

NA7-Hf-QGP “STRONG 2020” Workshop

Fluid dynamics of multiple conserved charges

Jan Fotakis

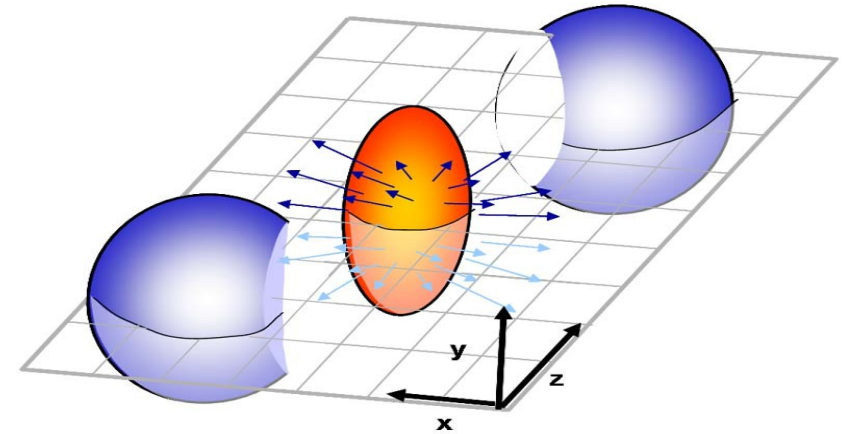
University of Frankfurt

Harri Niemi, Etele Molnár, Gabriel Denicol, Dirk Rischke, Carsten Greiner

Traditionally:

Viewed as 'blob' of one type of matter (single component) with one velocity field

- usually 'blob' of energy without charge

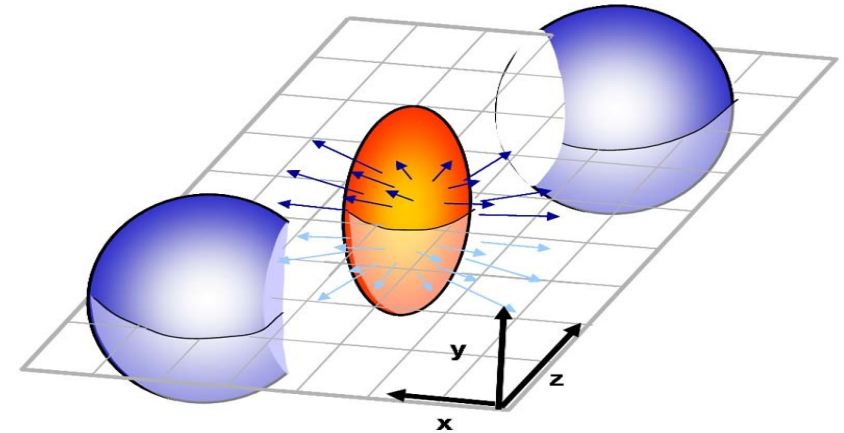


<https://www.quantumdiaries.org/wp-content/uploads/2011/02/FlowPr.jpg>

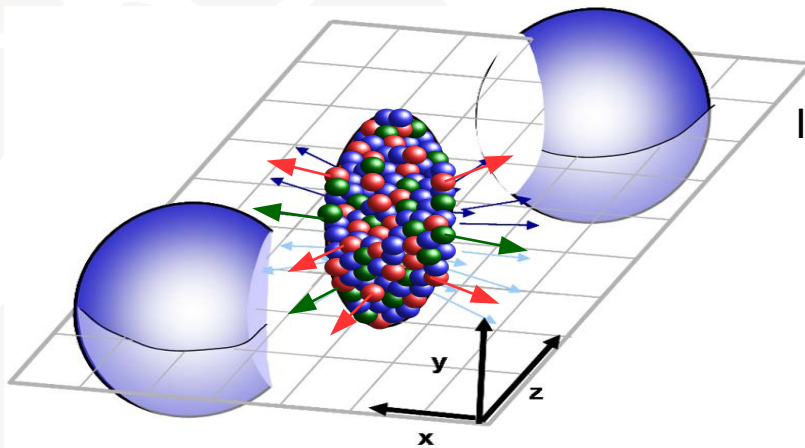
Traditionally:

Viewed as 'blob' of one type of matter (single component) with one velocity field

- usually 'blob' of energy without charge



<https://www.quantumdiaries.org/wp-content/uploads/2011/02/FlowPr.jpg>



In general:

Consists of multiple components with various properties with multiple velocity fields

- with **multiple conserved quantities** (e.g. energy, electric charge, baryon number, strangeness, ...)
- mixed chemistry → **coupled charge currents!**

Hydrodynamics: macroscopic effective field theory of thermal matter close to local equilibrium

Here: single fluid velocity field u^μ

$$T^{\mu\nu} = \sum_i T_i^{\mu\nu} = \epsilon u^\mu u^\nu - (P_0 + \Pi) \Delta^{\mu\nu} + \pi^{\mu\nu}$$

$$N_q^\mu = \sum_i q_i N_i^\mu = n_q u^\mu + V_q^\mu$$

Hydrodynamics: macroscopic effective field theory of thermal matter close to local equilibrium

Here: single fluid velocity field u^μ

$$T^{\mu\nu} = \sum_i T_i^{\mu\nu} = \epsilon u^\mu u^\nu - (P_0 + \Pi) \Delta^{\mu\nu} + \pi^{\mu\nu}$$

Conservation of Energy and Momentum: $\partial_\mu T^{\mu\nu} = 0$

q-th conserved charge (eg. B, Q, S)

$$N_{[q]}^\mu = \sum_i [q_i] N_i^\mu = n_{[q]} u^\mu + V_{[q]}^\mu$$

Conservation of charge: $\partial_\mu N_{[q]}^\mu = 0$

Fluid dynamics of multi-component systems

Hydrodynamics: macroscopic effective field theory of thermal matter close to local equilibrium

Here: single fluid velocity field u^μ

$$T^{\mu\nu} = \sum_i T_i^{\mu\nu} = \epsilon u^\mu u^\nu - (P_0 + \Pi) \Delta^{\mu\nu} + \pi^{\mu\nu}$$

Conservation of Energy and Momentum: $\partial_\mu T^{\mu\nu} = 0$

q-th conserved charge (eg. B, Q, S)

$$N_{[q]}^\mu = \sum_i [q_i] N_i^\mu = n_{[q]} u^\mu + V_{[q]}^\mu$$

Conservation of charge: $\partial_\mu N_{[q]}^\mu = 0$

$10 + 4N_{\text{ch}}$ degrees of freedom, $4 + N_{\text{ch}}$ equations \rightarrow $6 + 3N_{\text{ch}}$ unknowns

Hydrodynamics: macroscopic effective field theory of thermal matter close to local equilibrium

Here: single fluid velocity field u^μ

$$T^{\mu\nu} = \sum_i T_i^{\mu\nu} = \epsilon u^\mu u^\nu - (P_0 + \Pi) \Delta^{\mu\nu} + \pi^{\mu\nu}$$

