

# Dileptons in HSD 2.5

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February 23, 2005

HSD home-page:

<http://www.th.physik.uni-frankfurt.de/~brat/hsd.html>

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# 1 Dilepton channels

The dilepton ( $e^+e^-$  or  $\mu^+\mu^-$ ) spectra in HSD are calculated perturbatively with the time integration method. For the details of the dilepton implementation see our review [1] and also Refs. [2, 3, 4, 5, 6, 7, 8, 9]. The time integration is performed over the actual dilepton emission rate during the full reaction time (contrary to the 'spontaneous decay' assumption which counts the dilepton radiation only at freeze-out).<sup>1</sup>

Table 1: Dilepton channels in HSD 2.5

i	Dilepton channel
1	Dalitz decay of $\pi^0$ : $\pi^0 \rightarrow \gamma e^+ e^-$
2	Dalitz decay of $\eta$ : $\eta \rightarrow \gamma e^+ e^-$ (or $\mu^+ \mu^-$ , also for channels below)
3	Dalitz decay of $\omega$ : $\omega \rightarrow \pi^0 e^+ e^-$
4	Dalitz decay of $\Delta$ : $\Delta \rightarrow N e^+ e^-$
5	direct decay of $\omega$ : $\omega \rightarrow e^+ e^-$
6	direct decay of $\rho$ : $\rho \rightarrow e^+ e^-$
7	direct decay of $\phi$ : $\phi \rightarrow e^+ e^-$
8	direct decay of $J/\Psi$ : $J/\Psi \rightarrow e^+ e^-$
9	direct decay of $\Psi'$ : $\Psi' \rightarrow e^+ e^-$
10	Dalitz decay of $\eta'$ : $\eta' \rightarrow \gamma e^+ e^-$
11	$pn$ bremsstrahlung: $pn \rightarrow p n e^+ e^-$
12	$\pi^\pm N$ bremsstrahlung: $\pi^\pm N \rightarrow \pi N e^+ e^-$ , where $N = p$ or $n$

All branching ratios, electromagnetic partial and total decay widths are taken from the PDG [10].

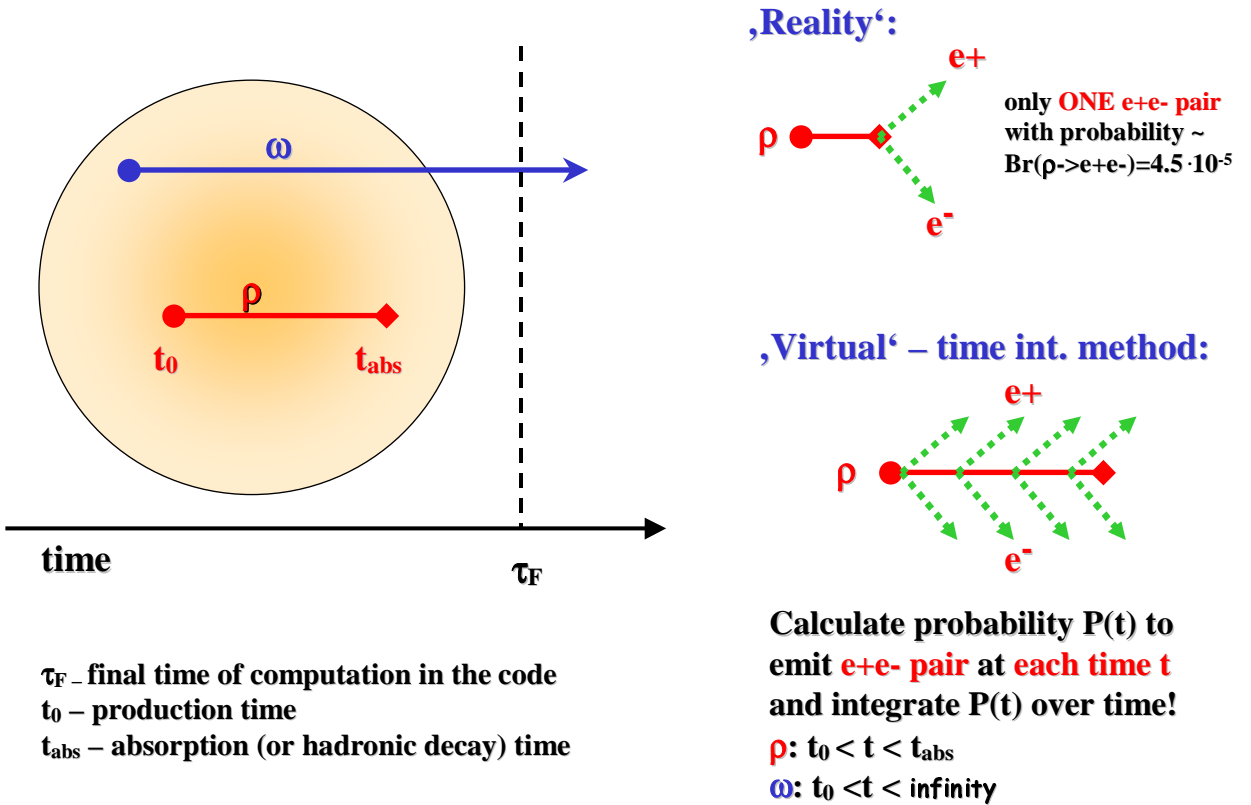
1. The  $pn$  and  $\pi^\pm N$  bremsstrahlungs are calculated in the soft-photon approximation (SPA). Only elastic  $pn$  and  $\pi^\pm N$  collisions are accounted in the bremsstrahlung (i.e.  $pn \rightarrow p n e^+ e^-$ ,  $\pi^\pm N \rightarrow \pi N e^+ e^-$ ). We stress that the SPA approximation might be considered as an upper limit for the bremsstrahlung contribution (especially for  $\pi N$  !). The bremsstrahlung channels are switched off for  $E_{lab} \geq 6$  GeV since it is very questionable to use the SPA at high energies.

