

# Exotic Phenomena in Nuclei



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# Shells, Clusters and Condensates



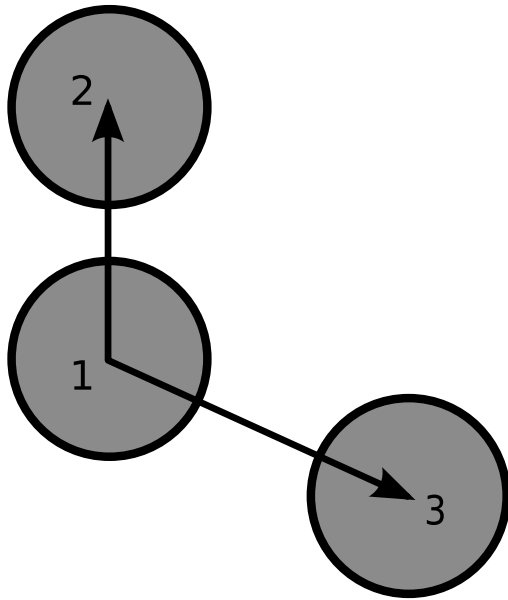
## Nucleus as a Mesoscopic System

- shell structure
- cluster structure
- $\alpha$ -cluster condensates close to the  $\alpha$ -threshold

## Hoyle State in $^{12}\text{C}$

- microscopic  $\alpha$ -cluster model
- microscopic  $\alpha$ -cluster model with condensate wave functions
- Fermionic Molecular Dynamics model that can describe shell model like configurations as well as  $\alpha$ -cluster configurations

# Microscopic $\alpha$ -Cluster Model



$$R_{12} = (2, 4, \dots, 10) \text{ fm}$$

$$R_{13} = (2, 4, \dots, 10) \text{ fm}$$

$$\cos(\vartheta) = (1.0, 0.8, \dots, -1.0)$$

altogether 165  
configurations

## Basis States

- describe  $^{12}\text{C}$  as a system of 3  $^4\text{He}$  nuclei
- Generator Coordinate Method:  
use Slater determinants to describe many-body state

$$|\Psi_{3\alpha}(\mathbf{R}_1, \mathbf{R}_2, \mathbf{R}_3)\rangle = \mathcal{A}\left\{|\psi_\alpha(\mathbf{R}_1)\rangle \otimes |\psi_\alpha(\mathbf{R}_2)\rangle \otimes |\psi_\alpha(\mathbf{R}_3)\rangle\right\}$$

## Volkov Interaction

- simple central interaction
  - parameters adjusted to reproduce  $\alpha$  binding energy and radius, exchange parameter adjusted to reproduce  $^{12}\text{C}$  ground state energy and radius
- ✗ only reasonable for  $^4\text{He}$ ,  $^8\text{Be}$  and  $^{12}\text{C}$  nuclei

# Condensate Wave Functions

## Basis States

- Funaki *et al* proposed a description based on condensate wave functions
- all  $\alpha$ -particles occupy the same spatially deformed center-of-mass orbit

$$|\Psi_{\text{BEC}}(\beta_x, \beta_y, \beta_z)\rangle = \int d^3R_1 d^3R_2 d^3R_3 \exp\left\{-\sum_{i=1}^3 \left(\frac{R_{ix}^2}{\beta_x^2} + \frac{R_{iy}^2}{\beta_y^2} + \frac{R_{iz}^2}{\beta_z^2}\right)\right\} |\Psi_{3\alpha}(\mathbf{R}_1, \mathbf{R}_2, \mathbf{R}_3)\rangle$$

- Generator Coordinate Method:  
consider the deformation parameters  $(\beta_x, \beta_y, \beta_z)$  as generator coordinates
- wave function has to be antisymmetrized, densities are not so low that antisymmetrization effects can be neglected
- detailed analysis by Matsumara *et al* finds 30%  $S$ -wave probability for the ground state and 70%  $S$ -wave probability for the Hoyle state