Conservation of Energy and Momentum: $\partial_\mu T^{\mu\nu} = 0$

q-th conserved charge (eg. B, Q, S)

$$N_{[q]}^\mu = \sum_i [q_i] N_i^\mu = n_{[q]} u^\mu + V_{[q]}^\mu$$

Conservation of charge: $\partial_\mu N_{[q]}^\mu = 0$

$10 + 4N_{\text{ch}}$ degrees of freedom, $4 + N_{\text{ch}}$ equations \rightarrow $6 + 3N_{\text{ch}}$ unknowns

What needs to be known:

- Equation of state
- Equations of motion for dissipative fields & transport coefficients
- Initial state
- Freeze-out and δf -correction

Hydrodynamics: macroscopic effective field theory of thermal matter close to local equilibrium

Here: single fluid velocity field u^μ

$$T^{\mu\nu} = \sum_i T_i^{\mu\nu} = \epsilon u^\mu u^\nu - (P_0 + \Pi) \Delta^{\mu\nu} + \pi^{\mu\nu}$$

Conservation of Energy and Momentum: $\partial_\mu T^{\mu\nu} = 0$

q-th conserved charge (eg. B, Q, S)

$$N_{[q]}^\mu = \sum_i [q_i] N_i^\mu = n_{[q]} u^\mu + V_{[q]}^\mu$$

Conservation of charge: $\partial_\mu N_{[q]}^\mu = 0$

$10 + 4N_{\text{ch}}$ degrees of freedom, $4 + N_{\text{ch}}$ equations \rightarrow $6 + 3N_{\text{ch}}$ unknowns

What needs to be known:

- Equation of state
- Equations of motion for dissipative fields & [transport coefficients](#)
- Initial state
- Freeze-out and δf -correction

Fluid dynamics with conserved baryon number:

Denicol et al., PRC 98, 034916 (2018)

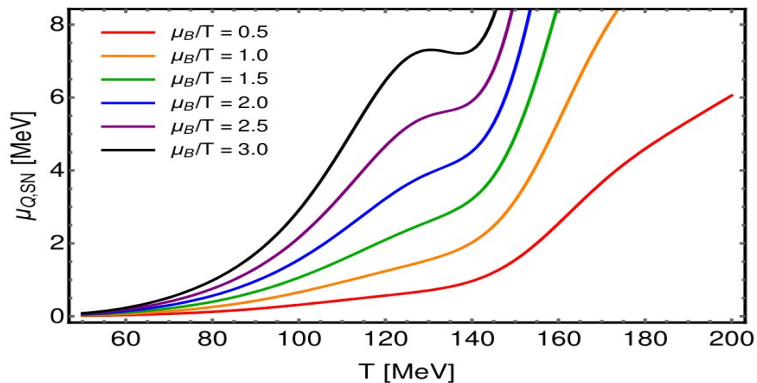
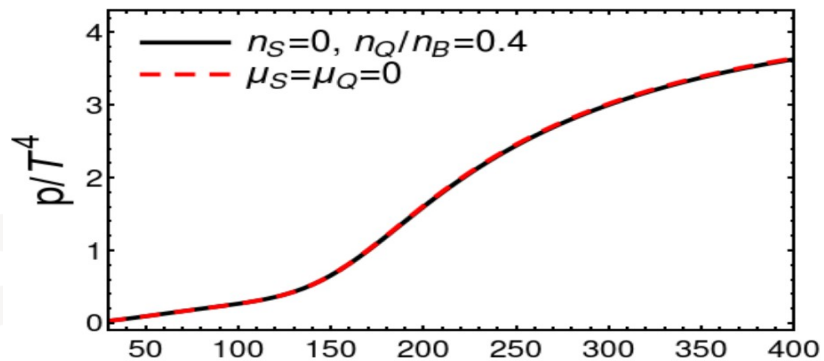
Du et al., Comp. Phys. Comm. 251, 107090 (2020)

Li et al., PRC 98, 064908 (2018)

Equation of state with multiple conserved charges

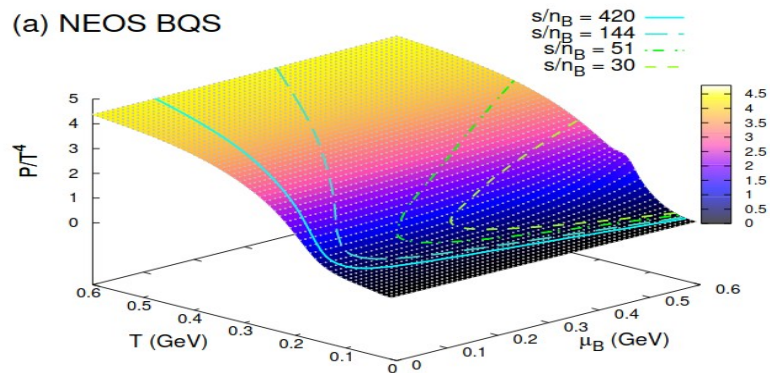
$$P_0(T) \rightarrow P_0(T, \mu_B, \mu_Q, \mu_S)$$

Noronha-Hostler et al., PRC 100, 064910 (2019)

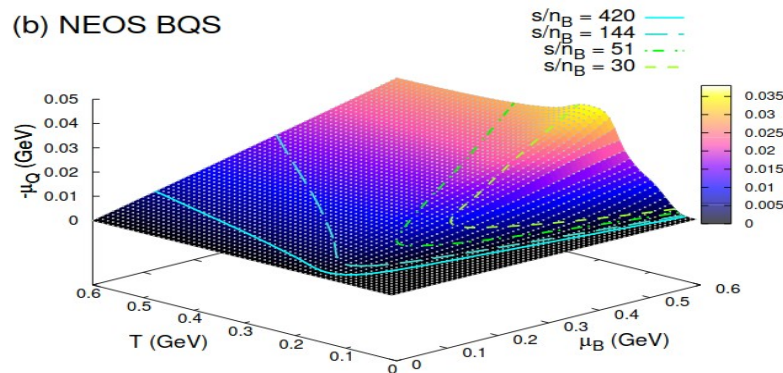


Monnai et al., PRC 100, 024907 (2019)

