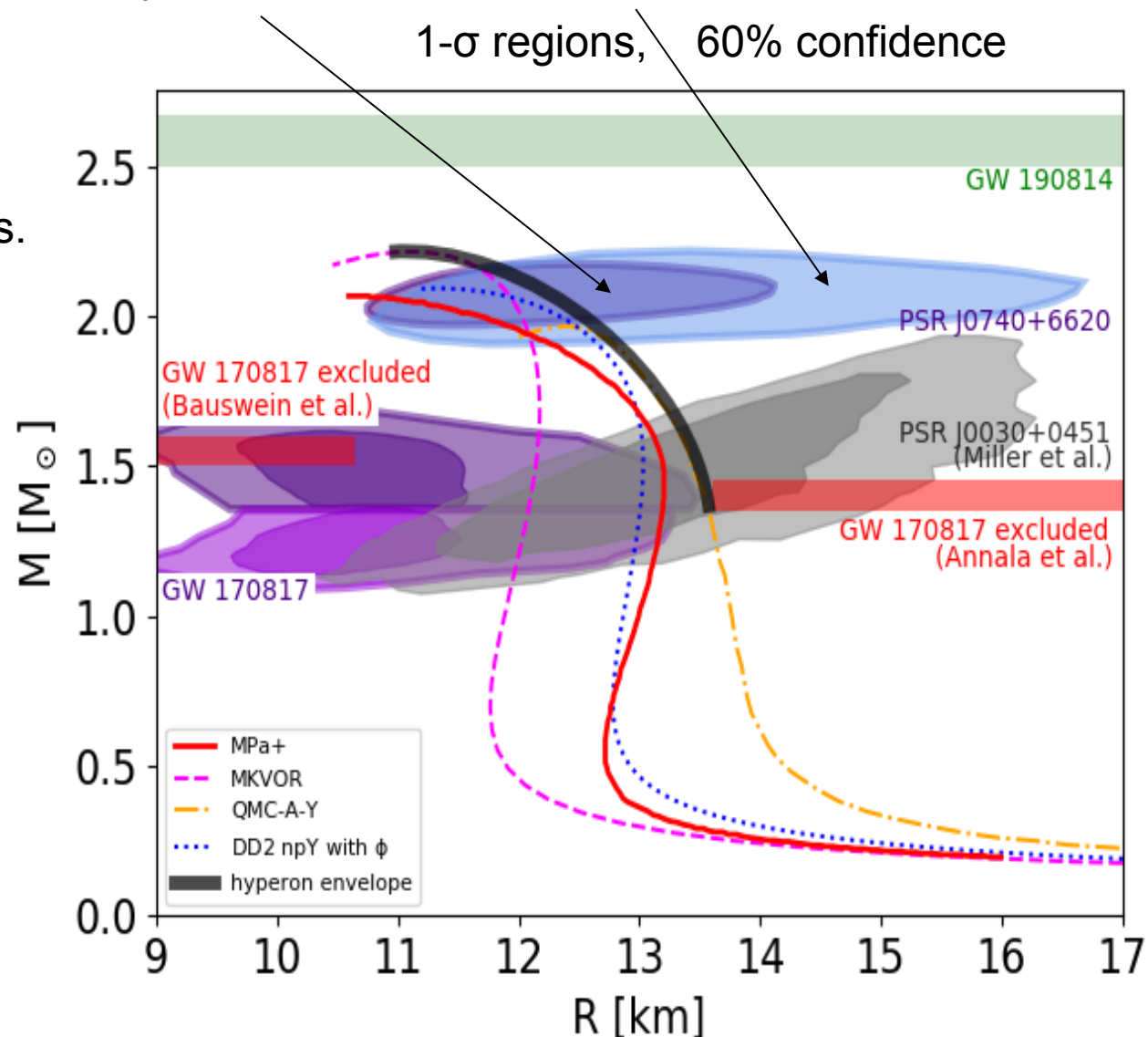
Microscopic EoS need high-density
Stiffening of the hypernuclear EoS,
e.g., by multi-pomeron interactions.

Yamamoto et al., PRC 96 (2017)

Relativistic mean-field EoS have a
Maximal NS radius $R_{2.0} \sim 13$ km

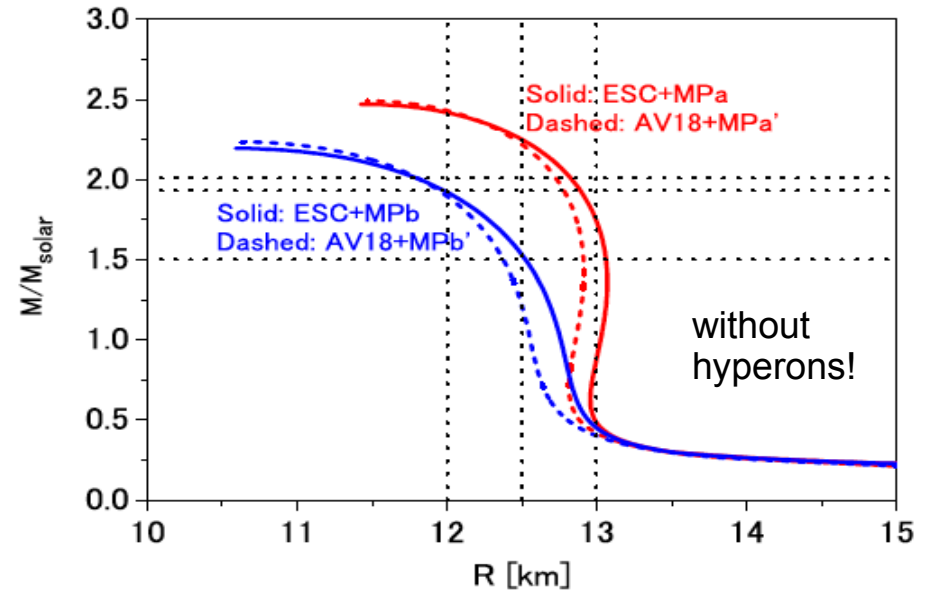
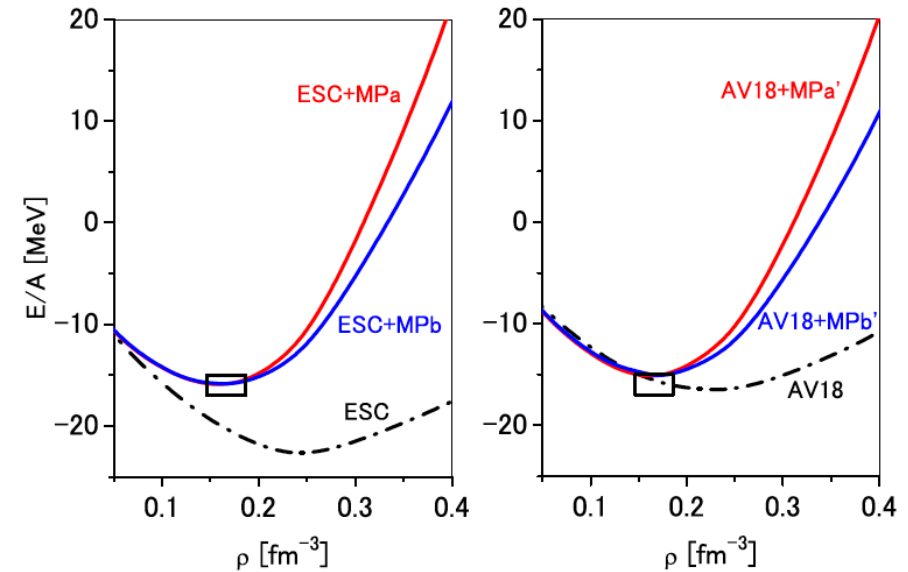
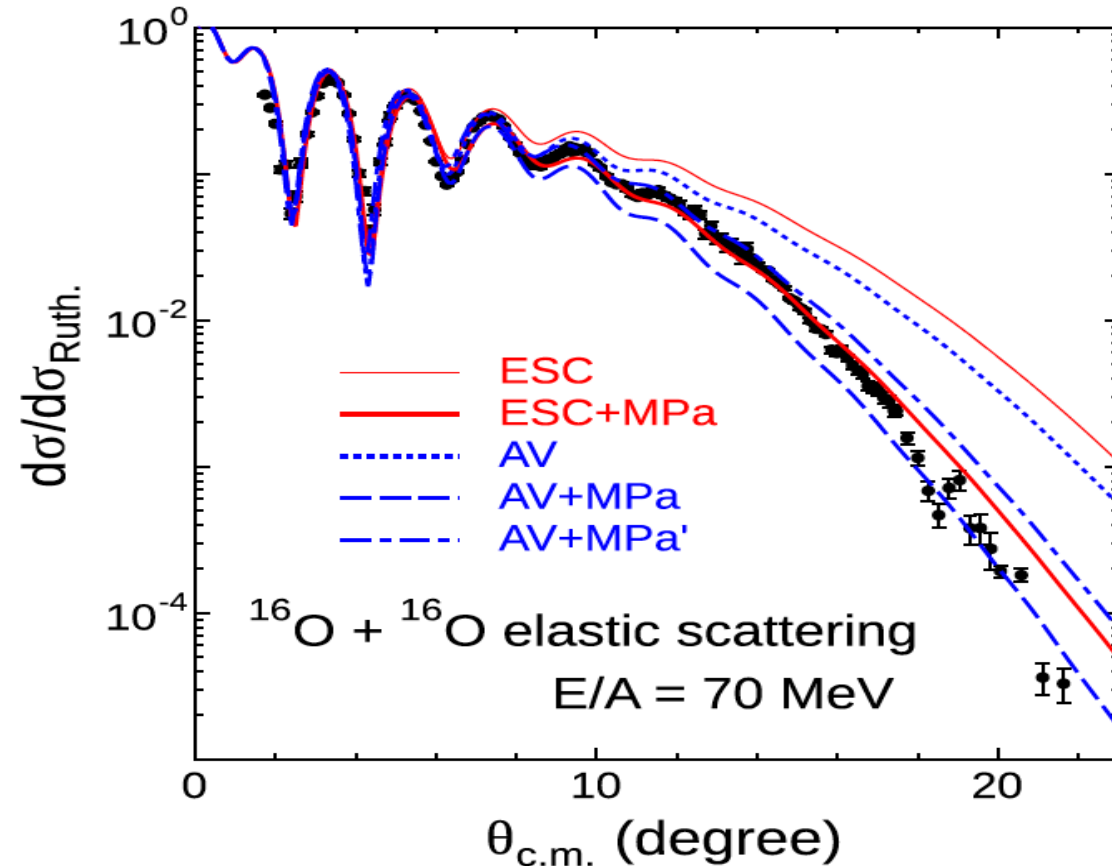
Way out:

early deconfinement to color
superconducting, stiff quark matter !



Shall the APR EoS be abandoned?

Y. Yamamoto, H. Togashi, T. Tamagawa, T. Furumoto, N. Yasutake, T. Rijken, PRC 96 (2017)



Short-range multipomeron exchange potential (MPP) added to AV18 potential gives significant improvement of large-angle scattering cross section (s.a.) and the Nuclear saturation properties, when compared to APR.
 → Neutron star radii $R(M < 2 M_{\text{sun}}) > 12 \text{ km} !!$

What is the special point? What are its properties?

The TOV equation

$$\frac{\partial P(r)}{\partial r} = - \frac{\epsilon(r)M(r) \left(1 + \frac{P(r)}{\epsilon(r)}\right) \left(1 + \frac{4\pi r^3 P(r)}{M(r)}\right)}{r^2 \left(1 - \frac{2M(r)}{r}\right)}, \quad \frac{\partial M(r)}{\partial r} = 4\pi r^2 \epsilon(r).$$

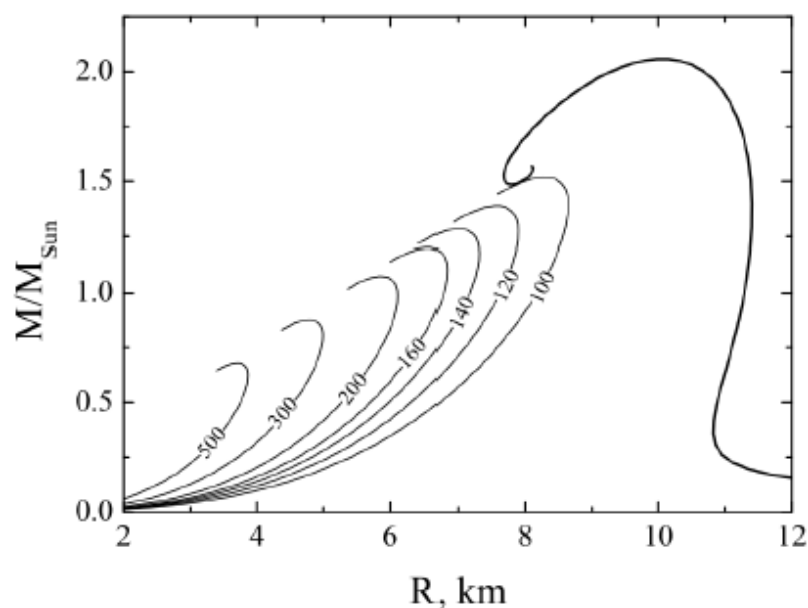


Fig. 1. Mass-radius diagram for a star made of ordinary matter (thick line) and purely quark stars (thin lines). The numbers at the lines indicate the parameter B .

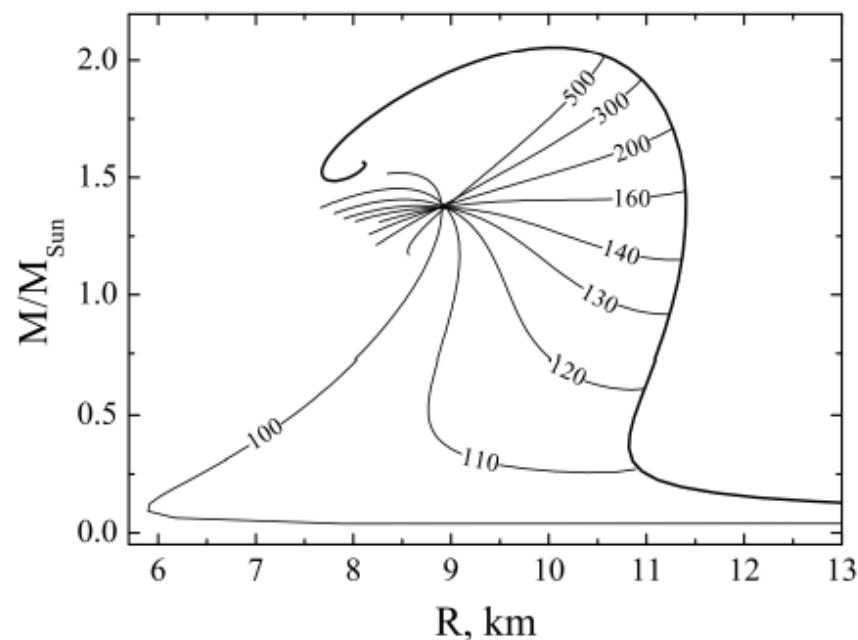


Fig. 2. Mass-radius diagram of hybrid stars for various values of the parameter B

What is the special point? What are its properties?

The constant-speed-of-sound (CSS) model:

– dimensionless baryochemical potential

$$\hat{\mu}_B = \frac{\mu_B}{\mu_{scale}} = \left(\frac{p+B}{A} \right)^{1/(1+\beta)},$$

– pressure

$$p(\mu_B) = A\hat{\mu}_B^{1+\beta} - B,$$

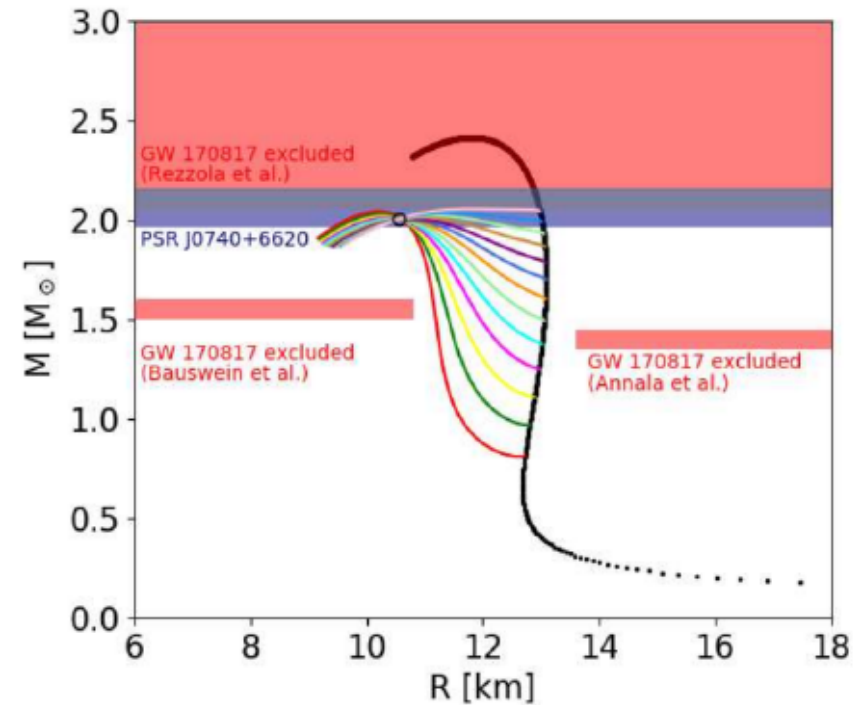
– baryon density

$$n_B(\mu_B) = (1 + \beta) \frac{A}{\mu_{scale}} \hat{\mu}_B^\beta,$$

– energy density

$$\epsilon = B + \beta A \hat{\mu}_B^{1+\beta},$$

– $p(\epsilon)$ relation: $\epsilon = \beta p + (1 + \beta)B$.



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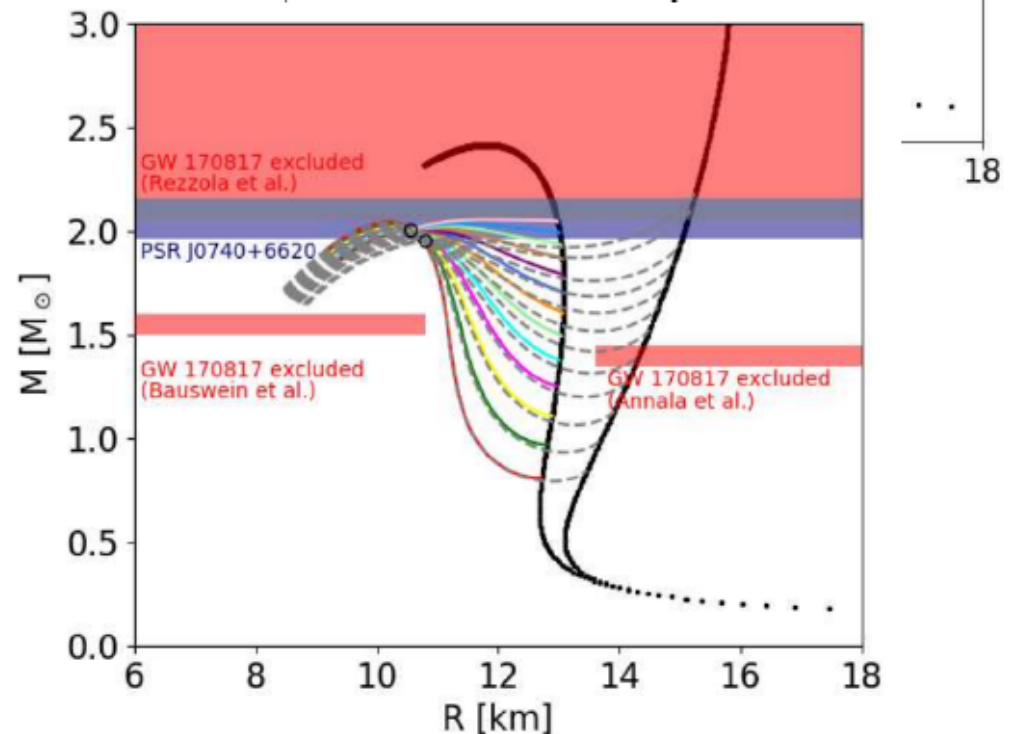
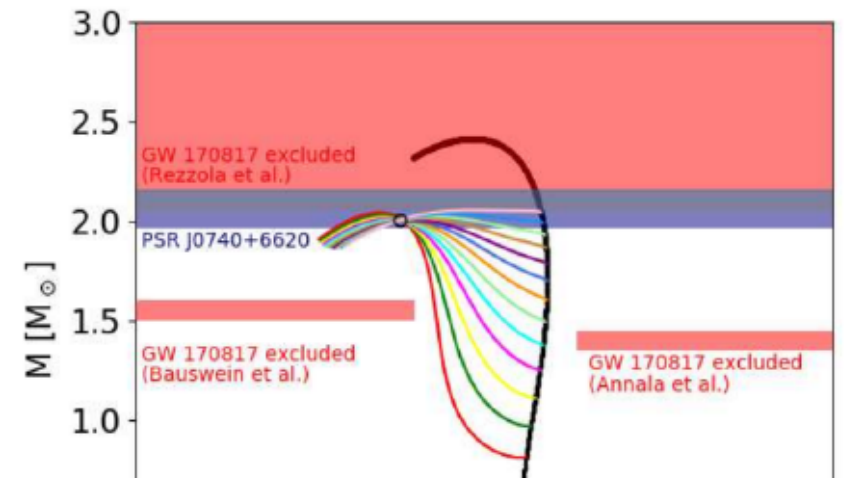
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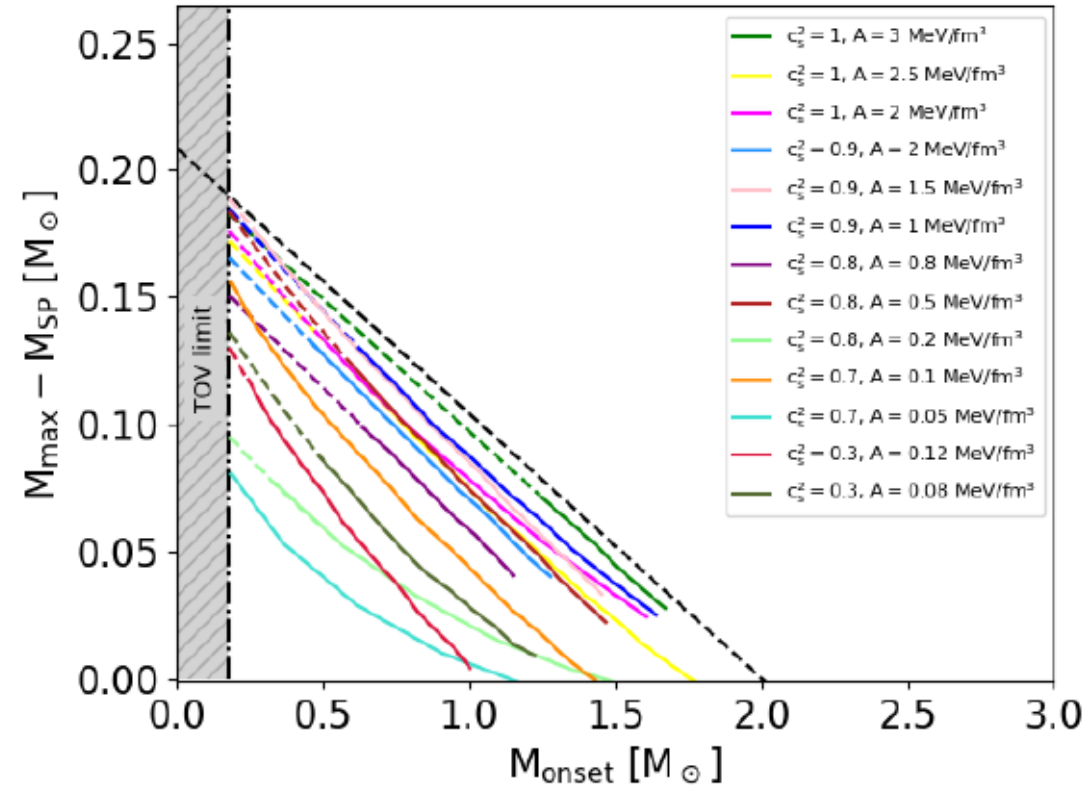
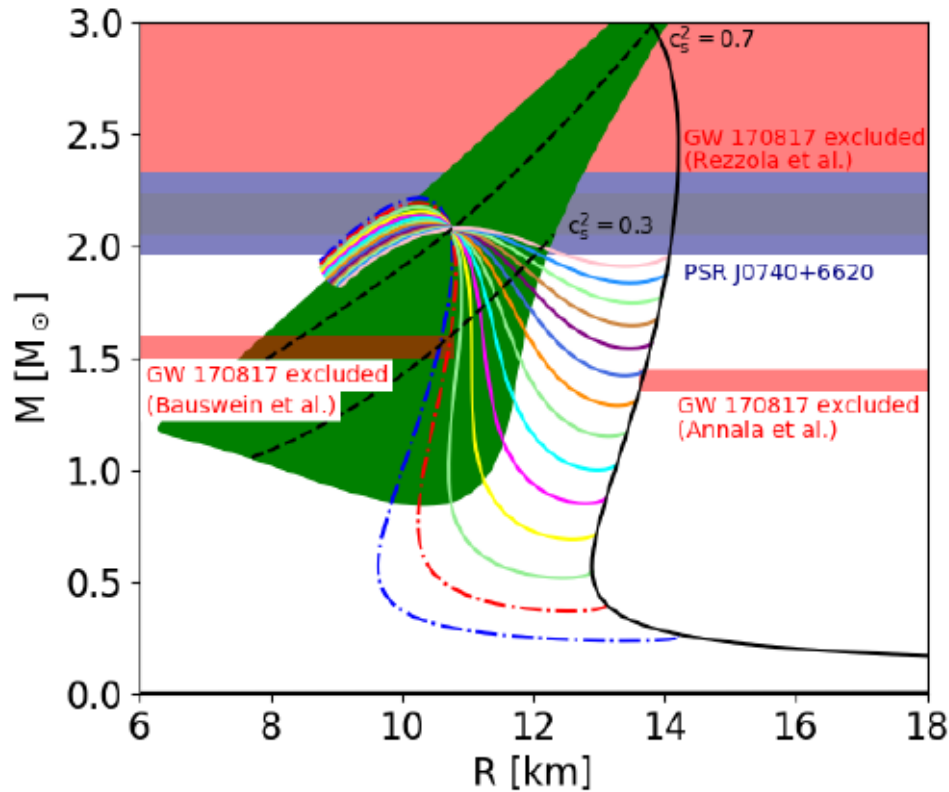
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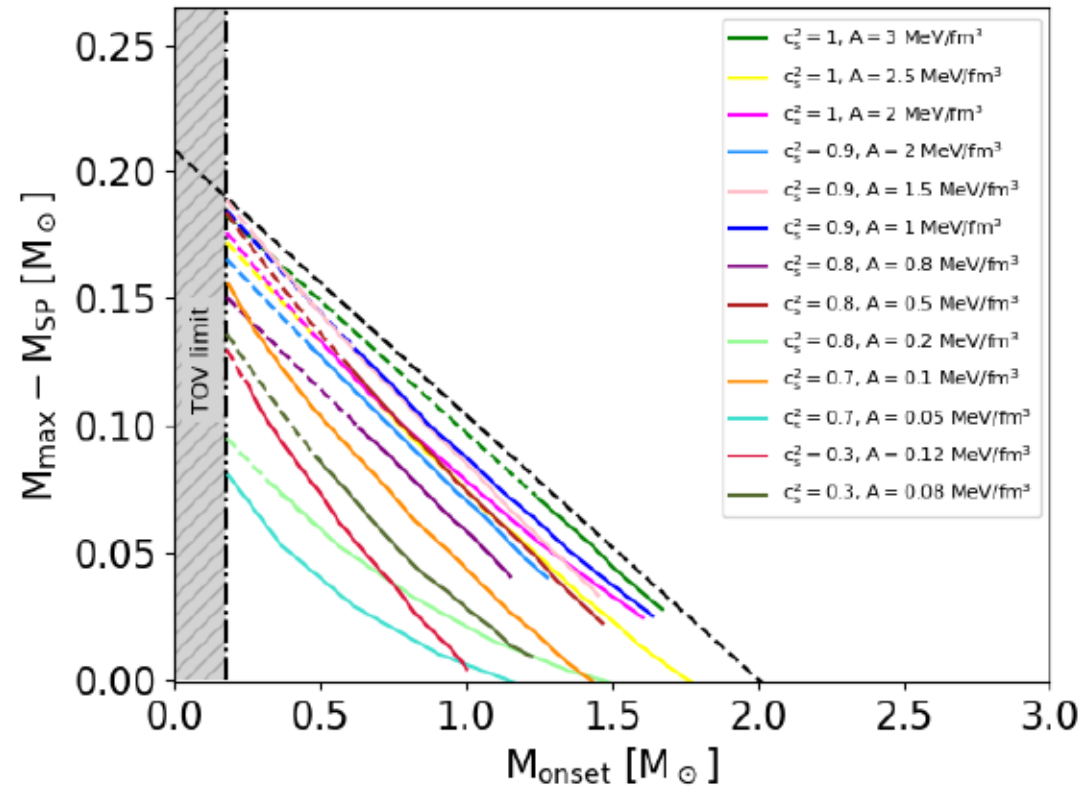
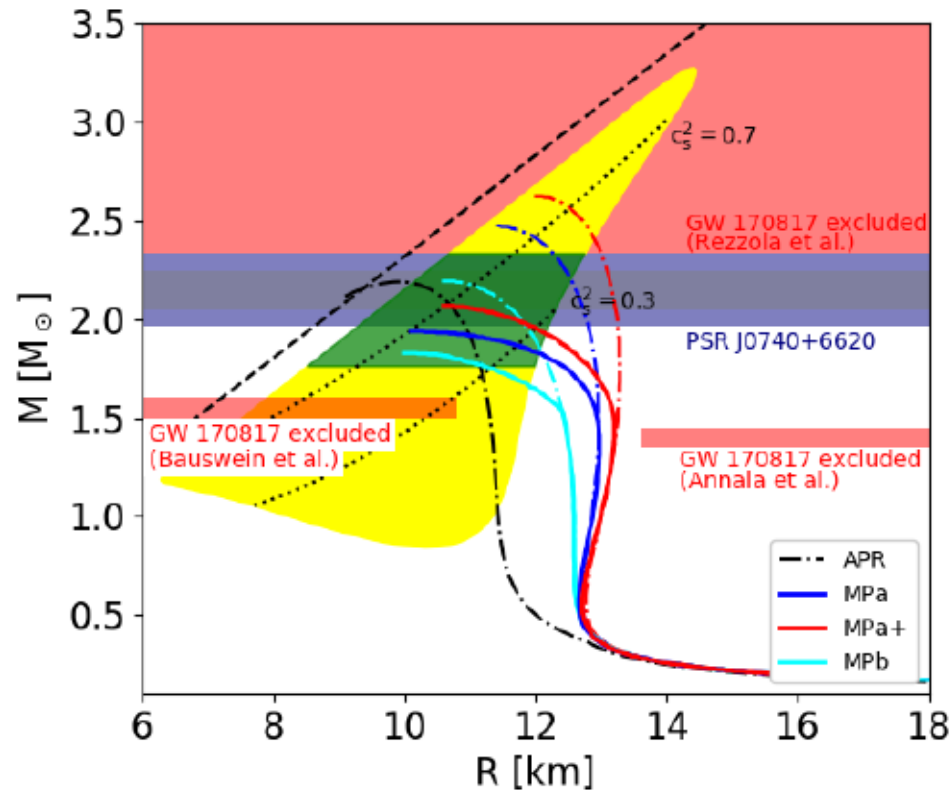


