

Time evolution of conserved-charge fluctuations near the QCD critical point

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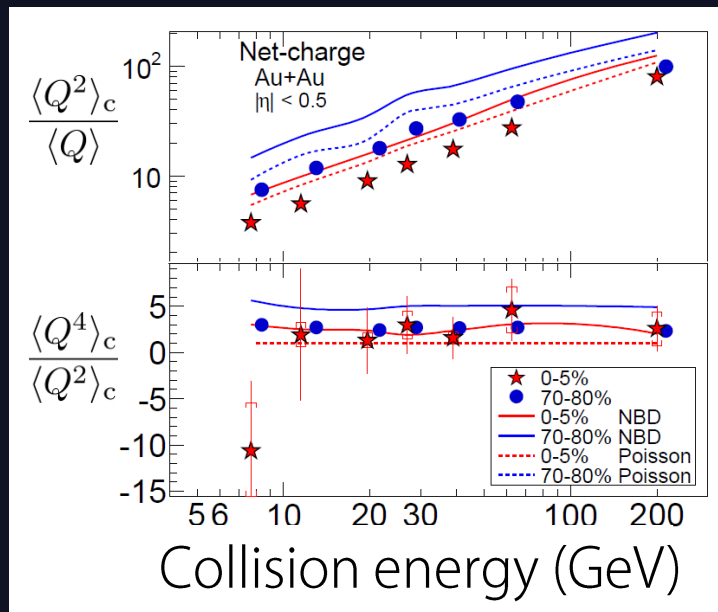
Fluctuation signature for QCD CP

$$\text{Near CEP, } \langle Q^2 \rangle_c^{\text{eq}} \sim \xi_{\text{eq}}^2 \rightarrow \infty$$

$$\langle Q^3 \text{ or } 4 \rangle_c^{\text{eq}} \sim \xi_{\text{eq}}^{4.5 \text{ or } 7} \rightarrow \infty + \text{sign change}$$

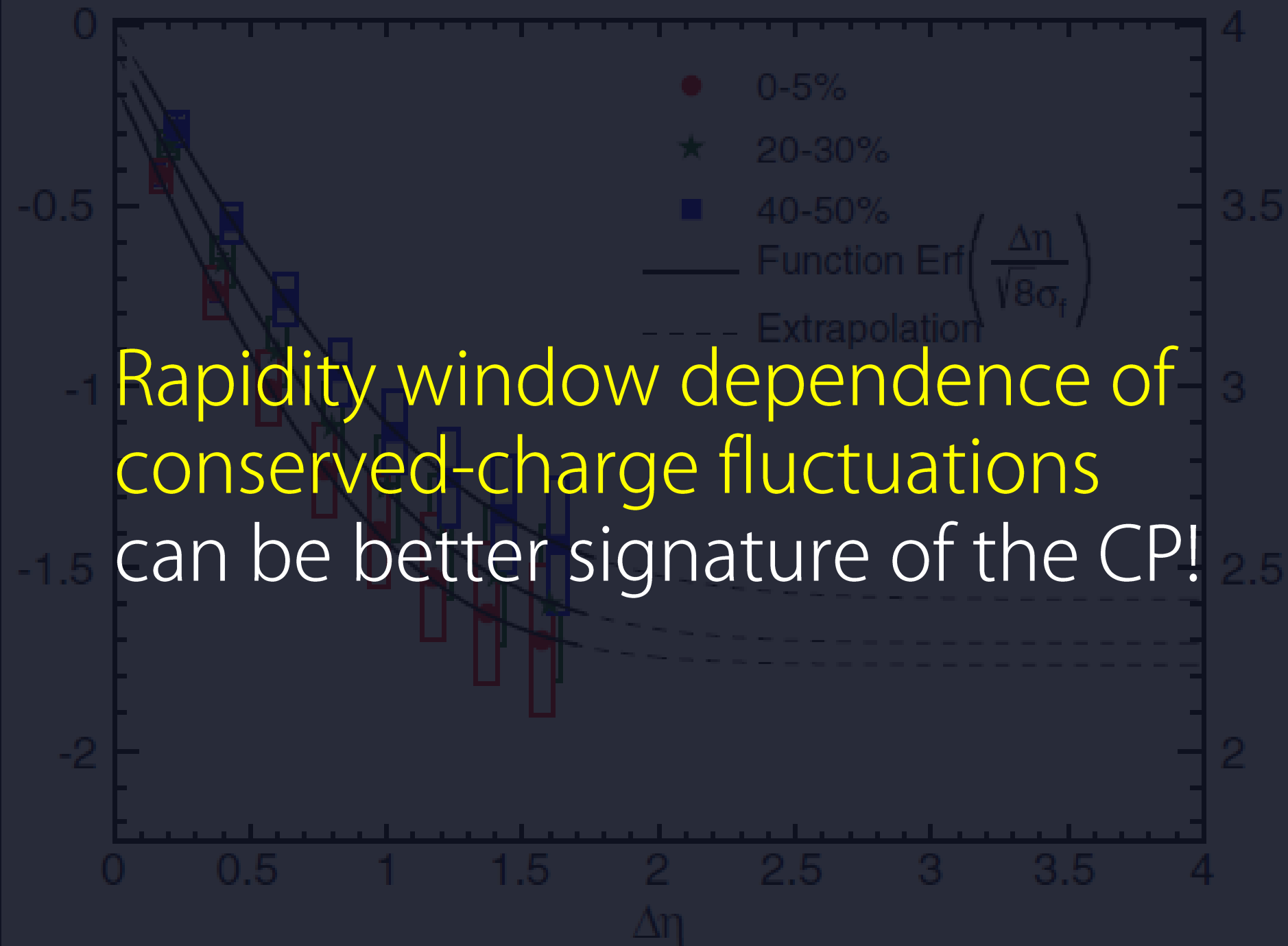
Asakawa et. al., PRL103,262301 (09), Stephanov, PRL102, 032301 (09)

singular behavior = signature!!



STAR, PRL113, 092301 (2014)

⚠ But due to critical slowing down & final state interactions, catching such signals is not so easy...