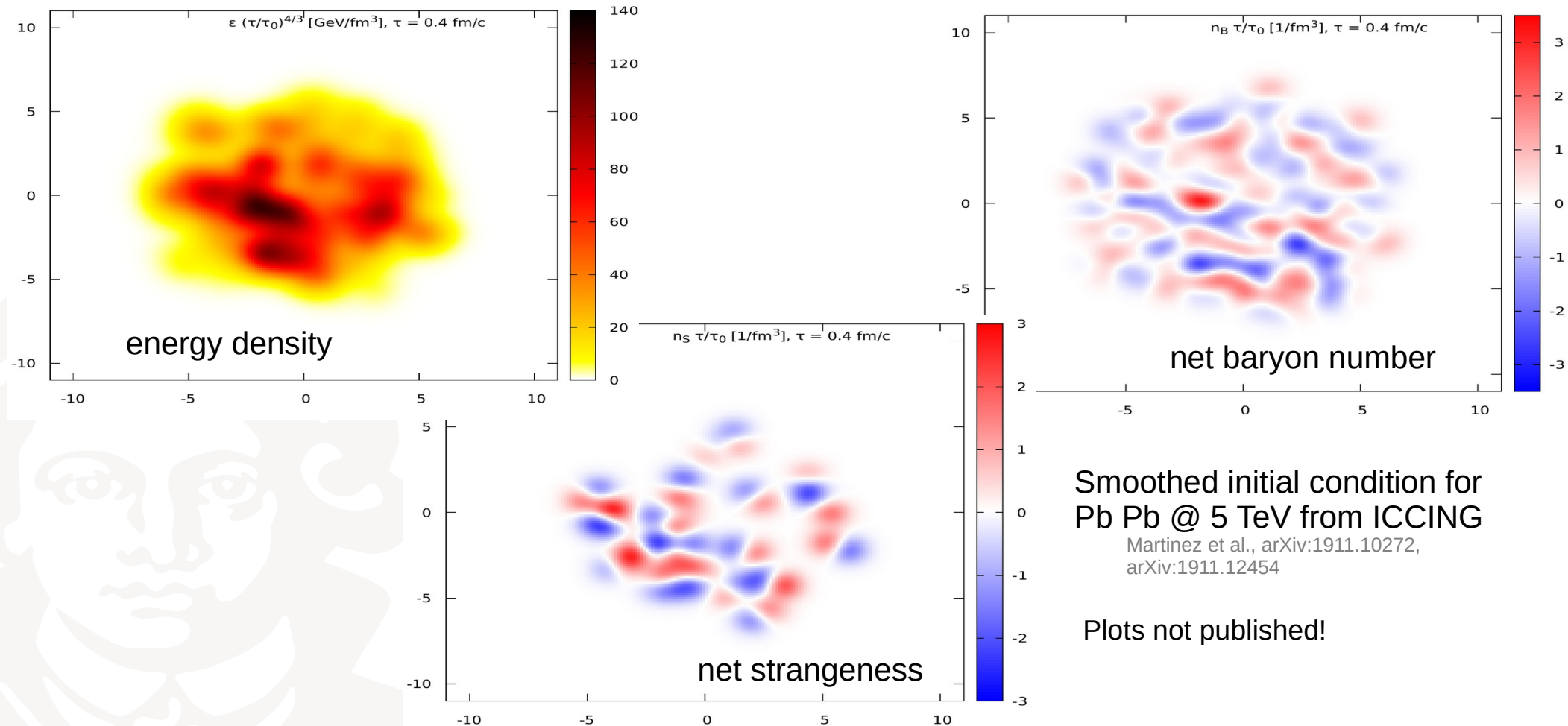
(a) NEOS BQS



(b) NEOS BQS



Initial state with multiple conserved charges



Deriving fluid dynamics from kinetic theory

Denicol et al., PRD 85, 114047 (2012)

On basis of DNMR theory: derivation from the Boltzmann equation

→ **upcoming publication!** (Fotakis, Molnár, Niemi, Denicol, Rischke, Greiner)

Also refer to: Monnai, Hirano, Nucl. Phys. A847:283-314 (2010) or Kikuchi et al. PRC 92, 064909 (2015)

relativistic Boltzmann eq.

$$k_i^\mu \partial_\mu f_{i,\mathbf{k}} = C_i[f_i]$$



2nd-order (multi-component) hydro

$$\dot{\Pi}, \dot{V}_q^{\langle\mu\rangle}, \dot{\pi}^{\langle\mu\nu\rangle}$$

Deriving fluid dynamics from kinetic theory

Denicol et al., PRD 85, 114047 (2012)

On basis of DNMR theory: derivation from the Boltzmann equation

→ upcoming publication! (Fotakis, Molnár, Niemi, Denicol, Rischke, Greiner)

Also refer to: Monnai, Hirano, Nucl. Phys. A847:283-314 (2010) or Kikuchi et al. PRC 92, 064909 (2015)

relativistic Boltzmann eq.

$$k_i^\mu \partial_\mu f_{i,\mathbf{k}} = C_i[f_i]$$



2nd-order (multi-component) hydro

$$\dot{\Pi}, \dot{V}_q^{\langle\mu\rangle}, \dot{\pi}^{\langle\mu\nu\rangle}$$

Idea:

Look at off-equilibrium moments:

$$\rho_{i,n}^{\mu\nu} = \sum_{i=1}^{N_{\text{species}}} \int \frac{d^3\mathbf{k}_i}{(2\pi)^3 E_{i,\mathbf{k}}} E_{i,\mathbf{k}}^n k_i^{\langle\mu} k_i^{\nu\rangle} \left(f_{i,\mathbf{k}} - f_{i,\mathbf{k}}^{(0)} \right)$$

Relate them to the fluid-dynamical fields with constituent's transport coefficients:

$$\rho_{i,n}^{\mu\nu} = \frac{\eta_{i,n}}{\eta} \pi^{\mu\nu} + \mathcal{O}(2)$$

Deriving fluid dynamics from kinetic theory

Denicol et al., PRD 85, 114047 (2012)

On basis of DNMR theory: derivation from the Boltzmann equation

→ **upcoming publication!** (Fotakis, Molnár, Niemi, Denicol, Rischke, Greiner)

Also refer to: Monnai, Hirano, Nucl. Phys. A847:283-314 (2010) or Kikuchi et al. PRC 92, 064909 (2015)

relativistic Boltzmann eq.

$$k_i^\mu \partial_\mu f_{i,\mathbf{k}} = C_i[f_i]$$



2nd-order (multi-component) hydro

$$\dot{\Pi}, \dot{V}_q^{\langle\mu\rangle}, \dot{\pi}^{\langle\mu\nu\rangle}$$

$$\tau_\Pi \dot{\Pi} + \Pi = -\zeta\theta + \mathcal{O}(2)[\Pi, V_q^\mu, \pi^{\mu\nu}, \theta, \sigma^{\mu\nu}, \nabla^\mu P_0, \nabla\alpha_q]$$

$$\sum_{q'} \tau_{qq'} \dot{V}_{q'}^{\langle\mu\rangle} + V_q^\mu = \sum_{q'} \kappa_{qq'} \nabla^\mu \alpha_{q'} + \mathcal{O}(2)[\Pi, V_q^\mu, \pi^{\mu\nu}, \theta, \sigma^{\mu\nu}, \omega^{\mu\nu}, \nabla^\mu P_0, \nabla\alpha_q]$$

$$\tau_\pi \dot{\pi}^{\langle\mu\nu\rangle} + \pi^{\mu\nu} = 2\eta\sigma^{\mu\nu} + \mathcal{O}(2)[\Pi, V_q^\mu, \pi^{\mu\nu}, \theta, \sigma^{\mu\nu}, \omega^{\mu\nu}, \nabla^\mu P_0, \nabla\alpha_q]$$

Deriving fluid dynamics from kinetic theory

Denicol et al., PRD 85, 114047 (2012)

On basis of DNMR theory: derivation from the Boltzmann equation

→ **upcoming publication!** (Fotakis, Molnár, Niemi, Denicol, Rischke, Greiner)

Also refer to: Monnai, Hirano, Nucl. Phys. A847:283-314 (2010) or Kikuchi et al. PRC 92, 064909 (2015)

relativistic Boltzmann eq.