What is the special point? What are its properties?



$$M_{max} = M_{SP} + 0.208M_{\odot} - 0.104M_{onset}$$

What is the special point? What are its properties?



$$M_{max} = M_{SP} + 0.208M_{\odot} - 0.104M_{onset}$$

Dependence on the phase transition construction?

The mixed phase parabolic ansatz:

$$P_M(\mu) = \alpha_2(\mu - \mu_c)^2 + \alpha_1(\mu - \mu_c) + (1 + \Delta_p)P_c,$$

Gibbs condition for phase equilibrium:

$$\begin{aligned} P_H(\mu_H) &= P_M(\mu_H), \\ P_Q(\mu_Q) &= P_M(\mu_Q), \\ \frac{\partial^k}{\partial \mu^k} P_H(\mu_H) &= \frac{\partial^k}{\partial \mu^k} P_M(\mu_H), \\ \frac{\partial^k}{\partial \mu^k} P_Q(\mu_Q) &= \frac{\partial^k}{\partial \mu^k} P_M(\mu_Q). \end{aligned}$$

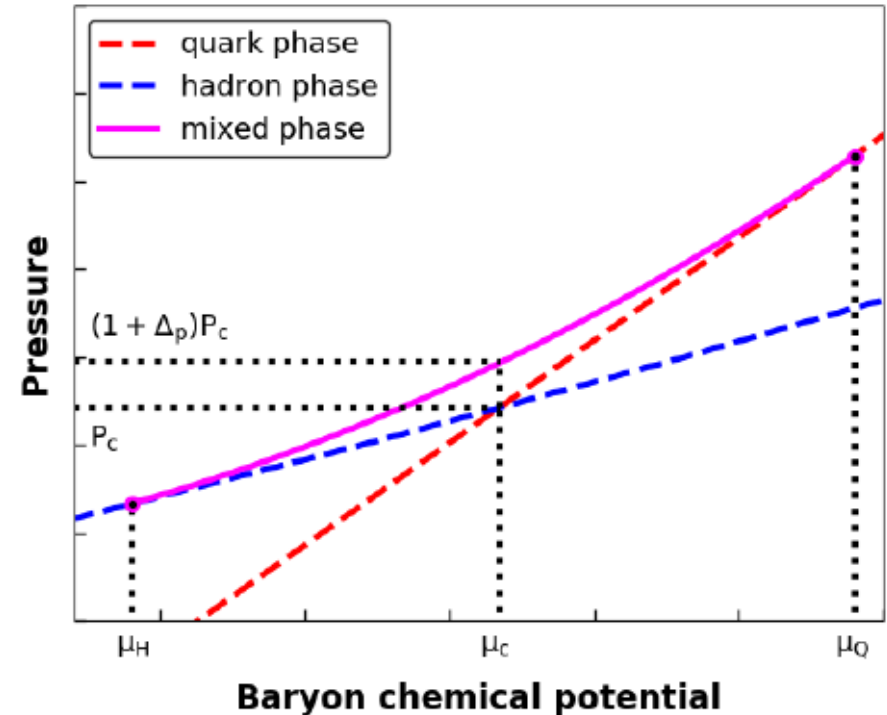
Derived parameters ($k = 1$):

$$\alpha_1 = \frac{-2\kappa_1 + \kappa_2(\mu_c - \mu_H)}{2(\mu_c - \mu_Q)(\mu_H - \mu_Q)},$$

$$\alpha_2 = \frac{-2\kappa_1 + \kappa_2(\mu_c - \mu_Q)}{2(\mu_c - \mu_H)(\mu_H - \mu_Q)},$$

$$\kappa_1 = n_Q(\mu_c - \mu_Q) - n_H(\mu_c - \mu_H) + P_Q - P_H,$$

$$\kappa_2 = n_Q - n_H.$$



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$$P_Q(\mu_Q) = P_M(\mu_Q),$$

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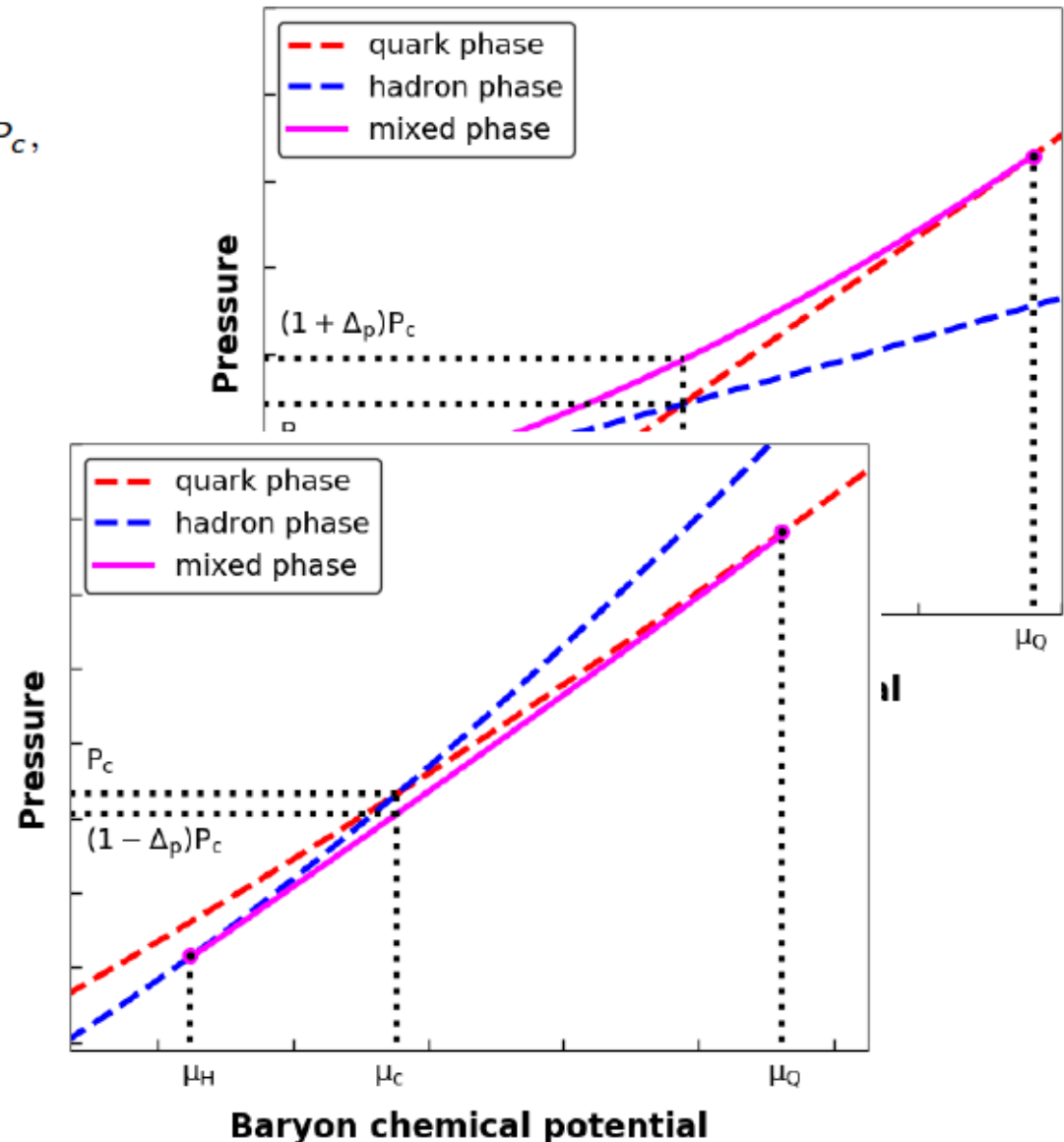
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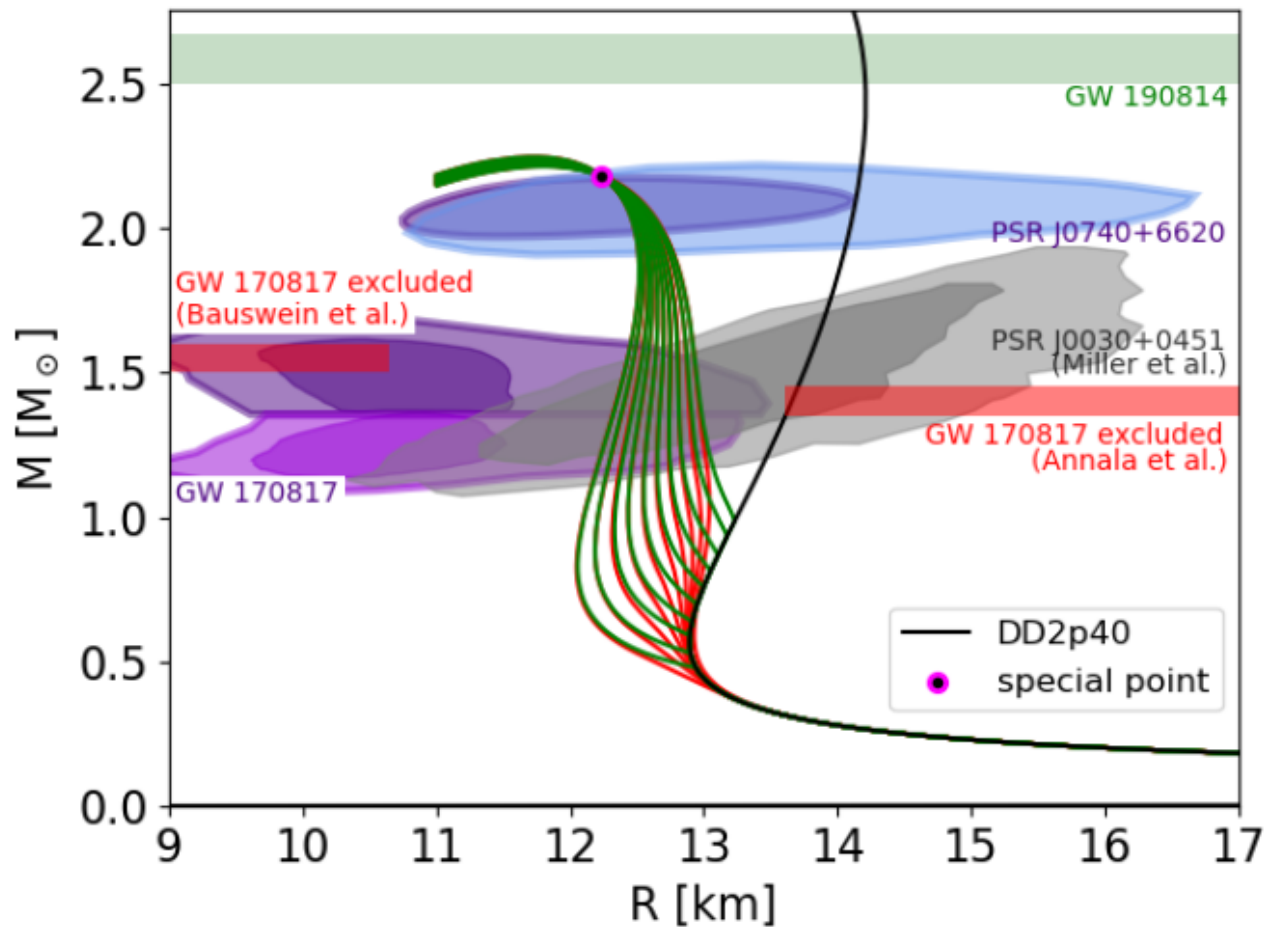
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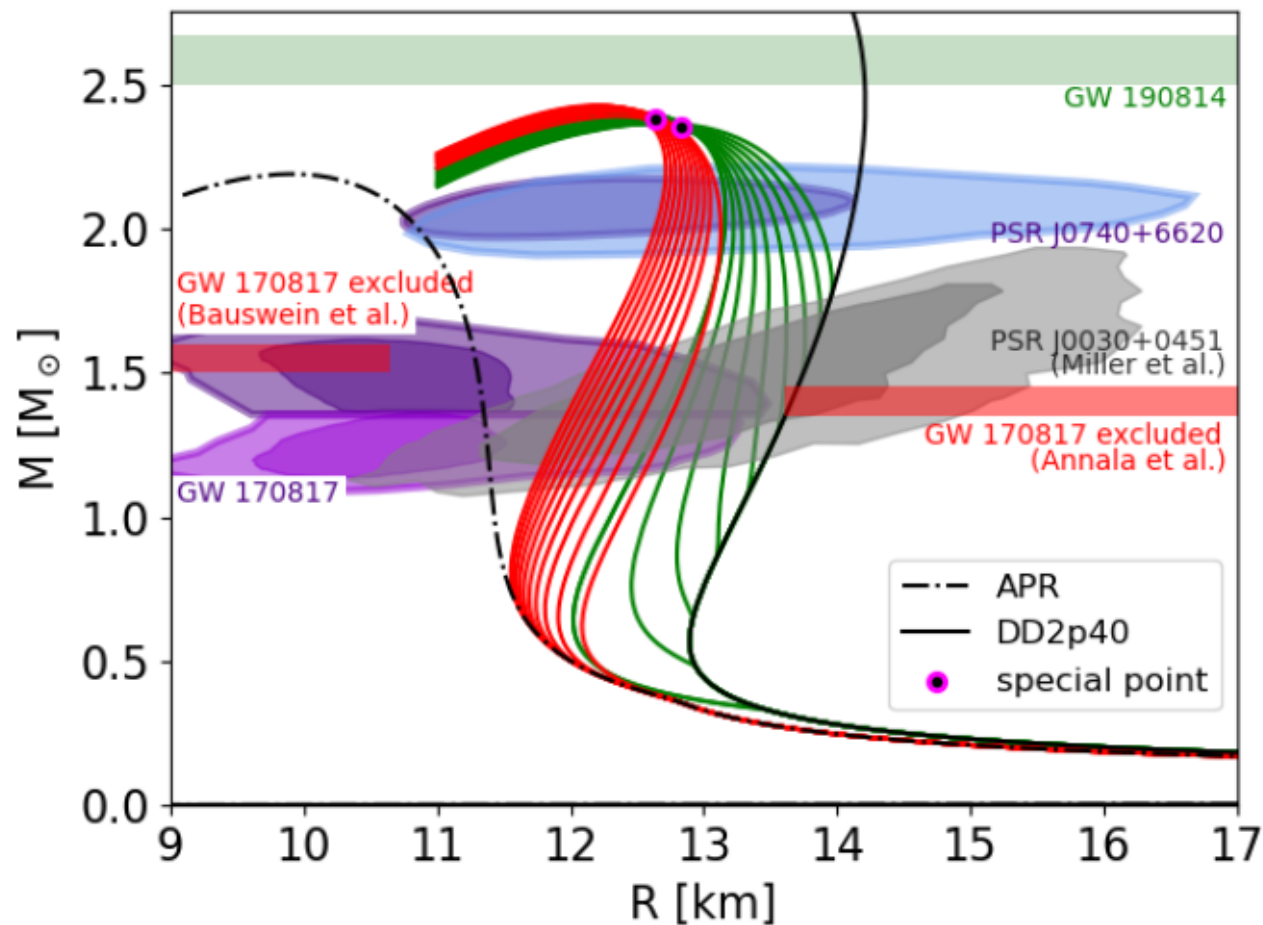
Dependence on the phase transition construction?

Invariance w.r.t. Maxwell \rightarrow mixed phase construction (pasta phases)



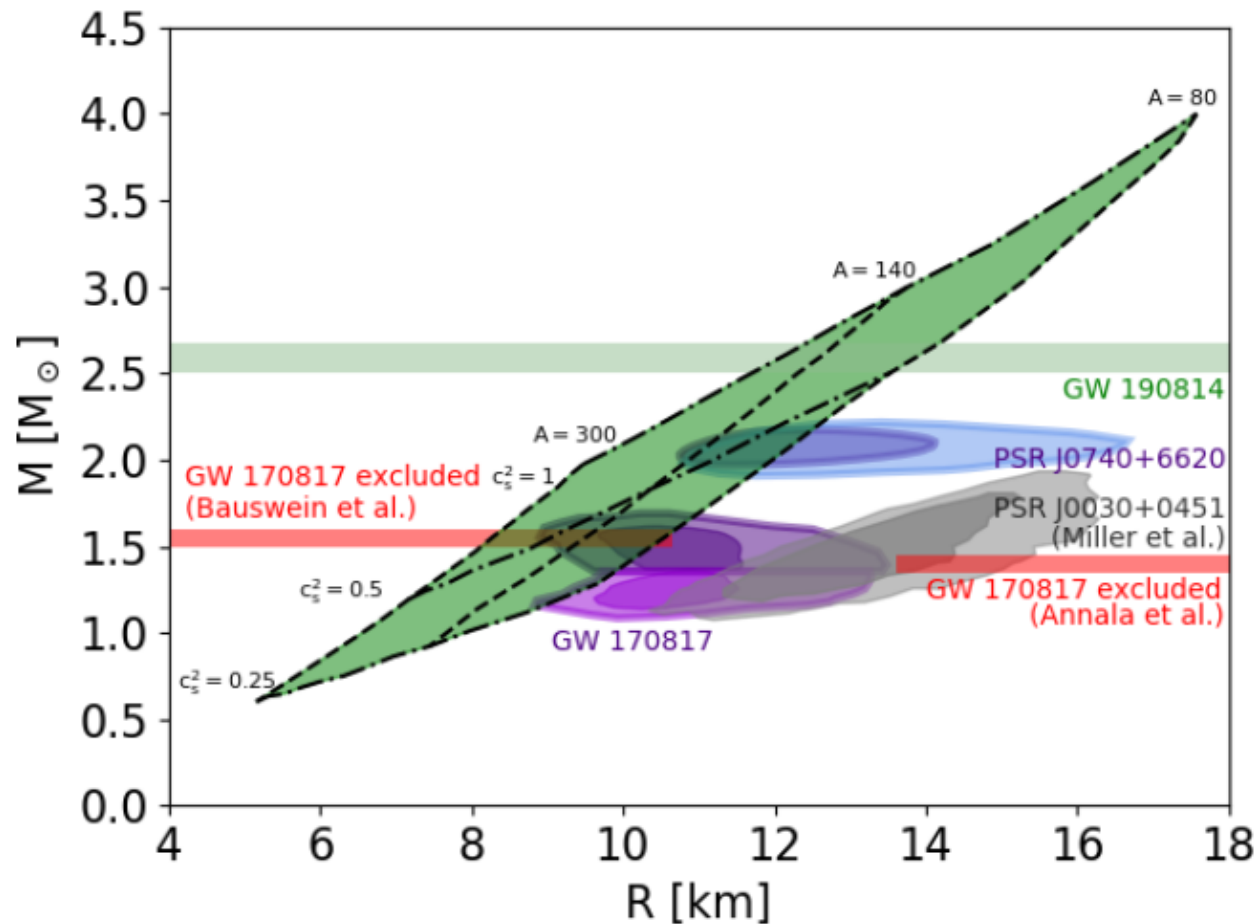
Dependence on the phase transition construction?