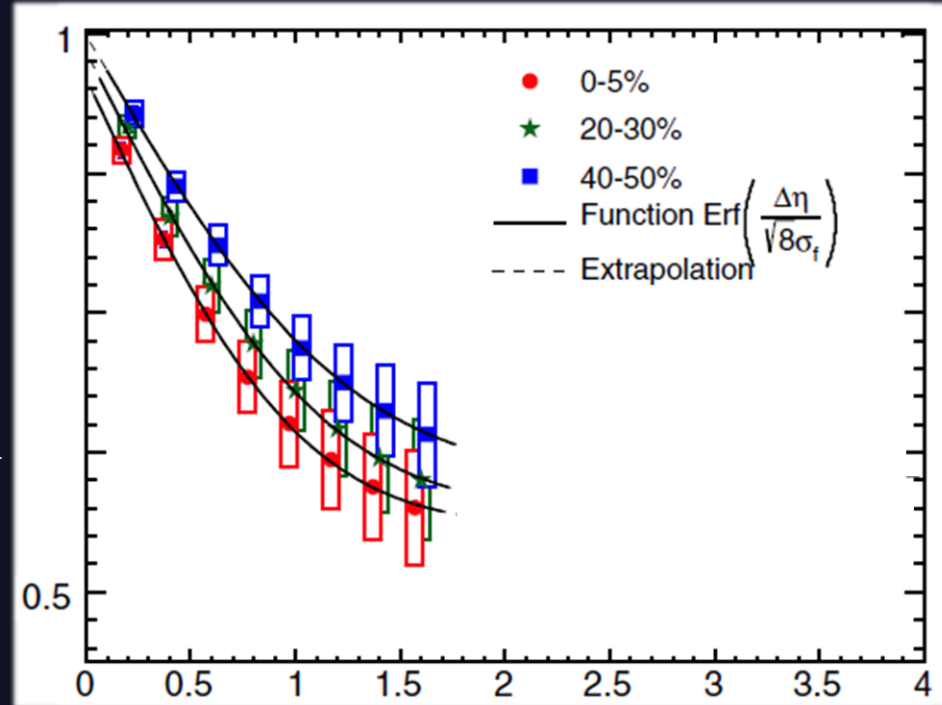
$\Delta\eta$ dependence of charge fluctuation

ALICE, PRL110, 152301 (2013)

Variance of net electric charge

$$\frac{\langle N_Q^{(\text{net})2} \rangle_c}{\langle N_Q^+ + N_Q^- \rangle_c}$$

Asakawa, Heinz, Muller (2000)
Jeon, Koch (2000)



$\Delta\eta$: rapidity window

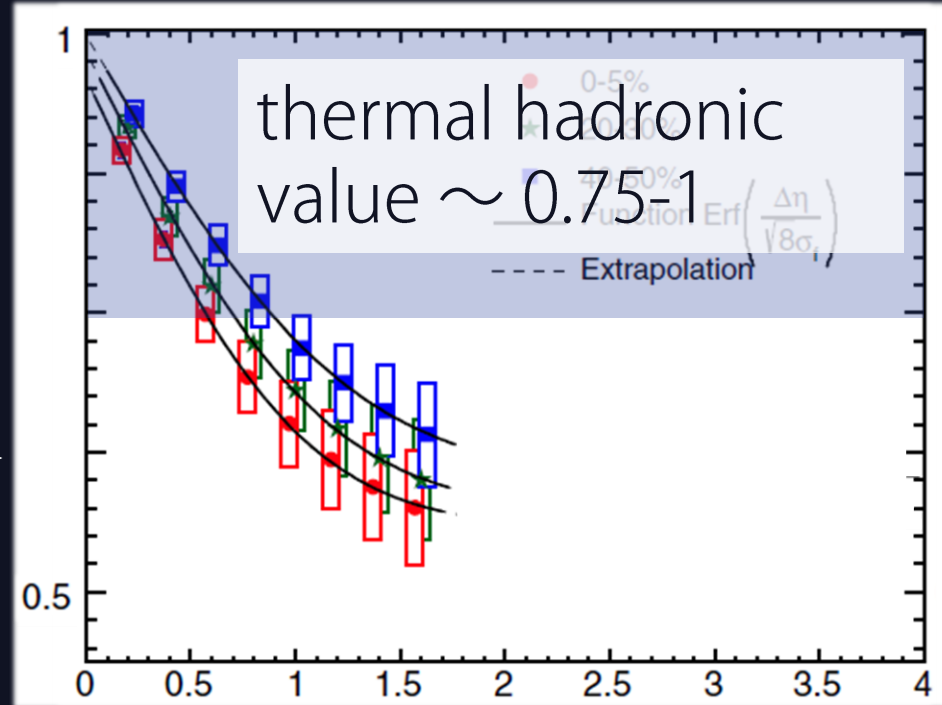
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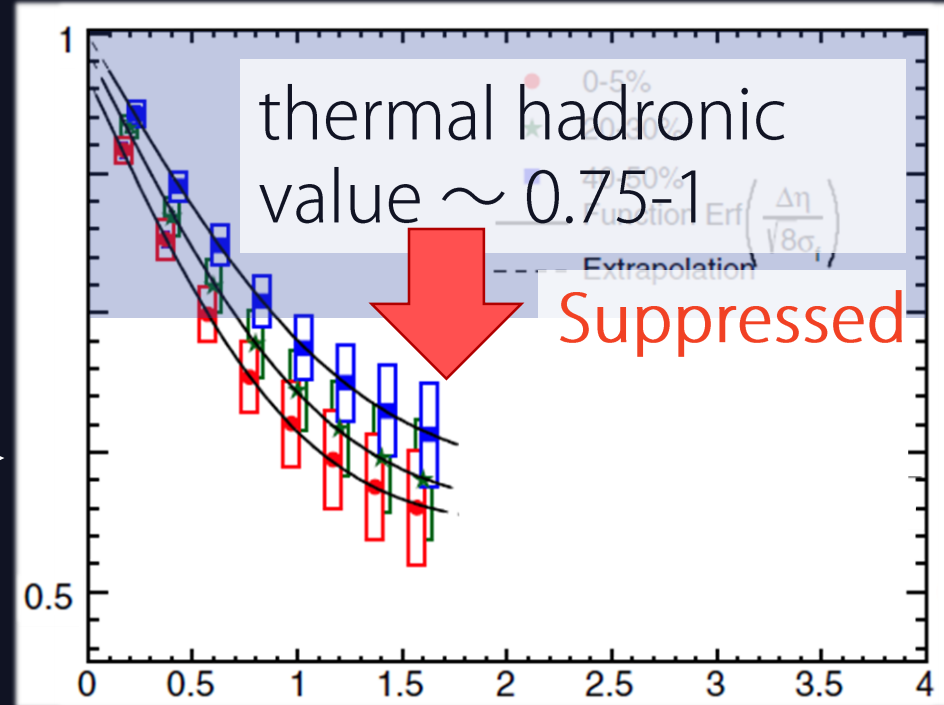
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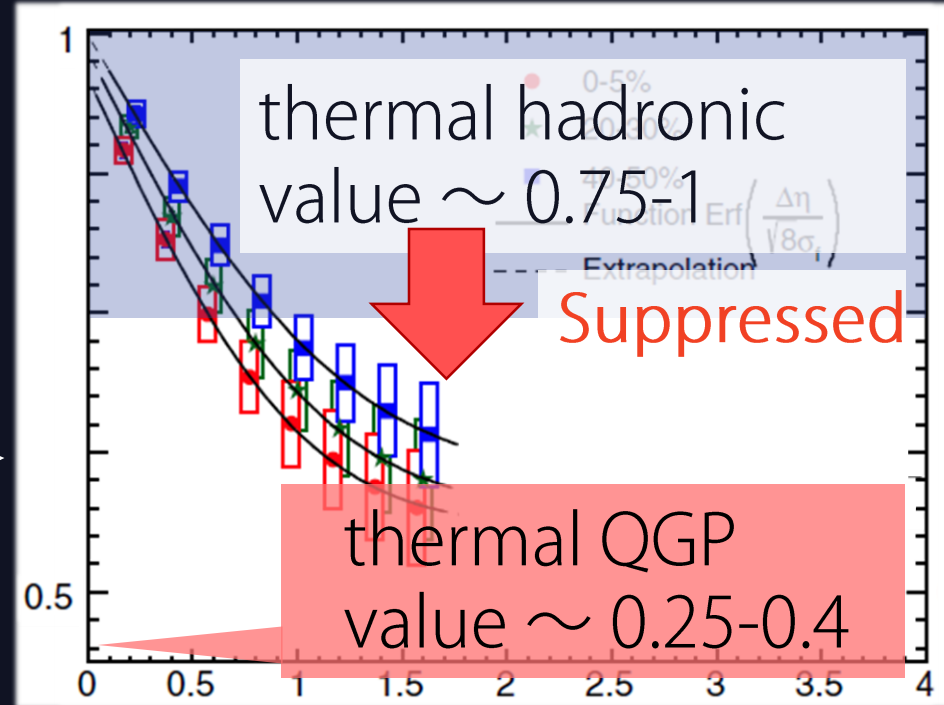
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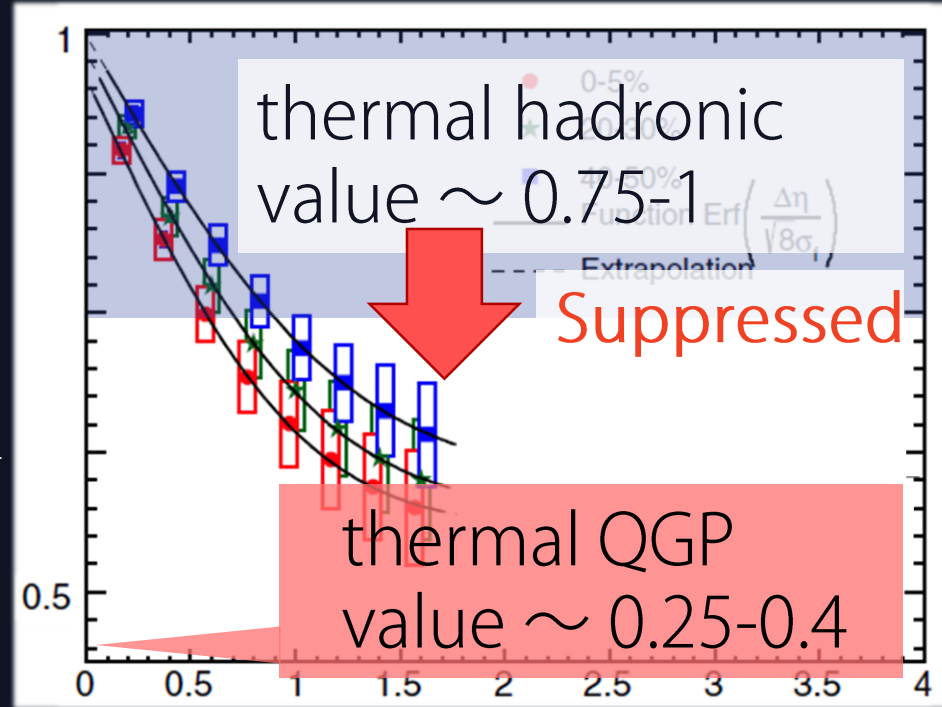
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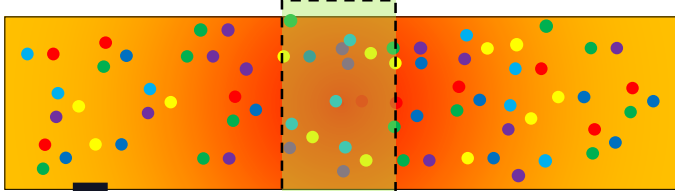