$$k_i^\mu \partial_\mu f_{i,\mathbf{k}} = C_i[f_i]$$

2nd-order (multi-component) hydro

$$\dot{\Pi}, \dot{V}_q^{\langle\mu\rangle}, \dot{\pi}^{\langle\mu\nu\rangle}$$

$$\tau_\Pi \dot{\Pi} + \Pi = \boxed{-\zeta\theta} + \boxed{\mathcal{O}(2)[\Pi, V_q^\mu, \pi^{\mu\nu}, \theta, \sigma^{\mu\nu}, \nabla^\mu P_0, \nabla\alpha_q]}$$

$$\sum_{q'} \tau_{qq'} \dot{V}_{q'}^{\langle\mu\rangle} + V_q^\mu = \boxed{\sum_{q'} \kappa_{qq'} \nabla^\mu \alpha_{q'}} + \boxed{\mathcal{O}(2)[\Pi, V_q^\mu, \pi^{\mu\nu}, \theta, \sigma^{\mu\nu}, \omega^{\mu\nu}, \nabla^\mu P_0, \nabla\alpha_q]}$$

$$\tau_\pi \dot{\pi}^{\langle\mu\nu\rangle} + \pi^{\mu\nu} = \boxed{2\eta\sigma^{\mu\nu}} + \boxed{\mathcal{O}(2)[\Pi, V_q^\mu, \pi^{\mu\nu}, \theta, \sigma^{\mu\nu}, \omega^{\mu\nu}, \nabla^\mu P_0, \nabla\alpha_q]}$$

1st order terms (Navier-Stokes): mixed chemistry already couples diffusion currents!

2nd order terms: couples all currents to each other; depend on all gradients!

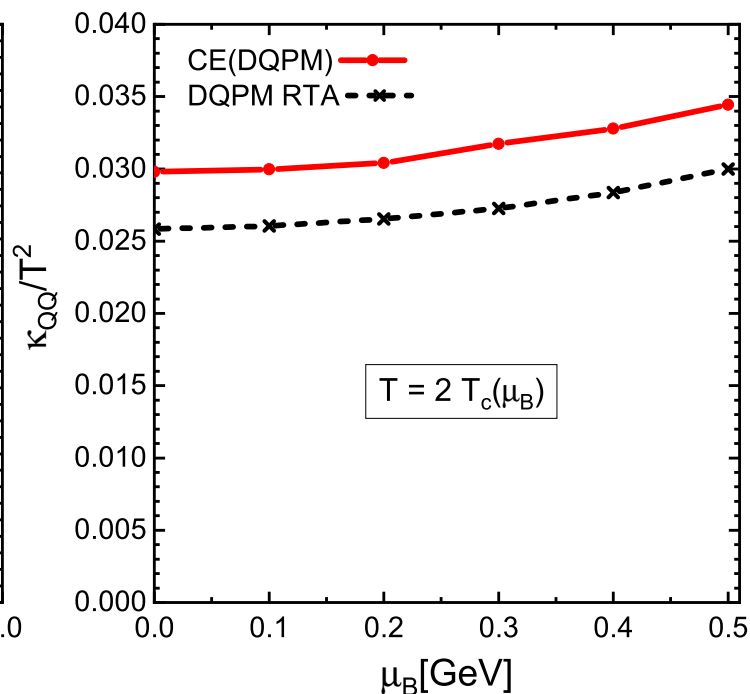
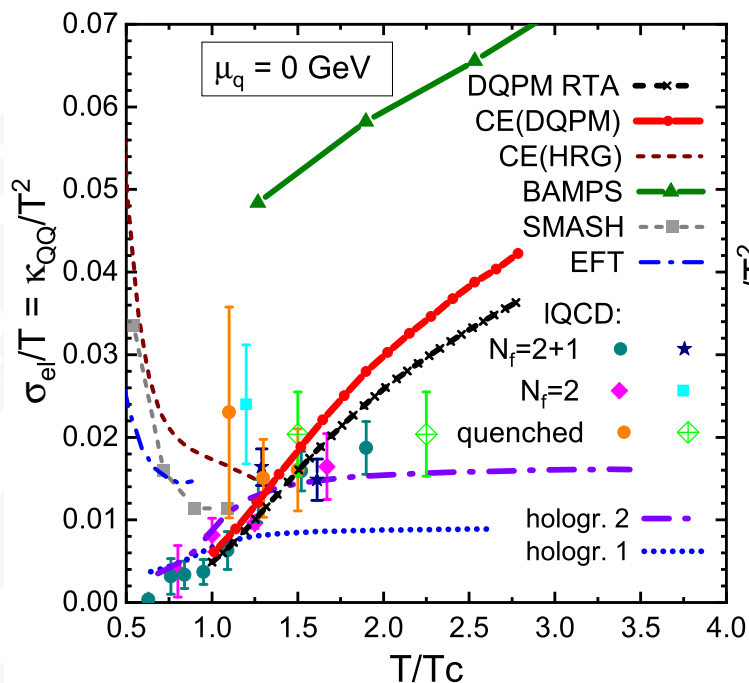
→ 3 conserved charges: 70+ transport coefficients (!!) with (T, μ_B, μ_Q, μ_S) -dependence

Computation of transport coefficients (Example: diffusion coefficients)

$$\kappa_{qq'} = \sum_{n=0}^{\infty} \sum_{i,j=1}^{N_{\text{species}}} (c^{-1})_{ji,0n}^{(1)} q_j \left(q'_i \mathcal{J}_{n+1,1}^{(i)} - \frac{n_{q'}}{\epsilon + P_0} \mathcal{J}_{n+2,1}^{(i)} \right)$$

Example: introduction of features from LQCD via the usage of DQPM

Fotakis, Soloveva et al, PRD 104, 034014 (2021)



A simplistic case study

Greif, Fotakis et al., PRL 120, 242301 (2018)
Fotakis, Greif et al., PRD 101, 076007 (2020)

- Investigate longitudinal evolution in Milne coordinates (transversally homogeneous)
- Conserved baryon number and strangeness, **neglect viscosity**, neglect 2nd order terms

~~$$\tau_{\Pi} \dot{\Pi} + \Pi = -\zeta \theta + \mathcal{O}(2)$$~~

~~$$\tau_{\pi} \dot{\pi}^{\langle \mu \nu \rangle} + \pi^{\mu \nu} = 2\eta \sigma^{\mu \nu} + \mathcal{O}(2)$$~~

~~$$\sum_{q'} \tau_{q'} \dot{V}_{q'}^{\langle \mu \rangle} + V_q^{\mu} = \sum_{q'} \kappa_{qq'} \nabla^{\mu} \alpha_{q'} + \mathcal{O}(2)$$~~

Only 3 transport coefficients left!

A simplistic case study

Greif, Fotakis et al., PRL 120, 242301 (2018)
Fotakis, Greif et al., PRD 101, 076007 (2020)

- Investigate **longitudinal** evolution in **Milne coordinates** (transversally homogeneous)
- Conserved **baryon number and strangeness**, **neglect viscosity**, neglect 2nd order terms
- **Equation of state**: Non-interacting, classical statistics,
Hadronic system with 19 lightest (stable) particle species

$$\tau_q \dot{V}_q^{\langle \mu \rangle} + V_q^\mu = \sum_{q'} \kappa_{qq'} \nabla^\mu \alpha_{q'}$$

Only 3 transport coefficients left!