Invariance w.r.t. Maxwell \rightarrow interpolation construction (soft - stiff transition)



Mapping the special point locations ...

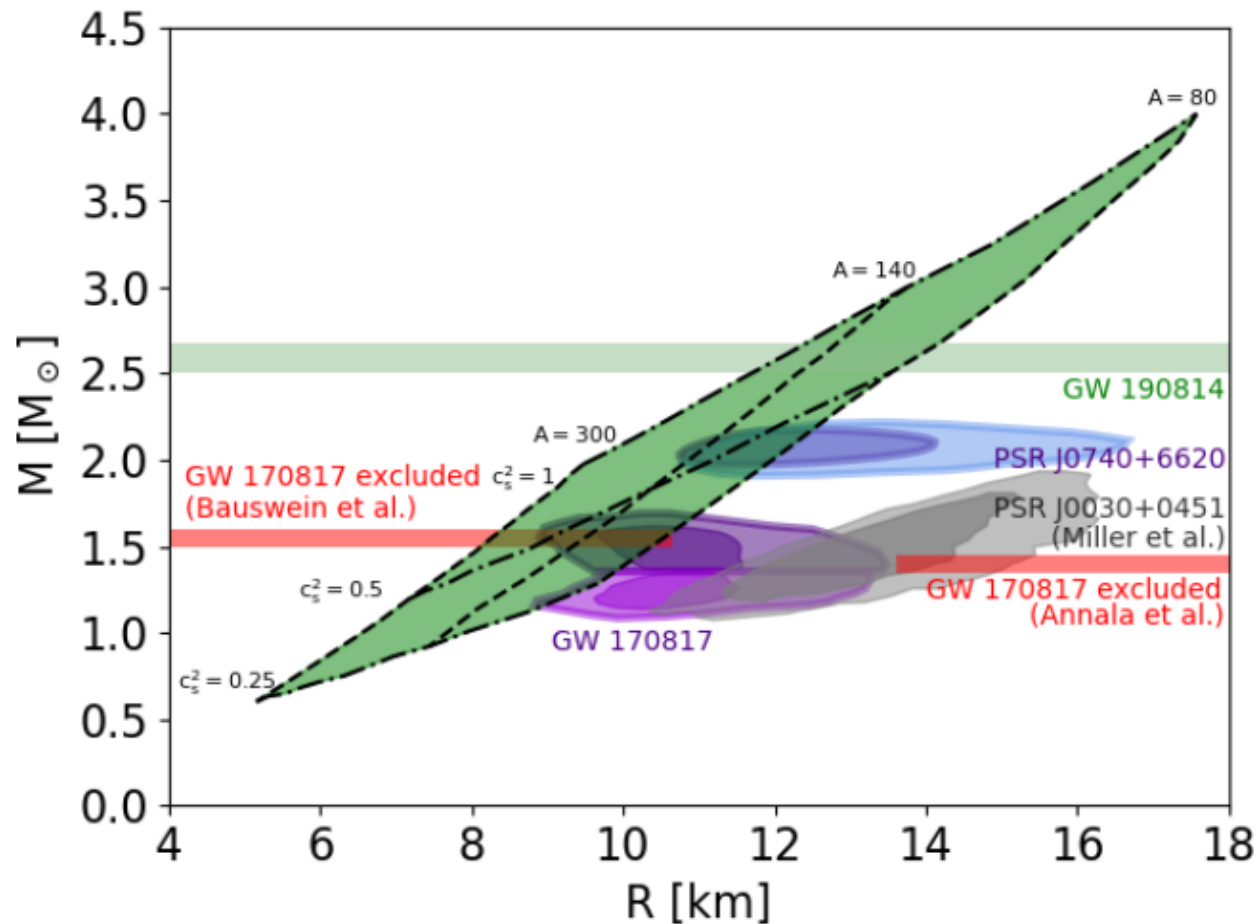
Special point locations for constant sound speed c_s



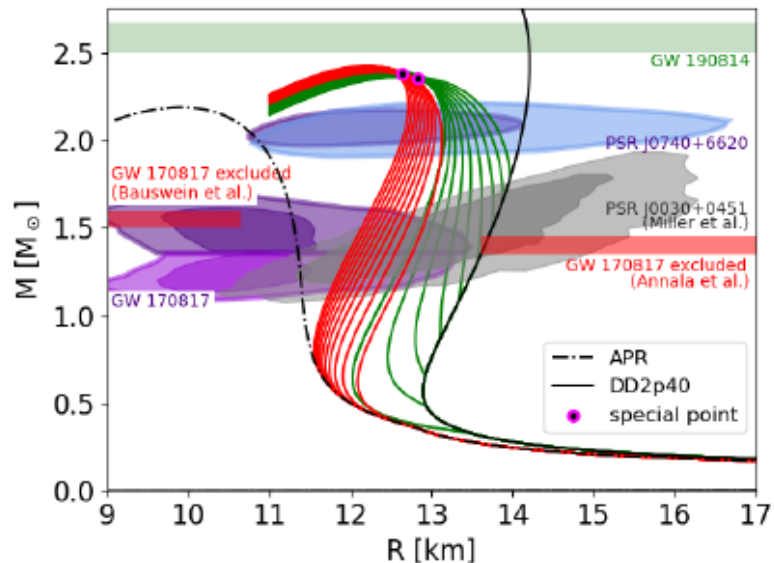
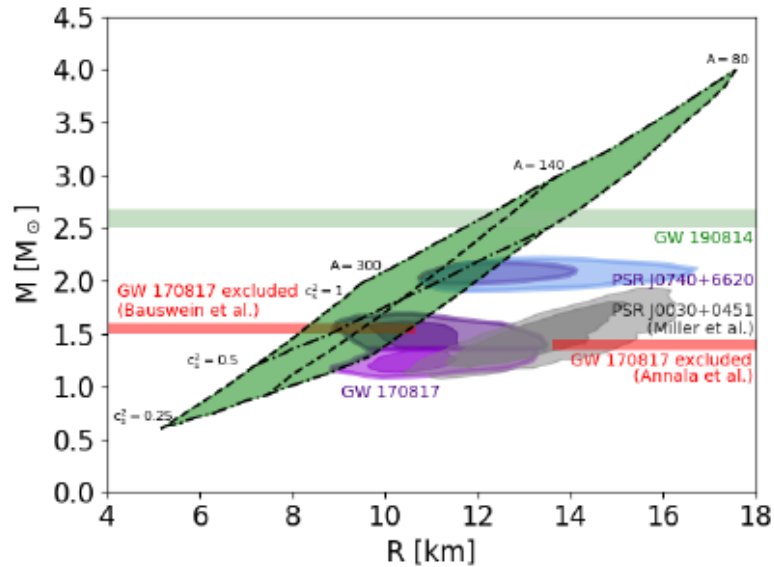
⁵Cierniak, Blaschke, Astron. Nachr. (2021), arXiv: 2106.06986

Mapping the special point locations ...

Special point locations for constant sound speed c_s
... and constant prefactor A



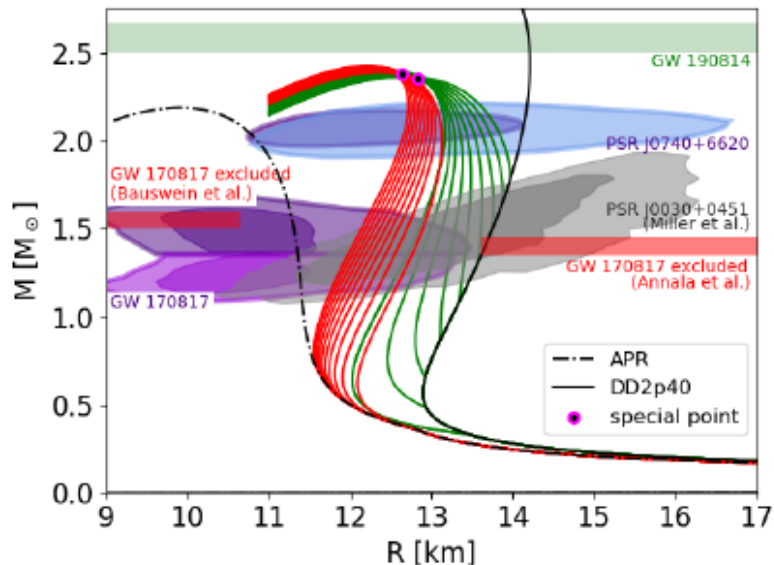
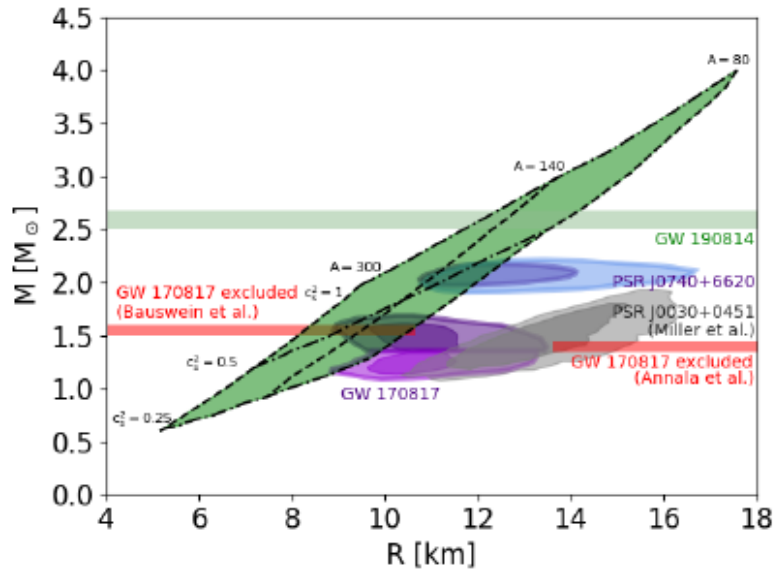
Mapping the special point locations ...



c_s^2	M_{SP} [M_\odot]	R_{min} [km]	R_{max} [km]
0.35	1.82	-	-
0.40	2.07	12.18	12.29
0.45	2.30	11.84	13.41
0.50	2.50	11.56	13.91
0.55	2.68	11.30	14.20
0.60	2.86	11.05	14.45
0.70	3.22	10.67	14.67
1.00	4.00	9.95	14.84

The values of c_s^2 , largest possible M_{SP} and the radii range ($R_{min} - R_{max}$) of a $2 M_\odot$ hybrid star.

Mapping the special point locations ...



c_s^2	M_{SP} [M_\odot]	R_{min} [km]	R_{max} [km]
0.35	1.82	-	-
0.40	2.07	12.18	12.29
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1.00	4.00	9.95	14.84

The values of c_s^2 , largest possible M_{SP} and the radii range ($R_{min} - R_{max}$) of a $2 M_\odot$ hybrid star. **Bold red rows correspond to the nINJL fit from [6].**

Constant sound speed (CSS) vs. nonlocal NJL model

$$\mathcal{L} = \bar{\psi} (-i\not{\partial} + m_c) \psi - \frac{G_S}{2} j_S^f j_S^f - \frac{G_D}{2} [j_D^a]^\dagger j_D^a + \frac{G_V}{2} j_V^\mu j_V^\mu, \quad \eta_D = G_D/G_S \text{ and } \eta_V = G_V/G_S$$

Nonlocal currents, formfactor $g(z)$

$$j_S^f(x) = \int d^4 z g(z) \bar{\psi}(x + \frac{z}{2}) \Gamma^f \psi(x - \frac{z}{2}),$$

$$j_D^a(x) = \int d^4 z g(z) \bar{\psi}_C(x + \frac{z}{2}) i\gamma_5 \tau_2 \lambda^a \psi(x - \frac{z}{2}),$$

$$j_V^\mu(x) = \int d^4 z g(z) \bar{\psi}(x + \frac{z}{2}) i\gamma_\mu \psi(x - \frac{z}{2}),$$

CSS equation of state

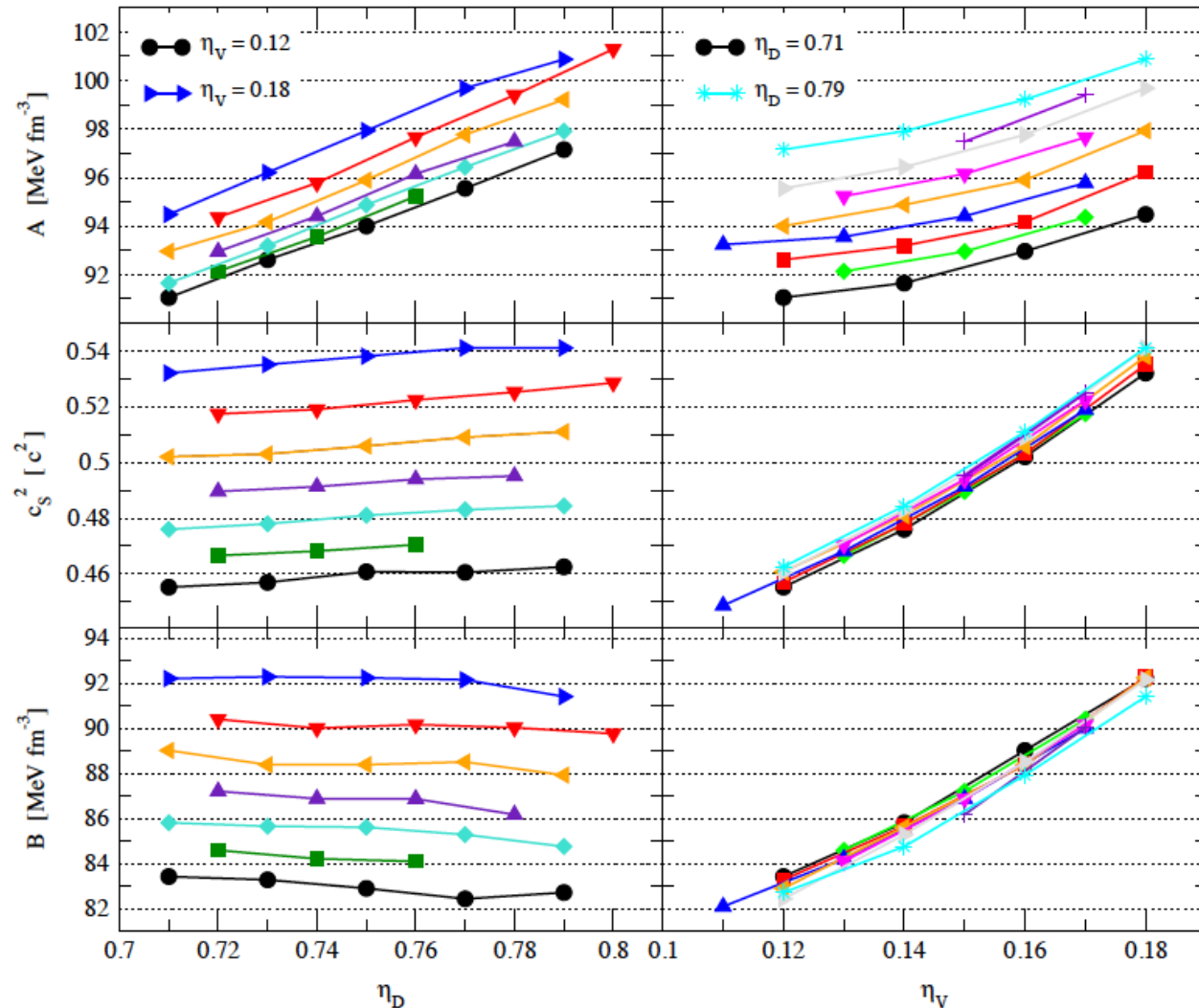
$$P(\mu) = A \left(\frac{\mu}{\mu_x} \right)^{1+\beta} - B,$$

Fitted relationship, see figure \rightarrow

$$A = a_1 \eta_D + b_1 \eta_V^2 + c_1$$

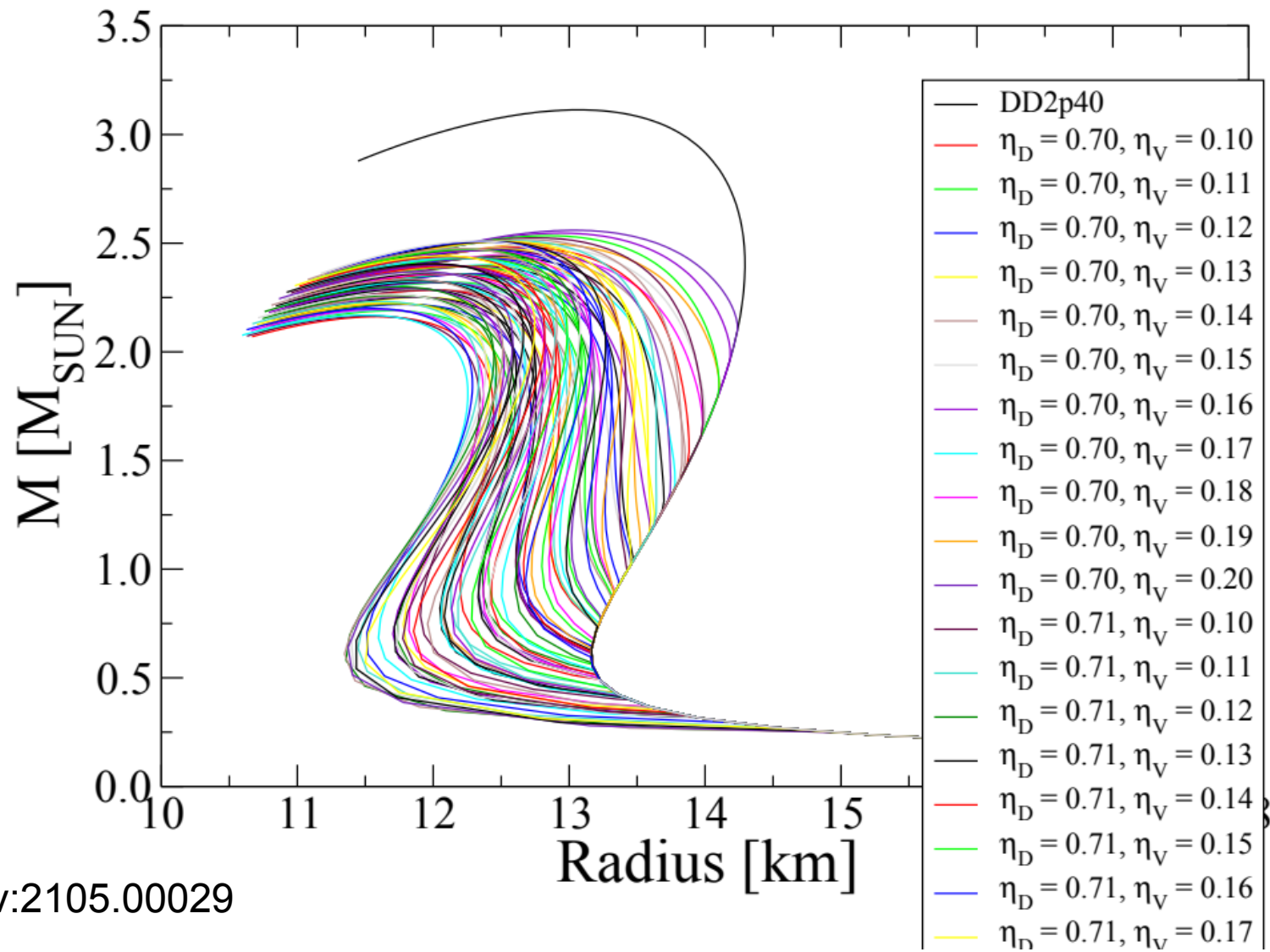
$$c_s^2 = a_2 \eta_D + b_2 \eta_V^2 + c_2$$

$$B = a_3 \eta_D + b_3 \eta_V^2 + c_3,$$



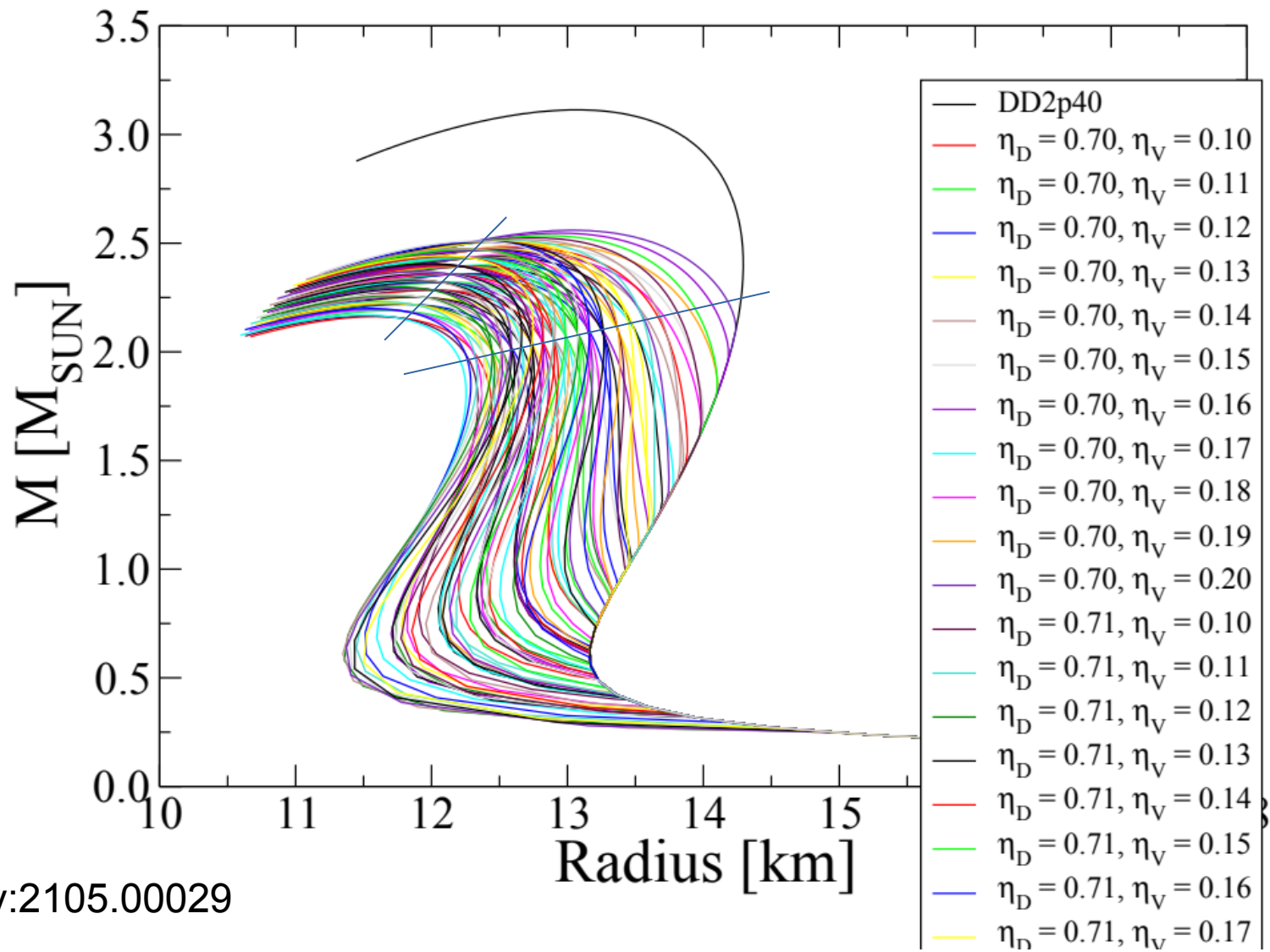
Constant sound speed (CSS) vs. nonlocal NJL model

“Trains” of special points, when η_D and η_V are varied systematically (grid)



Constant sound speed (CSS) vs. nonlocal NJL model

“Trains” of special points, when η_D and η_V are varied systematically (grid)



Old paradigm: hybrid stars smaller and lighter

Works on **Special Point** with M. Cierniak: 2012.15785 & 2009.12353; EPJ ST 229, 3663 (2020)

Dense quark plasma in color
superconducting phase: nINJL mode

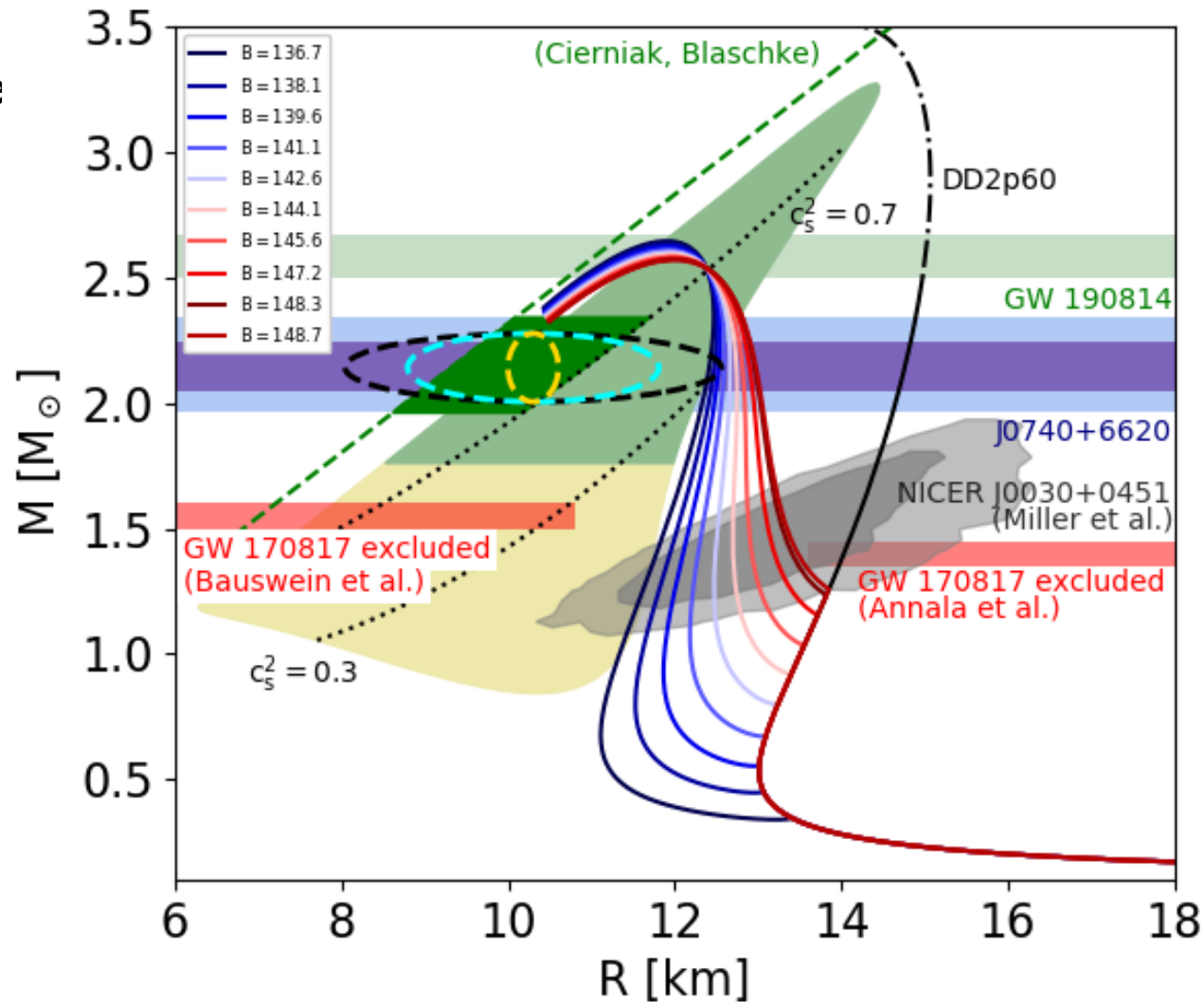
Constant-speed-of-sound (CSS)
Equation of state (EoS)

$$p(\mu) = A(\mu/\mu_0)^{1+c_s^{-2}} - B,$$

$$p = c_s^2 \varepsilon - (1 + c_s^2)B$$

Perfect mapping nINJL \rightarrow CSS,
Antic et al., arxiv:2105.00029

Maxwell construction with
(1st order phase transition)
Relativistic Density Functional
EoS “DD2pxy” by S. Typel
With density-dependent coupling
And excluded volume $v=x.y \text{ fm}^3$



2.6 M_{sun} object can be a hybrid neutron star! With early onset of deconfinement and twins!
NICER radius measurement on PSR J0740+6620 will put constraints on this too!