2. The channel  $\rho \rightarrow e^+ e^-$  includes the dilepton radiation by all rho mesons produced in baryon-baryon, meson-baryon or meson-meson (e.g.  $\pi^+ \pi^-$  annihilation) collisions. The same holds for the other mesons –  $\rho, \eta, \omega, \phi, J/\Psi, \Psi'$ .

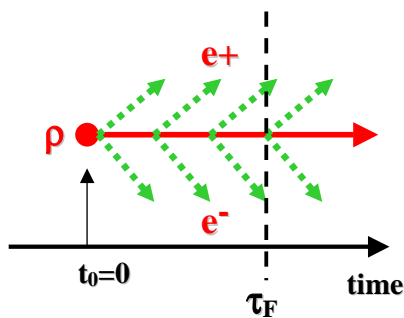
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<sup>1</sup>Useful link: HSD-home page - talk at HADES Collaboration Meeting (Feb. 2005).

# Time integration method for dileptons



## The time integration method for dileptons in HSD



### Dilepton emission rate:

$$\frac{dN(\rho \rightarrow e^+e^-)}{dt}(t) = -\frac{1}{\gamma\tau} N_\rho \text{Br}(\rho \rightarrow e^+e^-), \quad \tau = \frac{\hbar c}{\Gamma_{\text{tot}}}$$

$$N(\rho \rightarrow e^+e^-)(t) = N_0 e^{-\frac{t}{\gamma\tau}}$$

$$\int_0^\infty dt N(t) = \int_0^{\tau_F} dt N(t) + \int_{\tau_F}^\infty dt N(t)$$

numerically :  $dt \rightarrow \Delta t$ ,  $\frac{dN}{dt} \rightarrow \frac{1}{\Delta t} [N(t+\Delta t) - N(t)]$

### Dilepton invariant mass spectra:

$$\frac{dN^{\rho \rightarrow e^+e^-}}{dM} = \sum_{\text{time}}^{\tau_F} \Gamma^{\rho \rightarrow e^+e^-}(M) \cdot \frac{\Delta t}{\gamma(\hbar c)} \cdot \frac{1}{\Delta M} + \text{Br}^{\rho \rightarrow e^+e^-}(M) \cdot e^{-\frac{\Gamma_{\text{tot}}\tau_F}{\gamma(\hbar c)}} \cdot \frac{1}{\Delta M}$$

$$0 < t < \tau_F$$

$$\tau_F < t < \text{infinity}$$

❖ **The time integration method allows to account for the in-medium dynamics of vector mesons!**

## 2 Dalitz decays $A \rightarrow B l^+l^-$

Here  $l^+l^-$  are electron  $e^+e^-$  or muon  $\mu^+\mu^-$  pairs, i.e.  $m_l = m_e = 0.511 \cdot 10^{-3}$  GeV or  $m_l = m_\mu = 0.105658389$  GeV.