A simplistic case study

Greif, Fotakis et al., PRL 120, 242301 (2018)
Fotakis, Greif et al., PRD 101, 076007 (2020)

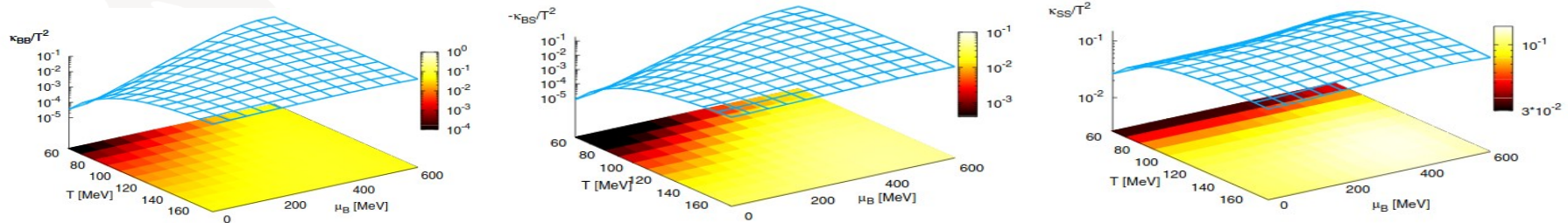
- Investigate **longitudinal** evolution in **Milne coordinates** (transversally homogeneous)
- Conserved **baryon number and strangeness**, **neglect viscosity**, neglect 2nd order terms
- **Equation of state**: Non-interacting, classical statistics, **Hadronic system** with 19 lightest (stable) particle species

$$\tau_q \dot{V}_q^{\langle \mu \rangle} + V_q^\mu = \sum_{q'} \kappa_{qq'} \nabla^\mu \alpha_{q'}$$

Only 3 transport coefficients left!

- **Diffusion coefficient matrix**:

- Assumed **elastic, isotropic, binary cross sections** from PDG, SMASH, GiBUU and UrQMD



A simplistic case study

Greif, Fotakis et al., PRL 120, 242301 (2018)
Fotakis, Greif et al., PRD 101, 076007 (2020)

- Investigate **longitudinal** evolution in **Milne coordinates** (transversally homogeneous)
- Conserved **baryon number and strangeness**, **neglect viscosity**, neglect 2nd order terms

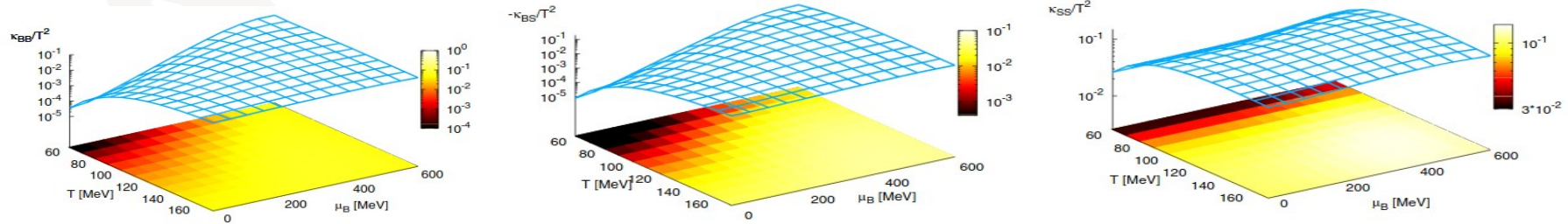
- **Equation of state**: Non-interacting, classical statistics,
Hadronic system with 19 lightest (stable) particle species

$$\tau_q \dot{V}_q^{\langle \mu \rangle} + V_q^\mu = \sum_{q'} \kappa_{qq'} \nabla^\mu \alpha_{q'}$$

Only 3 transport coefficients left!

- **Diffusion coefficient matrix**:

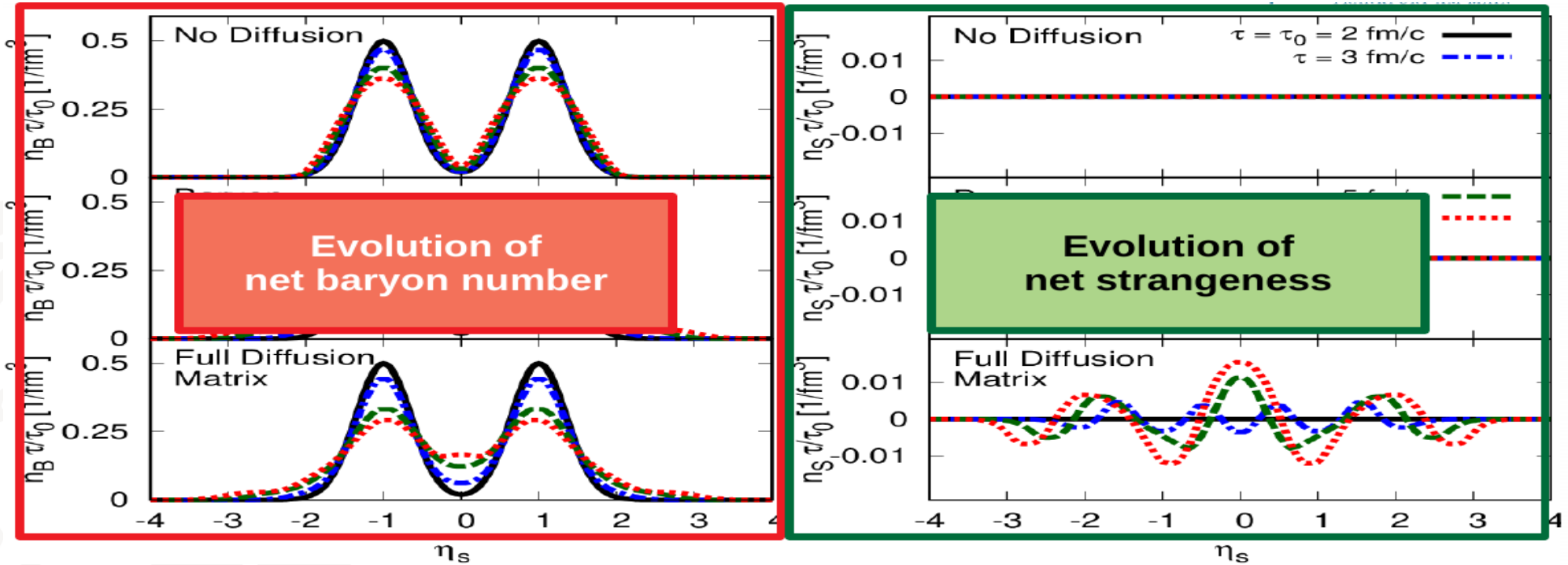
- Assumed **elastic, isotropic, binary cross sections** from PDG, SMASH, GiBUU and UrQMD



- **Simple initial state**: $T = 160$ MeV, **no initial net strangeness**, longitudinal double-gaussian profile in net baryon number, no initial dissipative currents