New paradigm: hybrid stars larger and heavier

Work based on **Special Point location** with M. Cierniak, in preparation

Dense quark plasma in color
superconducting phase: nINJL model

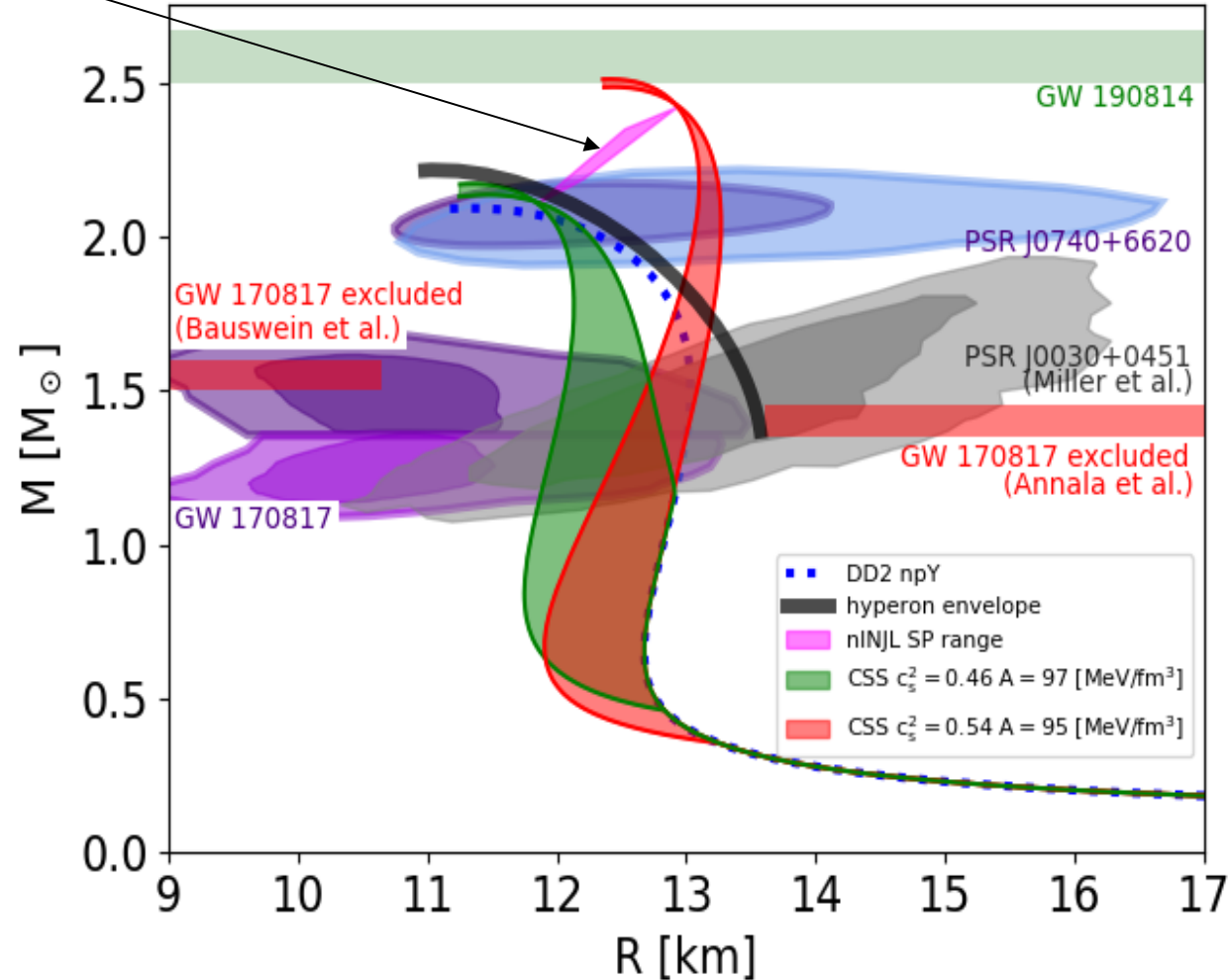
Constant-speed-of-sound (CSS)
Equation of state (EoS)

$$p(\mu) = A(\mu/\mu_0)^{1+c_s^{-2}} - B,$$

$$p = c_s^2 \varepsilon - (1 + c_s^2)B$$

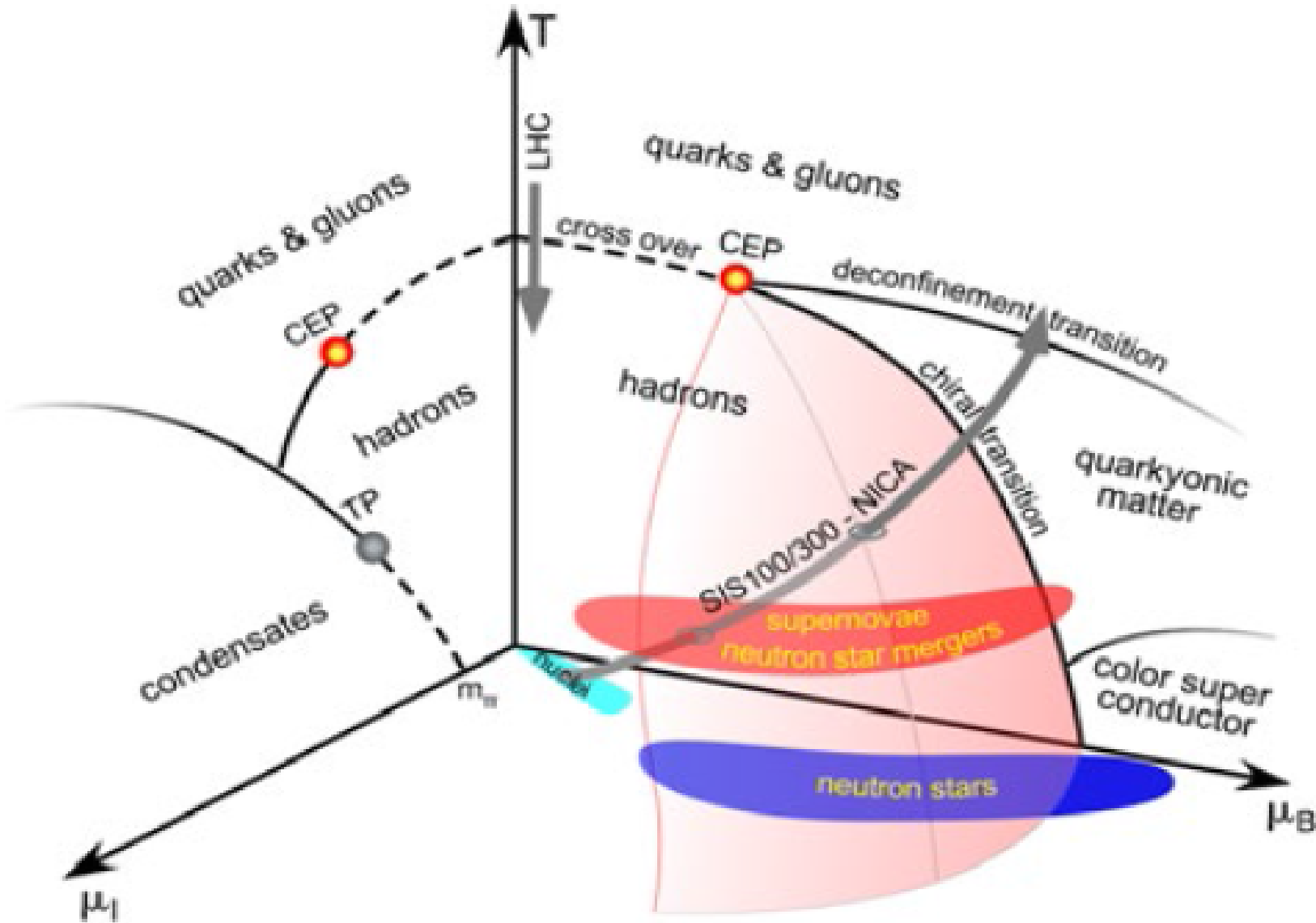
Perfect mapping nINJL \rightarrow CSS,
Antic et al., arxiv:2105.00029

Maxwell construction with
(1st order phase transition)
Relativistic Density Functional
EoS “DD2-Y-T” by S. Typel
With density-dependent coupling



2.5 M_{sun} object can be a hybrid neutron star! With early onset of deconfinement!
NICER radius measurement on PSR J0740+6620 best described by hybrid stars!

CEP in the QCD phase diagram: HIC vs. Astrophysics

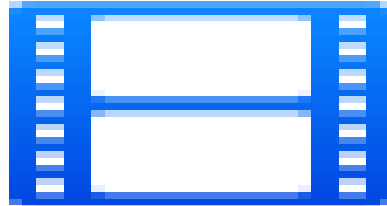


Binary neutron star merger simulation

S. Blacker & A. Bauswein (GSI Darmstadt), $1.35 M_{\text{sun}} + 1.35 M_{\text{sun}}$

<https://www.gsi.de/fileadmin/theorie/simulation-neutron-star-merger.mp4>

Population of the QCD phase diagram with mixed phase, 6... 25 ms

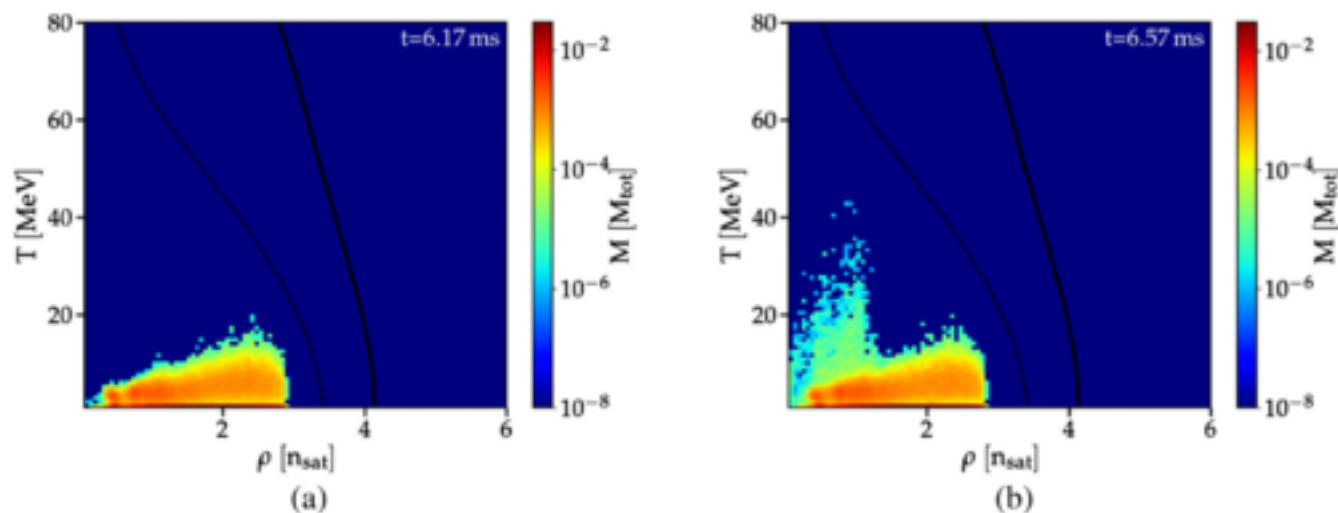


Binary neutron star merger simulation

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Population of the QCD phase diagram with mixed phase, 6... 25 ms

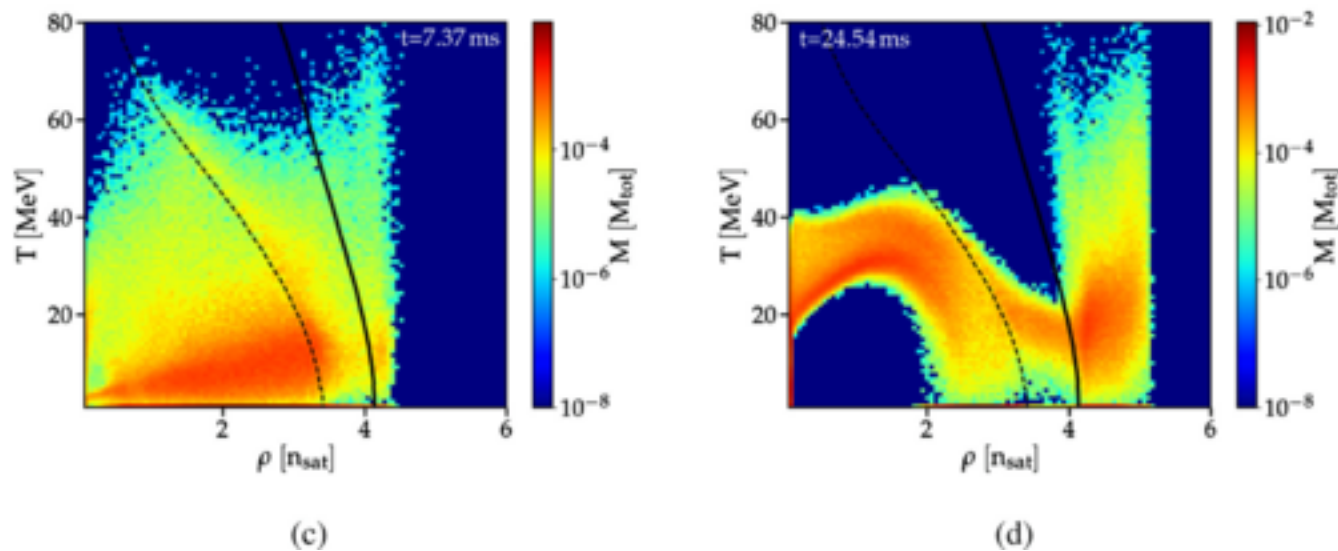


EoS for applications to
supernova and merger
Simulation:

CompOSE

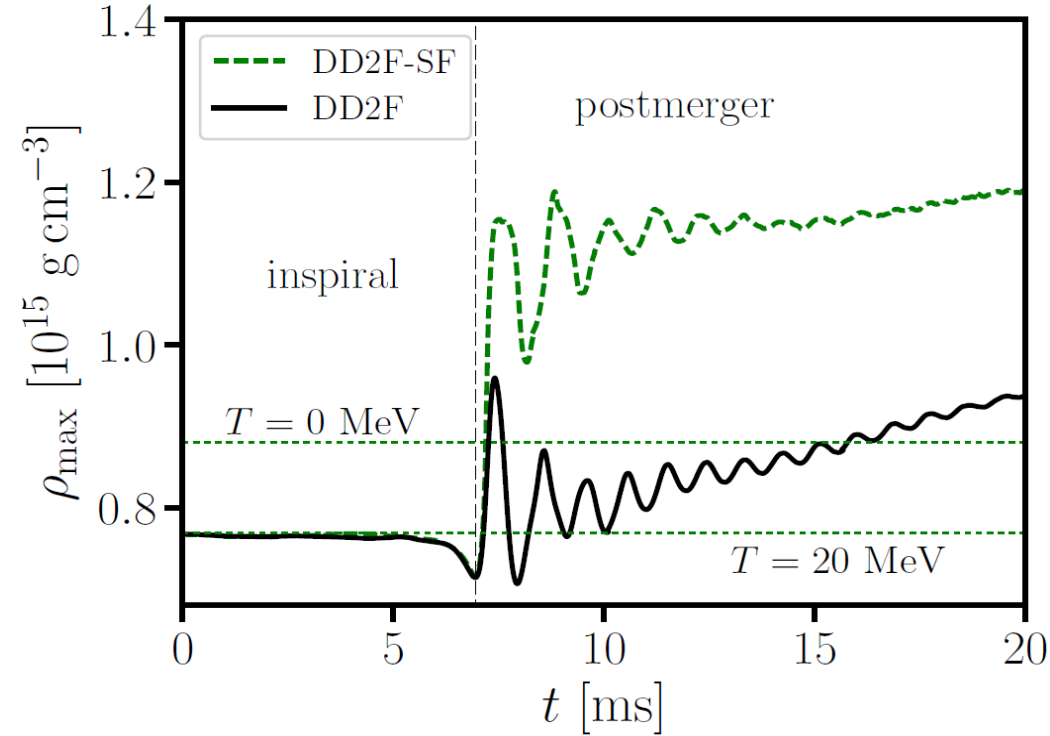
With deconfinement:

<https://compose.obspm.fr/eos/166>

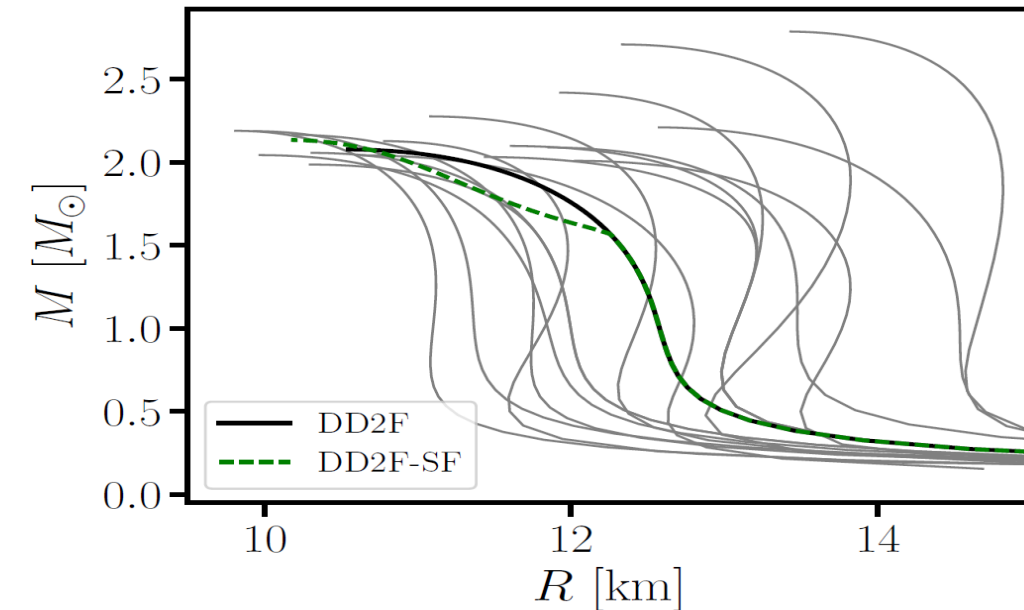
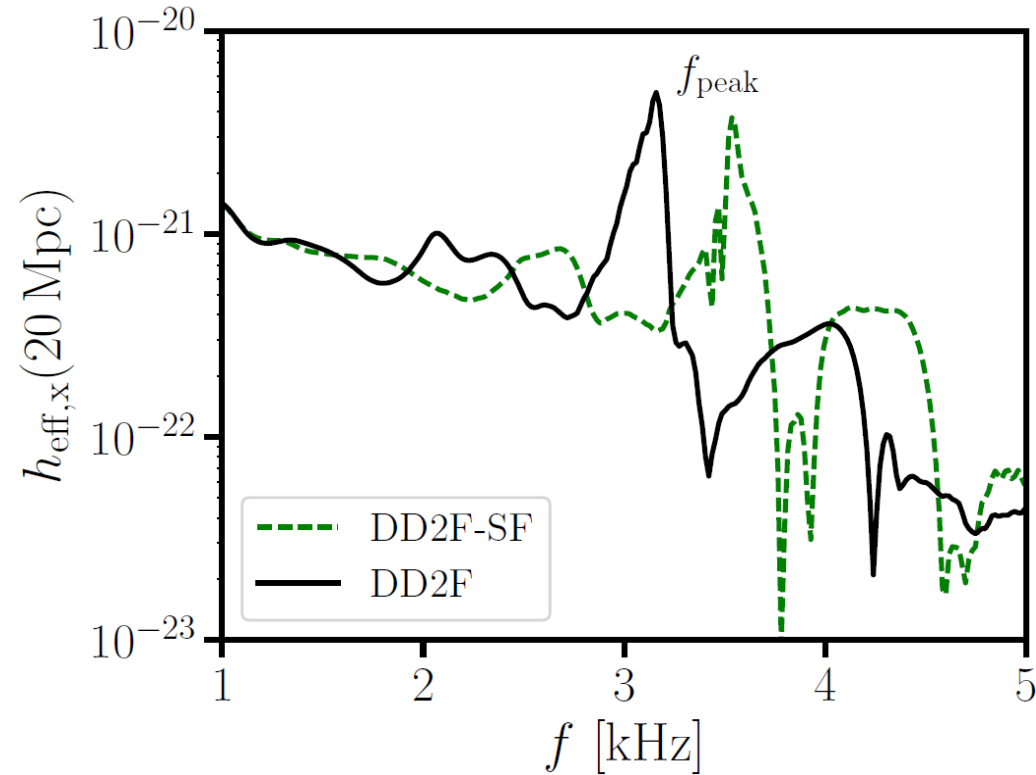