## Volkov Interaction

- use same interaction as for the microscopic  $\alpha$ -cluster model

# Fermionic Molecular Dynamics

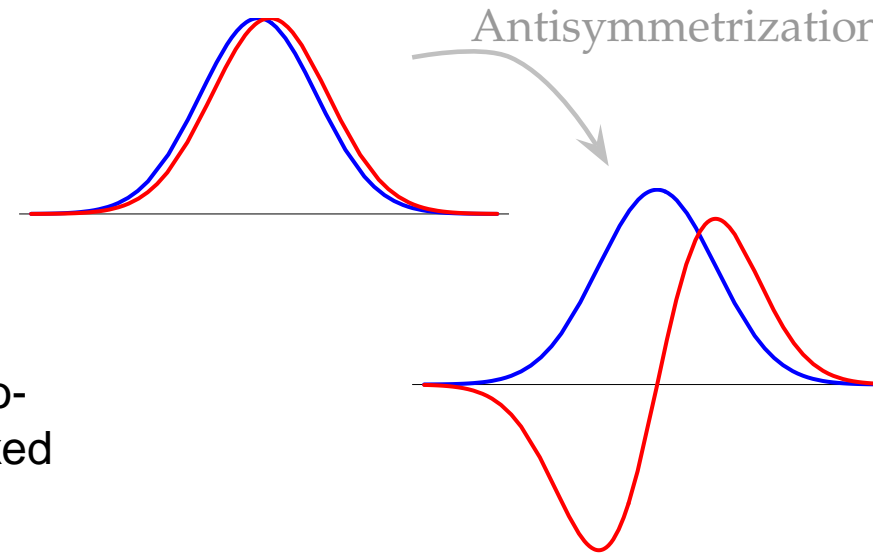
## Wave Functions

- describe many-body system by Slater determinants

$$|Q\rangle = \mathcal{A}\left(|q_1\rangle \otimes \cdots \otimes |q_A\rangle\right)$$

- single-particle states are Gaussian wave-packets in phase-space (complex parameter  $\mathbf{b}_i$  encodes mean position and mean momentum), spin is free, isospin is fixed

$$\langle \mathbf{x} | q \rangle = \exp\left\{-\frac{(\mathbf{x} - \mathbf{b})^2}{2a}\right\} \otimes |\chi^\uparrow, \chi^\downarrow\rangle \otimes |\xi\rangle$$



## Interaction

- derive effective interaction from Argonne V18 with the **Unitary Correlation Operator Method** which introduces short-range central and tensor correlations
- augment with phenomenological two-body correction that contains momentum-dependent central and spin-orbit terms fitted to doubly-magic nuclei to mimic effects of three-body correlations and genuine three-body forces

# Fermionic Molecular Dynamics

## Basis States

- 20 FMD states obtained in **Variation after Angular Momentum and Parity Projection** on  $0^+$  and  $2^+$  with constraints on the radius
- 42 FMD states obtained in **Variation after Parity Projection** with constraints on radius and quadrupole deformation
- 165 explicit  $\alpha$ -cluster configurations

$$P^{\mathbf{P}} = \frac{1}{(2\pi)^3} \int d^3X \exp\{-i(\mathbf{P} - \mathbf{P}) \cdot \mathbf{X}\}$$