### 2.1 Dalitz decay $\pi^0 \rightarrow \gamma l^+l^-$

$$\frac{d\Gamma^{\pi^0 \rightarrow \gamma l^+l^-}}{dM} = \frac{4\alpha}{3\pi} \frac{\Gamma^{\pi^0 \rightarrow \gamma\gamma}}{M} \left(1 - \frac{4m_l^2}{M^2}\right)^{1/2} \left(1 + \frac{2m_l^2}{M^2}\right) \left(1 - \frac{M^2}{m_\pi^2}\right)^3 |F^{\pi^0 \rightarrow \gamma\gamma}(M)|^2, \quad (1)$$

where

$$\begin{aligned} F^{\pi^0 \rightarrow \gamma\gamma}(M) &= 1 + B_{\pi^0} M^2, & B_{\pi^0} &= 5.5 \text{ GeV}^{-2} \\ \Gamma^{\pi^0 \rightarrow \gamma\gamma} &= 7.8 \cdot 10^{-9} \text{ GeV} \\ \Gamma_{\pi^0}^{tot} &\simeq \Gamma^{\pi^0 \rightarrow \gamma\gamma} \\ Br^{\pi^0 \rightarrow \gamma\gamma} &= 0.988 \end{aligned} \quad (2)$$

### 2.2 Dalitz decay $\eta \rightarrow \gamma l^+l^-$

$$\frac{d\Gamma^{\eta \rightarrow \gamma l^+l^-}}{dM} = \frac{4\alpha}{3\pi} \frac{\Gamma^{\eta \rightarrow \gamma\gamma}}{M} \left(1 - \frac{4m_l^2}{M^2}\right)^{1/2} \left(1 + \frac{2m_l^2}{M^2}\right) \left(1 - \frac{M^2}{m_\eta^2}\right)^3 |F^{\eta \rightarrow \gamma\gamma}(M)|^2, \quad (3)$$

where

$$\begin{aligned} F^{\eta \rightarrow \gamma\gamma}(M) &= \left(1 - \frac{M^2}{\Lambda_\eta^2}\right)^{-1} & \Lambda_\eta &= 0.72 \text{ GeV} \\ \Gamma^{\eta \rightarrow \gamma\gamma} &= 4.6 \cdot 10^{-7} \text{ GeV} \\ \Gamma_\eta^{tot} &= 1.18 \cdot 10^{-6} \text{ GeV} \\ Br^{\eta \rightarrow \gamma\gamma} &= 0.3933 \end{aligned} \quad (4)$$

### 2.3 Dalitz decay $\omega \rightarrow \gamma l^+l^-$

$$\begin{aligned} \frac{d\Gamma^{\omega \rightarrow \pi^0 l^+l^-}}{dM} &= \frac{2\alpha}{3\pi} \frac{\Gamma^{\omega \rightarrow \pi^0\gamma}}{M} \left(1 - \frac{4m_l^2}{M^2}\right)^{1/2} \left(1 + \frac{2m_l^2}{M^2}\right) \\ &\times \left[ \left(1 + \frac{M^2}{(m_\omega^2 - m_\pi^2)}\right)^2 - \frac{4m_\omega^2 M^2}{(m_\omega^2 - m_\pi^2)^2} \right]^{3/2} |F^{\omega \rightarrow \pi^0 l^+l^-}(M)|^2, \end{aligned} \quad (5)$$

where

$$\begin{aligned} |F^{\omega \rightarrow \pi^0 l^+l^-}(M)|^2 &= \frac{\Lambda_\omega^4}{(\Lambda_\omega^2 - M^2)^2 + \Lambda_\omega^2 \Gamma_\omega^2} & \Lambda_\omega &= 0.65 \text{ GeV}, & \Gamma_\omega &= 0.075 \text{ GeV} \\ \Gamma^{\omega \rightarrow \pi^0\gamma} &= 7.17 \cdot 10^{-4} \text{ GeV} \\ \Gamma_\omega^{tot} &= 8.44 \cdot 10^{-3} \text{ GeV} \\ Br^{\omega \rightarrow \pi^0\gamma} &= 0.085 \end{aligned} \quad (6)$$