S. Blacker, A. Bauswein, et al.,
Phys. Rev. D 102 (2020) 123023

Hybrid star formation in postmerger phase



Strong phase transition in postmerger GW,
A. Bauswein et al. arxiv:1809.01116

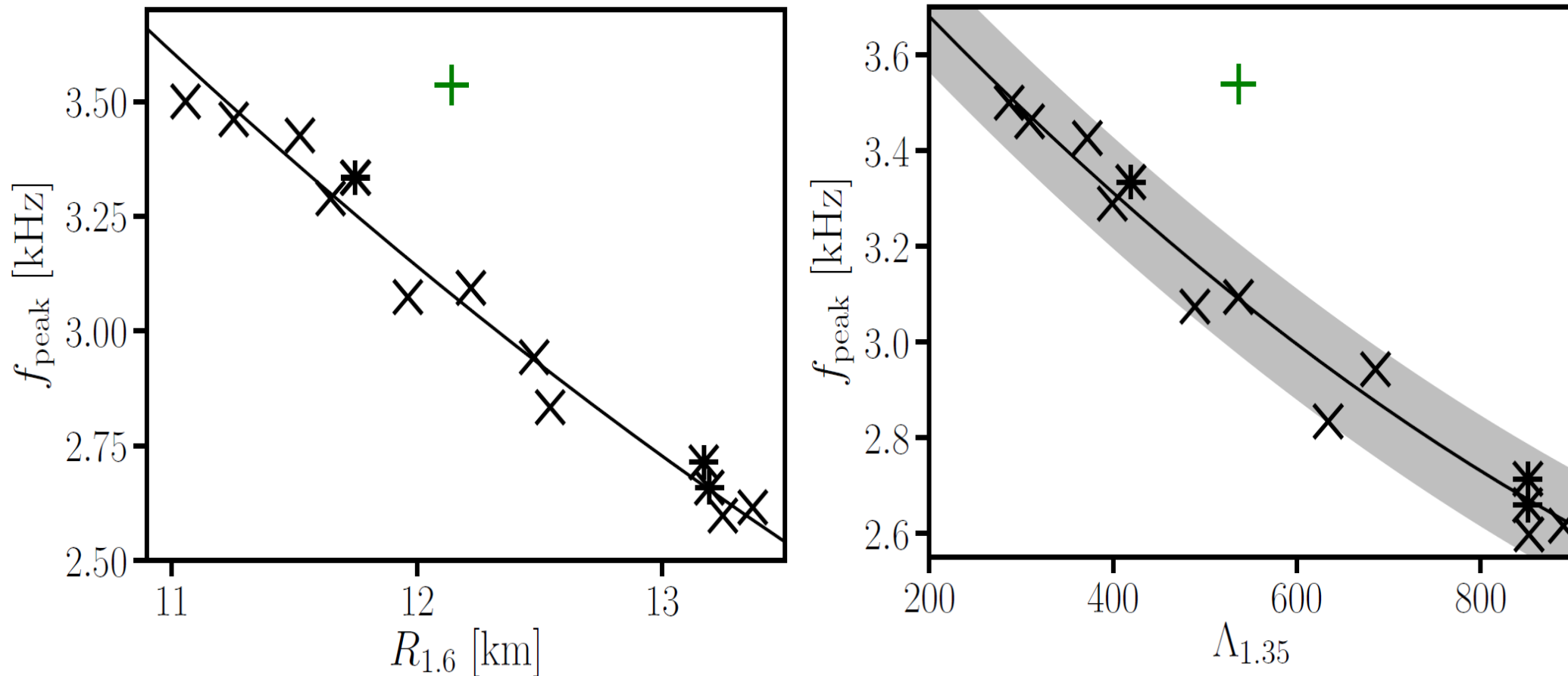


Hybrid star formation during NS merger
→ higher densities and compacter star
→ higher peak frequency of the GW

A. Bauswein et al., PRL 122 (2019) 061102

Hybrid star formation in postmerger phase

Strong phase transition in postmerger GW signal,
A. Bauswein et al., PRL 122 (2019) 061102; [arxiv:1809.01116]

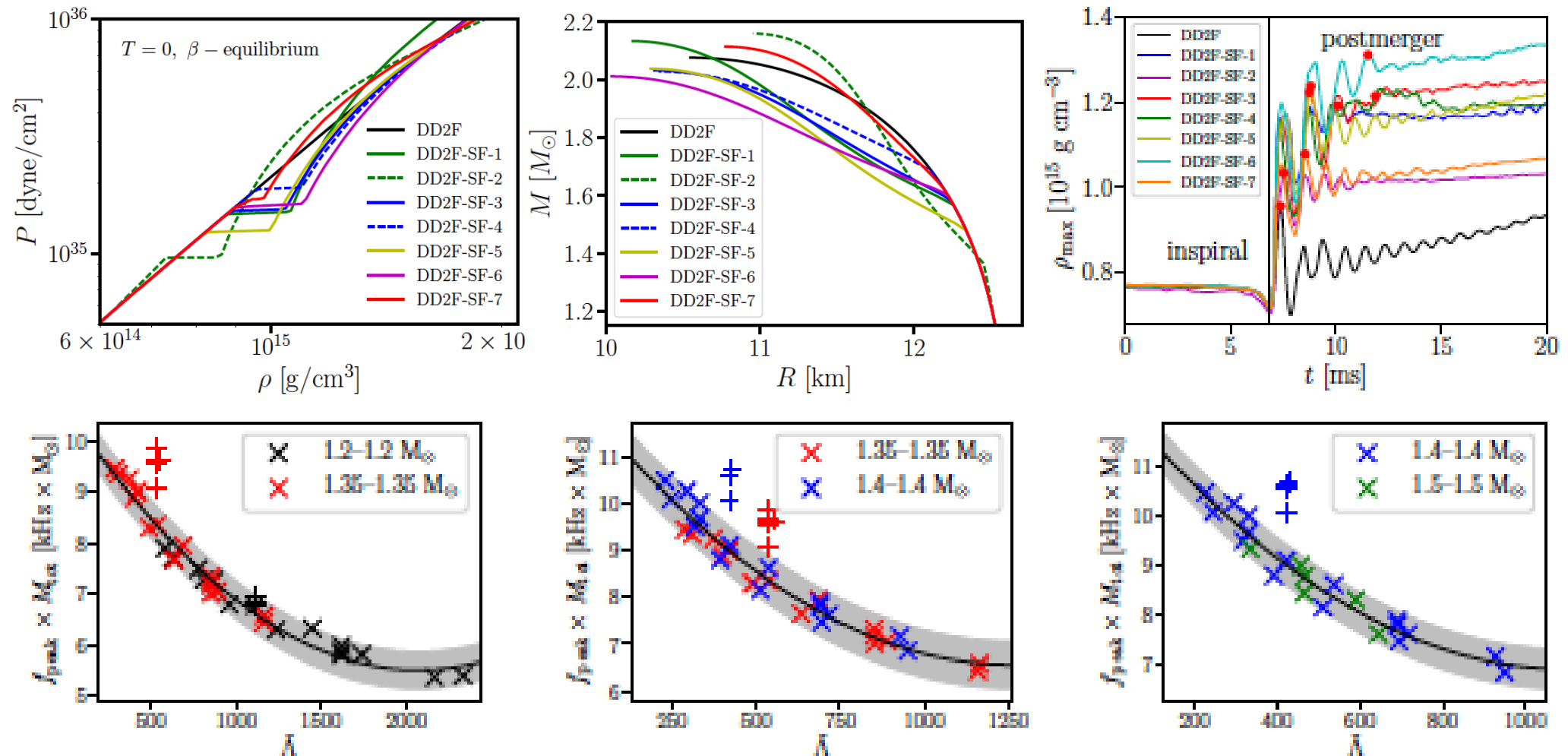


Strong deviation from $f_{\text{peak}} - R_{1.6}$ relation signals **strong phase transition** in NS merger!

Complementarity of f_{peak} from **postmerger** with tidal deformability $\Lambda_{1.35}$ from **inspiral phase**.

Hybrid star formation in postmerger phase

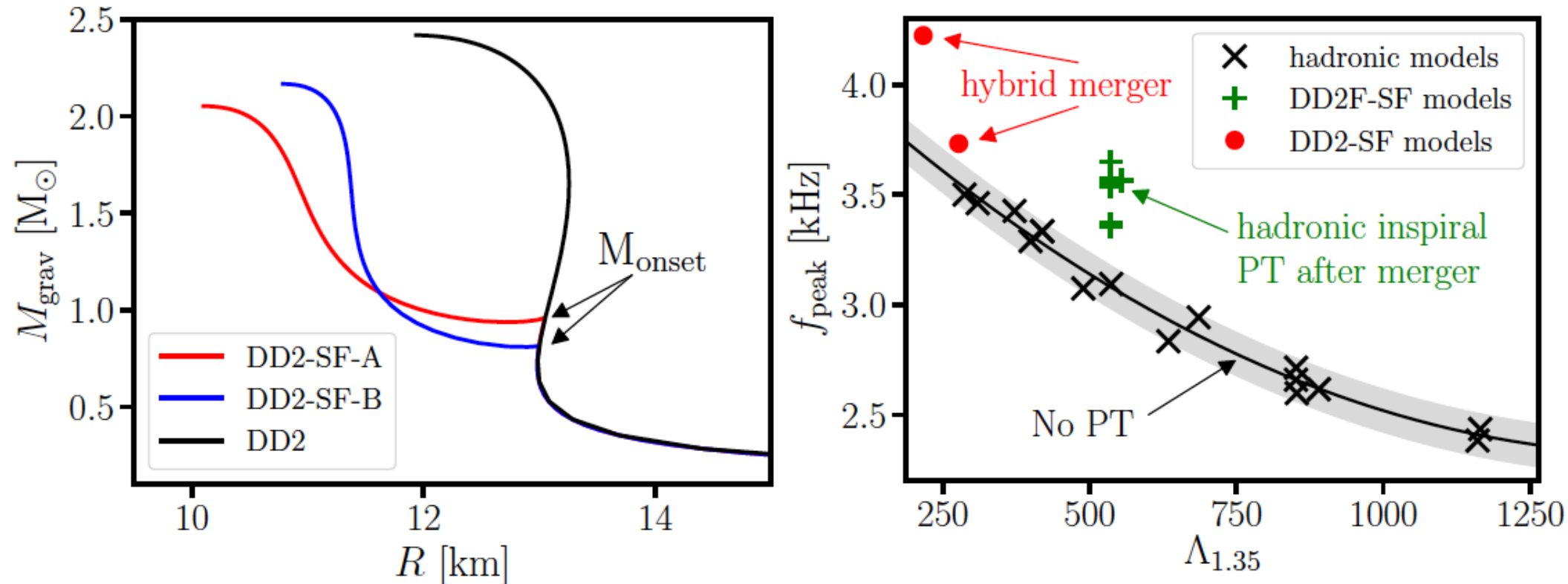
Strong PT in postmerger GW signal, S. Blacker et al., arxiv:2006.03789, PRD102 (2020) 123023



Dominant **postmerger** frequency f_{peak} vs. tidal deformability $\Lambda_{1.35}$ from **inspiral** phase:
 Results from hybrid models appear as **outliers** of the grey band (maximal deviation of purely hadronic models from a least squares fit) = signalling a **strong phase transition in NS** !

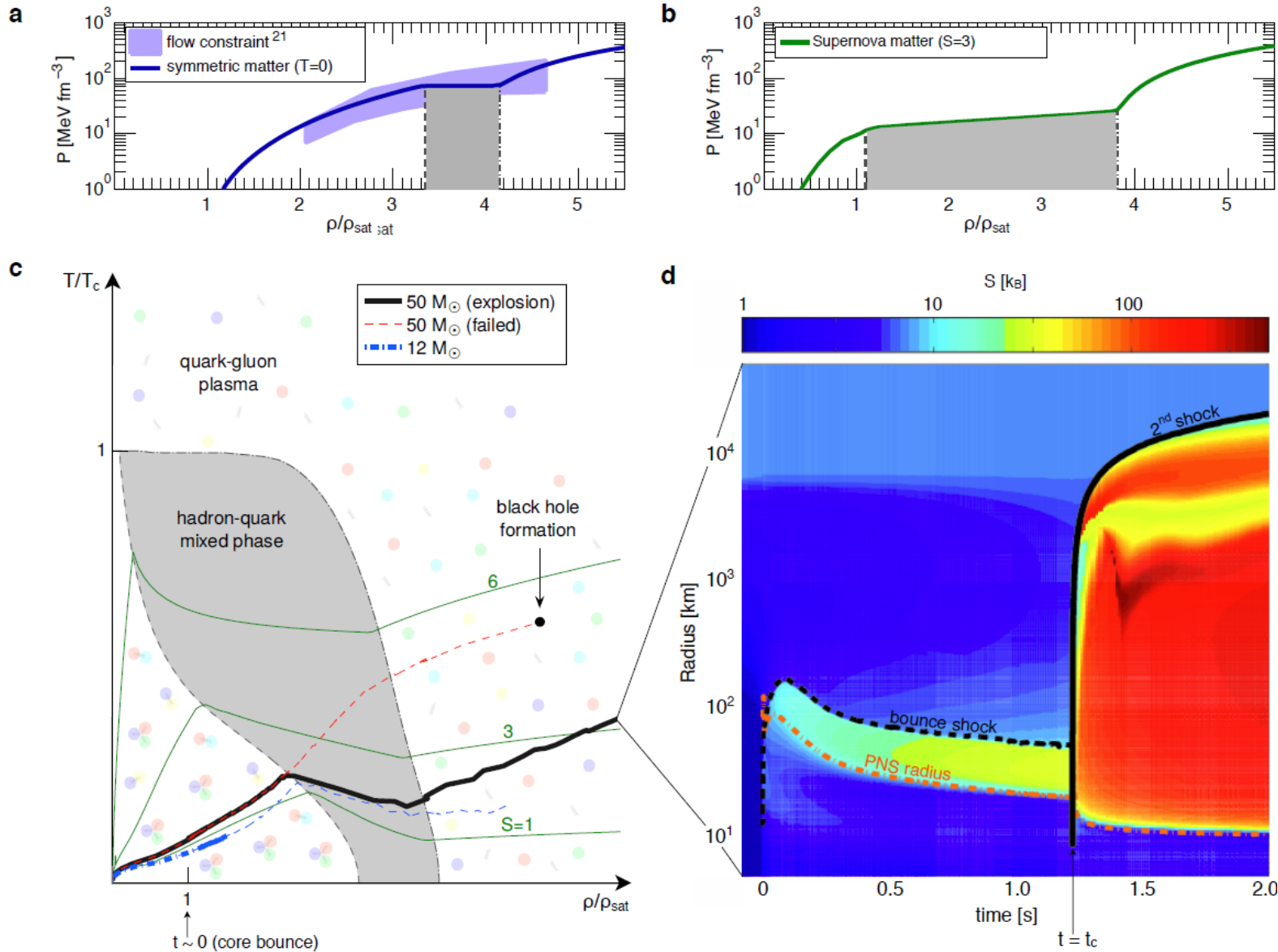
GW signal of deconfinement in merger of hybrid stars

Merger of hybrid stars with early phase transition: Bauswein & Blacker, EPJ ST 229 (2020)



The combination of stiff hadronic EoS (DD2) and string-flip (SF) model allows for early onset of deconfinement in low-mass neutron stars and even third-family solutions (mass twins). For these cases, the event GW170817 could have been a **merger of two hybrid stars!** Also in these cases (red dots in above figure) a **significant deviation** from the grey band of Purely hadronic star mergers without a phase transition is obtained!

Deconfinement transition as SN explosion mechanism



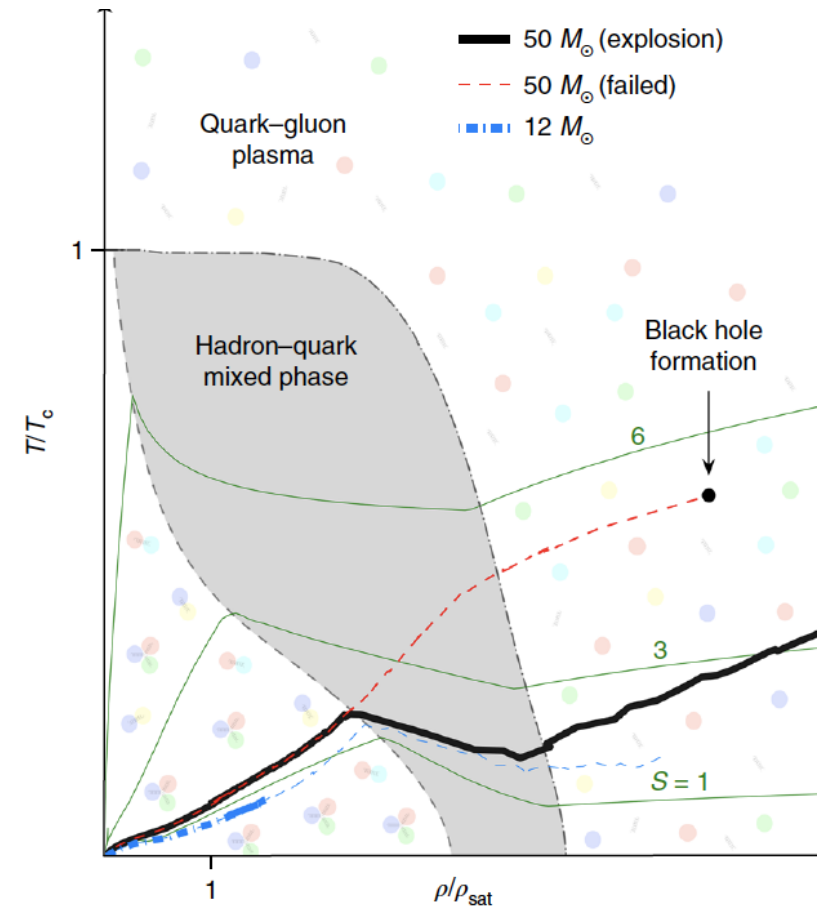
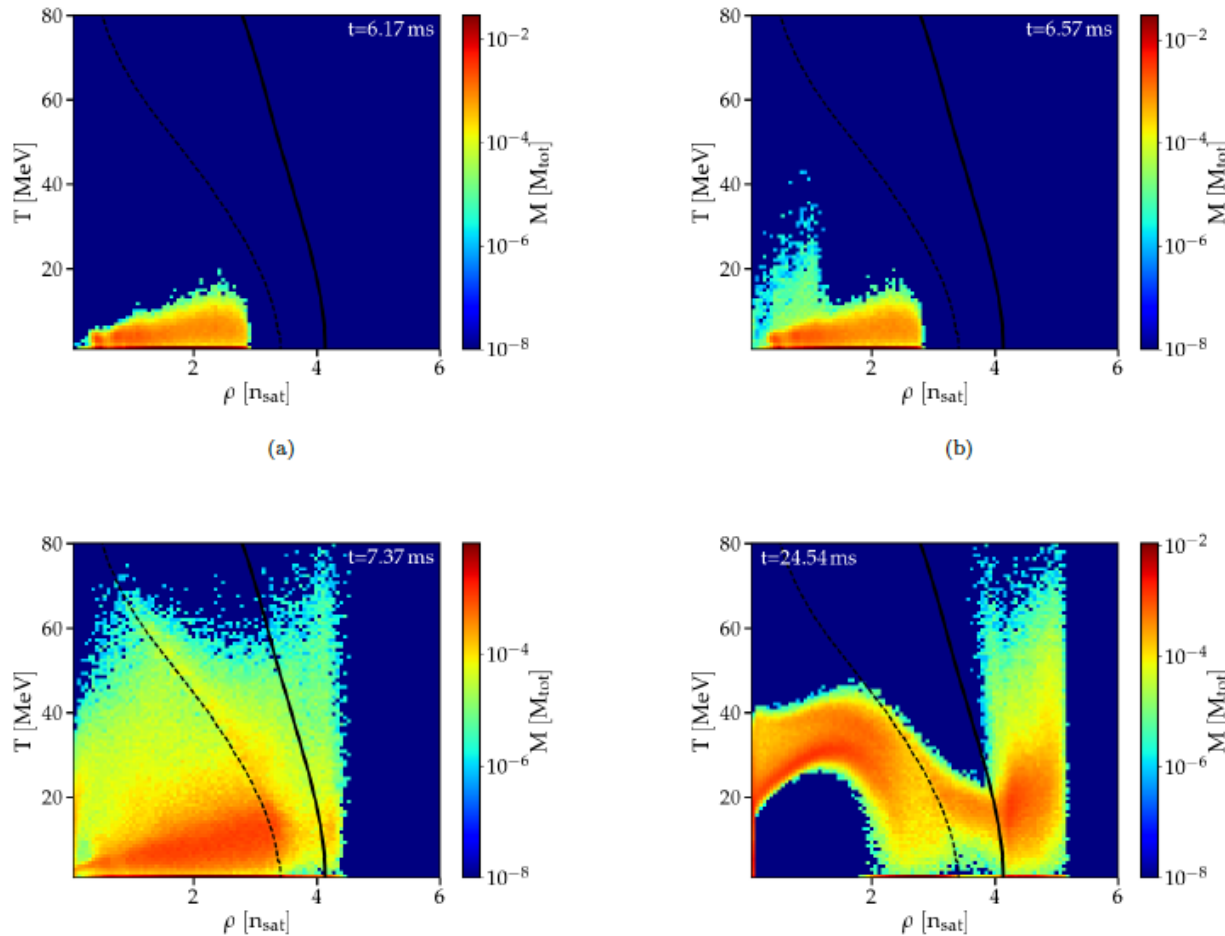
Progenitor:
M = 50 M_⊙

T. Fischer, N.-U. Bastian et al., Quark deconfinement as supernova engine of massive blue Supergiant star explosions, Nature Astronomy 2 (2018) 980-986; arxiv:1712.08788

Population of the QCD Phase Diagram in Mergers & SNe

Binary NS merger, 1.35 M_{sun} + 1.35 M_{sun}

SN explosion, 50 M_{sun}

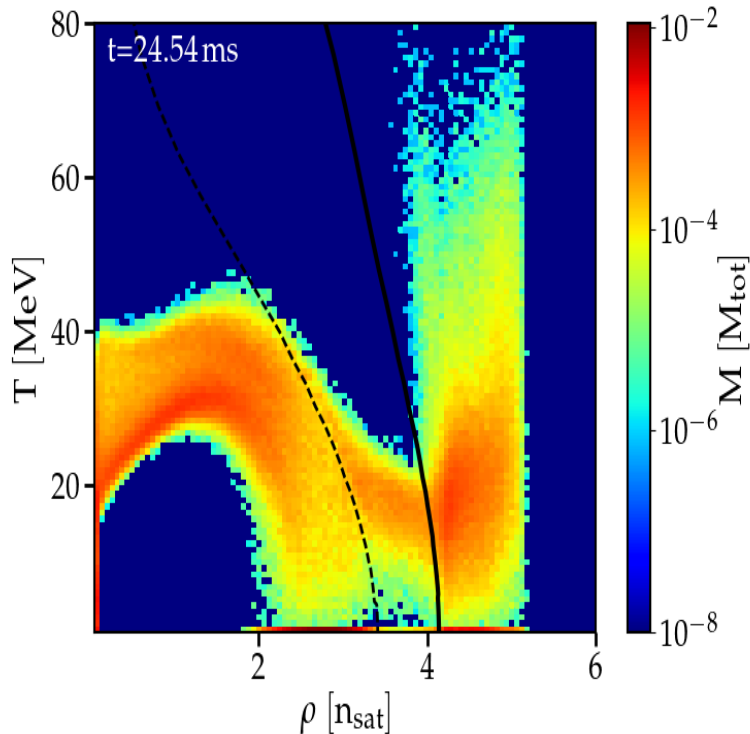


S. Blacker, A. Bauswein et al.,
 Phys. Rev. D102 (2020) 123023; arxiv:2006.03789

T. Fischer et al.,
 Nat. Astron. 2 (2018) 980;
 arxiv:1712.08788

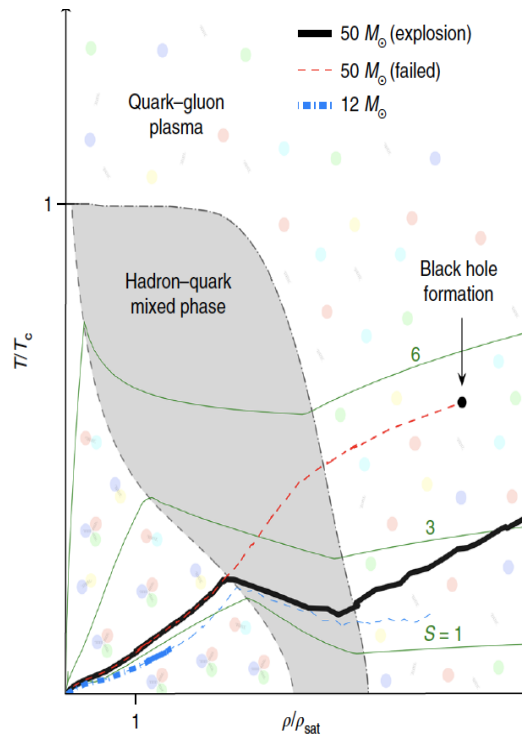
Population of the QCD Phase Diagram

Binary NS merger,
1.35+1.35 M_{sun}



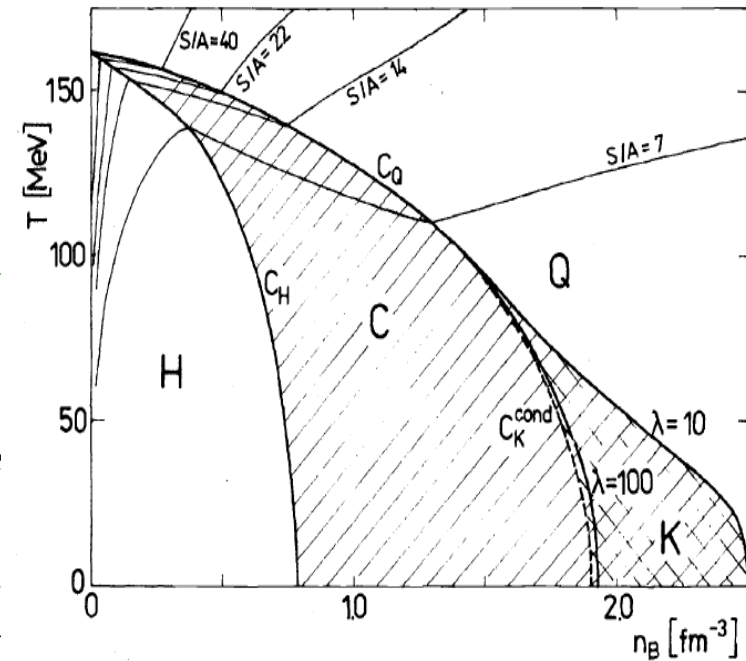
S. Blacker, A. Bauswein et al.,
PRD 102 (2020) 123023
arXiv:2006.03789