$\Delta\eta$: rapidity window

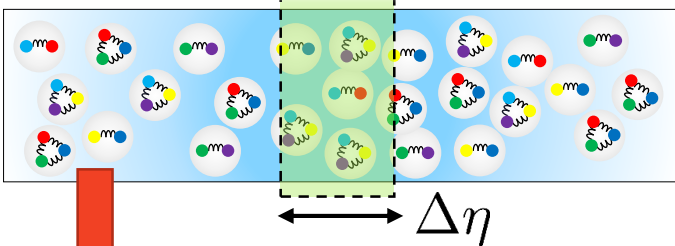
Fluctuations for larger $\Delta\eta$ are not thermal at kinetic freeze-out (?)

Time evolution of fluctuations

① QGP (thermal)

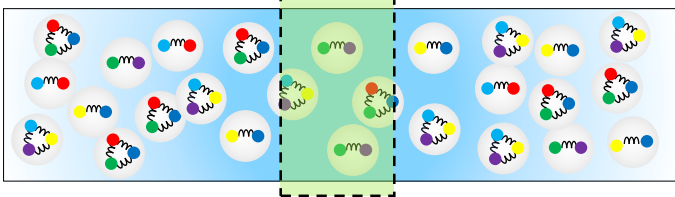


② Hadrons (non-thermal)



③ relaxation

④ Hadrons (thermal)

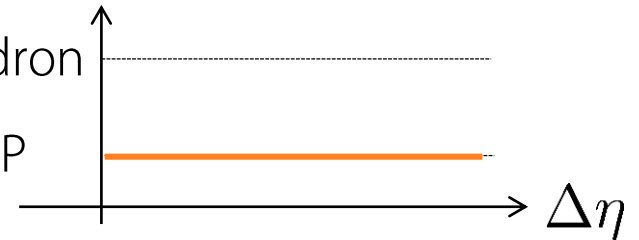


①, ②

2nd order fluctuation

thermal hadron

thermal QGP

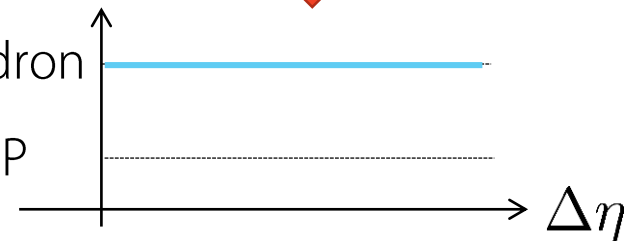


③ relaxation

④

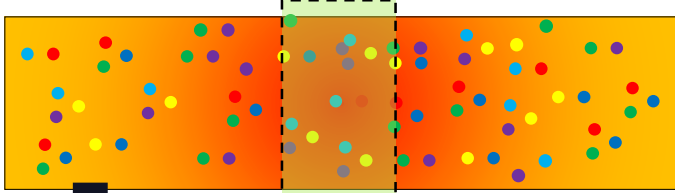
thermal hadron

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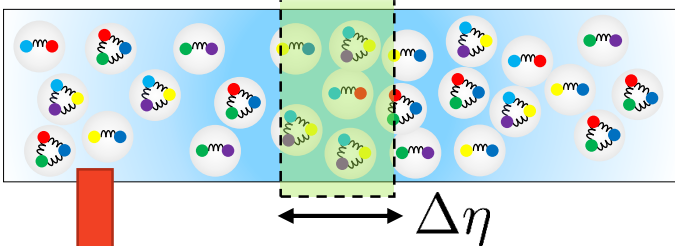