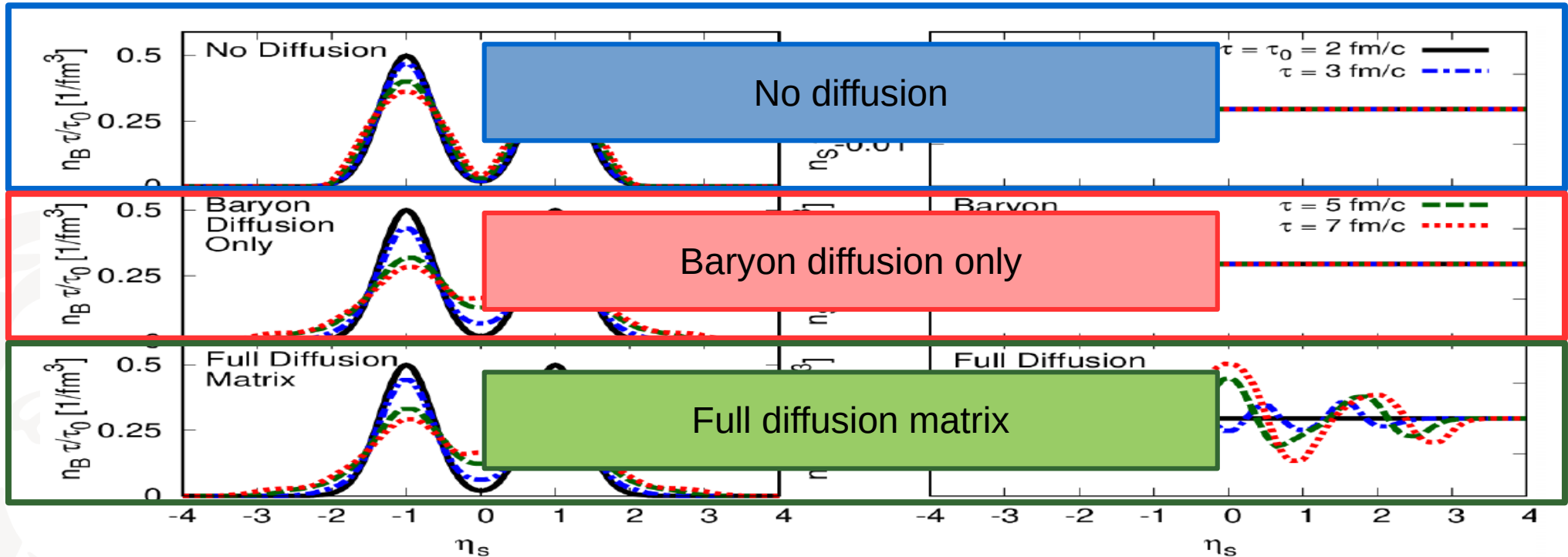
A simplistic case study

Greif, Fotakis et al., PRL 120, 242301 (2018)
Fotakis, Greif et al., PRD 101, 076007 (2020)



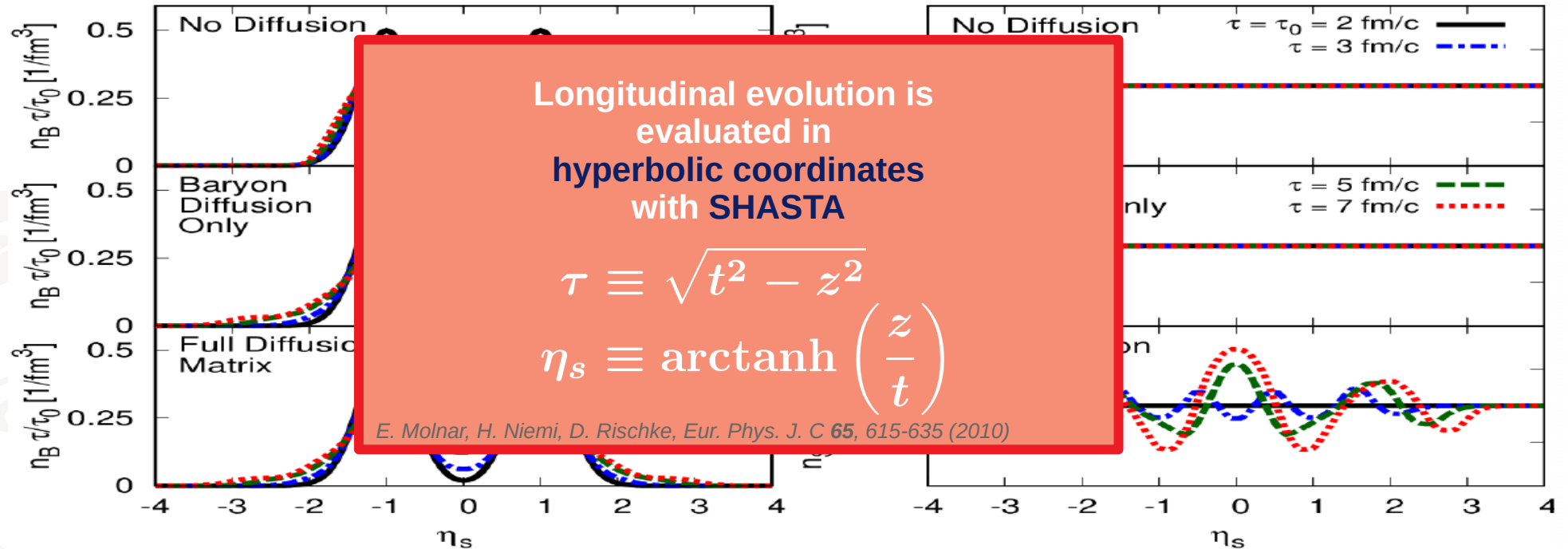
A simplistic case study

Greif, Fotakis et al., PRL 120, 242301 (2018)
Fotakis, Greif et al., PRD 101, 076007 (2020)