## 2.4 Dalitz decay $\eta' \rightarrow \gamma l^+ l^-$

$$\frac{d\Gamma^{\eta' \rightarrow \gamma l^+ l^-}}{dM} = \frac{4\alpha}{3\pi} \frac{\Gamma^{\eta' \rightarrow \gamma\gamma}}{M} \left(1 - \frac{4m_l^2}{M^2}\right)^{1/2} \left(1 + \frac{2m_l^2}{M^2}\right) \left(1 - \frac{M^2}{m_{\eta'}^2}\right)^3 |F^{\eta' \rightarrow \gamma\gamma}(M)|^2, \quad (7)$$

where

$$\begin{aligned} |F^{\eta' \rightarrow \gamma\gamma}(M)|^2 &= \frac{\Lambda_{\eta'}^4}{(\Lambda_{\eta'}^2 - M^2)^2 + \Lambda_{\eta'}^2 \Gamma_{\eta'}^2} & \Lambda_{\eta'} &= 0.75 \text{ GeV} & \Gamma_{\eta'} &= 0.14 \text{ GeV} \\ \Gamma^{\eta' \rightarrow \gamma\gamma} &= 4.28 \cdot 10^{-6} \text{ GeV} \\ \Gamma_{\eta'}^{\text{tot}} &= 0.202 \cdot 10^{-3} \text{ GeV} \\ Br^{\eta' \rightarrow \gamma\gamma} &= 0.0212 \\ m_{\eta'} &= 0.95778 \text{ GeV} \end{aligned} \quad (8)$$

## 2.5 Dalitz decay $\Delta \rightarrow N l^+ l^-$

$$\frac{d\Gamma^{\Delta \rightarrow N l^+ l^-}}{dM} = \frac{2\alpha}{3\pi} \frac{\Gamma_0(M, M_\Delta)}{M}, \quad (9)$$

$$\begin{aligned} \Gamma_0(M, M_\Delta) &= \frac{\lambda^{1/2}(M^2, m_N^2, M_\Delta^2)}{16\pi M_\Delta^2} \cdot m_N \cdot [2m_T(M, M_\Delta) + m_L(M, M_\Delta)] \\ m_L(M, M_\Delta) &= (efg)^2 \frac{M_\Delta^2}{9m_N} M^2 \cdot 4(M_\Delta - m_N - q_0), \quad e^2 = 4\pi\alpha, \quad g = 5.44 \\ m_T(M, M_\Delta) &= (efg)^2 \frac{M_\Delta^2}{9m_N} [q_0^2(5M_\Delta - 3(q_0 + m_N)) - M^2(M_\Delta + m_N + q_0)] \\ f &= -1.5 \frac{M_\Delta + m_N}{m_N((m_N + M_\Delta)^2 - M^2)} \\ q_0 &= (M^2 + p_f^2)^{1/2} \\ p_f^2 &= \frac{(M_\Delta^2 - (m_N + M)^2)(M_\Delta^2 - (m_N - M)^2)}{4M_\Delta^2} \\ \lambda(M^2, m_N^2, M_\Delta^2) &= M^4 + m_N^4 + M_\Delta^4 - 2(M^2 m_N^2 + M^2 M_\Delta^2 + m_N^2 M_\Delta^2). \end{aligned} \quad (10)$$

Here  $M_\Delta$  is the current mass of the  $\Delta$ -resonance - calculated in HSD according to the spectral function with the total width from Ref. [12] (cf. also Ref. [11]) :

$$\Gamma_\Delta^{\text{tot}}(M_\Delta) = \Gamma_R \frac{M_{\Delta 0}}{M_\Delta} \cdot \left(\frac{q}{q_r}\right)^3 \cdot F^2(q) \quad (11)$$

$$F(q) = \frac{\beta_r^2 + q_r^2}{\beta^2 + q^2}$$

$$q^2 = \frac{(M_\Delta^2 - (m_N + m_\pi)^2)(M_\Delta^2 - (m_N - m_\pi)^2)}{4M_\Delta^2}$$

$$q_r^2 = 0.051936, \quad \beta_r^2 = 0.09$$

$$\Gamma_R = 0.11 \text{ GeV}, \quad M_{\Delta 0} = 1.232 \text{ GeV} \quad (12)$$

### 3 Direct decay of vector mesons $V \rightarrow l^+l^-$

The dilepton decay width of vector meson  $V$  with the mass  $M$  (calculated in HSD according to the spectral function) is

$$\Gamma^{V \rightarrow l^+l^-}(M) = C_V \frac{m_V^4}{M^3}, \quad (13)$$

where  $C_V = \frac{\Gamma^{V \rightarrow l^+l^-}(m_V)}{m_V}$ ,  $m_V$  is the pole mass of the vector meson  $V$ .