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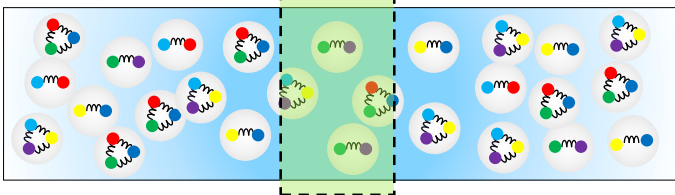


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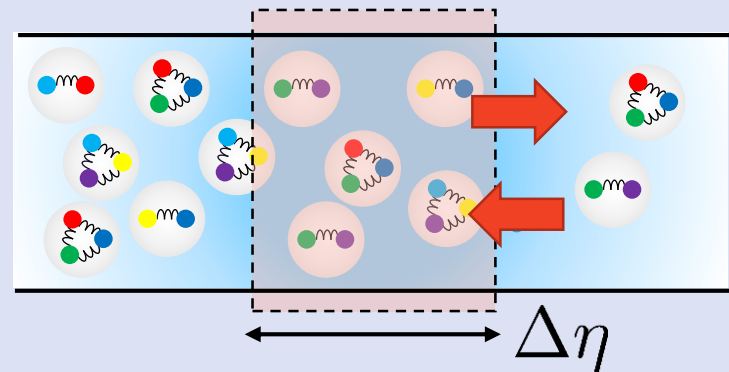
③ relaxation

④ Hadrons (thermal)



Relaxation can **only** proceed by **charge diffusion**!!

Shuryak, Stephanov (2001)

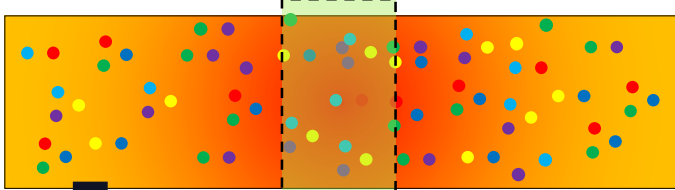


$\Delta\eta \rightarrow$ larger

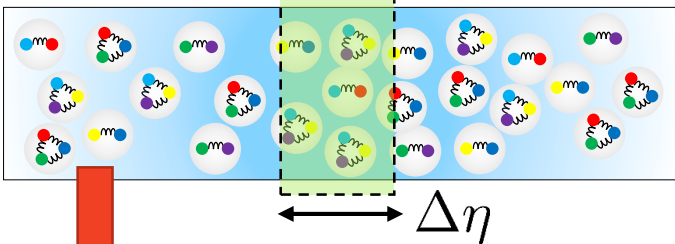
\rightarrow relaxation time \rightarrow longer
(more QGP value)

Time evolution of fluctuations

① QGP (thermal)

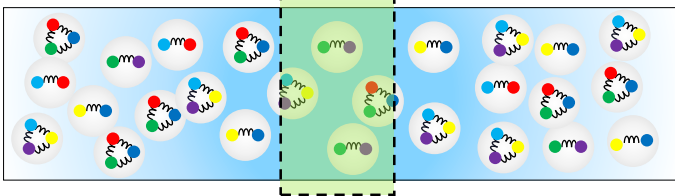


② Hadrons (non-thermal)



③ relaxation

④ Hadrons (thermal)

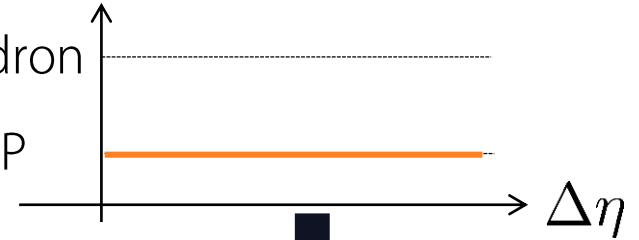


①, ②

thermal hadron

thermal QGP

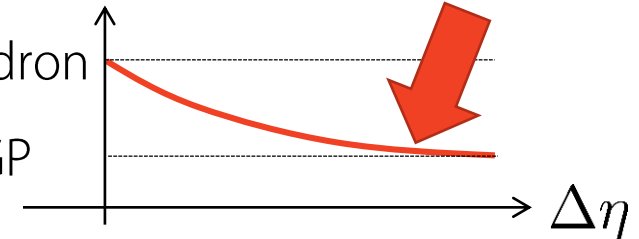
2nd order fluctuation



③

thermal hadron

thermal QGP



④

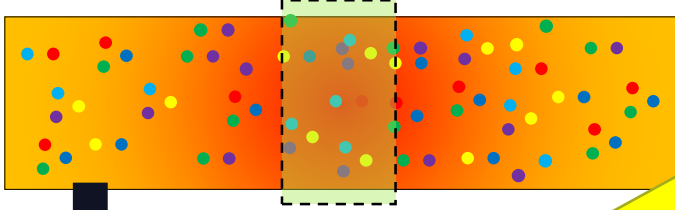
thermal hadron

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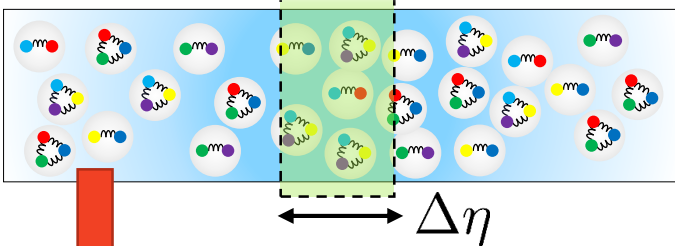


Time evolution of fluctuations

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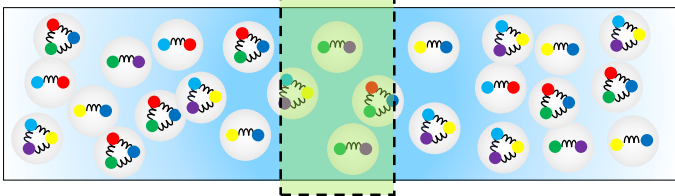


② Hadrons (non-thermal)



③ relaxation

④ Hadrons (thermal)

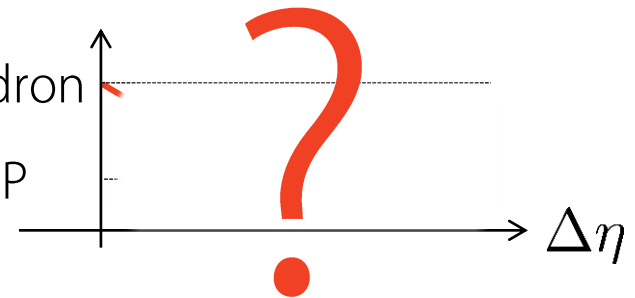


If system passes near the CP,

③

thermal hadron

thermal QGP



criticality should appear for larger $\Delta\eta$!!