$$P_{MK}^J = \frac{2J+1}{8\pi^2} \int d^3\Omega D_{MK}^J(\Omega) R(\Omega)$$

## Solve the Many-Body Schrödinger Equation

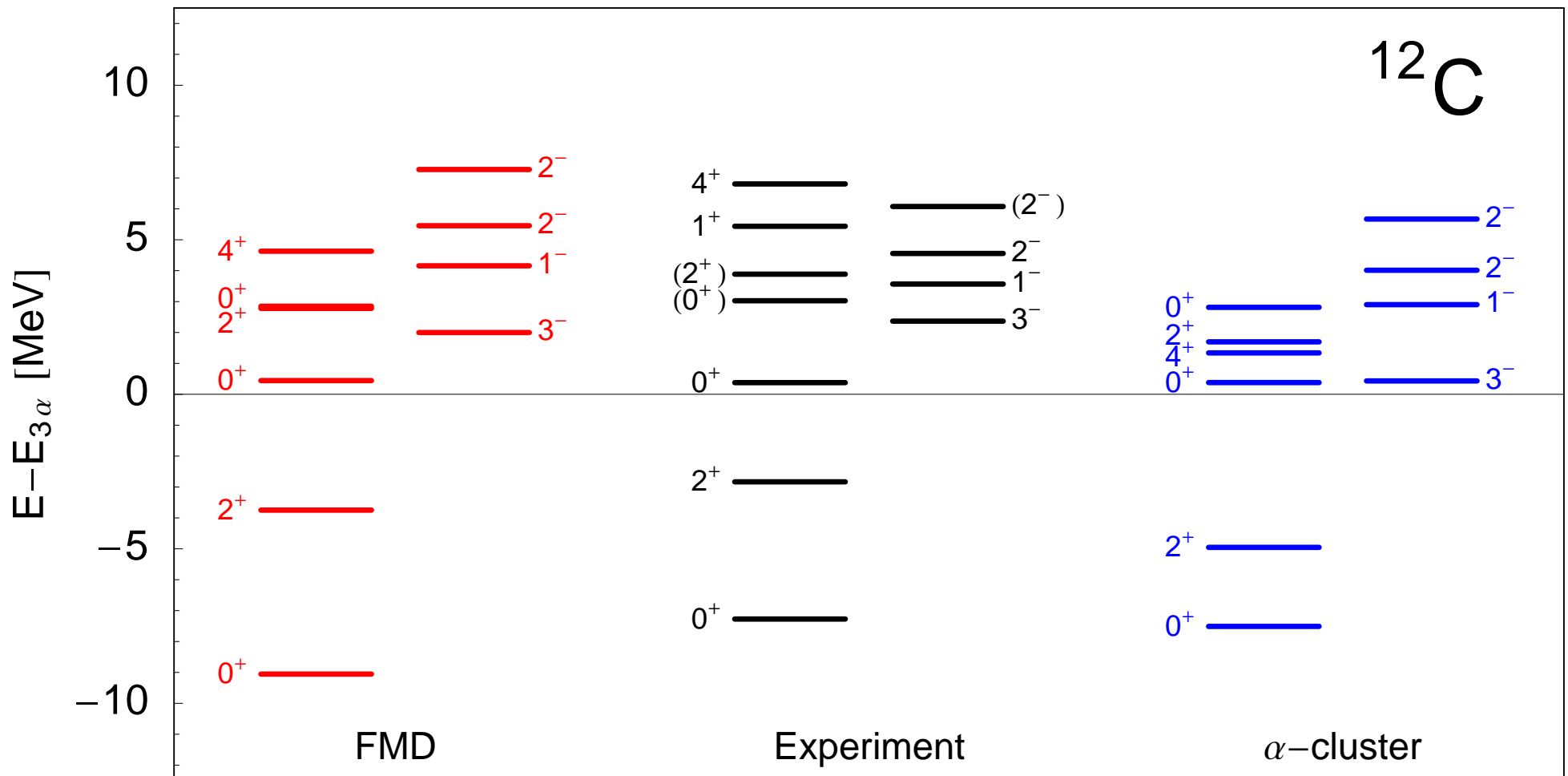
- Generalized Eigenvalue problem
- matrix elements are projected on parity, angular momentum and linear momentum

$$\sum_{K'b} \langle Q^{(a)} | HP_{KK'}^{J^\pi} P^{\mathbf{P}=0} | Q^{(b)} \rangle \cdot c_{K'b}^{(i)} = E^{J^\pi(i)} \sum_{K'b} \langle Q^{(a)} | P_{KK'}^{J^\pi} P^{\mathbf{P}=0} | Q^{(b)} \rangle \cdot c_{K'b}^{(i)}$$

# Spectra

FMD and  $\alpha$ -cluster reproduce Hoyle state close to threshold

differences for  $0_3^+$  and  $2_2^+$  states



# Comparison

	Exp <sup>1</sup>	Exp <sup>2</sup>	Exp <sup>3</sup>	FMD	$\alpha$ -cluster	'BEC'
$E(0_1^+)$	-92.16			-92.64	-89.56	-89.52
$E^*(2_1^+)$	4.44			5.31	2.56	2.81
$E(3\alpha)$	-84.89			-83.59	-82.05	-82.05
<b><math>E(0_2^+) - E(3\alpha)</math></b>	<b>0.38</b>			<b>0.43</b>	<b>0.38</b>	<b>0.26</b>
$E(0_3^+) - E(3\alpha)$	(3.0)	2.7(3)	3.96(5)	2.84	2.81	
$E(2_2^+) - E(3\alpha)$	(3.89)	2.6(3)	6.63(3)	2.77	1.70	
$r_{\text{charge}}(0_1^+)$	2.47(2)			2.53	2.54	
$r(0_1^+)$				2.39	2.40	2.40
<b><math>r(0_2^+)</math></b>				<b>3.38</b>	<b>3.71</b>	<b>3.83</b>
$r(0_3^+)$				4.62	4.75	
$r(2_1^+)$				2.50	2.37	2.38
$r(2_2^+)$				4.43	4.02	
$M(E0, 0_1^+ \rightarrow 0_2^+)$	5.4(2)			6.53	6.52	6.45
$B(E2, 2_1^+ \rightarrow 0_1^+)$	7.6(4)			8.69	9.16	
$B(E2, 2_1^+ \rightarrow 0_2^+)$	2.6(4)			3.83	0.84	