For broad resonances such as  $\rho$  meson, the branching ratio to dileptons depends on the mass  $M$ :

$$Br^{V \rightarrow l^+l^-}(M) = \frac{\Gamma_{V \rightarrow l^+l^-}(M)}{\Gamma_{tot}^V(M)}. \quad (14)$$

Here the total width of the  $\rho$  meson is

$$\begin{aligned} \Gamma_{tot}^\rho(M) &\simeq \Gamma_{\rho \rightarrow \pi\pi} = \Gamma_0 \left(\frac{m_V}{M}\right)^2 \left(\frac{q}{q_V}\right)^3 \\ q &= \frac{(M^2 - 4m_\pi^2)^{1/2}}{2}, \quad q_V = \frac{(m_V^2 - 4m_\pi^2)^{1/2}}{2} \end{aligned} \quad (15)$$

For narrow resonances such as  $\omega, \phi, J/\Psi, \Psi'$  a constant total width and branching ratio are used:  $\Gamma_{tot}^V \equiv \Gamma_{tot}(m_V)$ ,  $Br_0^{V \rightarrow l^+l^-} \equiv Br^{V \rightarrow l^+l^-}(m_V)$ .

Table 2: The parameters for dilepton decay of vector mesons used in HSD 2.5

meson V	electron pair: $e^+e^-$	muon pair $\mu^+\mu^-$
$\rho$ -meson	$Br^{\rho \rightarrow e^+e^-}(m_\rho) = 4.49 \cdot 10^{-5}$ $\Gamma^{\rho \rightarrow e^+e^-}(m_\rho) = 6.77 \cdot 10^{-6} \text{ GeV}$ $C^{\rho \rightarrow e^+e^-}(m_\rho) = 8.814 \cdot 10^{-6}$	$Br^{\rho \rightarrow \mu^+\mu^-}(m_\rho) = 4.6 \cdot 10^{-5}$ $\Gamma^{\rho \rightarrow \mu^+\mu^-}(m_\rho) = 6.9 \cdot 10^{-6} \text{ GeV}$ $C^{\rho \rightarrow \mu^+\mu^-}(m_\rho) = 8.96 \cdot 10^{-6}$
$\omega$ -meson	$Br^{\omega \rightarrow e^+e^-}(m_\omega) = 7.07 \cdot 10^{-5}$ $\Gamma^{\omega \rightarrow e^+e^-}(m_\omega) = 0.6 \cdot 10^{-6} \text{ GeV}$ $C^{\omega \rightarrow e^+e^-}(m_\omega) = 0.767 \cdot 10^{-6}$	$Br^{\omega \rightarrow \mu^+\mu^-}(m_\omega) = 8.06 \cdot 10^{-5} \text{ GeV}$ $\Gamma^{\omega \rightarrow \mu^+\mu^-}(m_\omega) = 0.68 \cdot 10^{-6}$ $C^{\omega \rightarrow \mu^+\mu^-}(m_\omega) = 0.863 \cdot 10^{-6}$
$\phi$ -meson	$Br^{\phi \rightarrow e^+e^-}(m_\phi) = 2.91 \cdot 10^{-4}$ $\Gamma^{\phi \rightarrow e^+e^-}(m_\phi) = 1.297 \cdot 10^{-6} \text{ GeV}$ $C^{\phi \rightarrow e^+e^-}(m_\phi) = 1.27 \cdot 10^{-6}$	$Br^{\phi \rightarrow \mu^+\mu^-}(m_\phi) = 3.7 \cdot 10^{-4}$ $\Gamma^{\phi \rightarrow \mu^+\mu^-}(m_\phi) = 1.649 \cdot 10^{-6} \text{ GeV}$ $C^{\phi \rightarrow \mu^+\mu^-}(m_\phi) = 1.618 \cdot 10^{-6}$
$J/\Psi$ -meson	$Br^{J/\Psi \rightarrow e^+e^-}(m_{J/\Psi}) = 5.93 \cdot 10^{-2}$ $\Gamma^{J/\Psi \rightarrow e^+e^-}(m_{J/\Psi}) = 5.26 \cdot 10^{-6} \text{ GeV}$ $C^{J/\Psi \rightarrow e^+e^-}(m_{J/\Psi}) = 1.698 \cdot 10^{-6}$	$Br^{J/\Psi \rightarrow \mu^+\mu^-}(m_{J/\Psi}) = 5.88 \cdot 10^{-2}$ $\Gamma^{J/\Psi \rightarrow \mu^+\mu^-}(m_{J/\Psi}) = 5.12 \cdot 10^{-6} \text{ GeV}$ $C^{J/\Psi \rightarrow \mu^+\mu^-}(m_{J/\Psi}) = 1.652 \cdot 10^{-6}$
$\Psi'$ -meson	$Br^{\Psi' \rightarrow e^+e^-}(m_{\Psi'}) = 8.8 \cdot 10^{-3}$ $\Gamma^{\Psi' \rightarrow e^+e^-}(m_{\Psi'}) = 2.12 \cdot 10^{-6} \text{ GeV}$ $C^{\Psi' \rightarrow e^+e^-}(m_{\Psi'}) = 0.575 \cdot 10^{-6}$	$Br^{\Psi' \rightarrow \mu^+\mu^-}(m_{\Psi'}) = 1.03 \cdot 10^{-2}$ $\Gamma^{\Psi' \rightarrow \mu^+\mu^-}(m_{\Psi'}) = 2.853 \cdot 10^{-6} \text{ GeV}$ $C^{\Psi' \rightarrow \mu^+\mu^-}(m_{\Psi'}) = 0.774 \cdot 10^{-6}$