Time evolution of fluc. near the CP

Q. How to describe time evolution of CC fluc. near the CP?

■ Previous study on dynamical critical fluc.

Berdnikov, Rajagopal (2000), Nonaka, Asakawa (2005), Mukherjee et. al. (2015)

☹do not take into account the conservation effect

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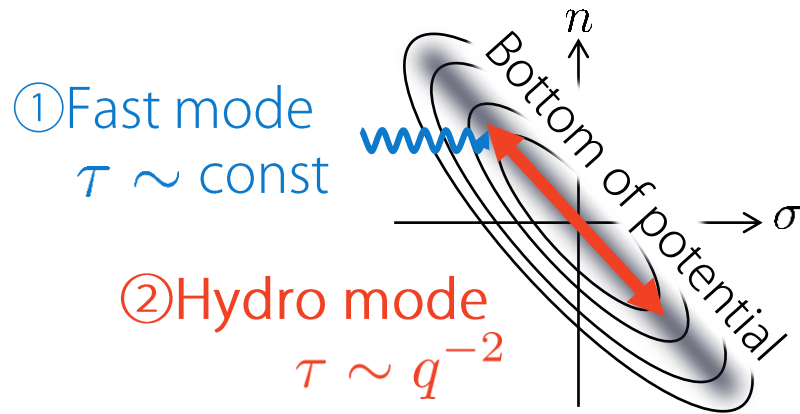
At QCD CP, n and σ are coupled
→ critical fluctuation is hydrodynamic mode

Fujii, Ohtani, PRD70, 014016 (2004), Son, Stephanov, PRD70, 056001(2004)

→ We must take into account the conservation effect in order to discuss dynamical CC fluc. near CP.

How to describe?

- Near the QCD CP \rightarrow Stochastic diffusion eq. (SDE)



Fujii, Ohtani, PRD70, 014016 (2004)
Son, Stephanov, PRD70, 056001(2004)

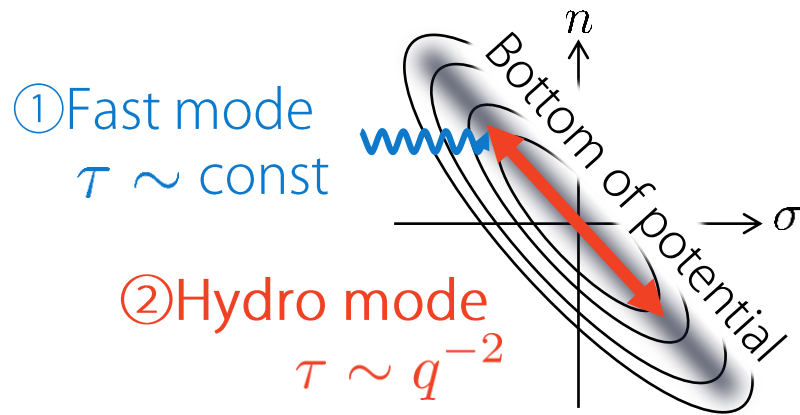
At longer time scale,
 σ can be forgotten.

\downarrow n relaxes solely
with diffusive time scale.

$$\frac{\partial}{\partial \tau} n = D \frac{\partial^2}{\partial \eta^2} n + \frac{\partial}{\partial \eta} \xi$$

How to describe?

- Near the QCD CP \rightarrow Stochastic diffusion eq. (SDE)



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n can evolve solely!

$$\frac{\partial}{\partial \tau} n = D \frac{\partial^2}{\partial \eta^2} n + \frac{\partial}{\partial \eta} \xi$$

- Non-critical region \rightarrow also SDE

Shuryak, Stephanov (2001), Kitazawa, Asakawa, Ono (2013), MS, Asakawa, Kitazawa (2014)

We can use SDE uniformly at all stage !

Our Strategy

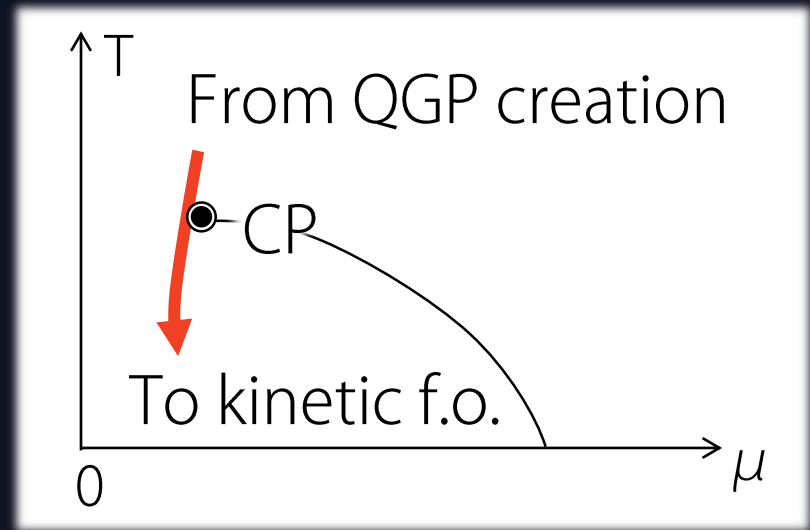
Stochastic diffusion eq.

$$\frac{\partial}{\partial \tau} n = D \frac{\partial^2}{\partial \eta^2} n + \frac{\partial}{\partial \eta} \xi$$



Solve(assumption: white noise)

$\Delta \eta$ dependence of conserved charge fluc.



Our Strategy

Stochastic diffusion eq. $\frac{\partial}{\partial \tau} n = D \frac{\partial^2}{\partial \eta^2} n + \frac{\partial}{\partial \eta} \xi$

■ Near the QCD CP

➤ Dynamical universality class → model H

susceptibility $\chi_B \sim \xi^2$, Diffusion coefficient $D \sim \xi^{-1}$

Hohenberg, Halperin, Rev. Mod. Phys. 49, 435 (1977)

➤ 3D Ising mapping

■ Non-critical region

➤ $\chi_B^{\text{hadron}} = 1, \chi_B^{\text{QGP}} = 0, D = \text{const}$

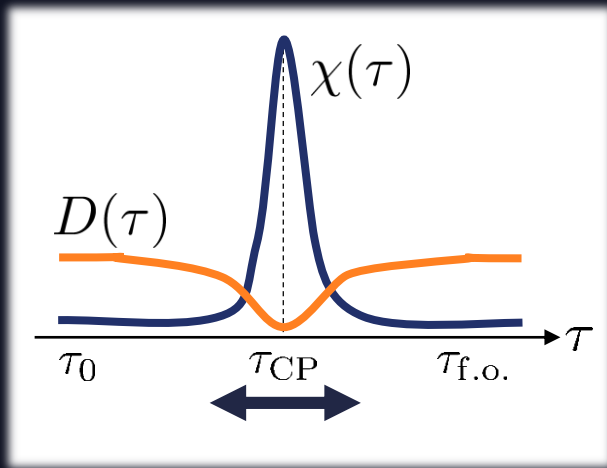
Solve(assumption: white noise)

$\Delta \eta$ dependence of conserved charge fluc.