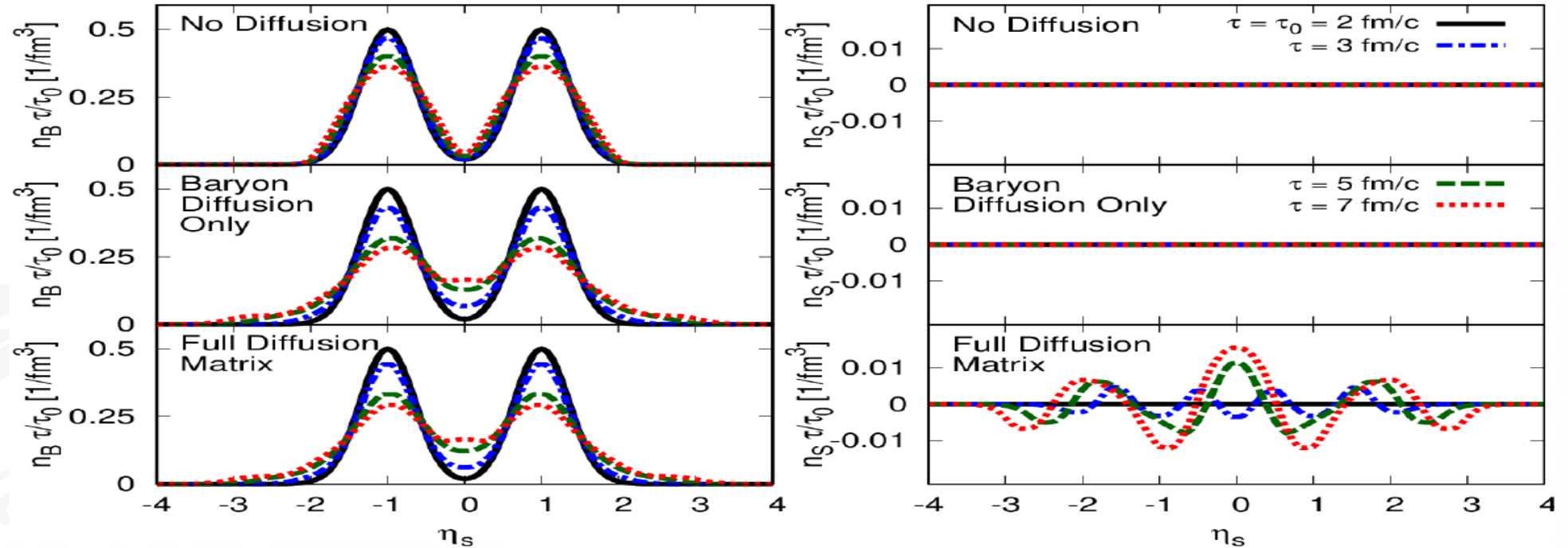
A simplistic case study

Greif, Fotakis et al., PRL 120, 242301 (2018)
 Fotakis, Greif et al., PRD 101, 076007 (2020)



A simplistic case study

Greif, Fotakis et al., PRL 120, 242301 (2018)
 Fotakis, Greif et al., PRD 101, 076007 (2020)



Mixed chemistry couples diffusion currents and introduces charge-correlation through EoS

$$\mu_S \equiv \mu_S(\epsilon, n_B, n_S)$$

e.g.: $\nabla^\mu \alpha_S \sim \nabla^\mu n_B$



Generation of domains of non-vanishing local net charge (here net strangeness)!

- Implement derived fluid dynamic theory in **existing (3+1)D-hydro code**
- Evaluate **2nd order transport coefficients** from linearized Boltzmann equation
- Use more realistic **initial state** and **equation of state** (see above)
- Apply **freeze-out routines**, take δf -correction
- Find **observables** sensitive to charge-coupling \rightarrow investigate impact

Backup



Computation of transport coefficients (Example: diffusion coefficients)

On basis of DNMR theory: derivation from the Boltzmann equation

→ **upcoming publication!** (Fotakis, Molnár, Niemi, Denicol, Rischke, Greiner)

relativistic Boltzmann eq.

$$k_i^\mu \partial_\mu f_{i,\mathbf{k}} = C_i[f_i]$$

2nd-order (multi-component) hydro

$$\dot{\Pi}, \dot{V}_q^{\langle\mu\rangle}, \dot{\pi}^{\langle\mu\nu\rangle}$$

$$C_{i,n-1}^{\langle\mu\rangle} \equiv \int \frac{d^3\mathbf{k}_i}{(2\pi)^3 E_{i,\mathbf{k}}} E_{i,\mathbf{k}}^{n-1} k_i^{\langle\mu\rangle} C_i[f_i]$$

$$= - \sum_{m=0}^{\infty} \sum_j C_{ij,nm}^{(1)} \rho_{j,m}^\mu + \text{non-linear terms}$$

Entries of „collision matrix“ (for diffusive moments)

Computation of transport coefficients (Example: diffusion coefficients)

On basis of DNMR theory: derivation from the Boltzmann equation

→ **upcoming publication!** (Fotakis, Molnár, Niemi, Denicol, Rischke, Greiner)

relativistic Boltzmann eq.

$$k_i^\mu \partial_\mu f_{i,\mathbf{k}} = C_i[f_i]$$

2nd-order (multi-component) hydro

$$\dot{\Pi}, \dot{V}_q^{\langle\mu\rangle}, \dot{\pi}^{\langle\mu\nu\rangle}$$

$$C_{i,n-1}^{\langle\mu\rangle} \equiv \int \frac{d^3\mathbf{k}_i}{(2\pi)^3 E_{i,\mathbf{k}}} E_{i,\mathbf{k}}^{n-1} k_i^{\langle\mu\rangle} C_i[f_i]$$

$$= - \sum_{m=0}^{\infty} \sum_j C_{ij,nm}^{(1)} \rho_{j,m}^\mu + \text{non-linear terms}$$

Entries of „collision matrix“ (for diffusive moments)

$$\kappa_{qq'} = \sum_{n=0}^{\infty} \sum_{i,j=1}^{N_{\text{species}}} (C^{-1})_{ji,0n}^{(1)} q_j \left(q_i' \mathcal{J}_{n+1,1}^{(i)} - \frac{n_{q'}}{\epsilon + P_0} \mathcal{J}_{n+2,1}^{(i)} \right)$$

Diffusion coefficient matrix! (equivalent to our PRL and PRD expression)

Greif, Fotakis et al., PRL 120, 242301 (2018)
Fotakis, Greif et al., PRD 101, 076007 (2020)
Fotakis, Soloveva et al, PRD 104, 034014 (2021)

- Hadronic system including lightest 19 species

$$\pi^\pm, \pi^0, K^\pm, K^0, \bar{K}^0, p, \bar{p}, n, \bar{n}, \Lambda^0, \bar{\Lambda}^0, \Sigma^0, \bar{\Sigma}^0, \Sigma^\pm, \bar{\Sigma}^\pm$$

- Assume classical statistics and non-interacting limit

$$P_0(T, \{\mu_q\}) \equiv \frac{1}{3} \sum_{i=1}^{N_{\text{species}}} \int \frac{dp^3}{(2\pi)^3 E_{i,p}} (E_{i,p}^2 - m_i^2) g_i \exp(-E_{i,p}/T + \sum_q q_i \alpha_q)$$

- Only assume **baryon number** and **strangeness**, **neglect electric charge**
- Tabulate state variables over energy density and net charge densities

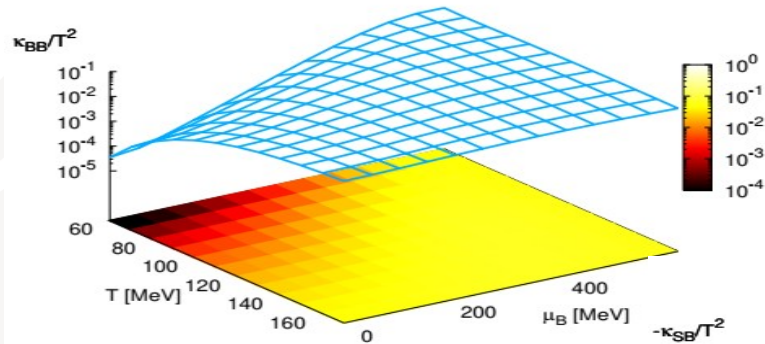
$$T \equiv T(\epsilon, \{n_q\}), \quad \mu_q \equiv \mu_q(\epsilon, \{n_q\}), \quad P_0 \equiv P_0(\epsilon, \{n_q\})$$