very large radii for  $0_2^+$ ,  
 $0_3^+$  and  $2_2^+$  states

experimental situation  
for  $0_3^+$  and  $2_2^+$  states  
still unsettled

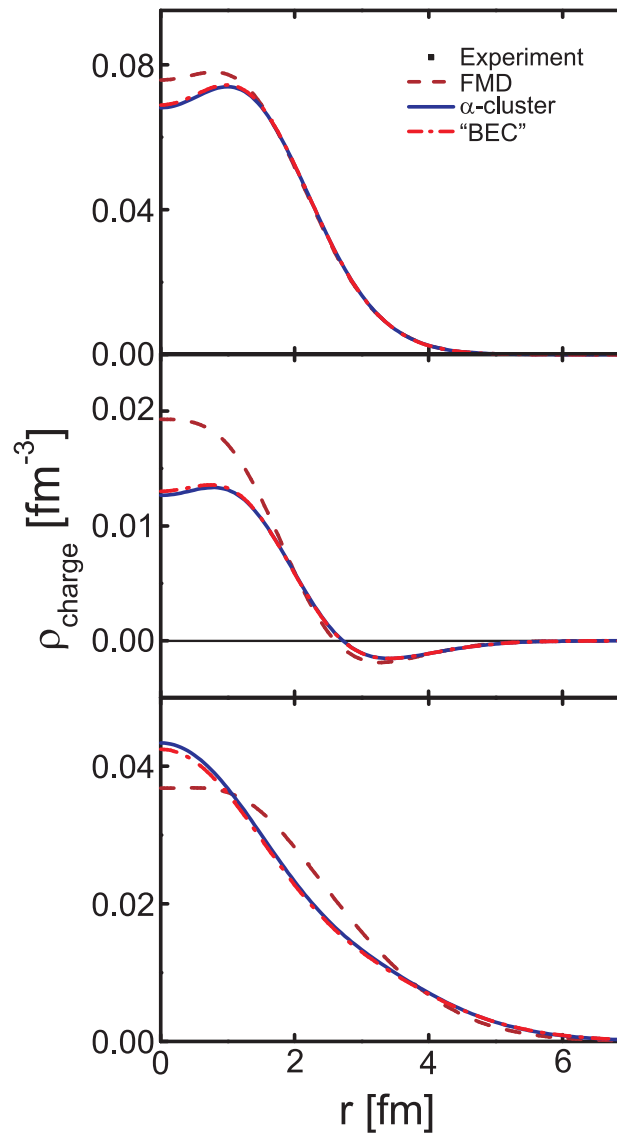
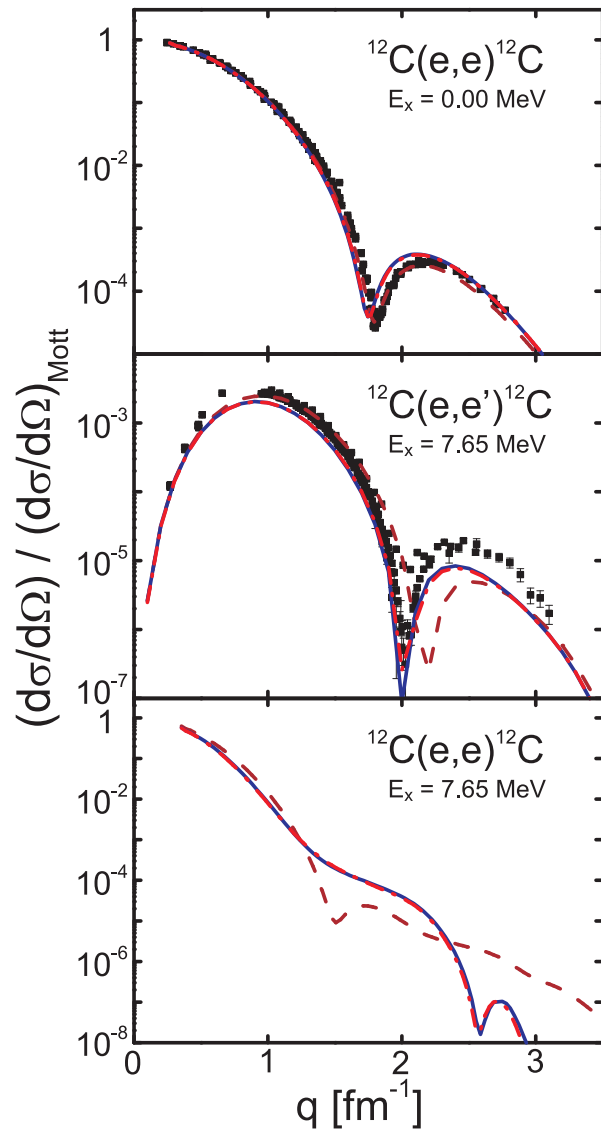
$2_2^+$  resonance at  
1.8 MeV above  
threshold included in  
NACRE compilation