## References

- [1] W. Cassing and E. L. Bratkovskaya,  
*'Hadronic and electromagnetic probes of hot and dense nuclear matter'*,  
Phys. Reports **308** (1999) 65.
- [2] E. L. Bratkovskaya,  
*' $\rho/\omega$  properties from dilepton spectra in  $pA$  reactions at 12 GeV'*,  
**Phys. Lett. B** **529** (2002) 26-35 [nucl-th/0108055].
- [3] E. L. Bratkovskaya,  
*' $e^+e^-$  production in  $pA$  reactions at SIS energies'*,  
**Nucl. Phys. A** **696** (2001) 761-787 [nucl-th/0101067].
- [4] E. L. Bratkovskaya, W. Cassing, and U. Mosel,  
*'Perspectives of  $e^+e^-$  production in  $pp$ ,  $pd$  and  $pBe$  reactions at SIS energies'*,  
**Nucl. Phys. A** **686** (2001) 568-588 [nucl-th/0008037].
- [5] E. L. Bratkovskaya, W. Cassing, M. Effenberger, and U. Mosel,  
*' $e^+e^-$  production from  $pp$  reactions at BEVALAC energies'*,  
**Nucl. Phys. A** **653** (1999) 301-317 [nucl-th/9903009].
- [6] E. L. Bratkovskaya and C. M. Ko,  
*'Low-mass dileptons and dropping rho meson mass'*,  
**Phys. Lett. B** **445** (1999) 265-270 [nucl-th/9809056].
- [7] E. L. Bratkovskaya, W. Cassing, R. Rapp, and J. Wambach,  
*'Dilepton production and  $m_T$ -scaling at BEVALAC/SIS energies'*,  
**Nucl. Phys. A** **634** (1998) 168-189 [nucl-th/9710043].
- [8] W. Cassing, E. L. Bratkovskaya, R. Rapp, and J. Wambach,  
*'Probing the  $\rho$  spectral function in hot and dense nuclear matter by dileptons'*,  
**Phys. Rev. C** **57** (1998) 916-921 [nucl-th/9708020].
- [9] E. L. Bratkovskaya and W. Cassing,  
*'Dilepton Production from AGS to SPS Energies within a Relativistic Transport Approach'*,  
**Nucl. Phys. A** **619** (1997) 413 [nucl-th/9611042].
- [10] K. Hagiwara *et al.*, (Review of Particle Properties), **Phys. Rev. D** **66**, 010001 (2002).
- [11] Gy. Wolf *et al.*, **Nucl. Phys. A** **517** (1990) 615.
- [12] J.H. Koch, E.J. Monitz and N. Ohtsuka, **Ann. Phys.** **154** (1984) 99.