Diffusion coefficient matrix - details

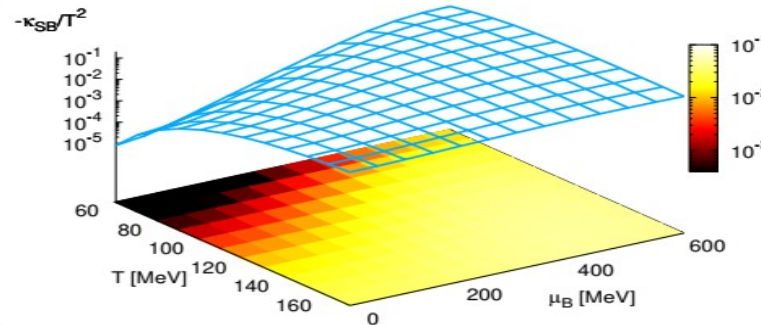
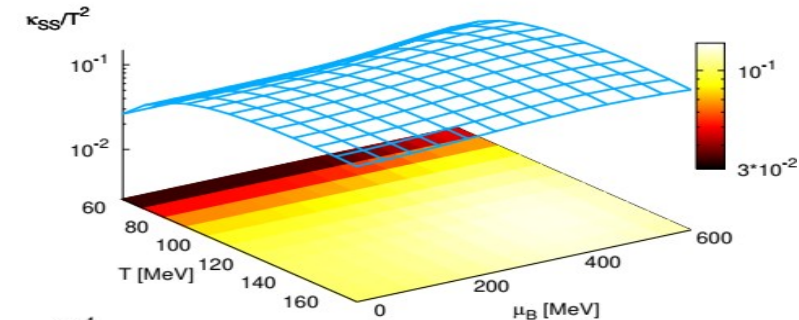
$$\begin{pmatrix} V_B^\mu \\ V_S^\mu \end{pmatrix} \sim \begin{pmatrix} \kappa_{BB} & \kappa_{BS} \\ \kappa_{SB} & \kappa_{SS} \end{pmatrix} \begin{pmatrix} \nabla^\mu \alpha_B \\ \nabla^\mu \alpha_S \end{pmatrix}$$

- Matrix is symmetric

L. Onsager, Phys. Rev. 37, 405 (1931) & Phys. Rev. 38, 2265 (1931)



- Elastic isotropic cross sections from PDG, SMASH, GiBUU, UrQMD



κ_{SB} is **negative** and has **similar magnitude** as κ_{BB}

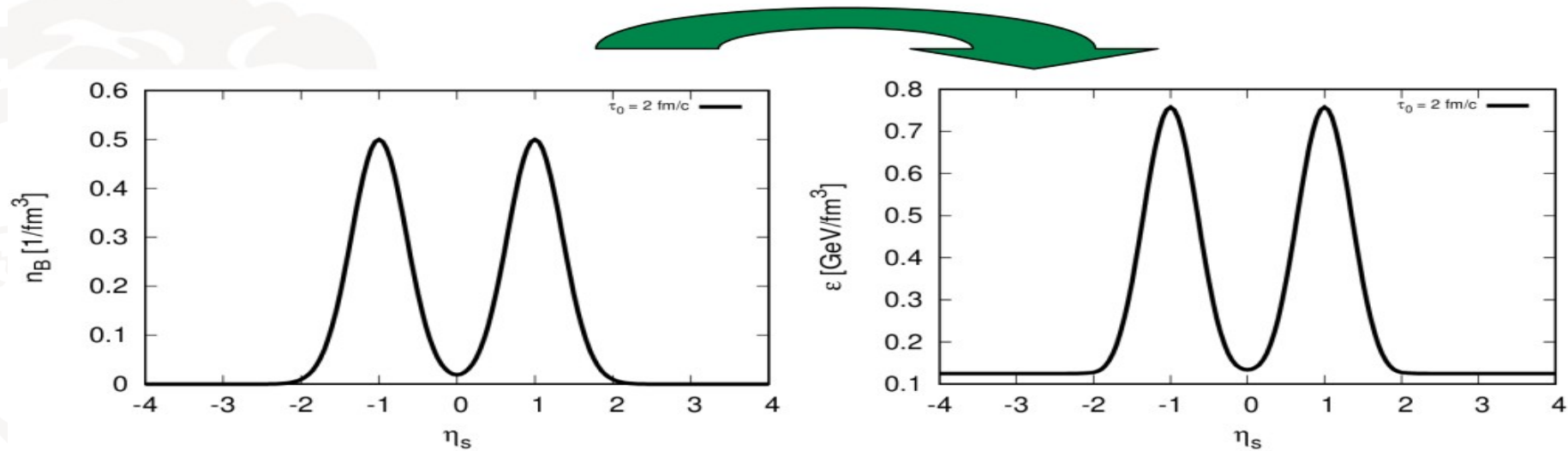
⇒ significant coupling?

- Tabulate coefficient matrix over T, μ_B, μ_S

- $\mu_Q = 0$

Initial conditions - details

- $\tau_0 = 2 \text{ fm}/c$
- Initially: no dissipation and only **Bjorken scaling flow**
- Temperature = 160 MeV
- Double-gaussian profile in net baryon number
- From **EoS**: get energy density



A simplistic case study

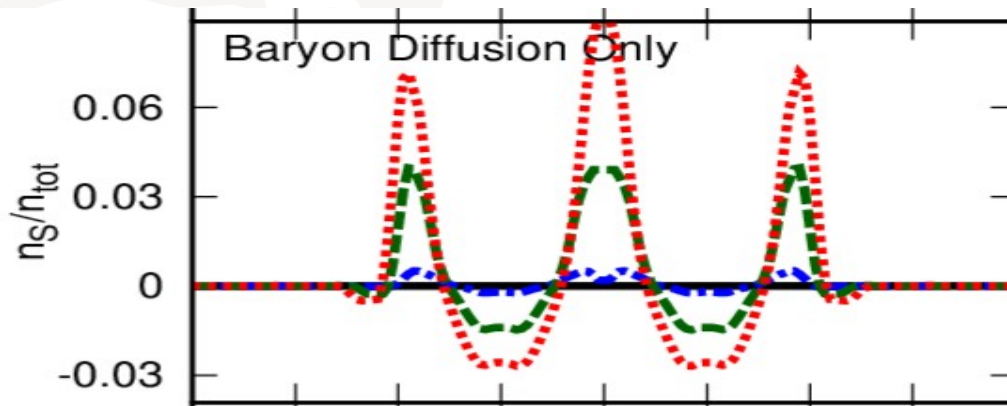
Greif, Fotakis et al., PRL 120, 242301 (2018)
Fotakis, Greif et al., PRD 101, 076007 (2020)

Let's **add shear** and also account for **second-order terms in the diffusion!**

Second-order terms in **ultrarelativistic limit**:

$$\tau_\pi \dot{\pi}^{\langle\mu\nu\rangle} + \pi^{\mu\nu} = 2\eta\sigma^{\mu\nu} - \frac{4}{3}\pi^{\mu\nu}\theta - 2\pi^{\langle\mu}_{\lambda}\omega^{\nu\rangle\lambda} - \frac{10}{7}\pi^{\lambda\langle\mu}\sigma^{\nu\rangle}_{\lambda}$$

With: $\frac{\eta}{s} = \frac{1}{4\pi}$



Second-order transport coefficients
not consistent with assumed (massive) system

→ generation of unphysical charge currents

Consistency is important in charge transport!