<sup>1</sup> Ajzenberg-Selove, Nuc. Phys. **A506**, 1 (1990)

<sup>2</sup> Itoh et al., Nuc. Phys. **A738**, 268 (2004)

<sup>3</sup> Fynbo et al, Nature **433**, 137 (2005). Diget et al., Nuc. Phys. **A738**, 760 (2005)

# Form factors and Densities



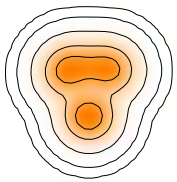
- compare with electron scattering data in Distorted Wave Born Approximation
- ➔ elastic form factor described very well by FMD
- ➔ transition form factor in first maximum better described by FMD, position of minimum and second maximum better described by cluster model

use intrinsic density

$$\rho(\mathbf{x}) = \sum_{k=1}^A \langle \Psi | \delta(\mathbf{x}_k - \mathbf{X} - \mathbf{x}) | \Psi \rangle$$

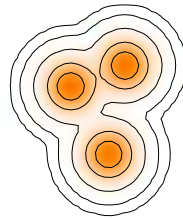
# Important Configurations

- Calculate the overlap with FMD basis states to find the most important contributions to the Hoyle state



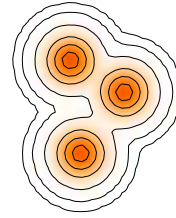
$$\langle \cdot | 0_1^+ \rangle = 0.94$$

$$\langle \cdot | 0_2^+ \rangle = 0.04$$



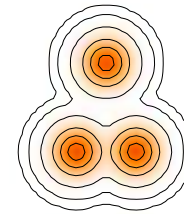
$$\langle \cdot | 0_1^+ \rangle = 0.30$$

$$\langle \cdot | 0_2^+ \rangle = 0.72$$



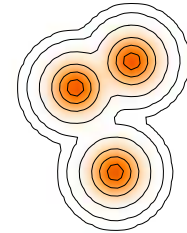
$$\langle \cdot | 0_1^+ \rangle = 0.25$$