Discussion of obtained formula

$$\langle n(q_1, \tau) n(q_2, \tau) \rangle_c = 4\pi q_1^2 \delta(q_1 + q_2) \int_{\tau_0}^{\tau} d\tau' \chi_B(\tau') D(\tau') \exp \left[-2q_1^2 \int_{\tau'}^{\tau} d\tau'' D(\tau'') \right]$$

World of τ



$D \rightarrow 0 =$ critical slowing down

Discussion of obtained formula

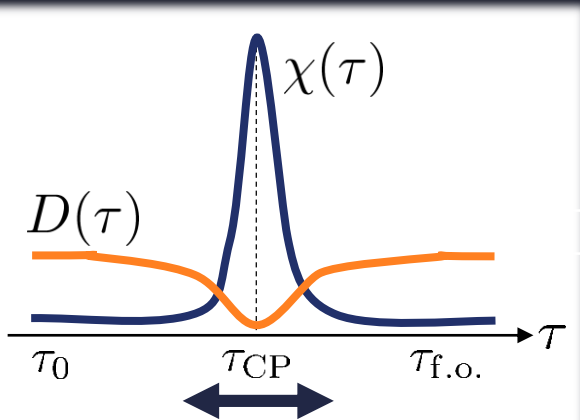
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Transform as

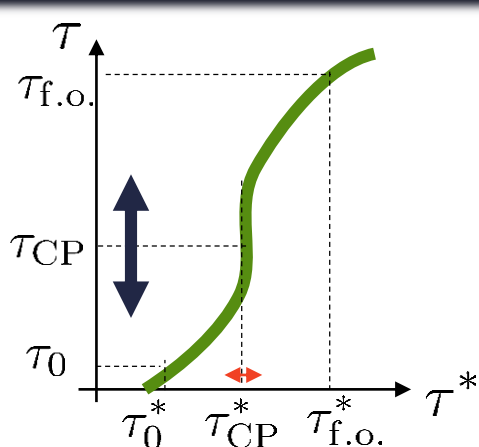
$$\int_{\tau_0}^{\tau} d\tau' \chi_B(\tau') D(\tau') \exp \left[-2q_1^2 \int_{\tau'}^{\tau} d\tau'' D(\tau'') \right] \rightarrow \int_{\tau_0^*}^{\tau^*} d\tau'^* \chi_B(\tau'(\tau'^*)) \exp \left[-2q_1^2 (\tau^* - \tau'^*) \right]$$

$d\tau^* = d\tau D(\tau)$

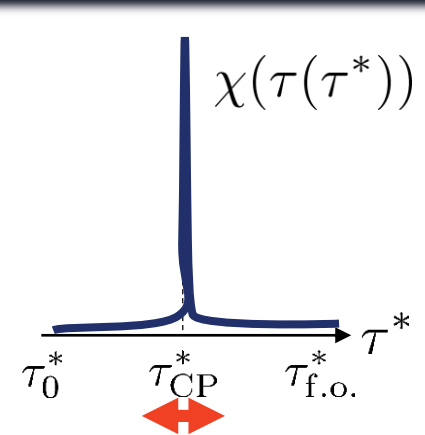
World of τ



Transform $\tau \rightarrow \tau^*$



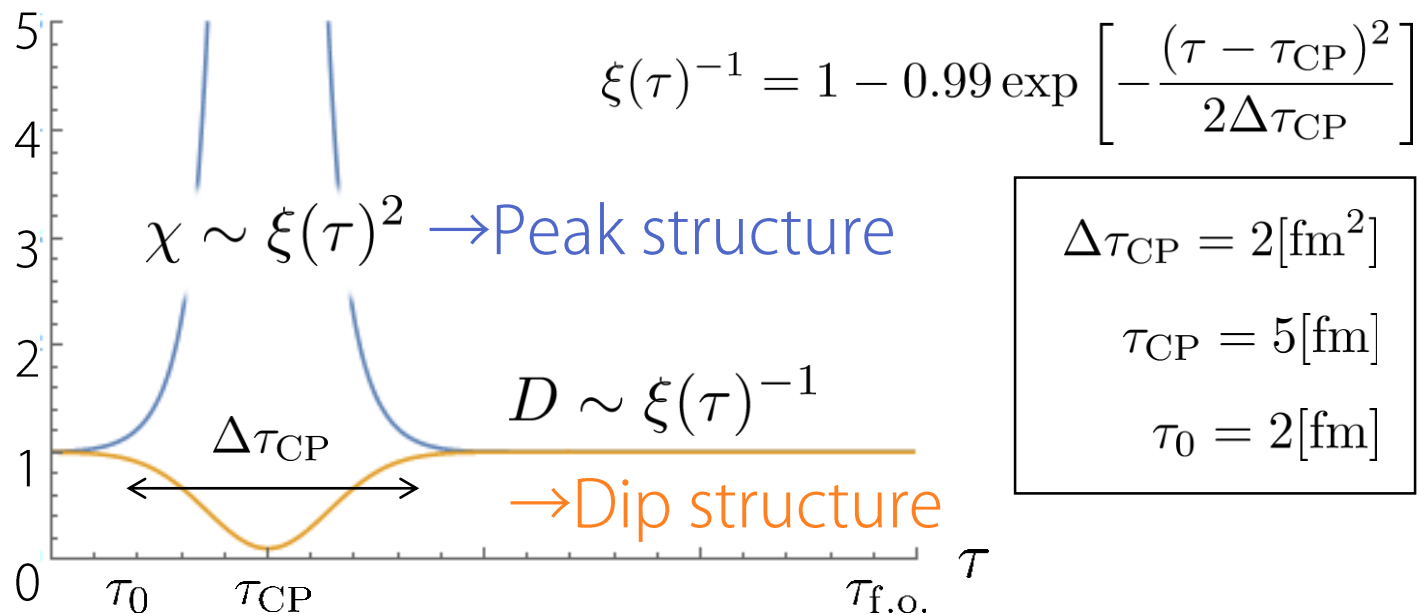
World of τ^*



Time scale is compressed near CP = critical slowing down

Our Strategy (Simple case)

Stochastic diffusion eq. $\frac{\partial}{\partial \tau} n = D \frac{\partial^2}{\partial \eta^2} n + \frac{\partial}{\partial \eta} \xi$