SN explosion,
Progenitor 50 M_{sun}



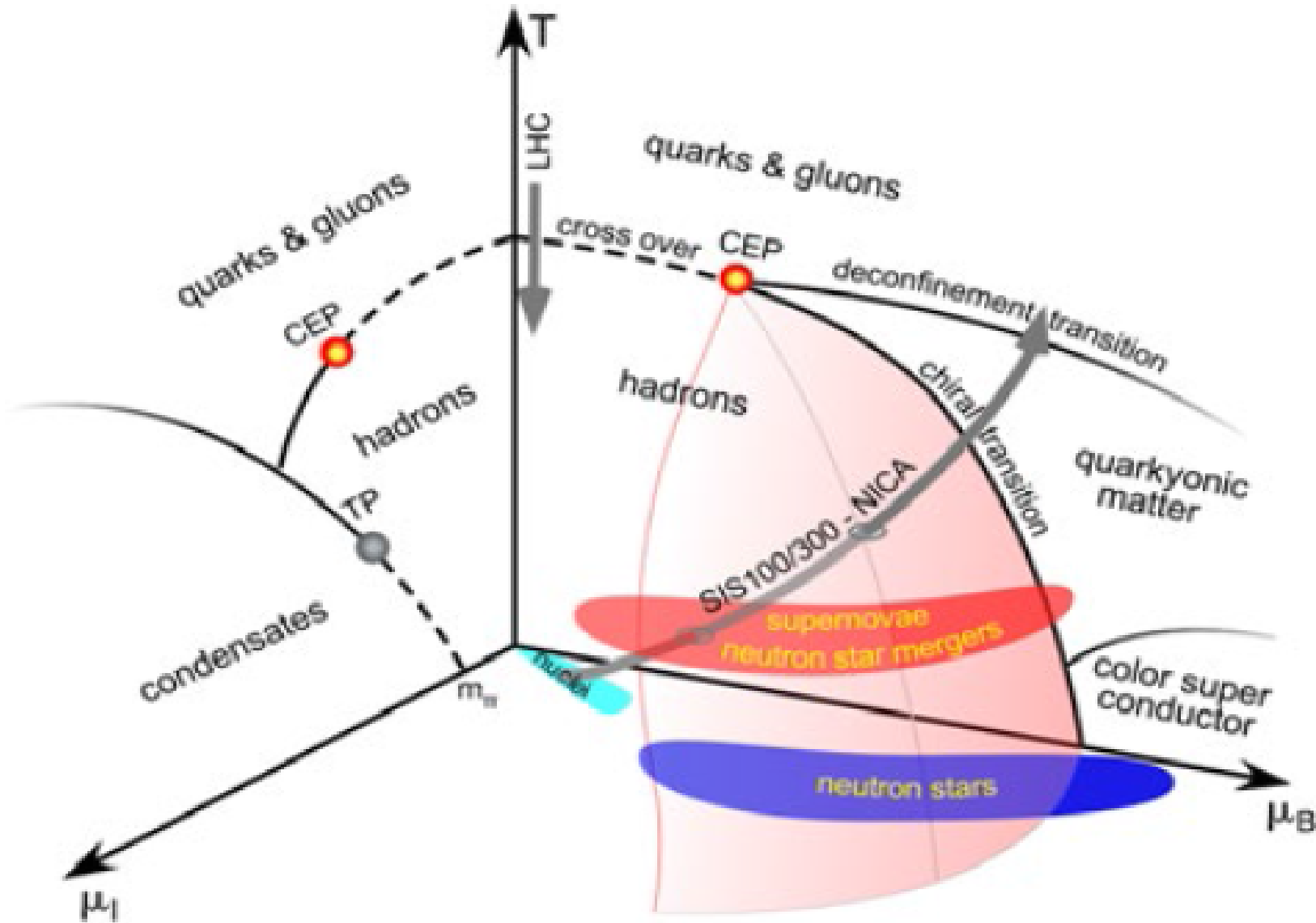
T. Fischer et al.,
Nat. Astron. 2 (2018) 980
arXiv:1712.08788

Ultrarelativistic HIC,
 \sqrt{s} [GeV]=16, 10, 7, 4



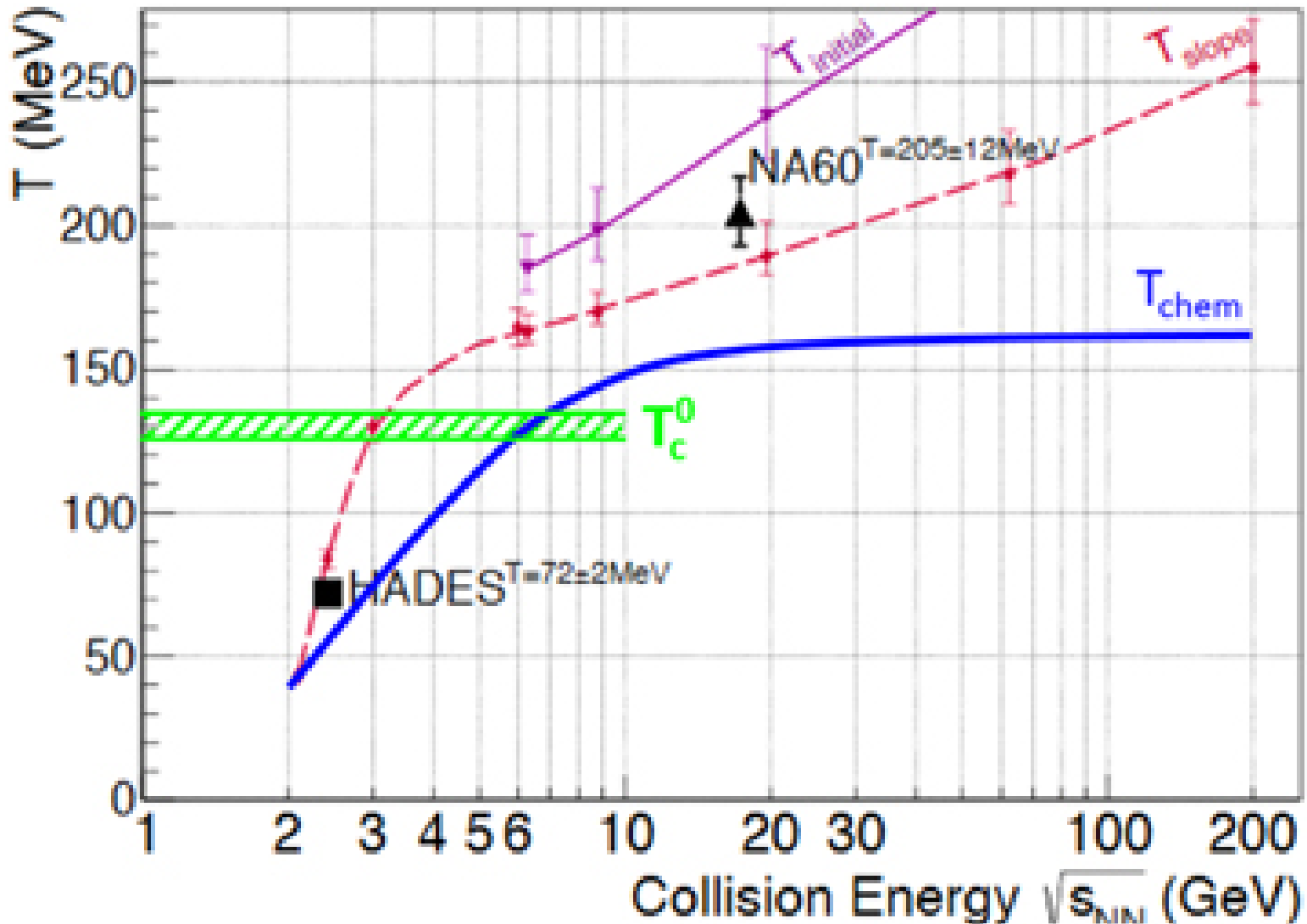
H.W. Barz, B. Friman et al.,
PRD 40 (1989) 157
GSI Preprint, GSI-89-13

CEP in the QCD phase diagram: HIC vs. Astrophysics



CEP in the QCD phase diagram: HIC vs. Astrophysics

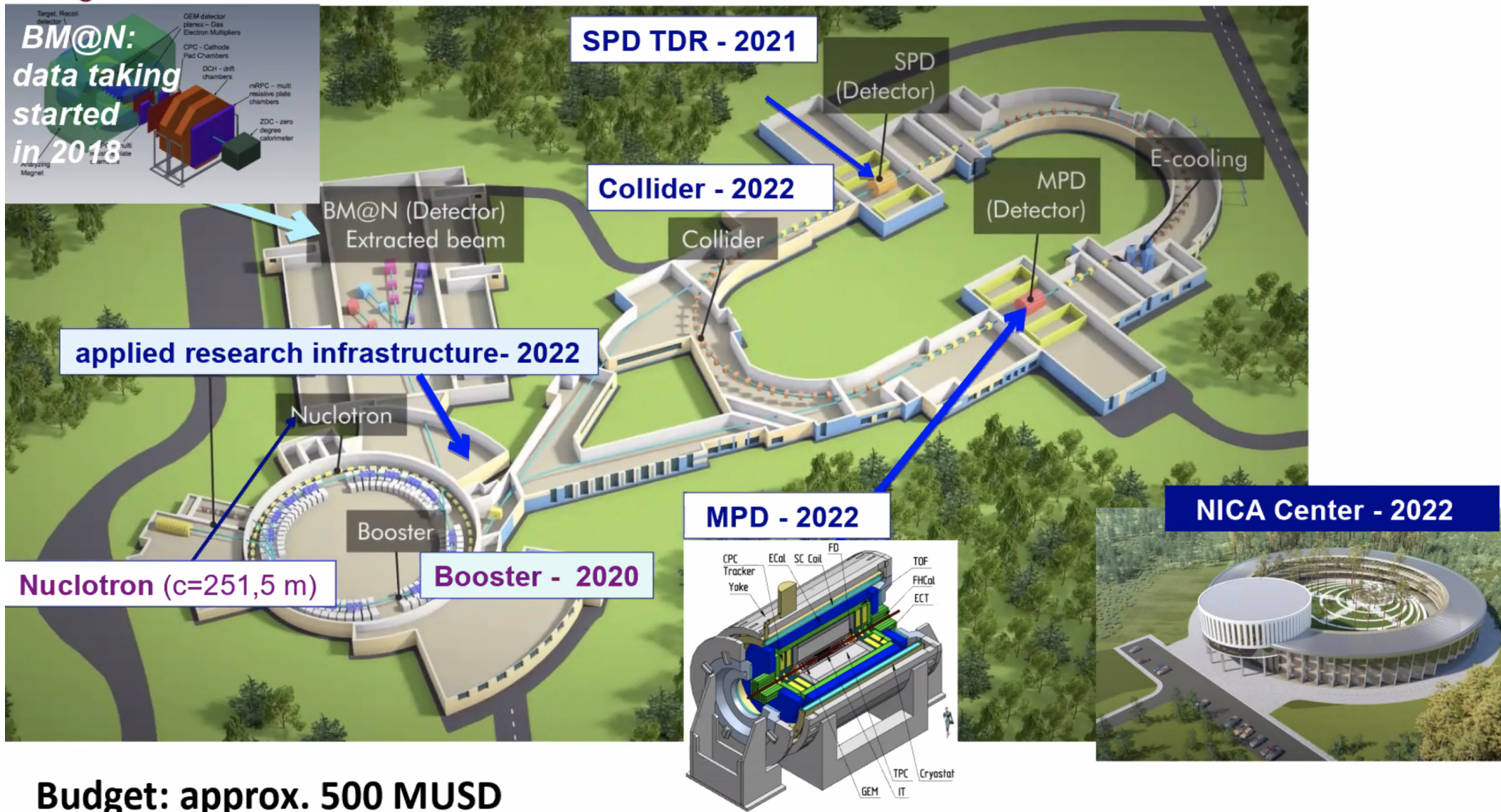
NICA experiments (BM @N, MPD)



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NICA Accelerator Complex in Dubna



The NICA Facility at JINR Dubna



NICA construction live



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NICA: Unique and complementary

Collider advantage:

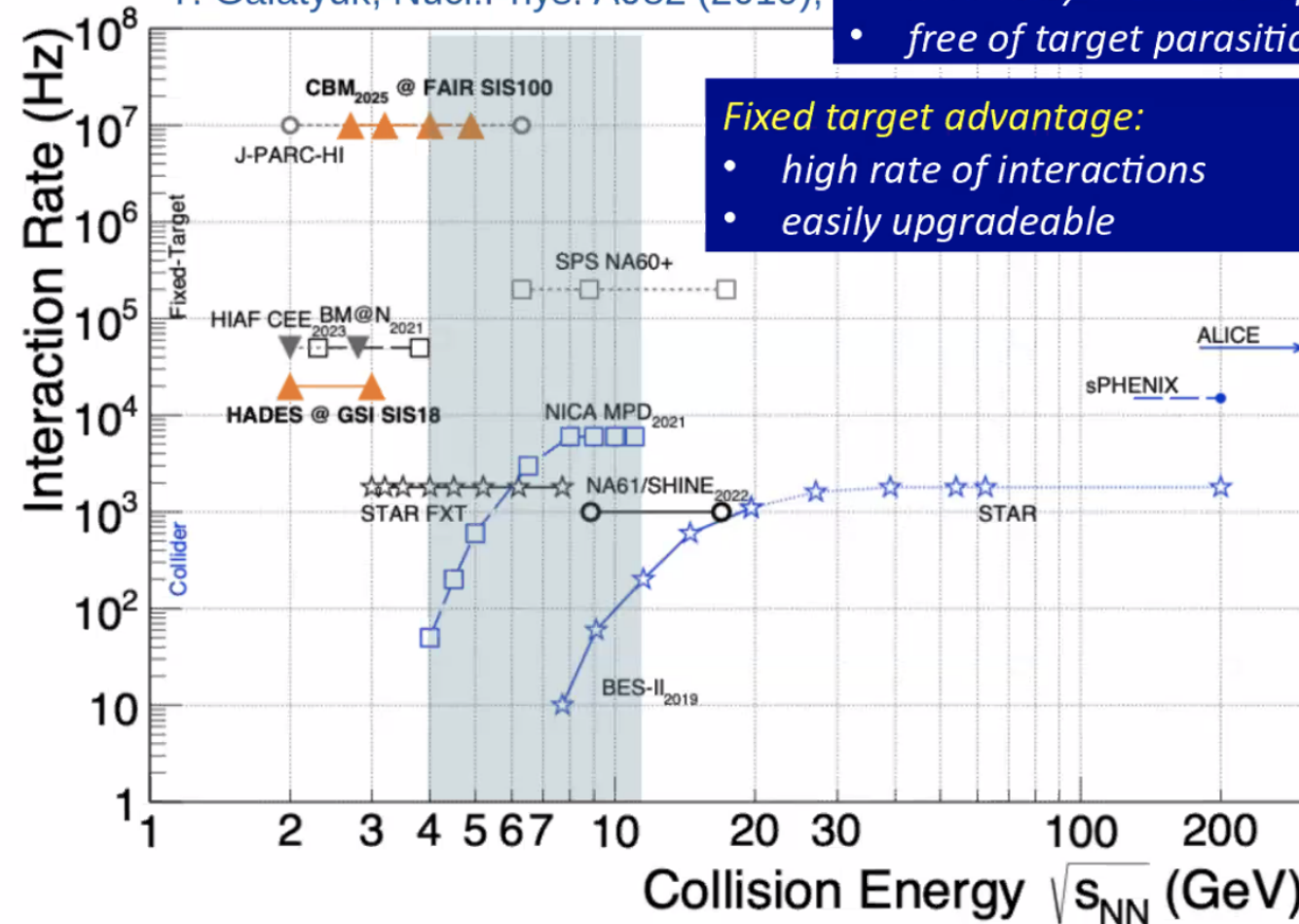
- coverage of max. phase space
- minimally biased acceptance
- free of target parasitic effects

In NICA Collider energy range maximum possible net-baryon density is reached

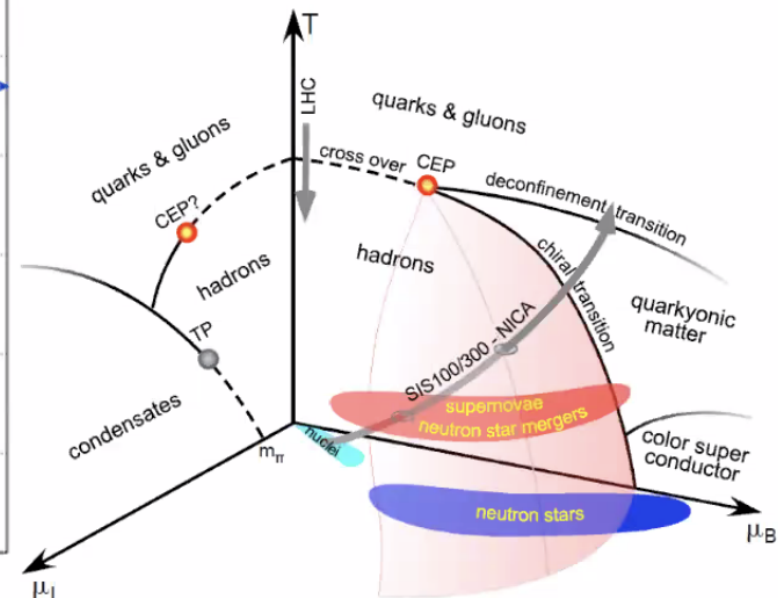
Fixed target advantage:

- high rate of interactions
- easily upgradeable

T. Galatyuk, Nucl.Phys. A982 (2019);



NUPECC Long Range Plan 2017



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Main parameters of accelerator complex

Nuclotron

Parameter	SC synchrotron
particles	$\uparrow p, \uparrow d, \text{nuclei (Au, Bi, ...)}$
max. kinetic energy, GeV/u	10.71 ($\uparrow p$); 5.35 ($\uparrow d$) 3.8 (Au)
max. mag. rigidity, Tm	38.5
circumference, m	251.52
vacuum, Torr	10^{-9}
intensity, Au /pulse	$1 \cdot 10^9$

Booster

	value
ion species	$A/Z \leq 3$
max. energy, MeV/u	600
magnetic rigidity, T m	1.6 – 25.0
circumference, m	210.96
vacuum, Tor	10^{-11}
intensity, Au /p	$1.5 \cdot 10^9$

The Collider

Design parameters, Stage II

45 T*m, 11 GeV/u for Au⁷⁹⁺

Ring circumference, m	503,04
Number of bunches	22
r.m.s. bunch length, m	0,6
β , m	0,35
Energy in c.m., GeV/u	4-11
r.m.s. $\Delta p/p$, 10^{-3}	1,6
IBS growth time, s	1800
Luminosity, $\text{cm}^{-2} \text{s}^{-1}$	1×10^{27}

Stage I:

- **without ECS in Collider, with stochastic cooling**
- **reduced number of RF**
- **reduced luminosity**

Collision system limited by source. Now Available:
C(A=12), N(A=14), Ne(A=20), Ar(A=40), Fe(A=56),
Kr(A=78-86), Xe(A=124-134), Bi(A=209)

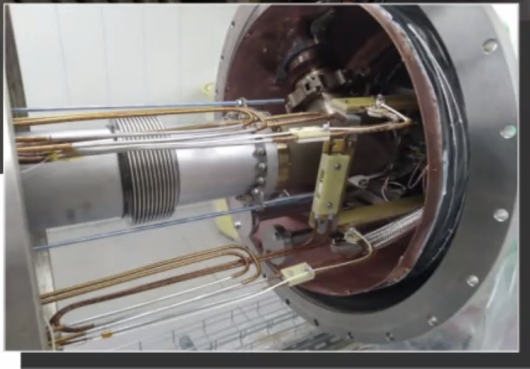
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Booster commissioning

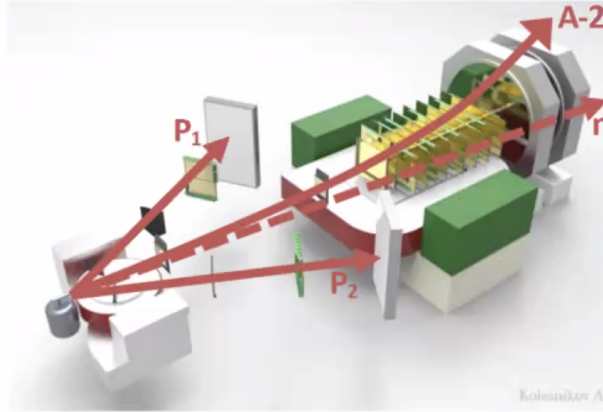
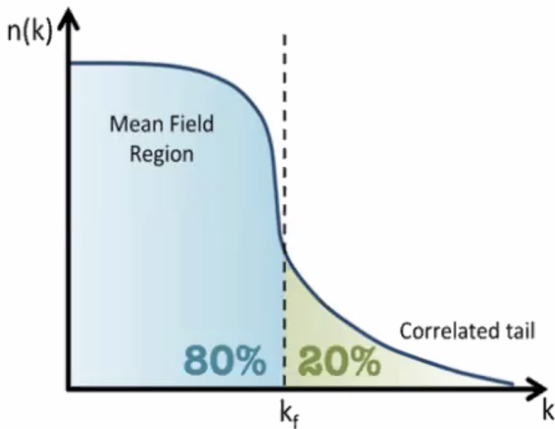


- ✓ Booster fully assembled in the tunnel
- ✓ Commissioning and test ongoing for beam diagnostics, beam acceleration, electron cooling, power supply, magnets, cryogenics



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Experiment with BM@N: Short-Range Correlations (SRC)



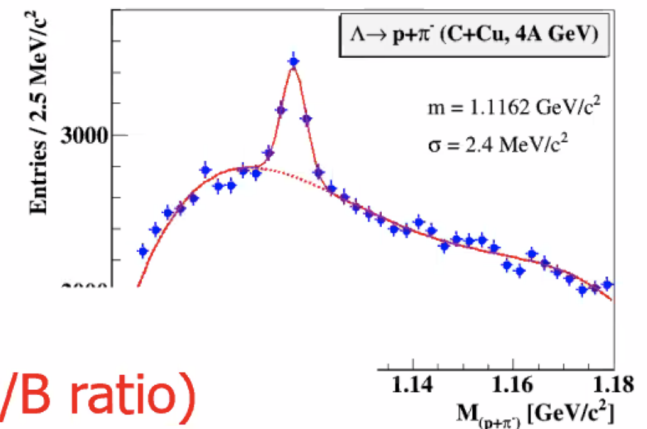
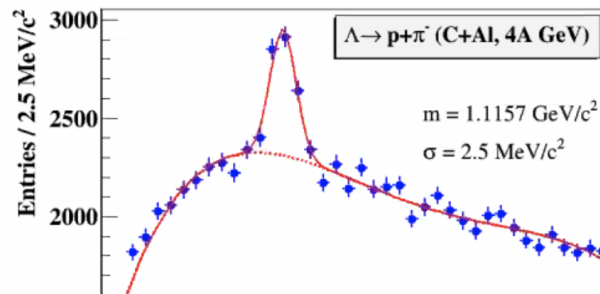
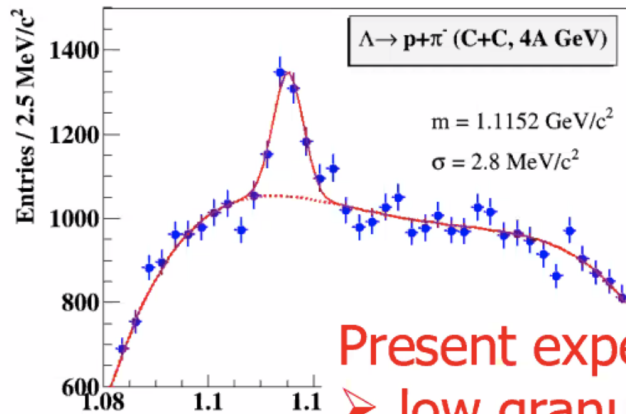
Experiment at BM@N with a 4A GeV C-beam:
 $^{12}\text{C} + p \rightarrow 2p + ^{10}_4\text{Be} + p$ (pp SRC)

First fully exclusive measurement in inverse kinematics probing the residual $A-2$ nuclear system!