$$\langle \cdot | 0_2^+ \rangle = 0.71$$



$$\langle \cdot | 0_1^+ \rangle = 0.15$$

$$\langle \cdot | 0_2^+ \rangle = 0.61$$



$$\langle \cdot | 0_1^+ \rangle = 0.08$$

$$\langle \cdot | 0_2^+ \rangle = 0.61$$

shell model like  
structure for the  
ground state

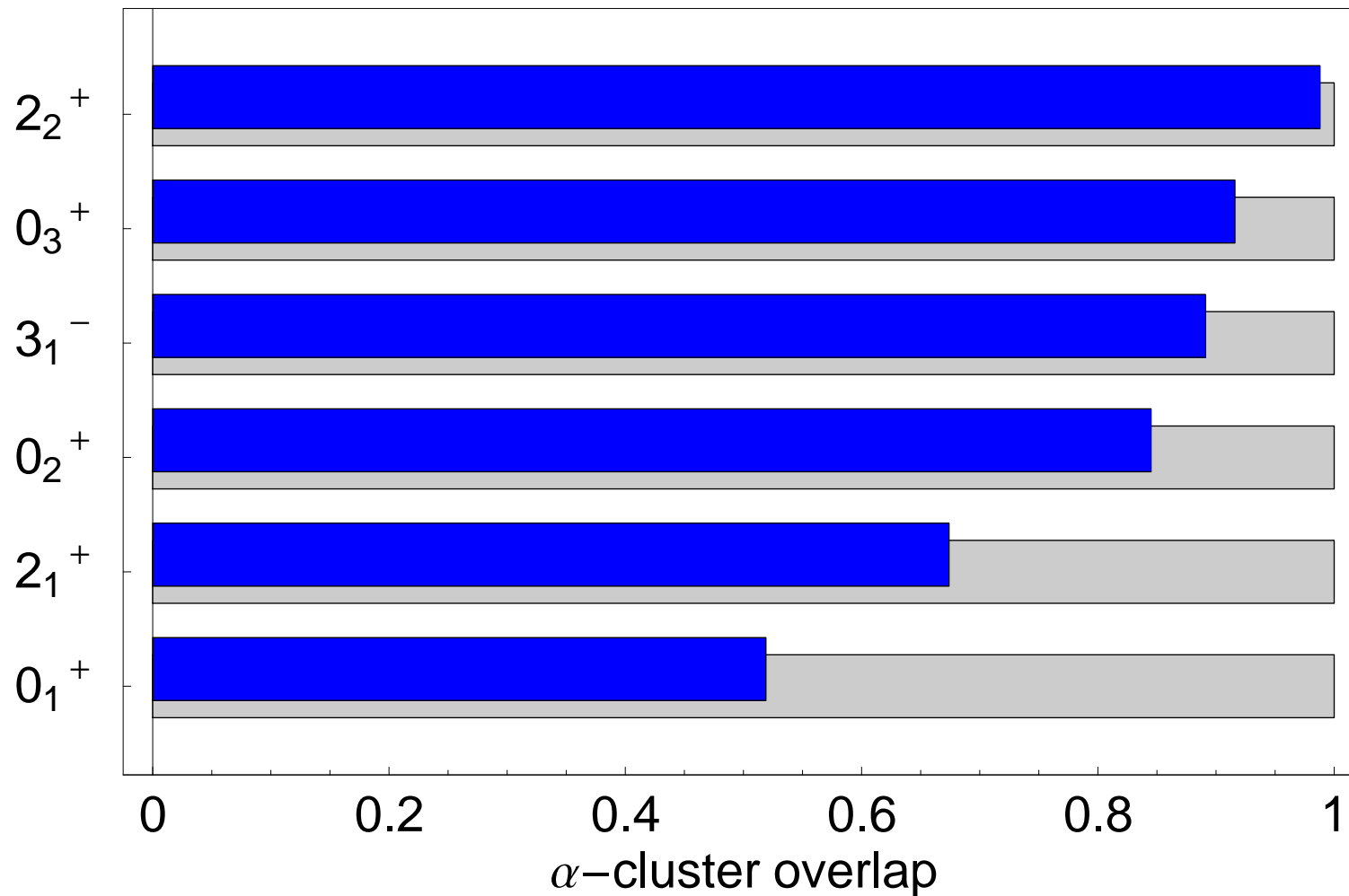
Hoyle state  
dominated by  
 $\alpha$ -cluster  
configurations

FMD basis states  
are not orthogonal!

# Overlap with Cluster Model Space

Calculate the overlap of FMD wave functions with pure  $\alpha$ -cluster model space