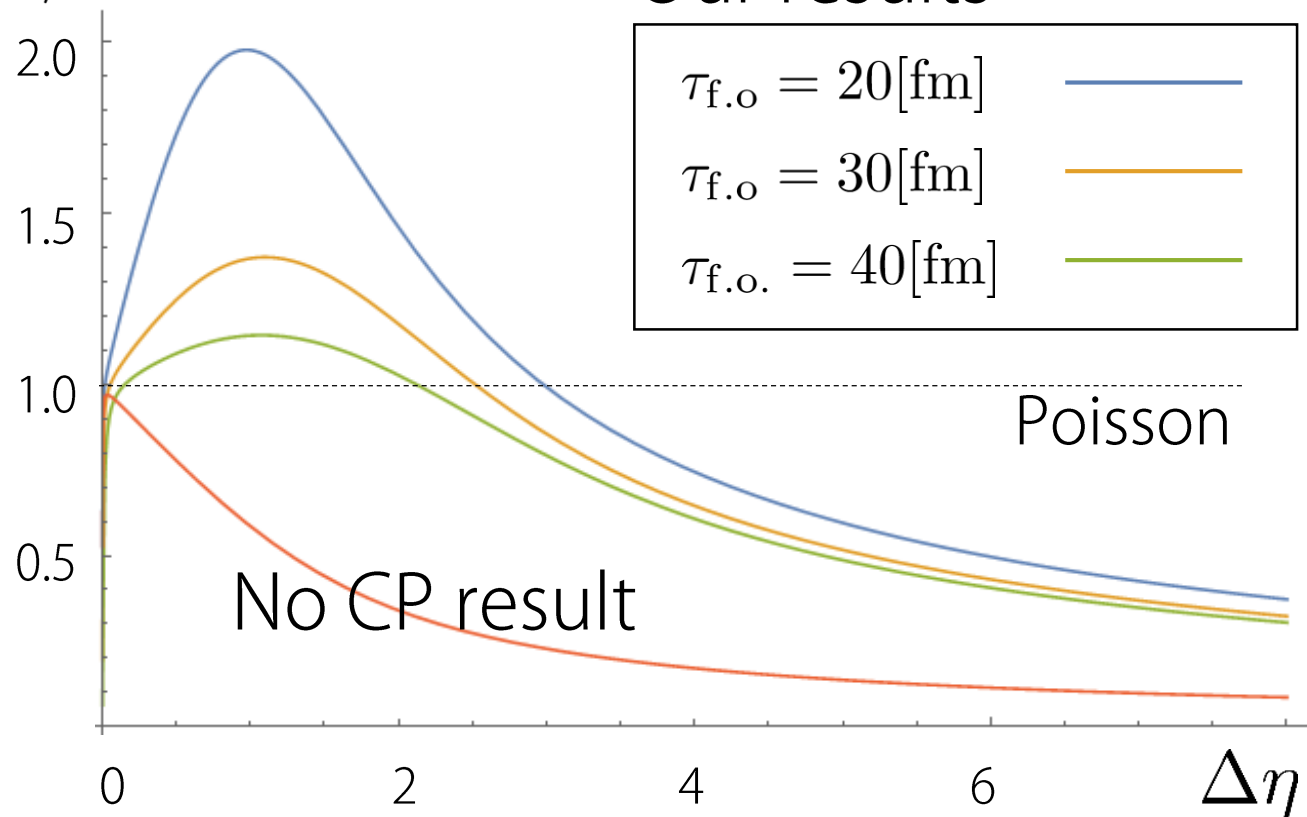


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$\Delta\eta$ dependence of conserved charge fluc.

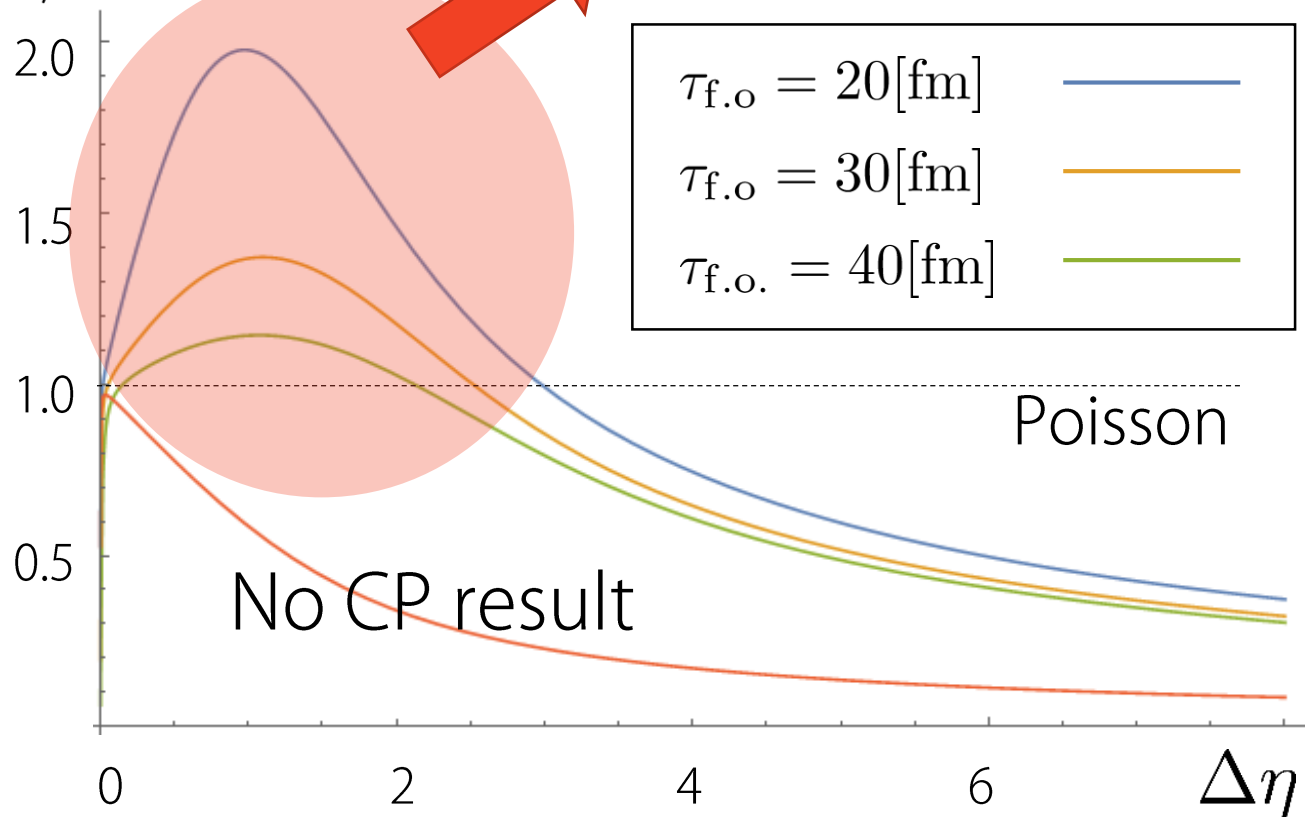
Result on $\Delta\eta$ dep. of CC fluctuations

$$\frac{\langle N_B^{(\text{net})2} \rangle_c}{\langle N_B^{(\text{tot})} \rangle}$$