M. Patsyuk et al., arXiv:2102.02626

Accepted for publication in *nature physics*

Experiment with BM@N: Λ 's in C + C, Al, Cu at 4A GeV



Present experimental limitations:

- low granularity tracking systems (small S/B ratio)
- air gaps in beam line from Nuclotron (low beam quality)
- no vacuum beam pipe in BM@N (large background)

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Interior of MPD Hall



*Opening of
solenoid
sarcophagus:
Mar. 23rd*

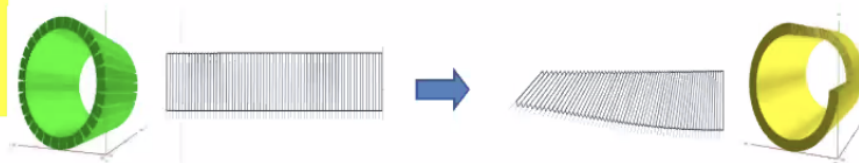
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Electromagnetic Calorimeter (ECAL)

- ❖ Pb+Sc “Shashlyk” read-out: WLS fibers + MAPD $L \sim 35 \text{ cm} (\sim 14 X_0)$
- ❖ Segmentation ($4 \times 4 \text{ cm}^2$) $\sigma(E)$ better than 5% @ 1 GeV time resolution $\sim 500 \text{ ps}$

Barrel ECAL = 38400 ECAL towers (2x25 half-sectors x 6x8 modules/half-sector x 16 towers/module)



Projective geometry

So far ~ 300 modules (16 towers each) = 3 sectors are produced
 Another 3 sectors are planned to be completed by May 2021
 Chinese collaborators will produce 8 sectors by the end of 2021
 25% of all modules are produced by JINR (production area in Protvino)
 75% produced in China, currently funding is secured for approx. 25%

Sectors in dedicated Containers

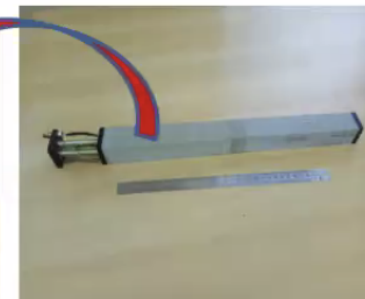
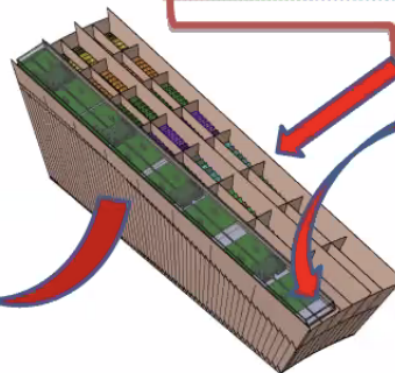
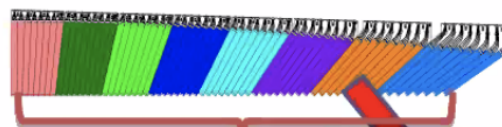
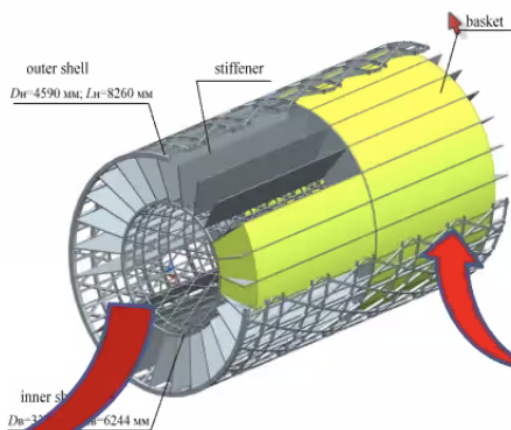
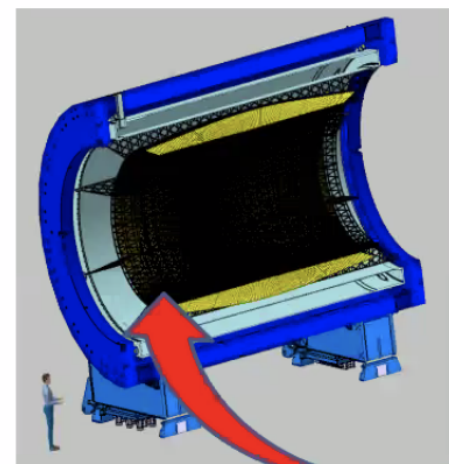


Photo of one element

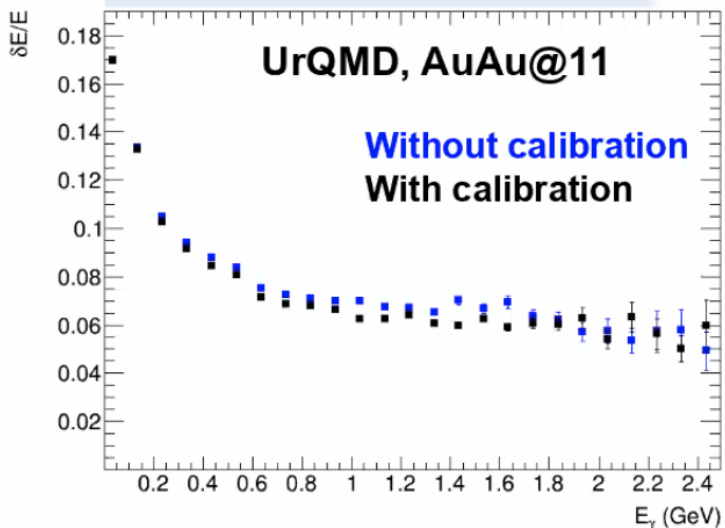
The NICA Facility at JINR Dubna



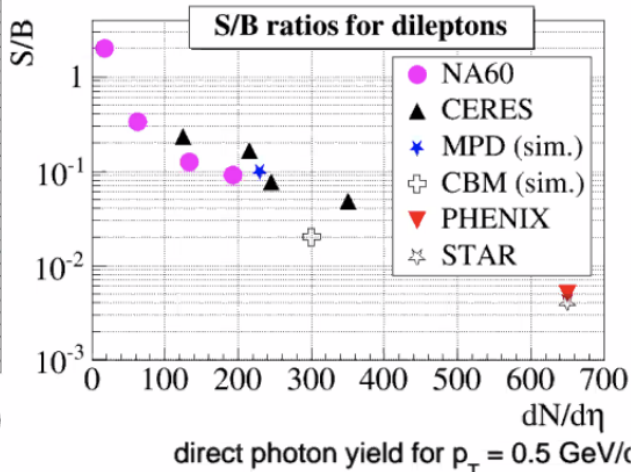
Electromagnetic probes in ECAL

- Realistic ECAL reconstruction & analysis – large acceptance ECAL with good energy resolution: ideal tool for measurement of neutral mesons in a wide momentum range

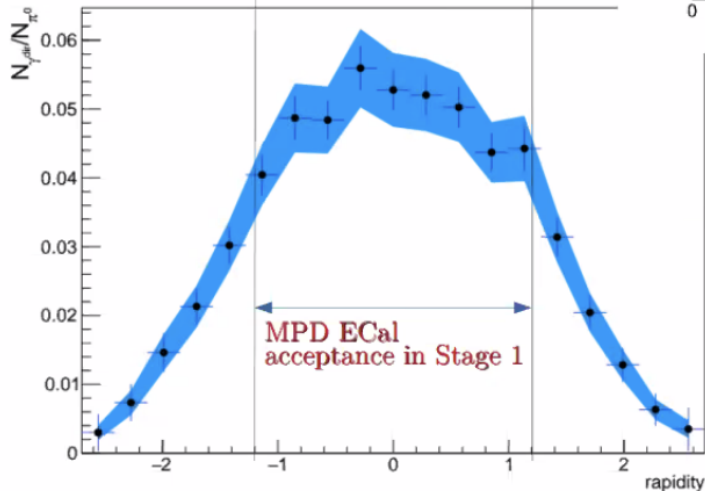
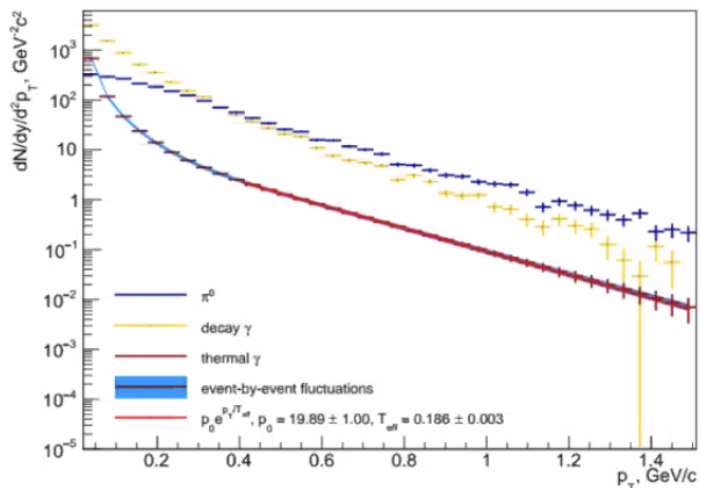
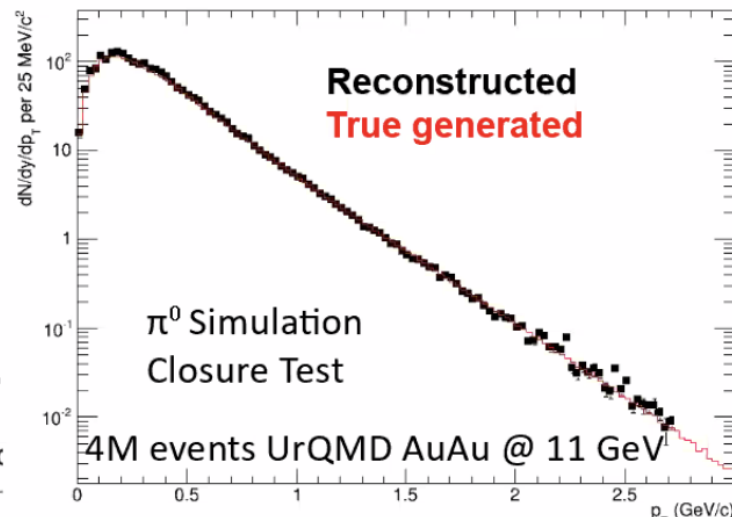
Photon energy resolution



direct γ and π^0 spectra. Au+Au $\sqrt{s_{NN}} = 11$ GeV. $b = 4.5$ fm



direct photon yield for $p_T = 0.5$ GeV/c



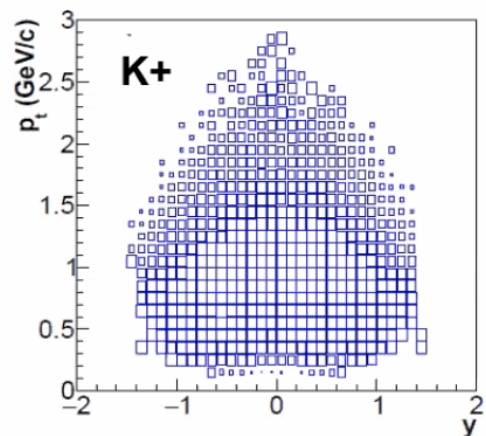
- Promising feasibility studies for prompt photon measurements in MPD

The NICA Facility at JINR Dubna



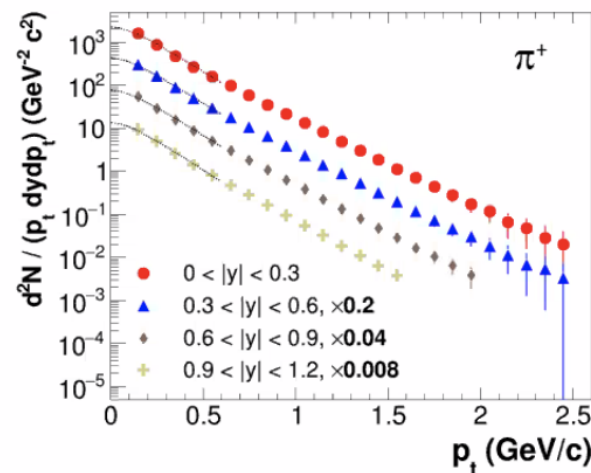
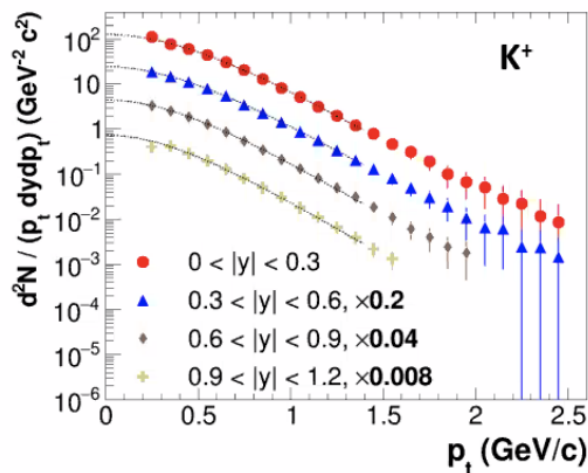
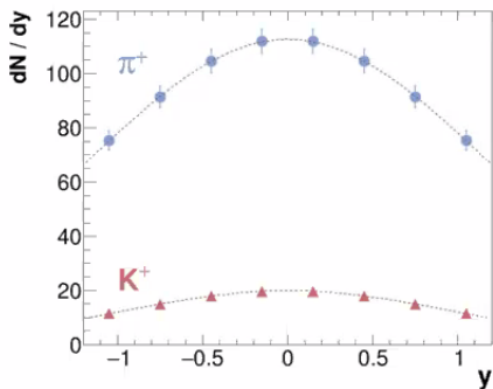
Hadroproduction with MPD

- Particle spectra, yields & ratios are sensitive to bulk fireball properties and phase transformations in the medium
- Uniform acceptance** and **large phase coverage** are crucial for precise mapping of the QCD phase diagram
- ✓ 0-5% central Au+Au at 9 GeV from the PHSD event generator, which implements partonic phase and CSR effects
- ✓ Recent reconstruction chain, combined dE/dx +TOF particle ID, spectra analysis



- MPD provides large phase-space coverage for identified pions and kaons (> 70% of the full phasespace at 9 GeV)
- Hadron spectra can be measured from $p_T=0.2$ to 2.5 GeV/c
- Extrapolation to full p_T -range and to the full phase space can be performed exploiting the spectra shapes (see BW fits for p_T -spectra and Gaussian for rapidity distributions)

Ability to cover full energy range of the „horn” with consistent acceptance

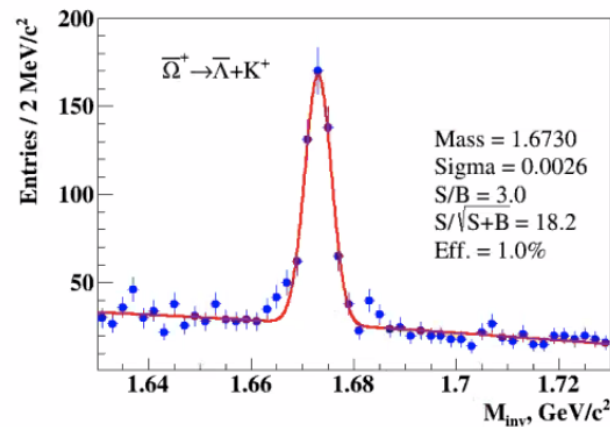
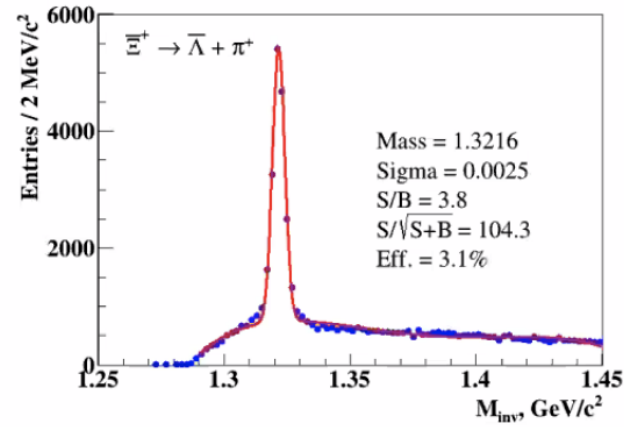
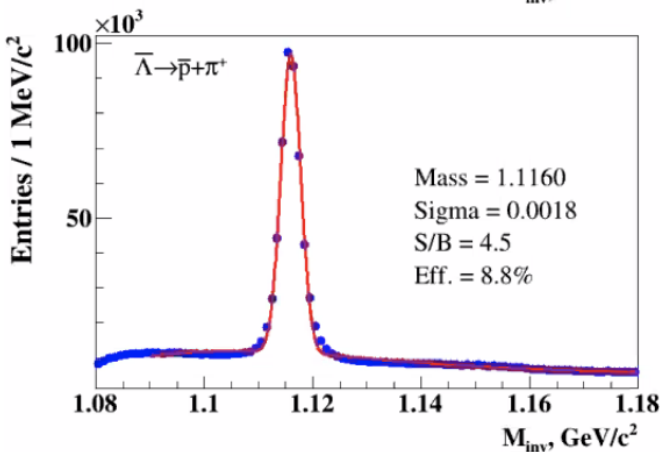
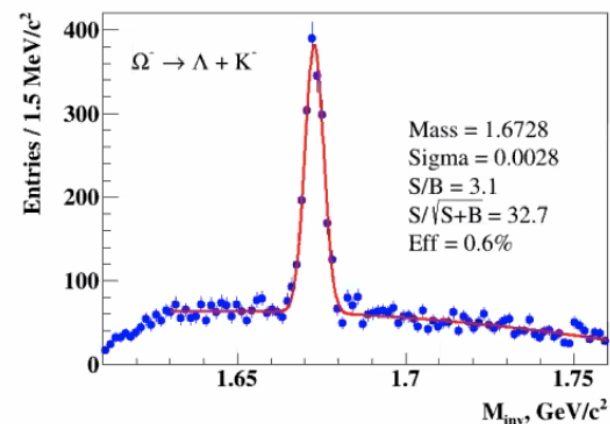
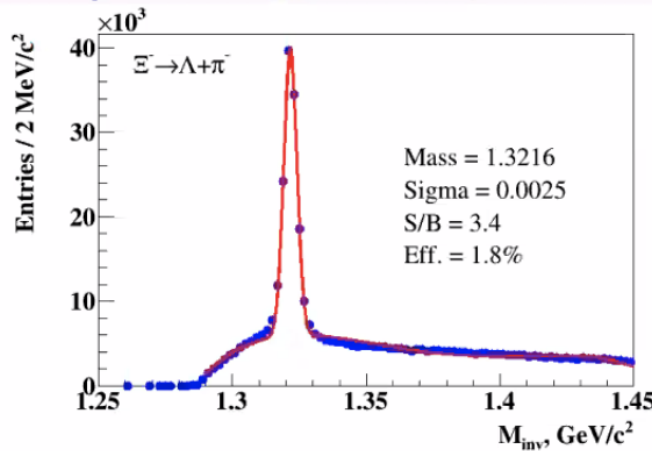
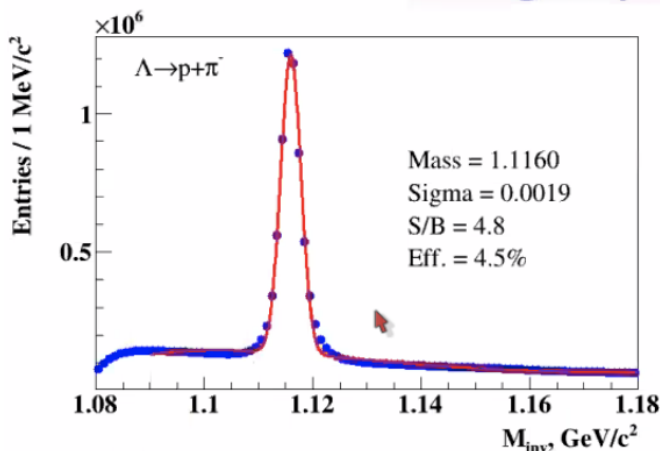


The NICA Facility at JINR Dubna



Strange and multi-strange baryons

Stage'1 (TPC+TOF): Au+Au @ 11 GeV, PHSD + MPDRoot reco.



particle	Λ	anti- Λ	Ξ^-	anti- Ξ^+	Ω^-	anti- Ω^+
yield in 10 weeks	$3 \cdot 10^8$	$3.5 \cdot 10^6$	$1.5 \cdot 10^6$	$8.0 \cdot 10^4$	$7 \cdot 10^4$	$1.5 \cdot 10^4$

The NICA Facility at JINR Dubna



NICA Facility running plan

- **Year 2021:**

- Extensive commissioning of Booster accelerator
- Heavy-ion (Fe/Kr/Xe) run of full Booster+Nucletron setup

- **Year 2022:**

- Completion of NICA Collider and transfer lines

- **Year 2023:**

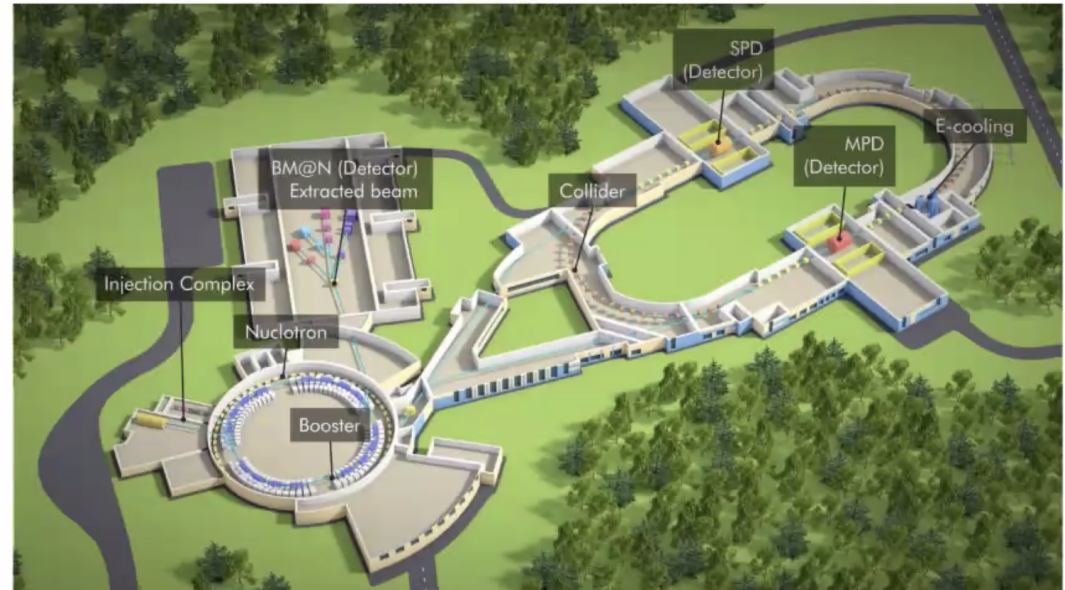
- Initial run of NICA with Bi+Bi @ 9.2 AGeV (other energies a second priority)
- Goal to reach luminosity of $10^{25} \text{ cm}^{-2}\text{s}^{-1}$

- **Year 2024:**

- Goal to have Au+Au collisions and acceleration in NICA (up to 11 AGeV)

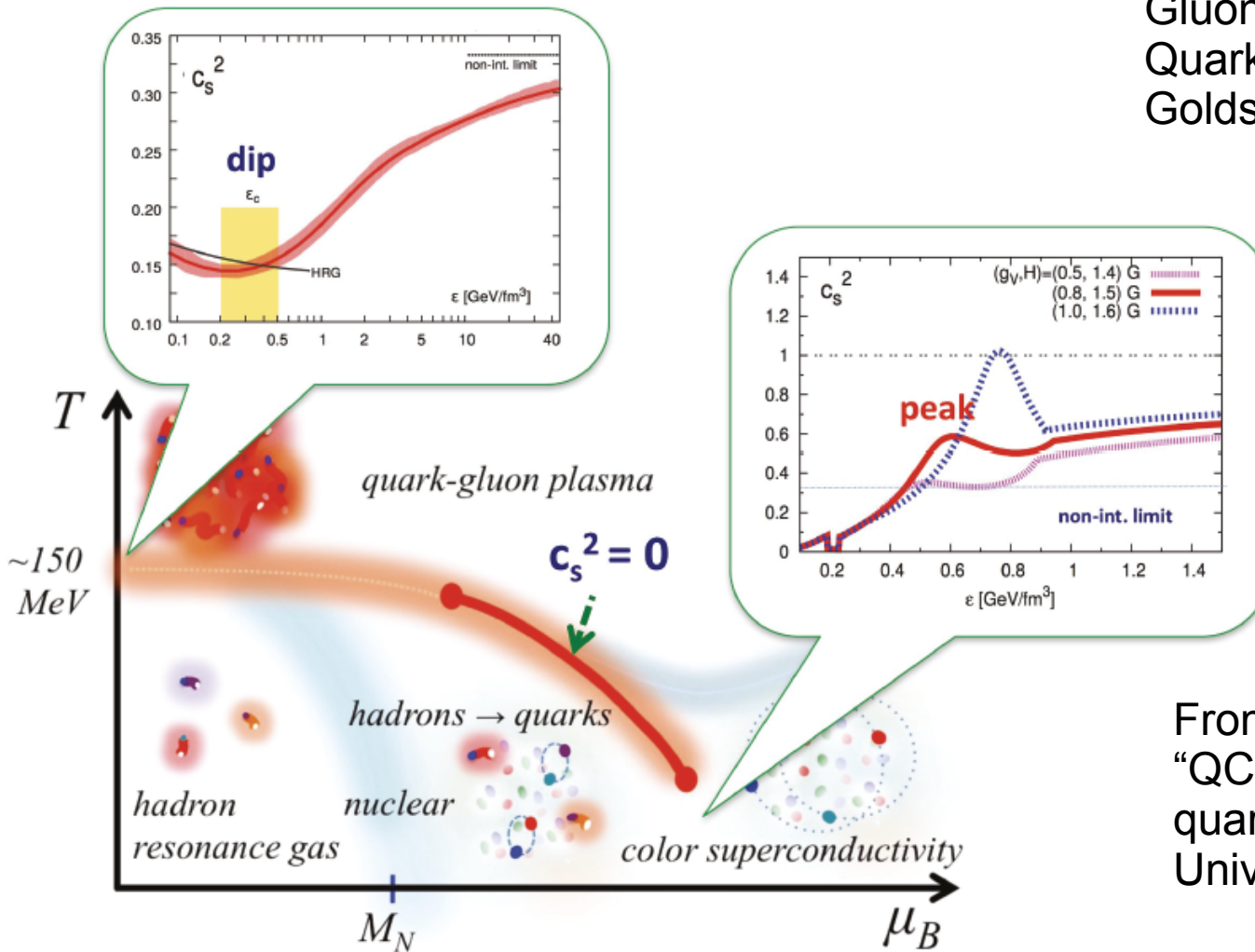
- **Beyond 2024:**

- Maximizing luminosity, possibility of collision energy and system size scan



2nd CEP in QCD phase diagram: Quark-Hadron Continuity?