$$N_\alpha = \langle \Psi | P_{3\alpha} | \Psi \rangle$$

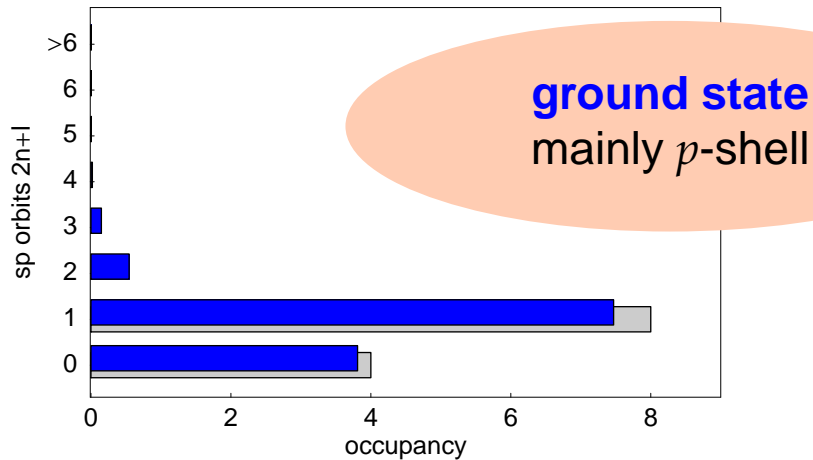


# Harmonic Oscillator Occupation Numbers

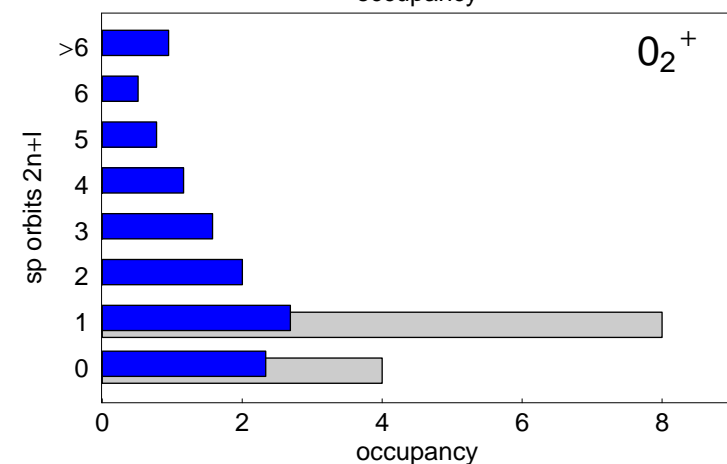
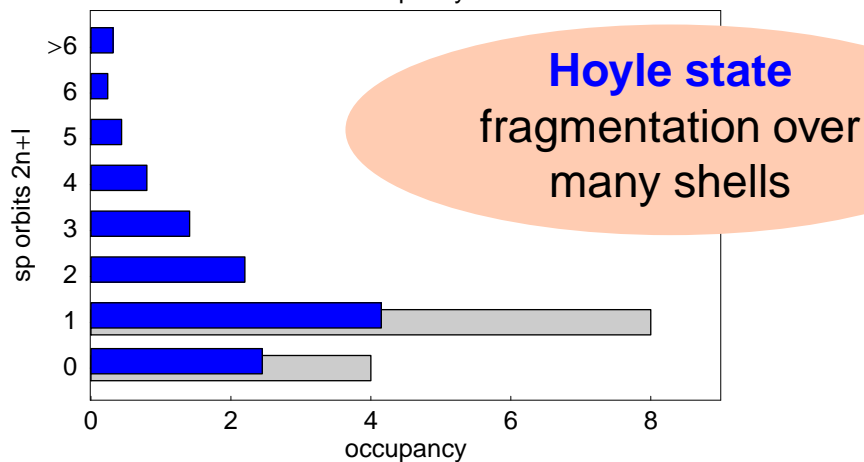
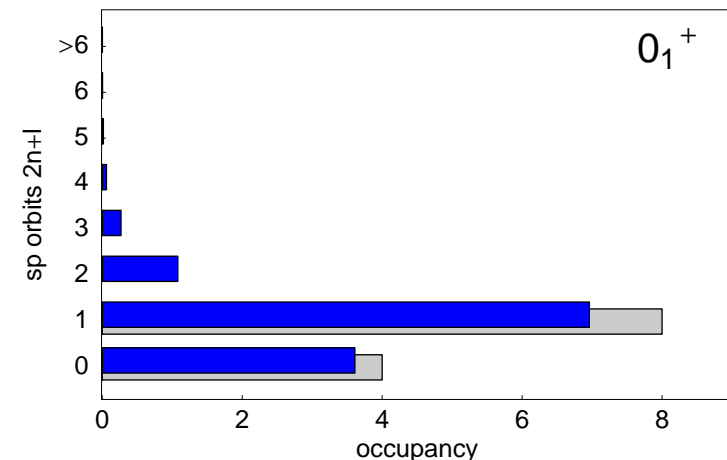
- calculate one-body density in harmonic oscillator basis