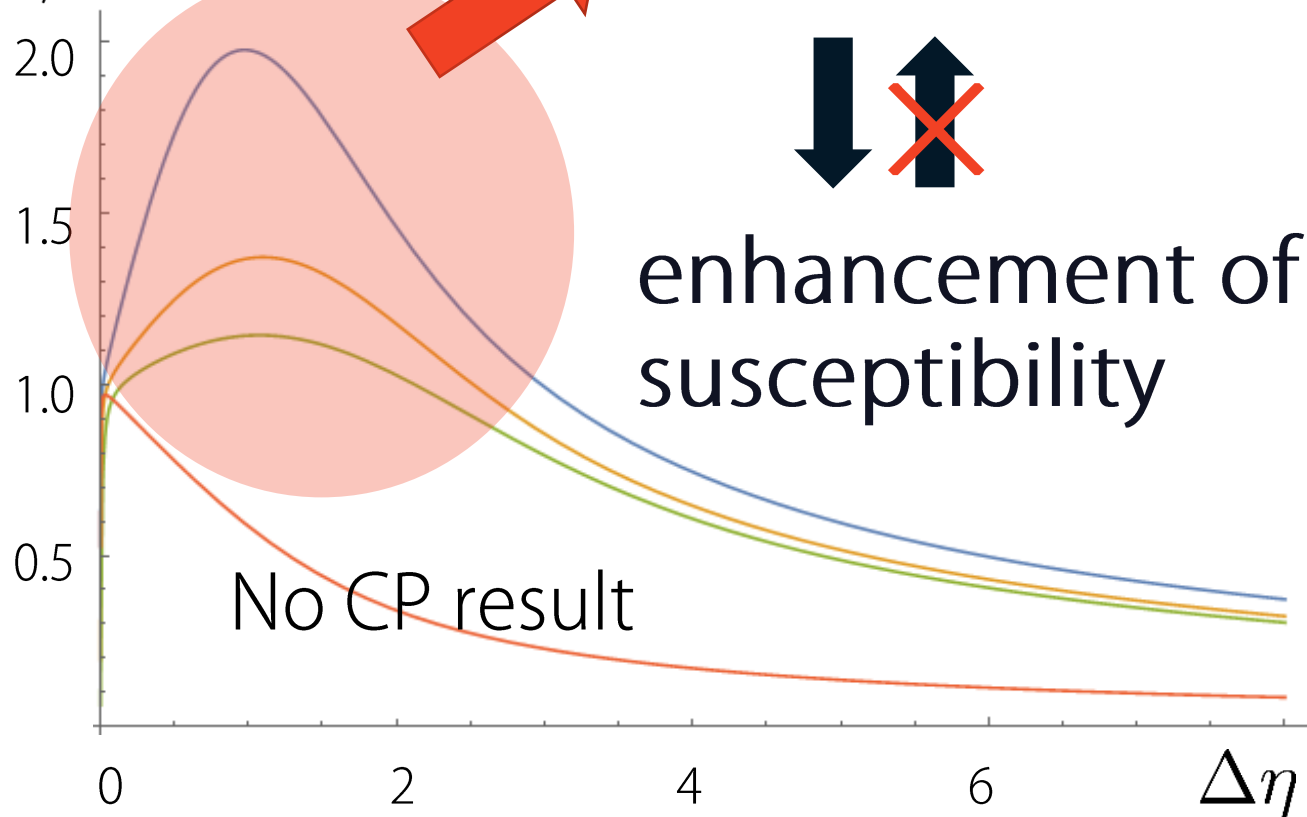
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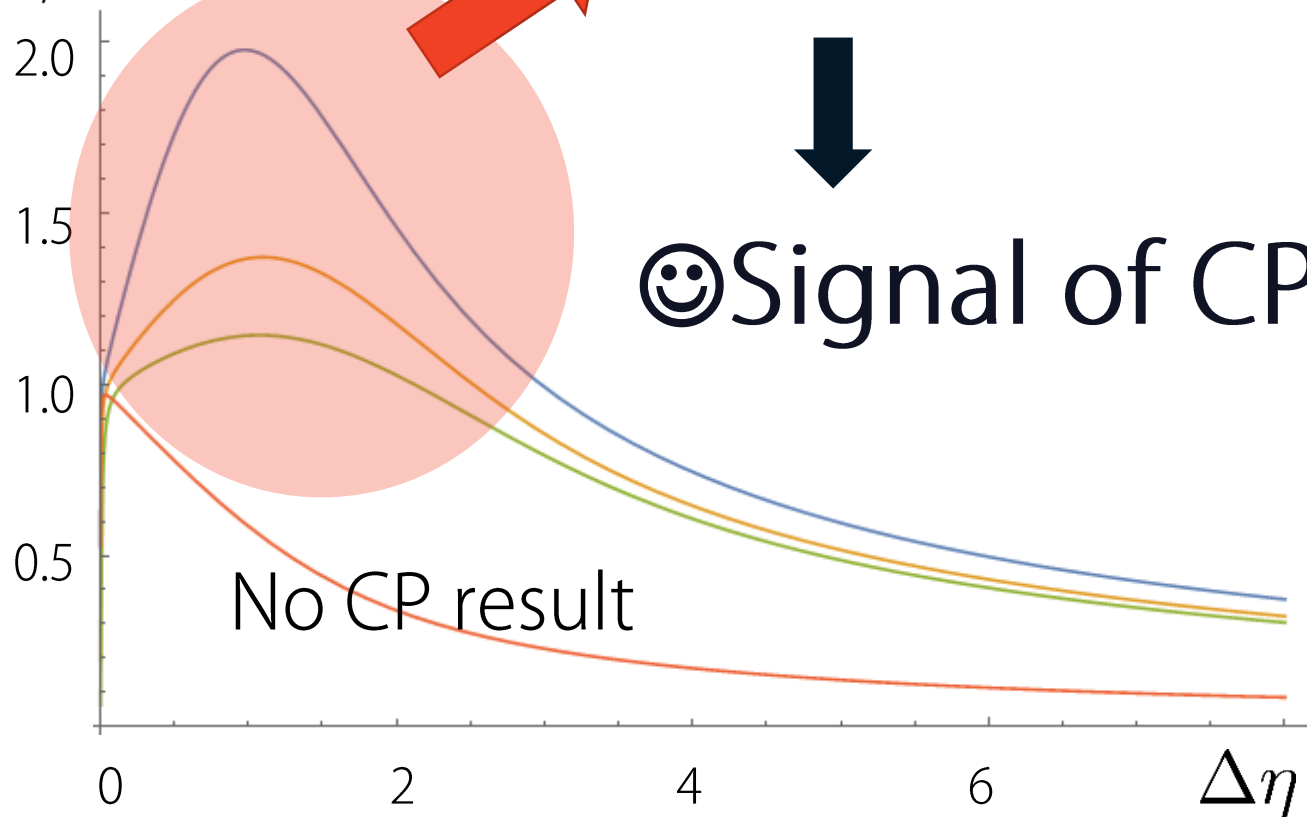
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Summary

We study the time evolution of CC fluctuations in the case that system pass near the CP.

- We use SDE from QGP creation to kinetic freeze out.
- If system passes near the CP,
 $\Delta\eta$ dep. of fluc. may show non-monotonic behavior.
→If the non-monotonic behavior can be observed in experiment, it is a signal of CP !!

$\Delta\eta$ dependence of CC fluctuations
can be better signature of the CP!