Gluons \leftrightarrow Vector mesons
 Quarks \leftrightarrow Baryons
 Goldstones \leftrightarrow Pseudoscalar mesons



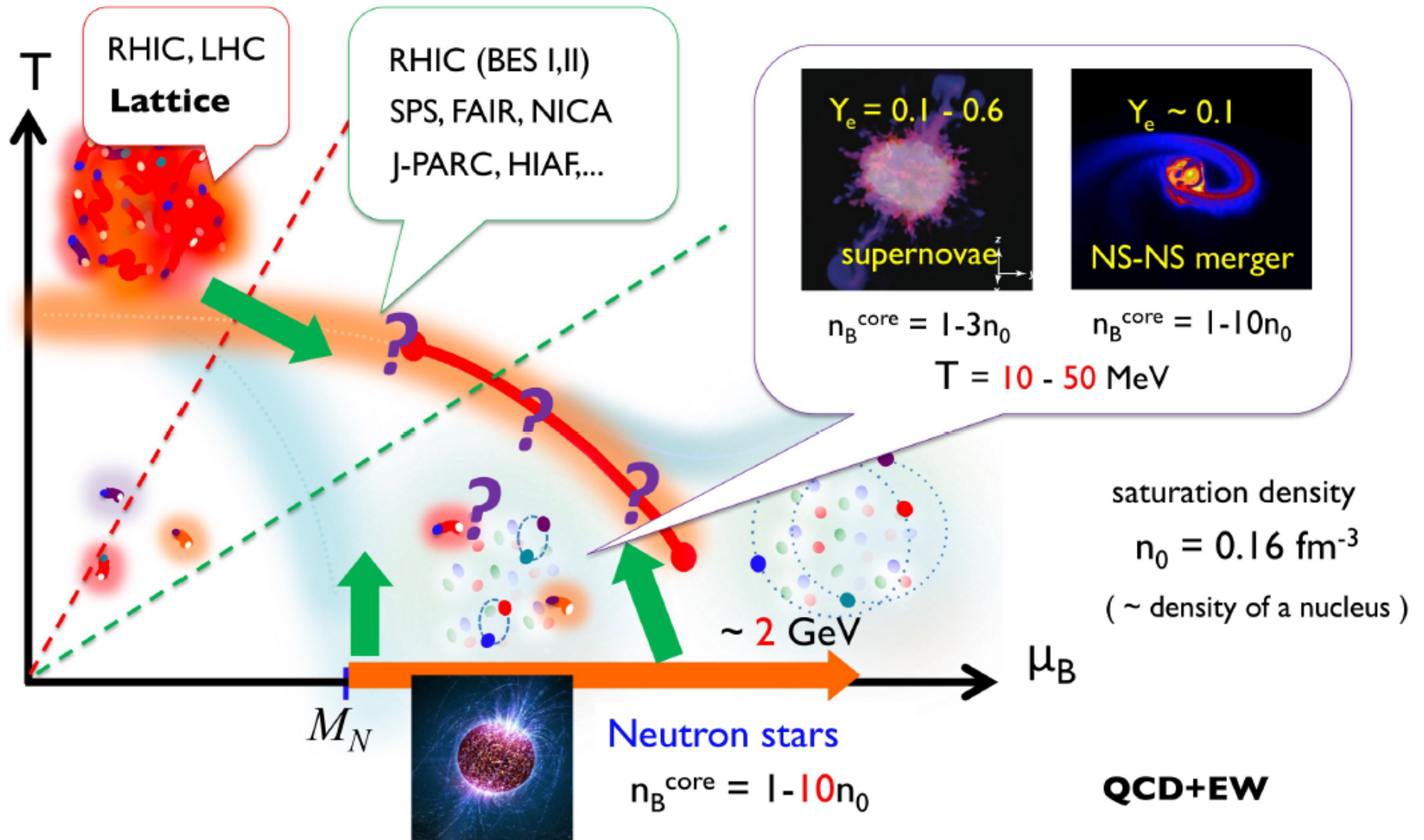
From: T. Kojo,
 "QCD equations of state in
 quark-hadron continuity",
 Universe 4 (2018) 42

T. Schaefer & F. Wilczek, Phys. Rev. Lett. 82 (1999) 3956

C. Wetterich, Phys. Lett. B 462 (1999) 164

T. Hatsuda, M. Tachibana, T. Yamamoto & G. Baym, Phys. Rev. Lett. 97 (2006) 122001

2nd or no CEP in QCD phase diagram: Crossover all over ?



Conclusions

- First observations of binary mergers open new possibilities to constrain properties of the Quark-gluon plasma at low temperatures and high baryon densities. Hybrid EoS are developed that allows to estimate quark plasma parameters in hypermassive (proto-) neutron stars
- GW170817: narrow window of small radii at $1.4 M_{\text{sun}}$ (Capano et al.: $10.4 < R_{1.4}[\text{km}] < 11.9$) strongly suggests an early onset of deconfinement with a critical density $n_c < 2 n_0$ and an onset mass $M_{\text{onset}} < 1.0 M_{\text{sun}}$ [Blaschke & Cierniak: 2012.15785]
- GW190814: the lighter object in the extremely asymmetric merger with its $2.6 M_{\text{sun}}$ can be either the heaviest neutron star or the lightest black hole. The central baryon density in such high-mass hybrid stars reaches $5.3 n_0$. Our EoS allows it to be a hybrid star ...
- NICER radius measurement on PSR J0740+6620 triggers a new paradigm: NS with $M > 2M_{\text{sun}}$ should have a deconfined quark matter core when $R_{2.0} > 13 \text{ km}$!

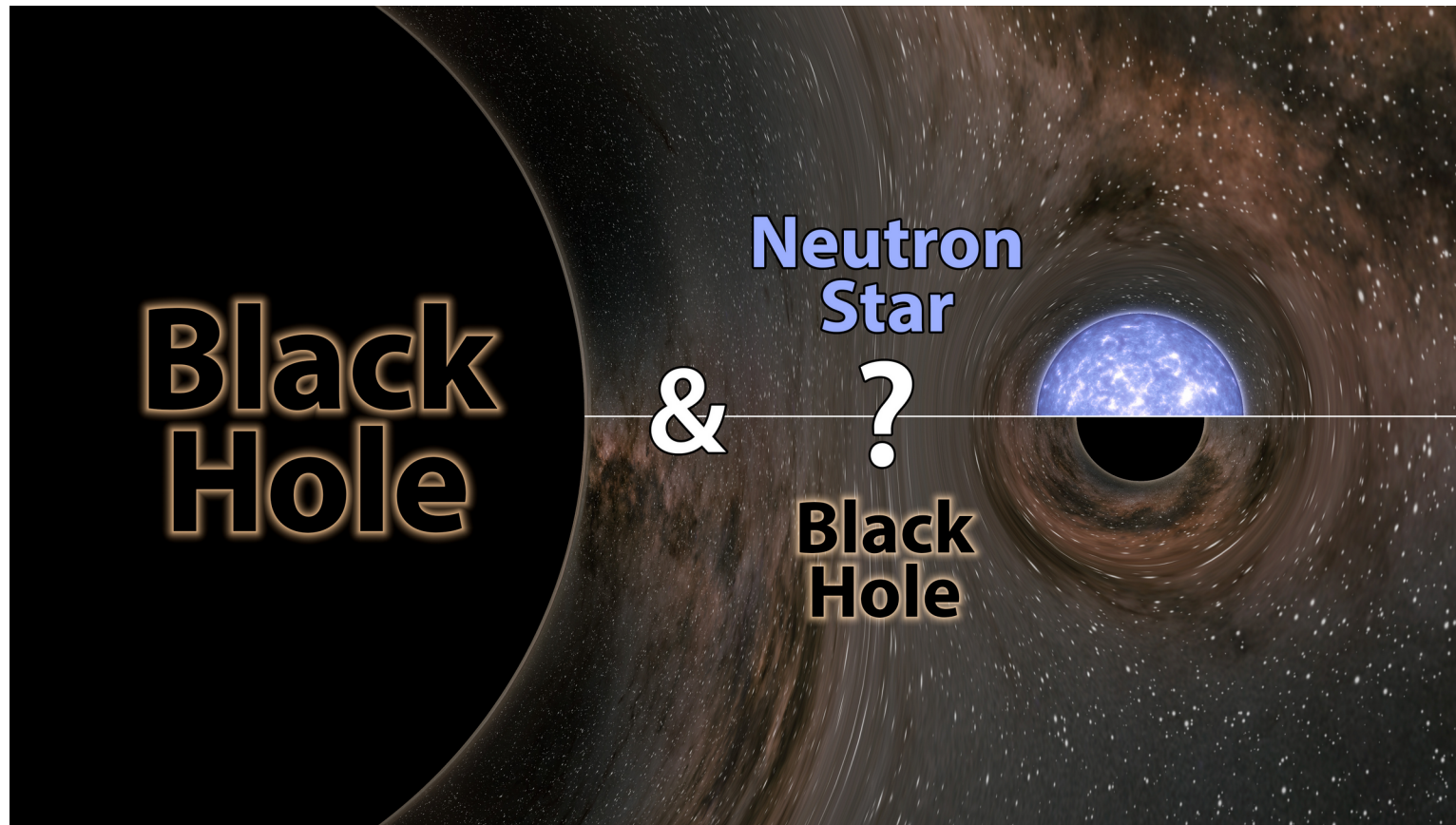
Such a result is similar to the “two families” scenario of Drago & Pagliara, PRD 102 (2020)
For the baryon density at the center of a star with $2.1 M_{\text{sun}}$ we find $n < 5 n_0$, $n_0 = 0.15 \text{ fm}^{-3}$.

- Consequences for supernova simulations: A new lower limit for onset of deconfinement?
- Consequences for merger simulations: Check the GW signal for deconfinement !
- Good news for entering a color superconducting quark matter phase at NICA (BMAN, MPD)

Backup Slides:

Limits of Neutron Star Physics

GW190814



What is the limiting
Mass of a neutron
Star?

Was GW190814 a
Merger of a 23- M_{sun}
Black hole with the

Lightest Black hole

Or

Heaviest Neutron star

at 2.6 M_{sun} ??

GW170817 – a merger of two compact stars

Neutron Star Merger Dynamics

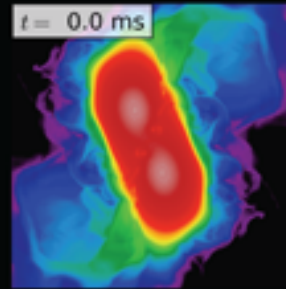
(General) Relativistic (Very) Heavy-Ion Collisions at ~ 100 MeV/nucleon

Simulations: Rezzola et al (2013)

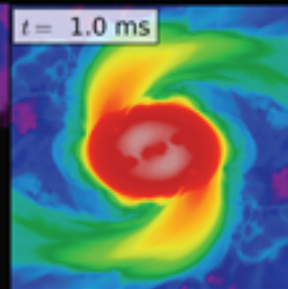
$t = -8.1$ ms



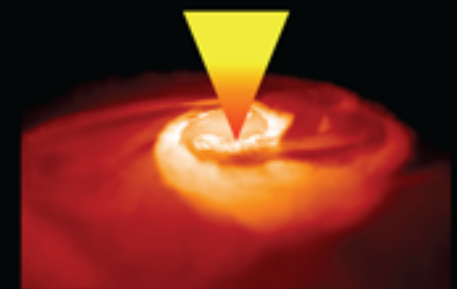
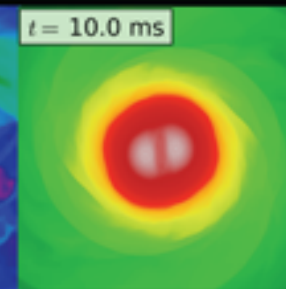
$t = 0.0$ ms



$t = 1.0$ ms



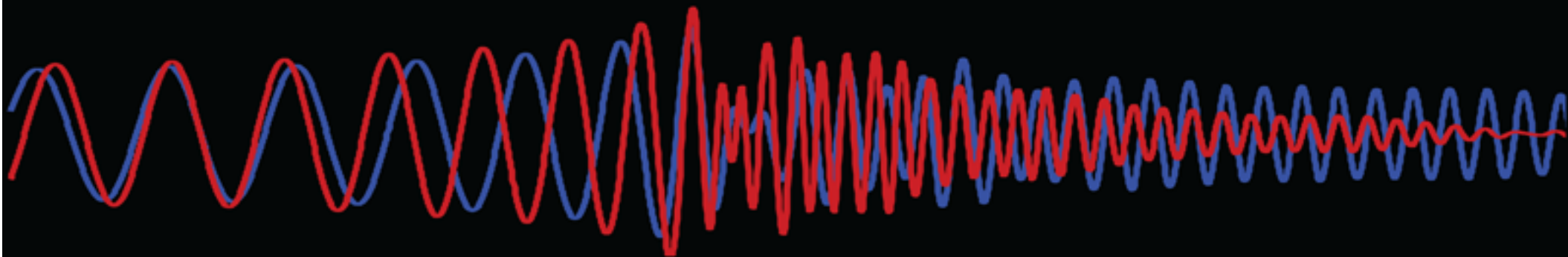
$t = 10.0$ ms



Inspiral:
Gravitational waves,
Tidal Effects

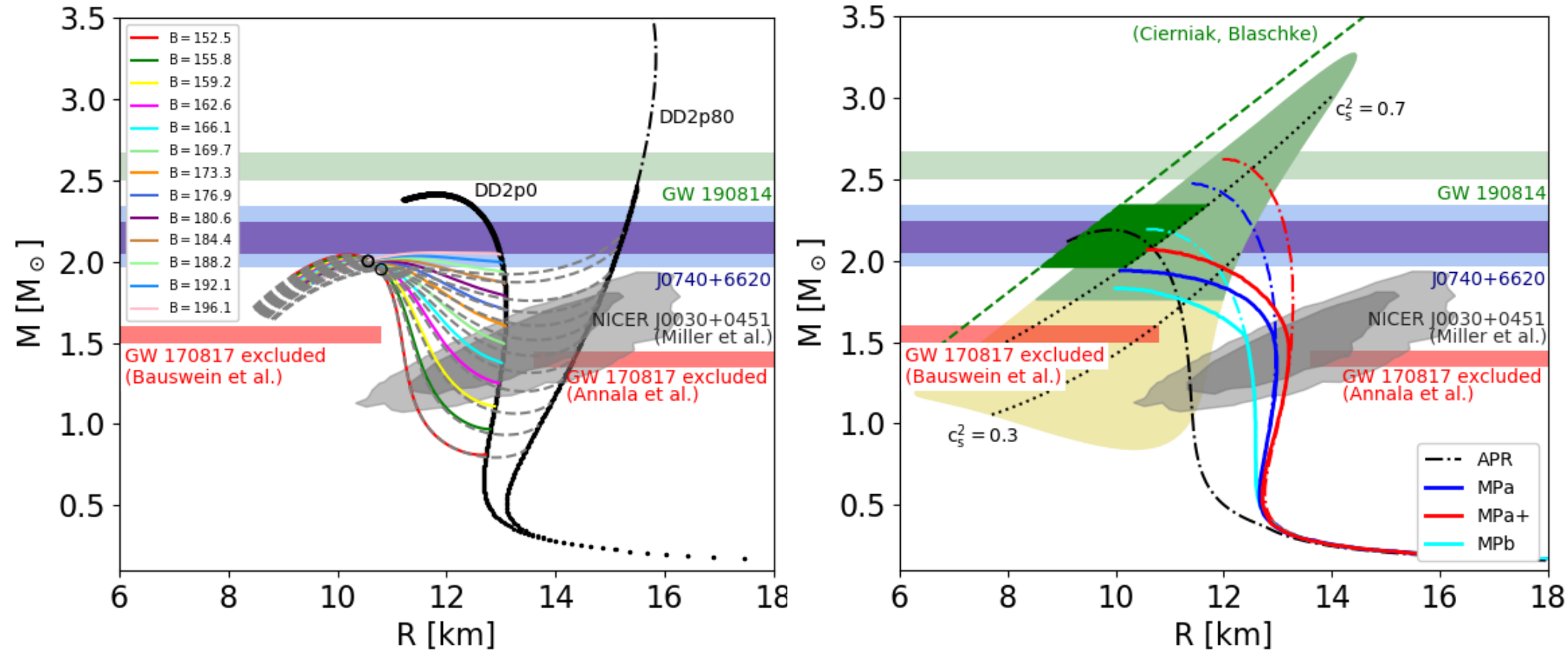
Merger:
Disruption, NS oscillations, ejecta
and r-process nucleosynthesis

Post Merger:
GRBs, Afterglows, and
Kilonova



Can NICER prove J0740+6620 to be a hybrid star?

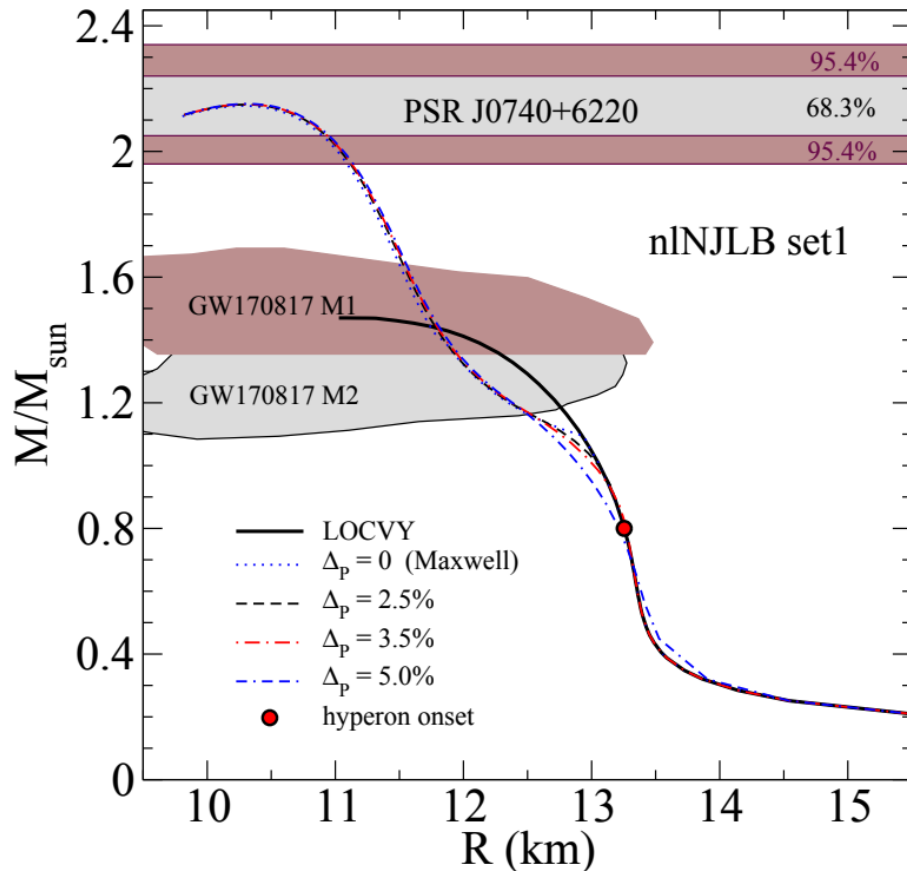
Work with Mateusz Cierniak, arxiv:2009.12353; EPJ ST 229 (2020) 3663
arxiv:2012.15785; AN (2021) accepted



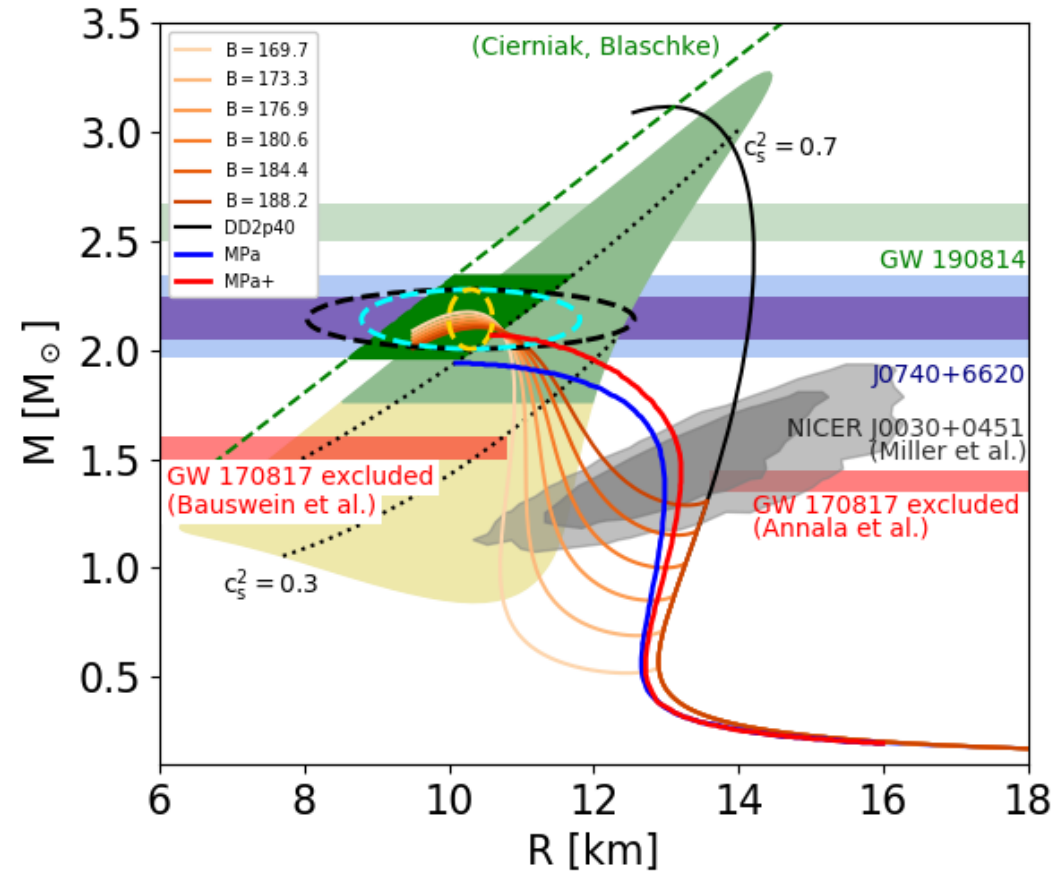
If radius of PSR J0740+6620 is measured in the dark-green region then it must harbor a core of superconducting quark matter!

Can NICER prove J0740+6620 to be a hybrid star?

Work with Mateusz Cierniak, arxiv:2009.12353; EPJ ST 229 (2020) 3663
 arxiv:2012.15785; AN (2021) accepted



If radius of PSR J0740+6620 is measured at ~ 10.5 km, then it is also compatible with the hybrid star solution of the hyperon puzzle; M. ShahrbaF et al., J. Phys. G 47 (2020) 115201



If radius of PSR J0740+6620 is measured at 10.2 km with the accuracy of the yellow ellipse, then it must harbor a core of superconducting quark matter!