$$n_{nlj} = \sum_m \langle \Psi | a_{nljm}^\dagger a_{nljm} | \Psi \rangle$$

FMD



$\alpha$ -cluster

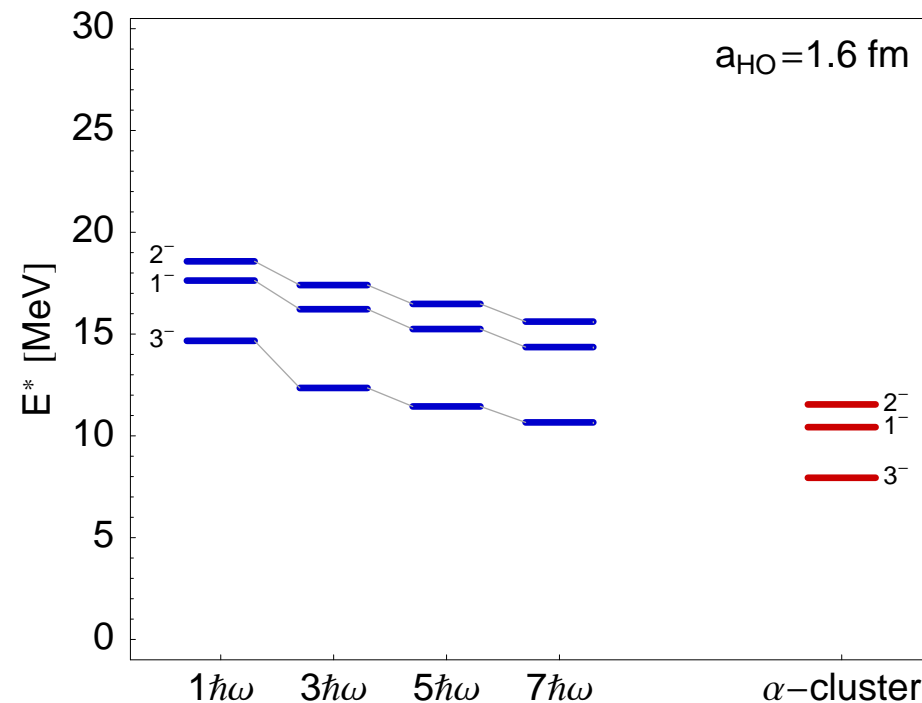
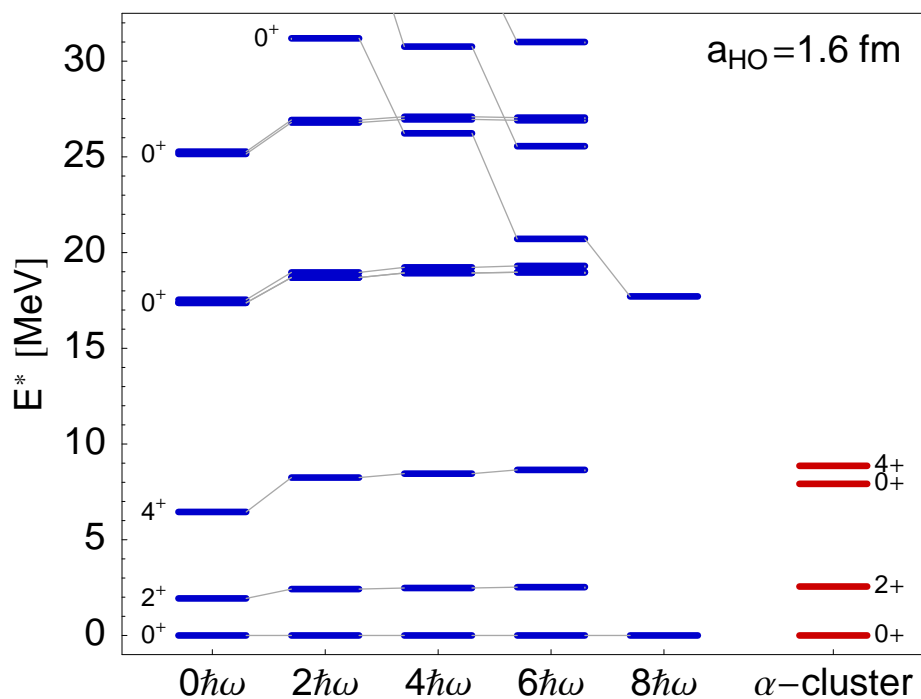


# $\alpha$ -cluster states in the No-Core Shell Model ?

- compare spectra in NCSM and  $\alpha$ -cluster model using the Volkov interaction
- bare interaction used in NCSM calculations

good agreement for  
ground state band  
( $0_1^+$ ,  $2_1^+$ ,  $4_1^+$ )

very slow convergence  
for cluster states



# Summary



- $\alpha$ -cluster model and “BEC” ansatz can describe  $^{12}\text{C}$  ground state and Hoyle state energies and form factors.
- Hoyle state is a dilute system of  $\alpha$ -particles
- results confirmed by FMD calculations that make no explicit assumption about the  $\alpha$ -cluster structure
- Harmonic Oscillator basis has a hard time describing the Hoyle state